

# Cosmology in the Unknown

## Particle Physics Inspired Ideas

### Limits of Experimental Physics

- We have very limited experience about particle physics above the electroweak scale
- For example, the most powerful collider in the world is the Large Hadron Collider
  - Protons hitting protons at 7 TeV + 7 TeV
- Protons are not fundamental
  - They contain a combination of three quarks plus some gluons
  - Optimistically, the quarks might carry  $\sim 2$  TeV of energy each
- As they collide, they can produce particles of energy up to  $\sim 4$  TeV
  - This corresponds to temperatures of  $k_B T$  up to  $\sim 1.3$  TeV
- Cosmic rays hit the Earth at much higher energy
- But these are rare and difficult to study directly
- Anything above temperatures of about 1.5 TeV have not been tested by experiment

# Sources of Inspiration

- No one thinks the standard model is the final answer on particle physics
- Many promising ideas have developed because of developments in particle physics
- Sometimes, these developments *should* have cosmological consequences
  - In some cases, we can learn about particle physics by studying these consequences
- The early universe can reach much higher energies than we can ever achieve in the laboratory
- There are also some unsolved problems in cosmology
- Sometimes, we can use the ideas *already* developed in particle physics to see if they can solve our cosmology problem
  - Sometimes, this seems very forced
- Other times, new ideas in particle physics are thought up in hopes of solving cosmological problems

# Supersymmetry

- In conventional particle physics, fermions and bosons are fundamentally different
  - And never the twain shall meet
- In a hypothesis called *supersymmetry*, fermions and bosons are interrelated

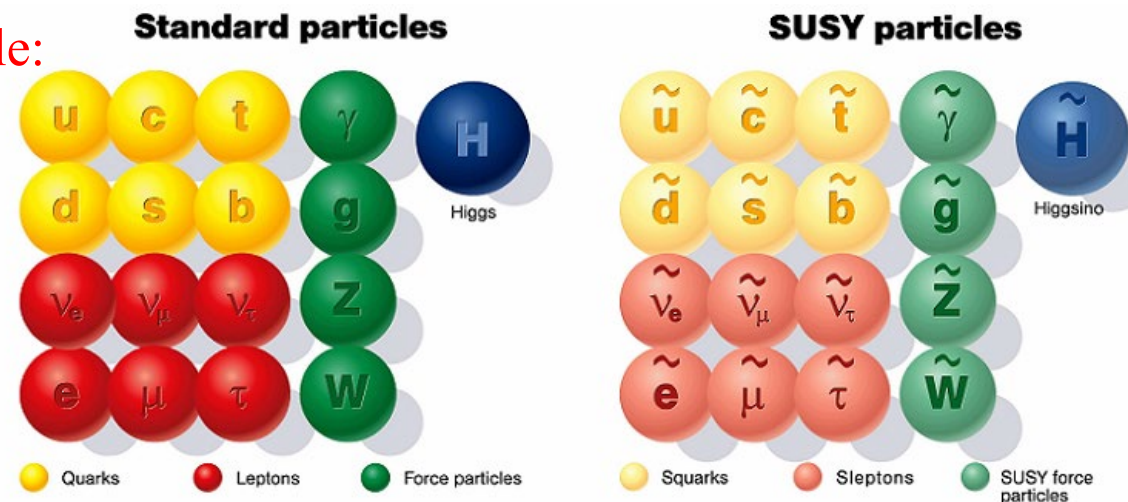
- There must be a *superpartner* for every particle:

- Supersymmetry also helps solve a problem called the *hierarchy problem*
  - But only if it doesn't happen at too high an energy
- If supersymmetry is right, then scale of supersymmetry breaking probably around

~~$k_B T = 500 \text{ GeV}$  or so.~~

$k_B T = 1.5 \text{ TeV}$  or so.

- If this is right, the LHC should discover it
- If supersymmetry is right, then we must revise upwards our estimate of where it happens



**Complete list of  
superpartners discovered at  
the Large Hadron Collider:**

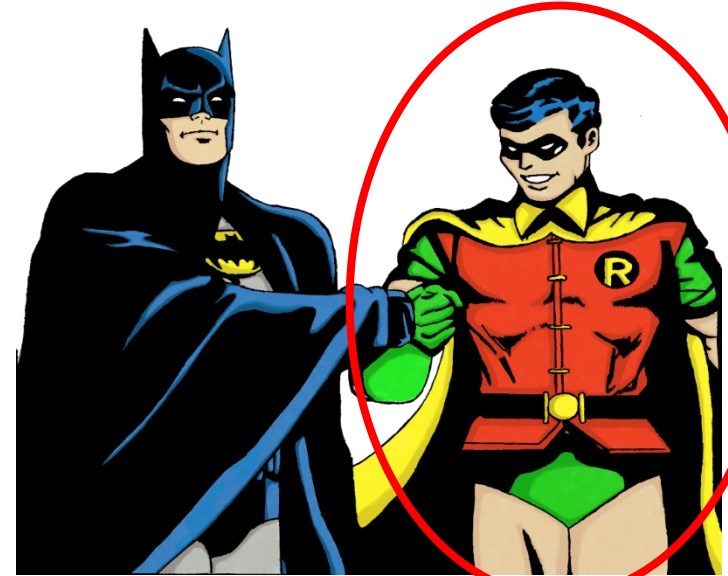
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# The Lightest Superpartner

- When would supersymmetry become relevant?
  - Assume supersymmetry breaking scale around 1.5 TeV

$$t = \frac{2.42 \text{ s}}{\sqrt{g_{\text{eff}}}} \left( \frac{\text{MeV}}{k_B T} \right)^2 = \frac{2.42 \text{ s}}{\sqrt{106.75}} \left( \frac{\text{MeV}}{1.5 \times 10^6 \text{ MeV}} \right)^2 \approx 10^{-13} \text{ s}$$

<u>Time</u>	<u><math>T</math> or <math>k_B T</math></u>	<u>Events</u>
$10^{-13} \text{ s}$	1.5 TeV	Supersymmetry Breaking
$10^{-11} \text{ s}$	160 GeV	Electroweak Symmetry Breaking



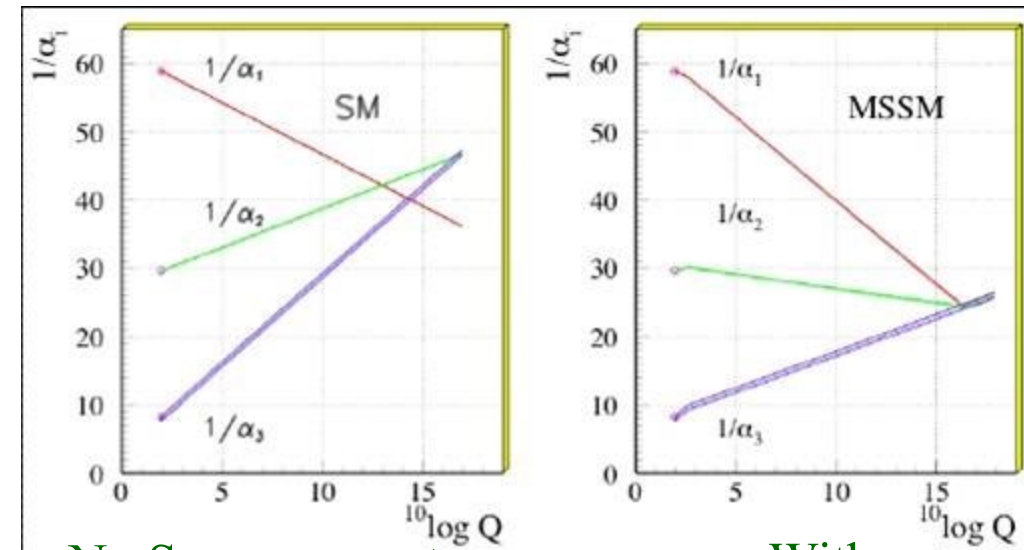
$$e^- + e^+ \leftrightarrow \tilde{\gamma} + \tilde{\gamma}$$

- Most variations of supersymmetry assume that supersymmetric particles can only be created or destroyed in pairs
  - Guaranteed by something called *R*-parity
- If this is true, then the lightest superpartner (LSP) must automatically be stable
- Such a particle could be the dark matter
  - More on this later

# Grand Unification Theories (GUT's)

- In the standard model, there are three fundamental forces, and three corresponding coupling constants
- These have rather different values
- But their strength changes as you change the energy of the experiment, theoretically
- How much they change depends on whether supersymmetry is right or not
- If supersymmetry is right, then at an energy of about  $10^{16}$  GeV, the three forces are equal in strength
- At  $k_B T = 10^{16}$  GeV, there will be another phase transition – the Grand Unification transition

$$t = \frac{2.42 \text{ s}}{\sqrt{210}} \left( \frac{\text{MeV}}{10^{19} \text{ MeV}} \right)^2 \approx 10^{-39} \text{ s}$$



