

Physics 780 – General Relativity
Solution Set B

5. [15] We can define the four-acceleration as $A^\mu = \frac{d}{d\tau}U^\mu = \frac{d^2}{d\tau^2}x^\mu$.

(a) [4] Show that if you are in a frame such that a particle is momentarily at rest, so

$U^\mu(\tau_0) = (1, 0, 0, 0)$, then $\eta_{\mu\nu}A^\mu A^\nu = \mathbf{a}^2$, where \mathbf{a} is the ordinary acceleration $\mathbf{a} = \frac{d\mathbf{v}}{dt}$.

Since $a = \sqrt{\eta_{\mu\nu}A^\mu A^\nu}$ is Lorentz invariant, this must be the magnitude of the acceleration as experienced by an object with arbitrary velocity U^μ .

The four-velocity is given by $U^\mu(\tau_0) = (\gamma, \gamma\mathbf{v})$, where $\gamma = (1 - v^2)^{-1/2}$. Taking the derivative, we have

$$A^\mu = \frac{d}{d\tau}U^\mu = \left(\frac{d\gamma}{d\tau}, \frac{d\gamma}{d\tau}\mathbf{v} + \gamma \frac{d\mathbf{v}}{d\tau} \right) = \frac{dt}{d\tau} \left(\frac{d\gamma}{dt}, \frac{d\gamma}{dt}\mathbf{v} + \gamma \frac{d\mathbf{v}}{dt} \right).$$

It is easy to show that $d\gamma/dt = (1 - v^2)^{-3/2} v(dv/dt) = \gamma^3 v(dv/dt)$. We also know that $dt/d\tau = \gamma$. So we have

$$A^\mu = \left(\gamma^4 v \frac{dv}{dt}, \gamma^4 \mathbf{v} v \frac{dv}{dt} + \gamma^2 \mathbf{a} \right).$$

If the instantaneous velocity is zero, then we have $v = 0$, $\gamma = 1$, and therefore $A^\mu = (0, \mathbf{a})$, so $\eta_{\mu\nu}A^\mu A^\nu = \mathbf{a}^2$. Hence $\sqrt{\eta_{\mu\nu}A^\mu A^\nu} = a$.

(b) [5] An object moves according to the formula $x = \sqrt{b^2 + t^2} - b$, $y = z = 0$. Rewrite these in terms of proper time τ (hint: this was almost done for you in class), then work out the four-velocity $U^\mu(\tau)$ the four-acceleration $A^\mu(\tau)$ and find a formula for the proper acceleration a , which should be constant.

This is almost identical with the problem done in class. We have $dx = t dt / \sqrt{b^2 + t^2}$, so

$$\tau = \int d\tau = \int \sqrt{-\eta_{\mu\nu}dx^\mu dx^\nu} = \int \sqrt{dt^2 - dx^2} = \int \sqrt{dt^2 - \frac{(t dt)^2}{b^2 + t^2}} = \int \frac{b dt}{\sqrt{b^2 + t^2}}.$$

This integral is identical to one we did in class, and is given by

$$\tau = b \sinh^{-1}(t/b).$$

We can invert this to get t in terms of τ , and then substitute to find x in terms of τ .

$$t = b \sinh(\tau/b),$$

$$x = \sqrt{b^2 + t^2} - b = \sqrt{b^2 + b^2 \sinh^2(\tau/b)} - b = b \cosh(\tau/b) - b.$$

We now start working out the four-velocity and the four-acceleration:

$$U^\mu = \frac{d}{d\tau}(b \sinh(\tau/b), b \cosh(\tau/b) - b, 0, 0) = (\cosh(\tau/b), \sinh(\tau/b), 0, 0)$$

$$A^\mu = \frac{d}{d\tau}(\cosh(\tau/b), \sinh(\tau/b), 0, 0) = \left(\frac{1}{b} \sinh\left(\frac{\tau}{b}\right), \frac{1}{b} \cosh\left(\frac{\tau}{b}\right), 0, 0\right).$$

The acceleration dotted into itself is therefore

$$\eta_{\mu\nu} A^\mu A^\nu = -(A^0)^2 + (A^1)^2 = -\frac{1}{b^2} \sinh^2\left(\frac{\tau}{b}\right) + \frac{1}{b^2} \cosh^2\left(\frac{\tau}{b}\right) = \frac{1}{b^2}.$$

We therefore conclude that $a^2 = 1/b^2$, so $b = 1/a$.