No Supersymmetry

With  
Supersymmetry

Baryogenesis might occur at this scale

Scale could be right for inflation

# Inflation

## How Flat Was the Universe?

- The universe today is very nearly flat:
- How flat was it in the past?
- Recall the Friedman equation:

$$\Omega = 0.9993(37)$$

$$H^2 = \frac{8\pi}{3} G \rho - \frac{k}{a^2}$$

- The radiation fraction  $\Omega_r$  is defined as
- The curvature (the last term) is given by

$$H^2 \Omega_r = \frac{8\pi}{3} G \rho_r$$

$$\frac{k}{a^2} = \frac{8\pi}{3} G \rho - H^2 = H^2 (\Omega - 1)$$

- Taking the ratio of these equations, we have
- Recall that radiation scales as  $a^{-4}$
- So it all scales approximately, as  $T^{-2}$ 
  - Not exactly, due to reheating

$$\frac{\Omega - 1}{\Omega_r} = \frac{3k}{8\pi G \rho_r a^2} \propto a^2 \propto T^{-2}$$

# The Flatness Problem

- When you include all the corrections for reheatings, you get  $\frac{\Omega-1}{\Omega_r} \propto \frac{1}{T^2}$

$$\frac{\Omega - 1}{\Omega_r} = \left( \frac{3.91}{g_{\text{eff}}} \right)^{2/3} \left( \frac{k_B T_0}{k_B T} \right)^2 \left( \frac{\Omega - 1}{\Omega_r} \right)_0$$

- For definiteness, pick  $g_{\text{eff}}=200$  and  $k_B T = 3 \times 10^{15}$  GeV
- Current values
- In the early universe expect  $\Omega_r = 1$

$$\Omega_0 = 0.9993(37)$$

$$\Omega_{r0} = 9 \times 10^{-5}$$

$$T_0 = 2.7255 \text{ K}$$

- So we find

[illegible]

- You should memorize this value to about 55 digits or so
- How did it happen to be so close to 1?

# The Horizon Problem

- The universe looks pretty uniform today
- The universe used to be smaller
  - Could it have been smoothed out in the past?
- Let's look at the universe at the time of Grand Unified scale ( $k_B T = 10^{16}$  GeV):
- At the time the universe was about  $10^{-39}$  seconds old
- Universe then *could be* smooth on scales of about

$$d = 2ct = 2(3 \times 10^8 \text{ m/s})(10^{-39} \text{ s}) = 6 \times 10^{-31} \text{ m}$$

- Since then, the universe has grown *immensely* in size
- Assume  $aT$  is roughly constant (not exactly)

$$\frac{a_0}{a} = \frac{k_B T}{k_B T_0} \left( \frac{g_{\text{eff}}}{3.91} \right)^{1/3} = \frac{10^{25} \text{ eV}}{(8.617 \times 10^{-5} \text{ eV/K})(2.725 \text{ K})} \left( \frac{200}{3.91} \right)^{1/3} \approx 1.6 \times 10^{29}$$

$$d_0 = \frac{a_0 d}{a} \approx 0.1 \text{ m}$$

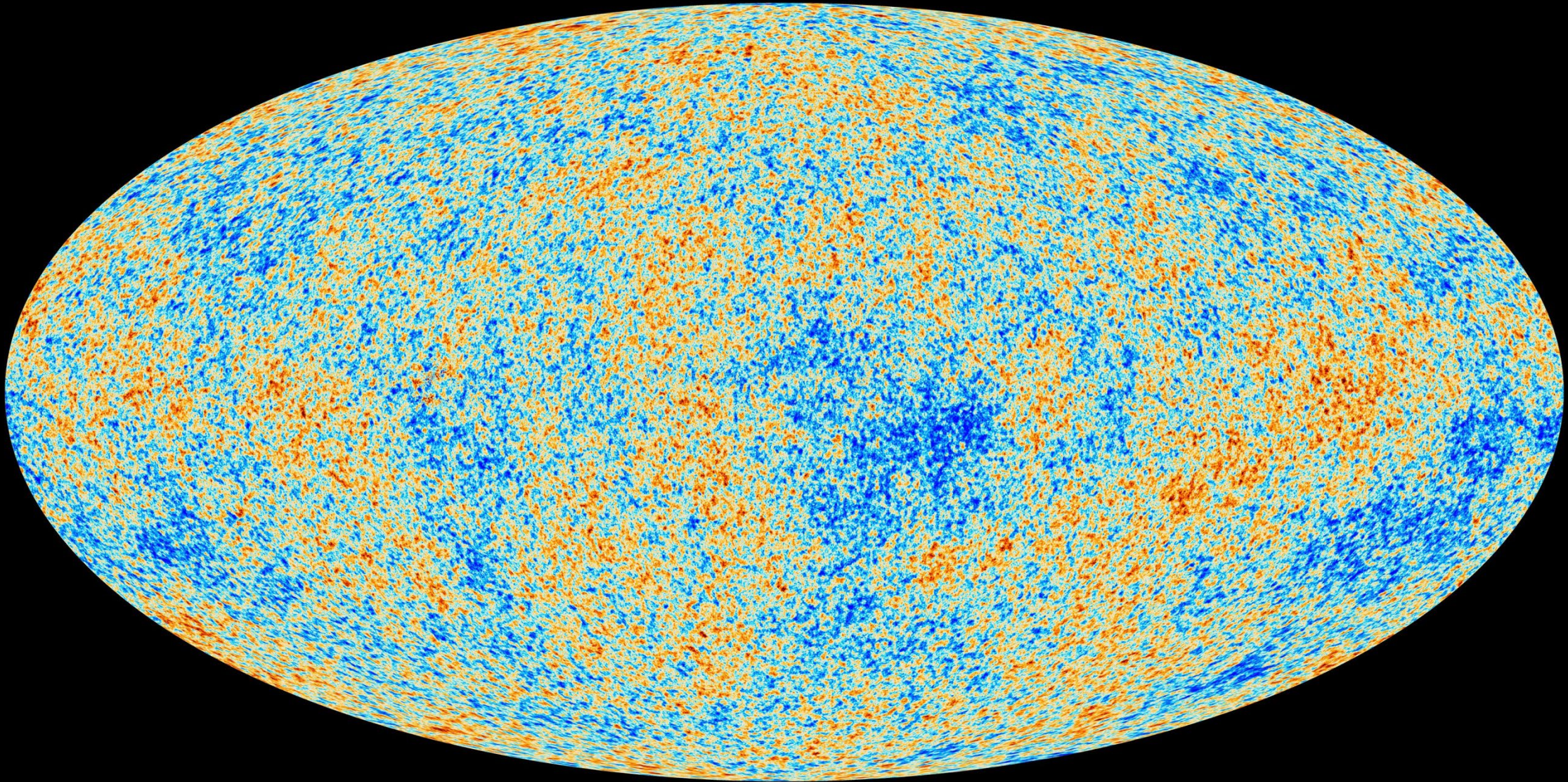
- So universe today could be smooth on scale of:
- Actual universe today smooth on scale of:

$$d \approx 3.3ct_0 \approx 4 \times 10^{26} \text{ m}$$



# The Origin of Structure Problem

- Where did all this variation come from?





# The Solution: Inflation

- How can we get  $\Omega$  to be close to one at some point in the past?

Recall,

$$\frac{\Omega - 1}{\Omega} \sim \frac{k/a^2}{\rho}$$

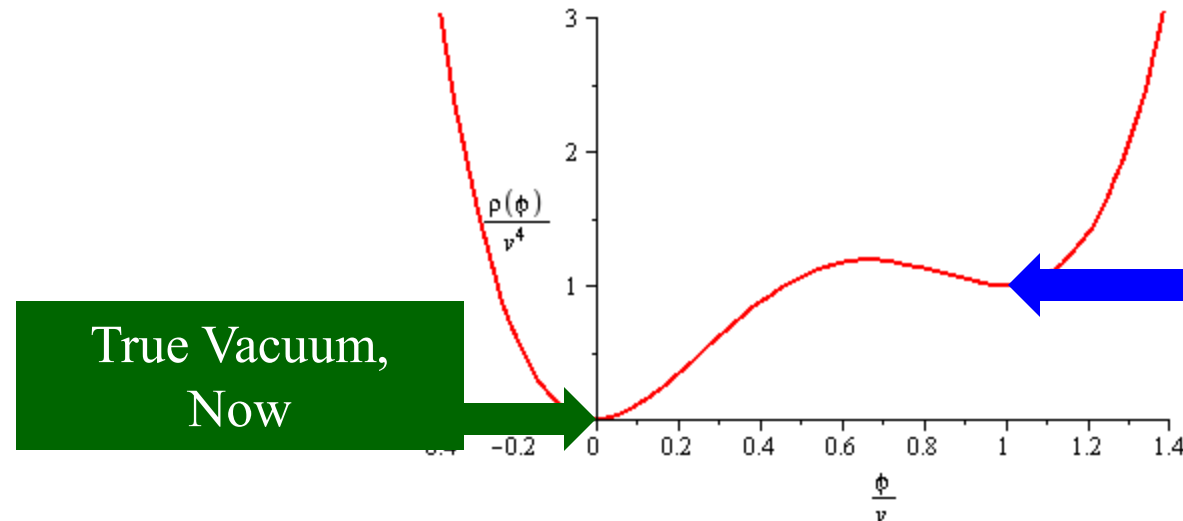
- Since  $\rho$  falls faster than  $1/a^2$ , as universe expands,  $\Omega$  is driven away from 1
- If  $\rho$  fell slower than  $1/a^2$ , as universe expands,  $\Omega$  would be driven *towards* 1
- Could there have been a period when universe dominated by vacuum in the past?
- Need some field that is stuck in “wrong place” – “false vacuum”

- Friedmann Equation:

$$\frac{\dot{a}^2}{a^2} = \frac{8\pi}{3} G \rho_{\Lambda} \equiv H_{\Lambda}^2$$

$$a(t) \propto e^{H_{\Lambda} t}$$

$$\begin{aligned} \rho_m &\propto a^{-3} \\ \rho_r &\propto a^{-4} \\ \rho_{\Lambda} &\propto 1 \end{aligned}$$



# The Inflationary Era

$$H_{\Lambda}^2 = \frac{8}{3} \pi G \rho_{\Lambda}$$

$$a(t) \propto e^{H_{\Lambda} t}$$

- Assume inflation occurs around the Grand unified scale,  $\nu = 10^{16}$  GeV

$$\rho_{\Lambda} \sim \frac{\nu^4}{10(\hbar c)^3 c^2} = \frac{(10^{25} \text{ eV})^4 (1.602 \times 10^{-19} \text{ J/eV})}{10(1.97 \times 10^{-7} \text{ eV} \cdot \text{m})^3 (3 \times 10^8 \text{ m/s})^2} = 2.33 \times 10^{83} \text{ kg/m}^3$$

- During inflation, Hubble's constant is:

$$H_{\Lambda} = \sqrt{\frac{8}{3} \pi G \rho_{\Lambda}} = \sqrt{\frac{8}{3} \pi (6.67 \times 10^{-11} \text{ m}^3/\text{kg/s}^2) (2.33 \times 10^{83} \text{ kg/m}^3)} = 1.1 \times 10^{37} \text{ s}^{-1}$$

- Universe increases by factor of  $e$  every  $10^{-37}$  s or so

- We are not radiation dominated, so we can't trust radiation era formulas

$$t \neq \frac{2.42 \text{ s}}{\sqrt{g_{\text{eff}}}} \left( \frac{\text{MeV}}{k_B T} \right)^2$$

# How Much Inflation Do We Need?

- During inflation,  $\Omega$  is driven towards one
- $\Omega - 1$  at GUT scale is less than  $10^{-55}$
- We need  $a$  to increase by about  $10^{27}$

$$\frac{\Omega - 1}{\Omega} \sim \frac{k/a^2}{\rho} \propto a^{-2}$$

- At GUT scale, universe was smooth on scale that becomes *at most* 0.1 m today
- Today universe is smooth on scale of *at least*  $10^{27}$  m
- We need  $a$  to increase by *at least*  $10^{28}$

$$a(t) \propto e^{H_{\Lambda} t} \quad e^{H_{\Lambda} t_I} > 10^{28} \quad H_{\Lambda} t_I > \ln(10^{28}) = 65$$

- Probably want 100 or so  $e$ -foldings of growth

$$t_I > 100 H_{\Lambda}^{-1} = 10^{-35} \text{ s}$$

- If it is any number significantly bigger than the minimum required, then
  - $\Omega = 1$  to *many* more digits than we need at GUT scale
  - $\Omega = 1$  today to *many* digits
  - Universe is smooth on scales *much* bigger than we can see

# More Comments on Inflation

- During inflation, the universe expands by at least factor of  $10^{28}$
- If there is anything *besides* vacuum, it is getting red shifted like crazy

$$T \propto a^{-1} \quad T_{\text{end}} < T_{\text{begin}} 10^{-28} \quad k_B T_{\text{end}} < 10^{-28} \cdot 10^{16} \text{ GeV} \quad T_{\text{end}} < 10 \text{ K}$$

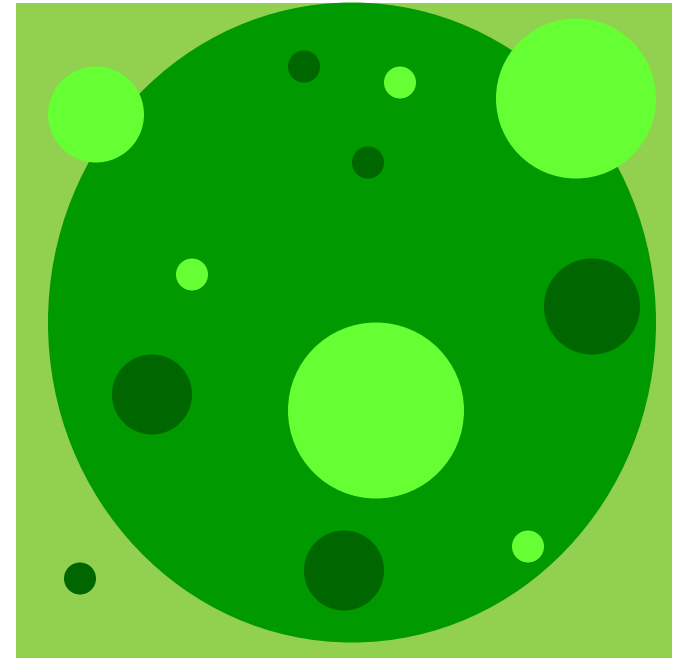
- Completely ignore radiation
- Anything else in the universe also is diluted and disappears

$$n \propto a^{-3} \quad n_{\text{end}} < 10^{-84} n_{\text{begin}}$$

- By the time inflation is over, any record of previous universe is effectively erased

# Density Fluctuations from Inflation

- During inflation, you could think universe is very smooth
- On microscopic scales, universe will have tiny density fluctuations, thanks to quantum mechanics
- The size and magnitude of these fluctuations can be calculated
  - Gaussian, distribution, for example
- In size, these are so small they would be irrelevant
  - Except the universe is inflating!
- Little fluctuations grow up to be big fluctuations
- Constant new fluctuations are appearing at all scales
- During inflation: predict “scale invariant” density perturbations
  - “Spectral index”  $n = 1$



# Exiting Inflation and Reheating

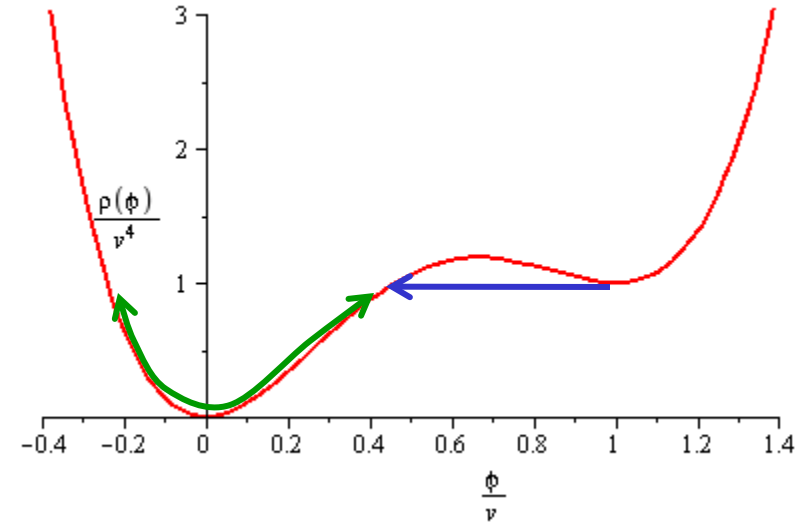
- Universe must make transition to true vacuum
- Old inflation picture:
  - It quantum tunnels from false vacuum to true
  - It oscillates back and forth around true
  - Oscillations generate all kinds of particles
    - Reheating

$$\rho_{\Lambda} \rightarrow \rho_r$$

$$\frac{\pi^2}{30} g_{\text{eff}} \frac{(k_B T)^4}{(\hbar c)^3 c^2} = \frac{v^4}{10(\hbar c)^3 c^2}$$

$$k_B T_r = v \left( \frac{10}{\pi^2 g_{\text{eff}}} \right)^{1/4}$$

$$k_B T_r \approx \frac{1}{4} v$$

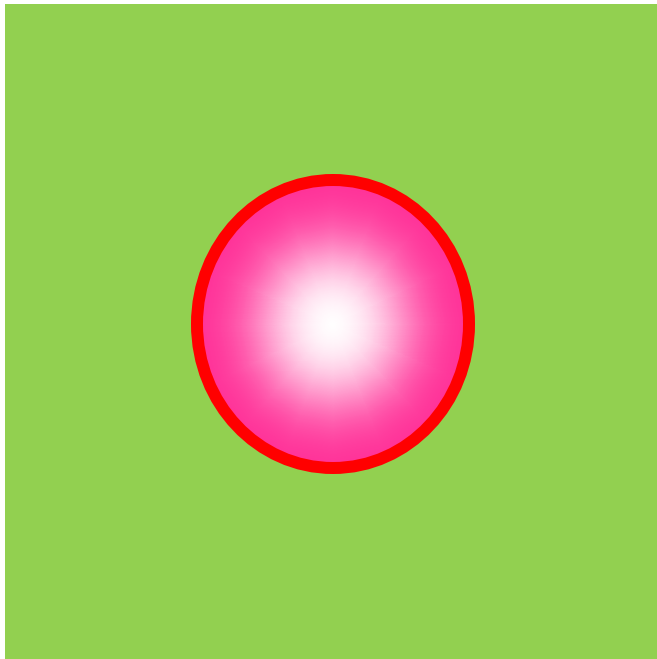
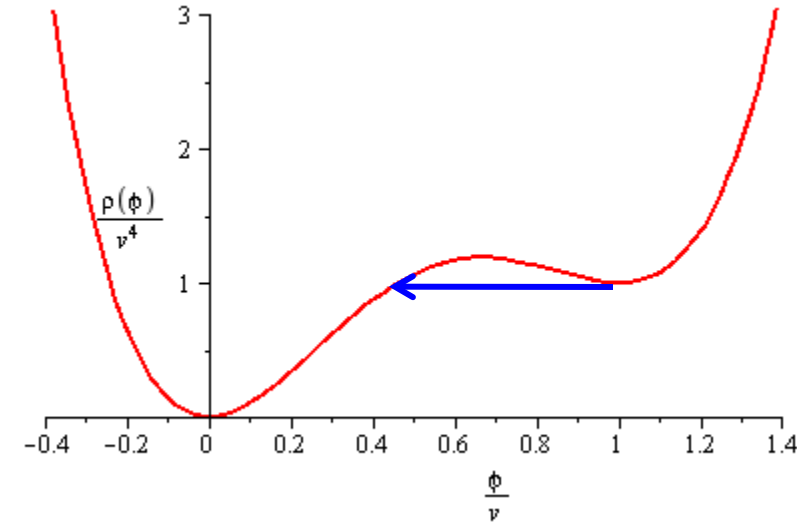


- Baryogenesis must occur after inflation ends
  - Otherwise any baryons would just inflate away
- We re-enter radiation era, but with  $t$  a little increased because of inflationary era

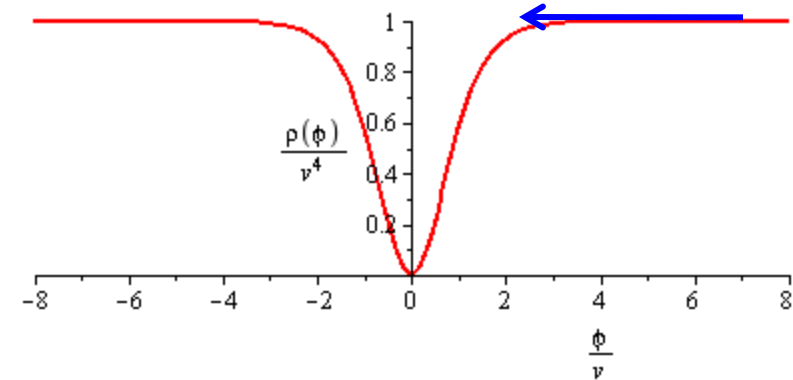
<u>Event</u>	<u><math>k_B T</math> or <math>T</math></u>	<u>Time</u>
Inflation begins	$10^{16}$ GeV	$10^{-39}$ s
Inflation ends, baryons created, etc.	$10^{15}$ GeV	$10^{-35}$ s

# Inflation is Dead. Long Live Inflation!

- The problem – it's hard to quantum tunnel
- It will only occur *very rarely*
- It will occur at one place
  - Bubble nucleation
- All the energy is on the walls of the expanding bubble
- If bubbles collided, we could get reheating
  - But universe expands so fast, bubbles never find each other



- The solution: change the potential
- No quantum tunneling needed
- “Slow roll” potentials
- Whole universe moves to minimum at same time



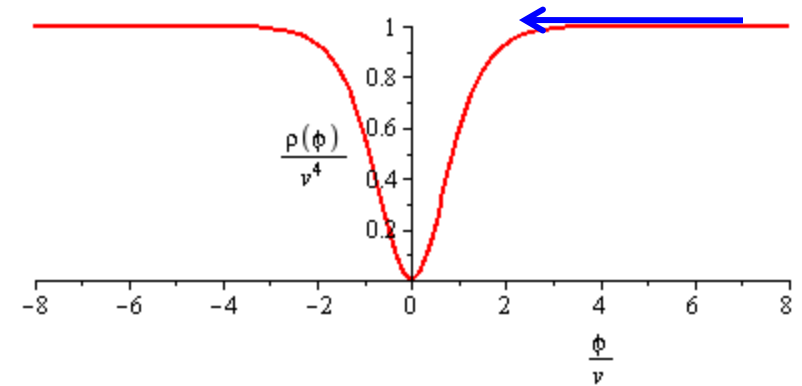


# Too Many Kinds of Inflation

- Many *many* variations on inflation have since been explored
  - Slow Roll inflation
  - Chaotic inflation
  - Designer inflation
- We don't know which (if any) of these are correct

Generic features of modern inflation:

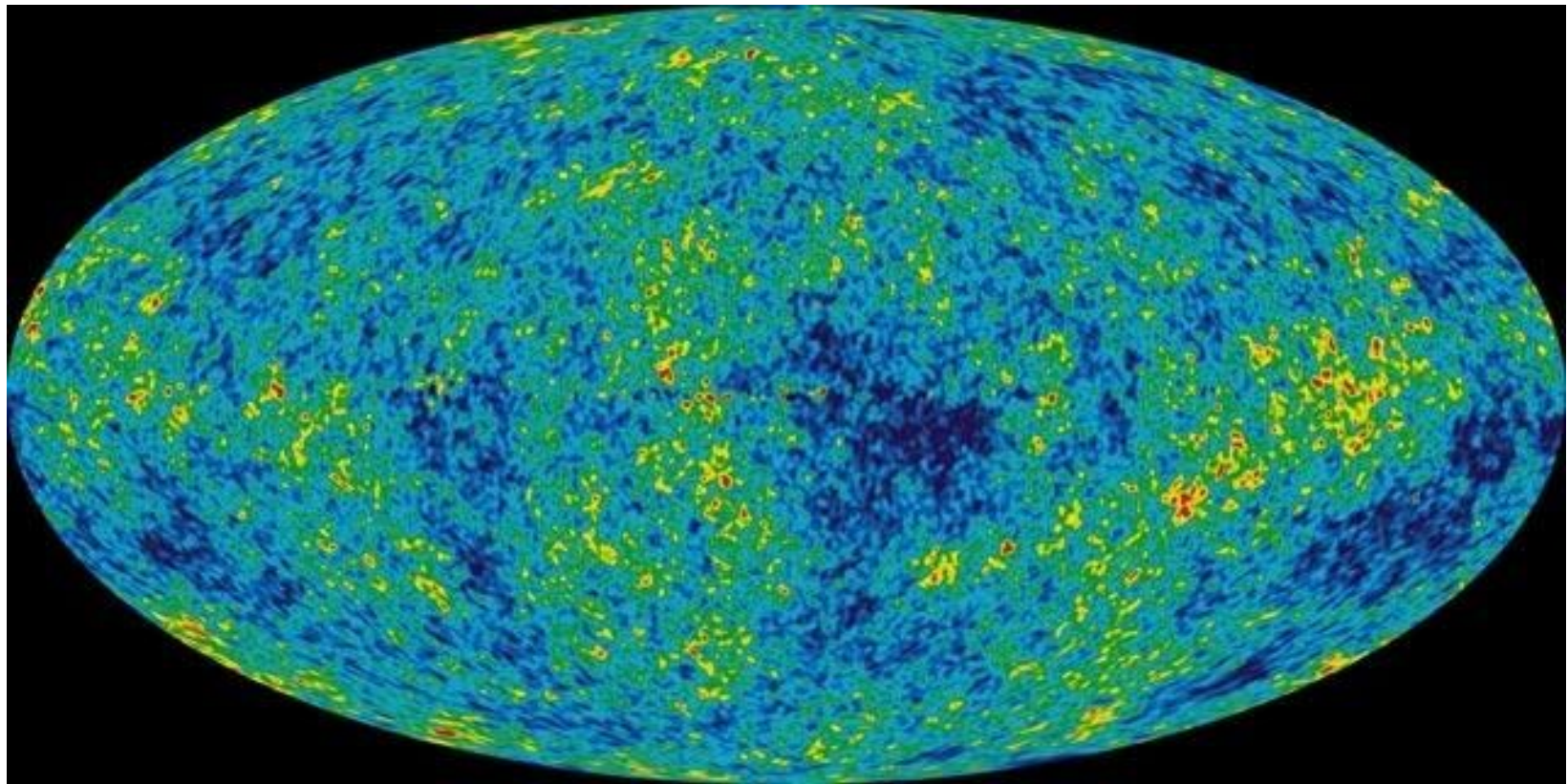
- Universe “slows down” just before exiting inflation
- Expansion not so fast at the end
- Small fluctuations don't grow quite so big
- Fluctuation spectrum has more power at smaller scales
  - Not quite scale invariant
  - Spectral index just under 1



# Density Fluctuations

## Why We Need Them

- If the universe were *perfectly* uniform, then how come the microwave background isn't uniform?
- Where did all the structure (galaxies, clusters, etc.) come from?
- Can we understand where these variations come from?
- Can we explain how these variations led to the structures we see?



# Review: Spherical Harmonics

- The temperature is a function of angle only
- The *spherical harmonics* are a series of functions of angle only
  - Taught about them in PHY 215 and other courses
- There are an infinite number of them
  - $l = 0, 1, 2, 3, \dots$
  - $m = -l, -l+1, \dots, l$

$$T(\theta, \phi)$$

$$Y_{lm}(\theta, \phi)$$

- The smaller indices represent functions that vary slowly with angle, the larger ones vary more quickly
- For example, the  $\phi$  dependence depends on  $m$
- The angular scale at which the function changes sign is, crudely,
  - This could be based on changes in  $\theta$  or  $\phi$

$$Y_{lm}(\theta, \phi) \propto e^{im\phi} = \cos(m\phi) + i \sin(m\phi)$$

$$\theta \approx \frac{\pi}{l}$$

- Any function that depends only on angle can be written as a linear combination of spherical harmonics
- The coefficients  $C_{lm}$  contain the same information as the function  $T(\theta, \phi)$

$$T(\theta, \phi) = \sum_{l=0}^{\infty} \sum_{m=-l}^l C_{lm} Y_{lm}(\theta, \phi)$$

# How to Describe the Fluctuations

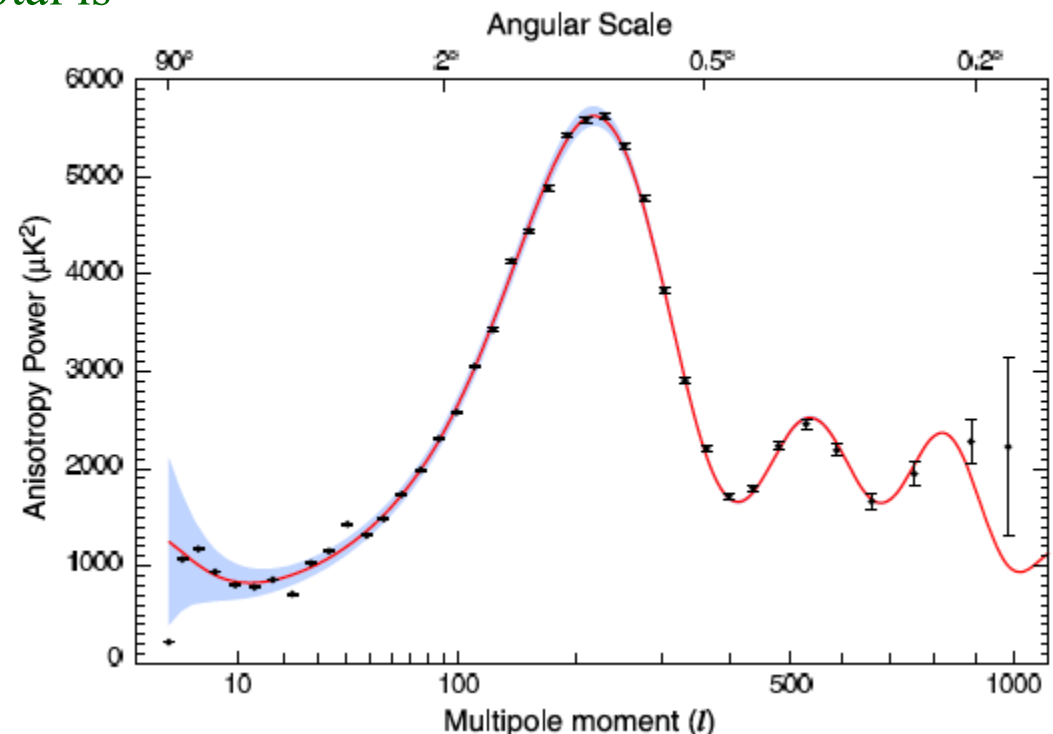
- When you rotate the coordinates, you change  $Y_{lm}$  spherical harmonics into others with the same value of  $l$  but different values of  $m$

$$T(\theta, \phi) = \sum_{l=0}^{\infty} \sum_{m=-l}^l C_{lm} Y_{lm}(\theta, \phi)$$

- The coefficients  $C_{lm}$  get mixed within different values of  $m$  when you rotate
  - Exact values of  $C_{lm}$  depends on choice of axes
- But, if you sum up the  $C_{lm}$ 's over fixed  $l$ , the total is independent of  $m$ , depends only on  $l$
- So the information about how much variation on each angular “scale” is contained in

$$\langle C_l^2 \rangle = \frac{1}{2l+1} \sum_{m=-l}^l |C_{lm}|^2$$

- Plot the result as a function of  $l$ 
  - Dots are the data, solid is the fit
- How do we explain all these features?



# Growth of Density Perturbations

- Assume inflation, or some other source, produces density perturbations at all scales
  - Spectral index = 1, or slightly less than 1
  - A bit more “power” at small scales than at large

Now, we need to discuss what happens at all possible scales during four eras:

- Radiation dominated era:  $z > 3000$
- Matter domination before recombination:  $1100 < z < 3000$
- Post recombination:  $10 < z < 1100$
- Structure formation:  $z < 10$

The cosmic microwave background shows us the universe at recombination

# Coupling in the Different Eras:

- The dark matter probably doesn't collide with anything
  - It will do its own thing
  - It is moving slowly and has little thermal pressure
- Neutrinos move nearly at the speed of light
  - As soon as they can, they simply erase any structure
  - Ignore them from now on

Everything else moves together before recombination:

- The baryons collide with electrons (Coulomb scattering)
- The photons collide with electrons (Thomson scattering)
- The baryons, photons, and electrons all move together
- The photons provide the lion's share of the pressure

After recombination, the atoms decouple from the photons

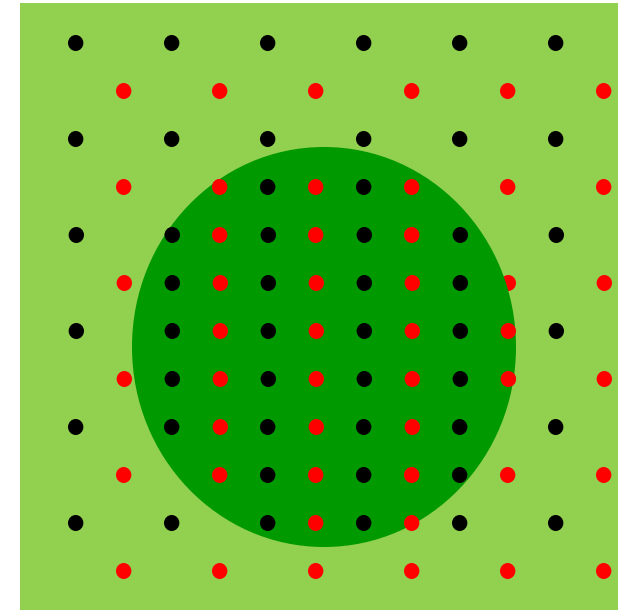
- The pressure drops effectively to zero

# The Radiation Era: What Happens

- Suppose we have a perturbation (say, a high density region) during this era:
- Inflation predicts: every type of particle will have higher density in this region

Radiation  
Dark Matter  
Baryons/Atoms

- Photons try to flow from high density to low density
- They move at speed of light
- When size of perturbation equals about  $ct$ , it gets wiped out
  - Actually, they oscillate, more on this later
- The baryons *also* get wiped out, because they are caught with the photons
- The dark matter stays behind, and gravity starts to gather it more
  - But very slowly, because radiation dominates



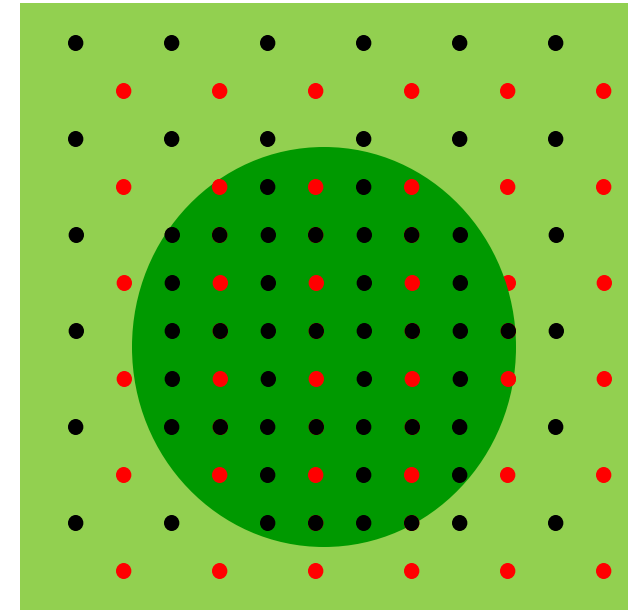
# Matter Domination, Before Recombination

- As before, the photons and **baryons** flow from high to low density as soon as they are able
  - Actually, they oscillate, as I'll explain soon

- However, now the dark matter dominates
- Gravity causes high density regions to get higher density
- They start to form structures
- These structures will become galaxies, etc.

- We see the microwave background from the *end* of this era
  - We don't see large density perturbations from this era
- We are looking only at the photons (coupled to baryons)

Radiation  
Dark Matter  
**Baryons/Atoms**





# Acoustic Modes

- Before recombination: The photons and baryons are all coupled together
- Consider regions of alternating high density/low density:

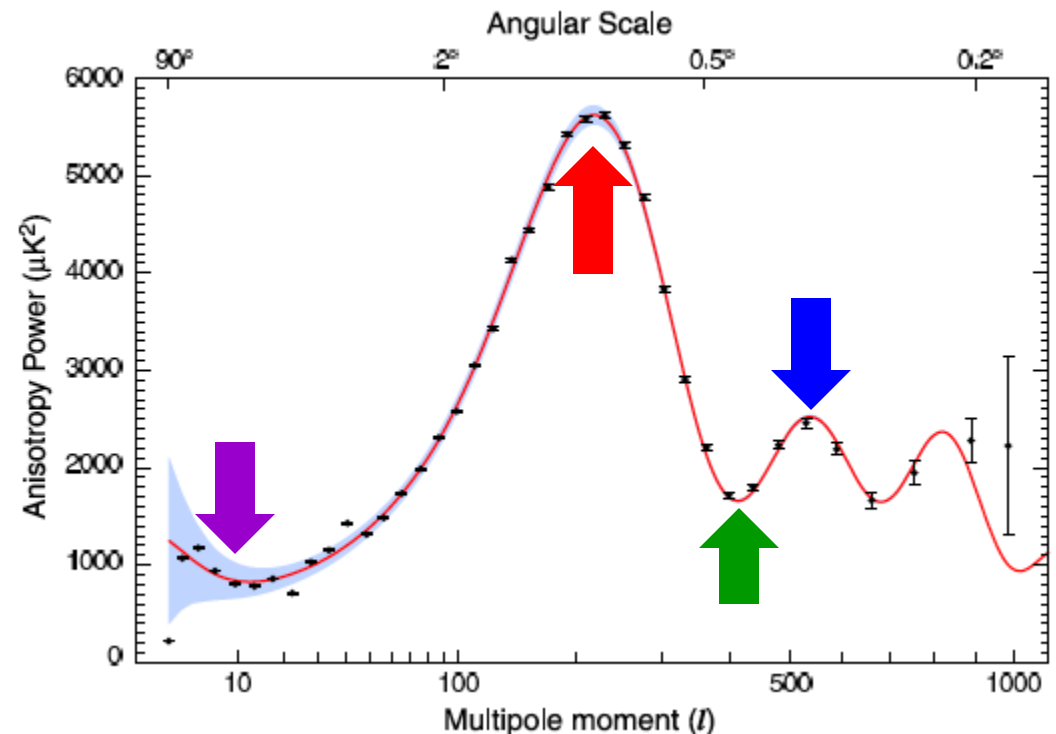


- High density regions try to flow to low density
- This creates a sound wave, just like an ordinary sound wave
- This process stops at recombination
  - We catch the wave, frozen at that moment of time

# Acoustic Modes and Scales

Where will we catch the wave at recombination? It depends on wavelength

- Longest wavelength: It hasn't had time to even cycle at all
  - Relatively little “power” at largest scales
- Next longest: It will be right at its first peak in oscillation cycle
  - This will be sensitive to  $\Omega$
- Slightly longer: Wave is at a node, or minimum
- Longer still: The second peak
  - Wave has oscillated, it is now at anti-node
  - Baryons feel gravity of dark matter, and resist (baryon drag)
  - Sensitive to  $\Omega_b$



# Where Will the First Peak Be? (1)

- This is a wave that has only gone through about a third of a cycle

$$\frac{1}{3}\lambda \approx v_s t \quad v_s = \sqrt{\frac{P}{\rho}} = \sqrt{\frac{\frac{1}{3}u}{u/c^2}} = \frac{c}{\sqrt{3}}$$

$$\lambda = \sqrt{3} (3 \times 10^8 \text{ m/s}) (3.8 \times 10^5 \text{ y}) (3.16 \times 10^7 \text{ s/y})$$

$$= 6.24 \times 10^{21} \text{ m}$$

- This scale has since grown to about

$$\lambda_0 = \lambda (1 + z_*) = 1092 (6.24 \times 10^{21} \text{ m}) = 6.81 \times 10^{24} \text{ m}$$

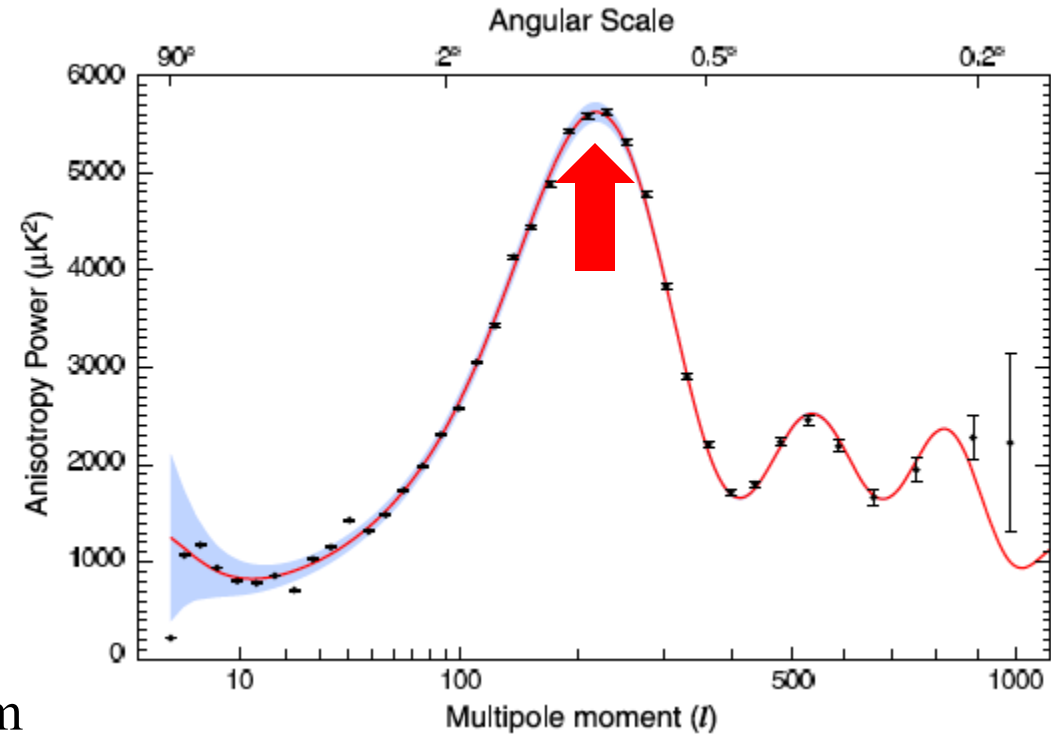
- The distance to this wave is now about
- The angular size of the wave is then

$$\theta = \frac{\lambda_0}{d} = \frac{6.81 \times 10^{24} \text{ m}}{4.4 \times 10^{26} \text{ m}} = 0.015$$

$$d \approx 3.3ct_0 \approx 4.4 \times 10^{26} \text{ m}$$

$$l = \frac{\pi}{\theta} \approx 210$$

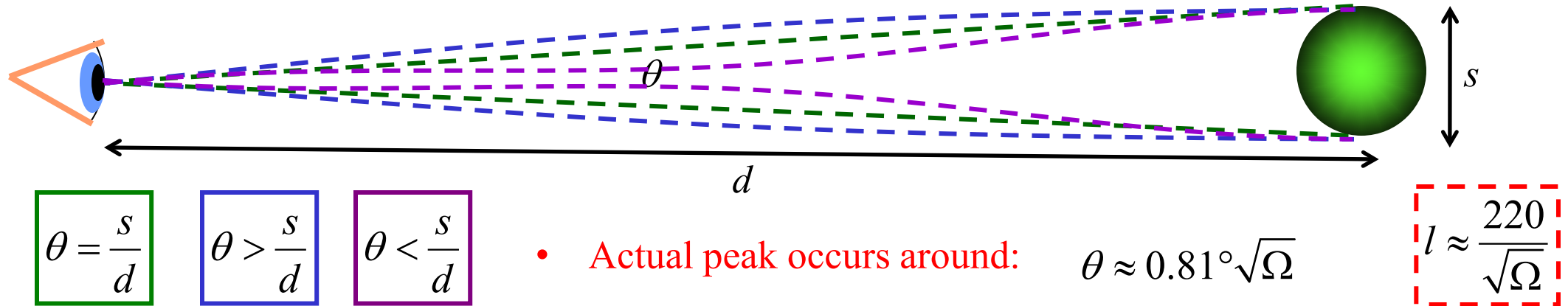
- More precise analysis says peak should be at  $l = 220$



# Where Will the First Peak Be? (2)

- All calculations assumed standard cosmological values
  - Including  $\Omega = 1$
- There are *many* effects that change this, most notably, the curvature of spacetime
- In flat space,  $\Omega = 1$ , we used the correct distance/angle relationship
- If  $\Omega > 1$ , then curvature of universe will make angular size appear bigger
- If  $\Omega < 1$ , then curvature of universe will make angular size appear smaller

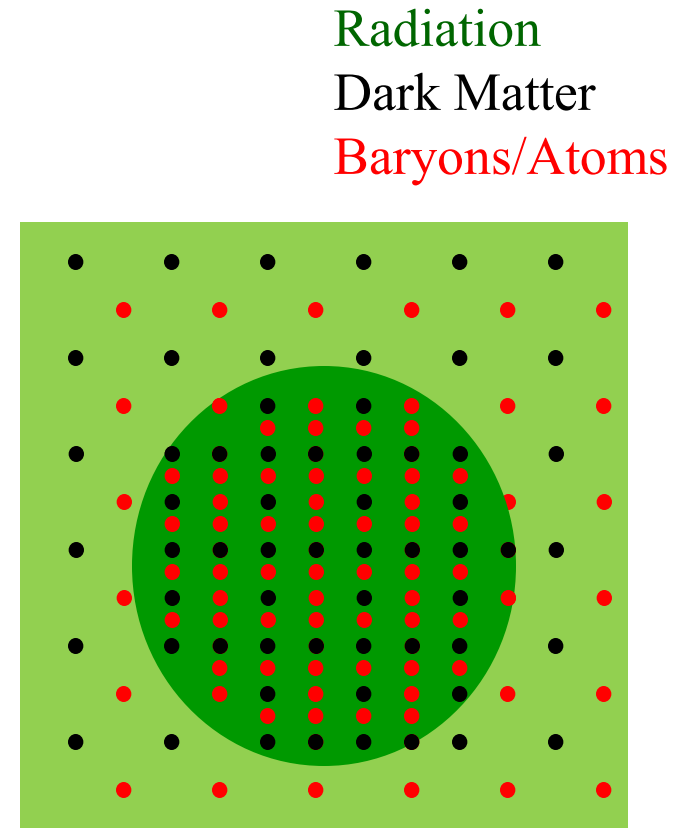
$$l = 220$$



- By measuring actual position of peak, we can estimate  $\Omega$ 
  - *This* is the primary source of the information that  $\Omega = 1.00$

# Matter Domination, After Recombination

- Photons completely decouple, they become completely smooth
- Any dark matter high density regions will gather baryons
  - These regions quickly become high density again
- These regions continue to grow more and more massive
- Throughout all these periods (up to now) these regions are continuing to grow with the rest of the universe in physical size



# Outline of History of Universe

<u>Time</u>	<u><math>T</math> or <math>k_B T</math></u>	<u>Events</u>
$10^{-43}$ s	$10^{18}$ GeV	Planck Era; time becomes meaningless?
$10^{-39}$ s	$10^{16}$ GeV	Inflation begins; forces unified
$10^{-35}$ s	$10^{15}$ GeV	Inflation ends; reheating; forces separate; baryosynthesis (?)
$10^{-13}$ s	1500 GeV	Supersymmetry breaking, LSP (dark matter)
$10^{-11}$ s	160 GeV	Electroweak symmetry breaking
14 $\mu$ s	150 MeV	Quark Confinement
0.4 s	1.5 MeV	Neutrino Decoupling
1.5 s	0.7 MeV	Neutron/Proton freezeout
20 s	170 keV	Electron/Positron annihilation
200 s	80 keV	Nucleosynthesis
57 ky	0.76 eV	Matter-Radiation equality
370 ky	0.26 eV	Recombination
600 My	30 K	First Structure/First Stars
13.8 Gy	2.725 K	Today

# Structure Formation

- Up to now, the density fluctuations have been assumed small
- As universe expands, can show in general that
- But the density itself is falling as
- Therefore,

$$\delta\rho/\rho \ll 1$$

$$\delta\rho \propto a^{-2}$$

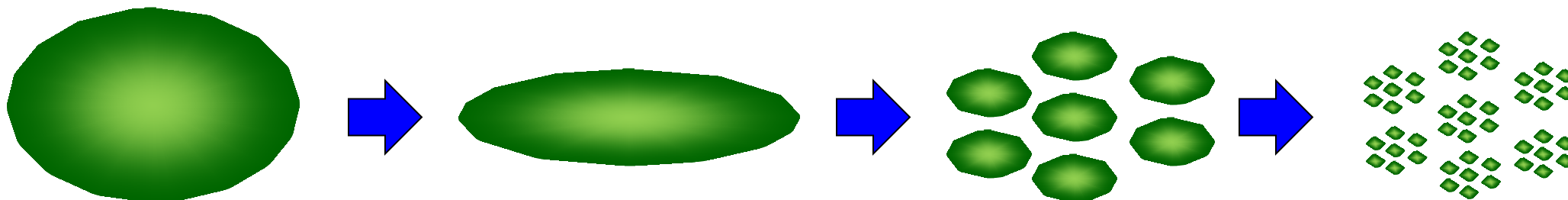
$$\rho \propto a^{-3}$$

$$\delta\rho/\rho \propto a$$

- Eventually, the perturbation becomes large (non-perturbative)
- To some extent, must rely on computer simulations

Qualitatively:

- Collapses in some direction first
- Then it fragments
- Fragments fragment further
- Form successively smaller structures



# Which Structures Form First?

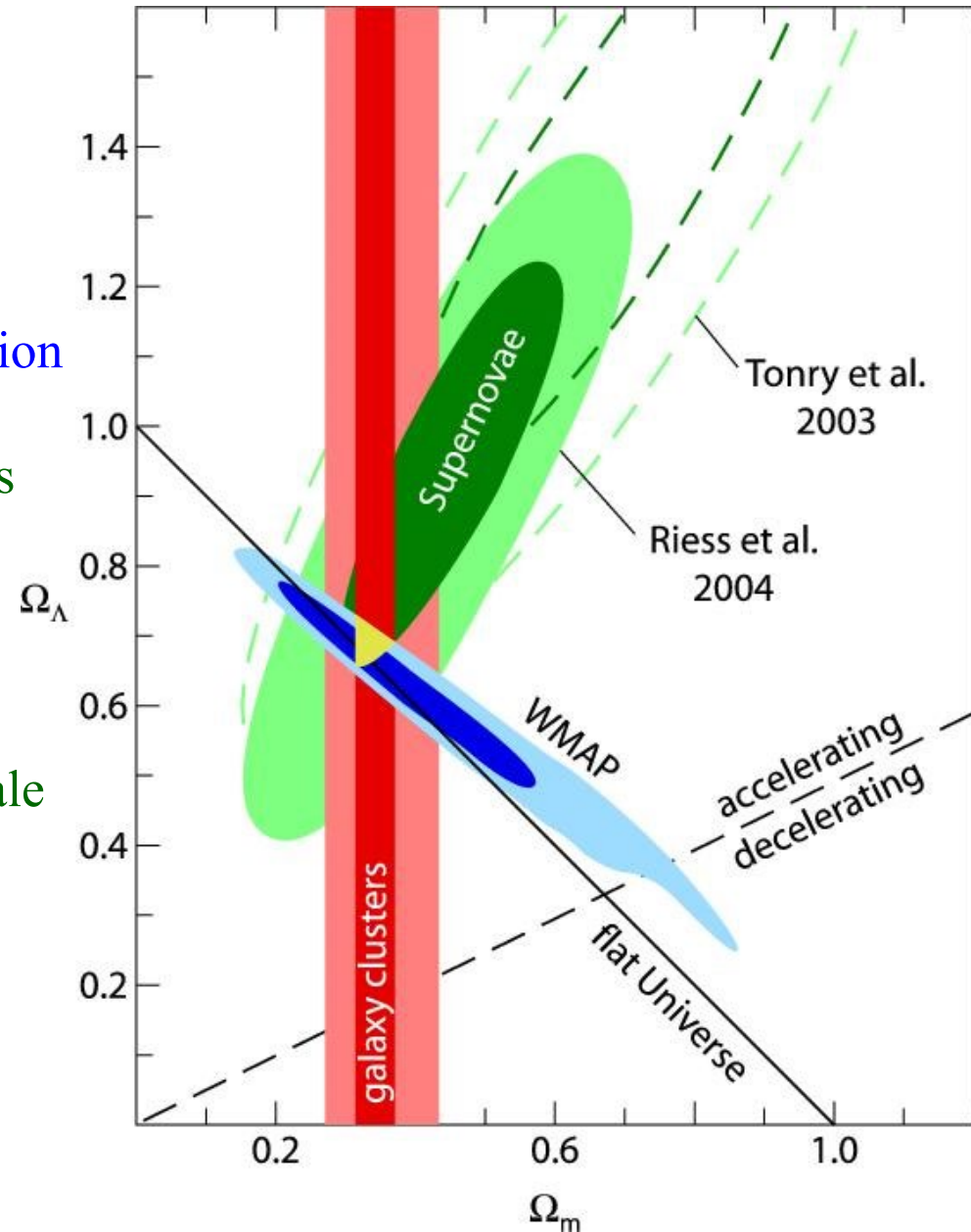
- For “scale invariant” density perturbations, all scales start about the same
  - But inflation implies smaller structures have a slight advantage
- During radiation era, smaller structures can grow a little
  - Another slight advantage for small structures
- Structures larger than the size of the universe  $ct$  at matter/radiation equality have a big disadvantage
  - This sets the largest scales at which structures form
  - This is a homework problem
- There is also a smallest structure that can form, set by the Jeans Mass
  - When pressure overcomes gravity
  - Homework problem
- Conclusion – smallest structures form first, but not by a lot
  - Globular Clusters/Dwarf Ellipticals – around  $z = 8 - 12$
  - Galaxies/quasars/etc. – around  $z = 5-8$
  - Galaxy Clusters – around  $z = 2-5$
  - Galaxy Superclusters – around  $z = 0-2$



# Cosmological Parameters (1)

- We have three independent ways of estimating cosmological parameters
  - Cosmic background
  - Structure in the current universe
  - Supernovae red shift
- Collectively, these give lots of independent information
- The fact that these all work *consistently* together tells us we are on the right track
  - And note that  $\Omega = 1$
- To extract the *best* values, do a full fit to *everything* simultaneously
- Indications are that inflation (especially not quite scale invariant structure) yields best fit
- *Most* cosmologists now think inflation is right

$$n_s = 0.9665 \pm 0.0038$$



# Cosmological Parameters (2)

Parameter	TT+lowE 68% limits	TE+lowE 68% limits	EE+lowE 68% limits	TT,TE,EE+lowE 68% limits	TT,TE,EE+lowE+lensing 68% limits	TT,TE,EE+lowE+lensing+BAO 68% limits
$\Omega_b h^2$ . . . . .	$0.02212 \pm 0.00022$	$0.02249 \pm 0.00025$	$0.0240 \pm 0.0012$	$0.02236 \pm 0.00015$	$0.02237 \pm 0.00015$	$0.02242 \pm 0.00014$
$\Omega_c h^2$ . . . . .	$0.1206 \pm 0.0021$	$0.1177 \pm 0.0020$	$0.1158 \pm 0.0046$	$0.1202 \pm 0.0014$	$0.1200 \pm 0.0012$	$0.11933 \pm 0.00091$
$100\theta_{\text{MC}}$ . . . . .	$1.04077 \pm 0.00047$	$1.04139 \pm 0.00049$	$1.03999 \pm 0.00089$	$1.04090 \pm 0.00031$	$1.04092 \pm 0.00031$	$1.04101 \pm 0.00029$
$\tau$ . . . . .	$0.0522 \pm 0.0080$	$0.0496 \pm 0.0085$	$0.0527 \pm 0.0090$	$0.0544^{+0.0070}_{-0.0081}$	$0.0544 \pm 0.0073$	$0.0561 \pm 0.0071$
$\ln(10^{10} A_s)$ . . . . .	$3.040 \pm 0.016$	$3.018^{+0.020}_{-0.018}$	$3.052 \pm 0.022$	$3.045 \pm 0.016$	$3.044 \pm 0.014$	$3.047 \pm 0.014$
$n_s$ . . . . .	$0.9626 \pm 0.0057$	$0.967 \pm 0.011$	$0.980 \pm 0.015$	$0.9649 \pm 0.0044$	$0.9649 \pm 0.0042$	$0.9665 \pm 0.0038$
$H_0$ [km s <sup>-1</sup> Mpc <sup>-1</sup> ] . .	$66.88 \pm 0.92$	$68.44 \pm 0.91$	$69.9 \pm 2.7$	$67.27 \pm 0.60$	$67.36 \pm 0.54$	$67.66 \pm 0.42$
$\Omega_\Lambda$ . . . . .	$0.679 \pm 0.013$	$0.699 \pm 0.012$	$0.711^{+0.033}_{-0.026}$	$0.6834 \pm 0.0084$	$0.6847 \pm 0.0073$	$0.6889 \pm 0.0056$
$\Omega_m$ . . . . .	$0.321 \pm 0.013$	$0.301 \pm 0.012$	$0.289^{+0.026}_{-0.033}$	$0.3166 \pm 0.0084$	$0.3153 \pm 0.0073$	$0.3111 \pm 0.0056$
$\Omega_m h^2$ . . . . .	$0.1434 \pm 0.0020$	$0.1408 \pm 0.0019$	$0.1404^{+0.0034}_{-0.0039}$	$0.1432 \pm 0.0013$	$0.1430 \pm 0.0011$	$0.14240 \pm 0.00087$
$\Omega_m h^3$ . . . . .	$0.09589 \pm 0.00046$	$0.09635 \pm 0.00051$	$0.0981^{+0.0016}_{-0.0018}$	$0.09633 \pm 0.00029$	$0.09633 \pm 0.00030$	$0.09635 \pm 0.00030$
$\sigma_8$ . . . . .	$0.8118 \pm 0.0089$	$0.793 \pm 0.011$	$0.796 \pm 0.018$	$0.8120 \pm 0.0073$	$0.8111 \pm 0.0060$	$0.8102 \pm 0.0060$
$S_8 \equiv \sigma_8 (\Omega_m/0.3)^{0.5}$ .	$0.840 \pm 0.024$	$0.794 \pm 0.024$	$0.781^{+0.052}_{-0.060}$	$0.834 \pm 0.016$	$0.832 \pm 0.013$	$0.825 \pm 0.011$
$\sigma_8 \Omega_m^{0.25}$ . . . . .	$0.611 \pm 0.012$	$0.587 \pm 0.012$	$0.583 \pm 0.027$	$0.6090 \pm 0.0081$	$0.6078 \pm 0.0064$	$0.6051 \pm 0.0058$
$z_{\text{re}}$ . . . . .	$7.50 \pm 0.82$	$7.11^{+0.91}_{-0.75}$	$7.10^{+0.87}_{-0.73}$	$7.68 \pm 0.79$	$7.67 \pm 0.73$	$7.82 \pm 0.71$
$10^9 A_s$ . . . . .	$2.092 \pm 0.034$	$2.045 \pm 0.041$	$2.116 \pm 0.047$	$2.101^{+0.031}_{-0.034}$	$2.100 \pm 0.030$	$2.105 \pm 0.030$
$10^9 A_s e^{-2\tau}$ . . . . .	$1.884 \pm 0.014$	$1.851 \pm 0.018$	$1.904 \pm 0.024$	$1.884 \pm 0.012$	$1.883 \pm 0.011$	$1.881 \pm 0.010$
Age [Gyr] . . . . .	$13.830 \pm 0.037$	$13.761 \pm 0.038$	$13.64^{+0.16}_{-0.14}$	$13.800 \pm 0.024$	$13.797 \pm 0.023$	$13.787 \pm 0.020$
$z_*$ . . . . .	$1090.30 \pm 0.41$	$1089.57 \pm 0.42$	$1087.8^{+1.6}_{-1.7}$	$1089.95 \pm 0.27$	$1089.92 \pm 0.25$	$1089.80 \pm 0.21$
$r_*$ [Mpc] . . . . .	$144.46 \pm 0.48$	$144.95 \pm 0.48$	$144.29 \pm 0.64$	$144.39 \pm 0.30$	$144.43 \pm 0.26$	$144.57 \pm 0.22$
$100\theta_*$ . . . . .	$1.04097 \pm 0.00046$	$1.04156 \pm 0.00049$	$1.04001 \pm 0.00086$	$1.04109 \pm 0.00030$	$1.04110 \pm 0.00031$	$1.04119 \pm 0.00029$
$z_{\text{drag}}$ . . . . .	$1059.39 \pm 0.46$	$1060.03 \pm 0.54$	$1063.2 \pm 2.4$	$1059.93 \pm 0.30$	$1059.94 \pm 0.30$	$1060.01 \pm 0.29$
$r_{\text{drag}}$ [Mpc] . . . . .	$147.21 \pm 0.48$	$147.59 \pm 0.49$	$146.46 \pm 0.70$	$147.05 \pm 0.30$	$147.09 \pm 0.26$	$147.21 \pm 0.23$
$k_D$ [Mpc <sup>-1</sup> ] . . . . .	$0.14054 \pm 0.00052$	$0.14043 \pm 0.00057$	$0.1426 \pm 0.0012$	$0.14090 \pm 0.00032$	$0.14087 \pm 0.00030$	$0.14078 \pm 0.00028$