(c) [3] Determine the value of b in years (or light-years) if $a = g = 9.80 \text{ m/s}^2$.

Keeping in mind that we are working in units where $c = 1$, we have

$$b = \frac{1}{g} = \frac{2.998 \times 10^8 \text{ m/s}}{9.80 \text{ m/s}^2} = \frac{3.059 \times 10^7 \text{ s}}{3.156 \times 10^7 \text{ s/yr}} = 0.969 \text{ yr}.$$

Since a light year is $c \cdot \text{yr}$, this is also the same in light years.

(d) [3] If you leave Earth, starting at rest, and accelerate at g , how much proper time in years would it take you go get to α -Centauri (4.3 $c \cdot \text{yr}$), the center of the galaxy (27,000 $c \cdot \text{yr}$) and the approximate edge of the Universe ($2 \times 10^{10} c \cdot \text{yr}$).

We note that our formulas have initially $x = 0$, so these formulas describe the motion as a function of proper time. Solving for τ in terms of x , we find

$$\tau = b \cosh^{-1}\left(\frac{x}{b} + 1\right).$$

We now simply substitute all the relevant distances

$$\tau_\alpha = (0.969 \text{ yr}) \cosh^{-1}\left(\frac{4.30 \text{ yr}}{0.969 \text{ yr}} + 1\right) = 2.30 \text{ yr}, \quad \tau_g = (0.969 \text{ yr}) \cosh^{-1}\left(\frac{27,000 \text{ yr}}{0.969 \text{ yr}} + 1\right) = 10.6 \text{ yr},$$

$$\tau_U = (0.969 \text{ yr}) \cosh^{-1}\left(\frac{2 \times 10^{10} \text{ yr}}{0.969 \text{ yr}} + 1\right) = 23.7 \text{ yr}.$$

Interestingly, you can get almost anywhere in the universe, but you have to use a substantial fraction of your lifetime to get there.

6. The electromagnetic field tensor $F_{\mu\nu}$ is given by equation (1.69). If the fields are given by the three components of \mathbf{E} and \mathbf{B} , what would be the new values of the electric and magnetic fields \mathbf{E}' and \mathbf{B}' if you

(a) Performed a rotation by an angle θ around the z -axis,

(b) Perform a boost in the x -direction by rapidity ϕ .

The corresponding inverse Lorentz matrices in each case are given below. Note that because $F_{\mu\nu}$ is anti-symmetric, you can calculate just six components of $F_{\mu\nu}$ to get \mathbf{E}' and \mathbf{B}' .

$$\Lambda(\theta)^{-1} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & \cos \theta & -\sin \theta & 0 \\ 0 & \sin \theta & \cos \theta & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}, \quad \Lambda(\phi)^{-1} = \begin{pmatrix} \cosh \phi & \sinh \phi & 0 & 0 \\ \sinh \phi & \cosh \phi & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

We simply start working them all out, using the inverse Lorentz transformations. For the rotations, we have

$$\begin{aligned} E'_1 = F'_{10} &= F_{\mu\nu} (\Lambda^{-1})^\mu_1 (\Lambda^{-1})^\nu_0 = F_{10} (\Lambda^{-1})^1_1 (\Lambda^{-1})^0_0 + F_{20} (\Lambda^{-1})^2_1 (\Lambda^{-1})^0_0 = E_1 \cos \theta + E_2 \sin \theta, \\ E'_2 = F'_{20} &= F_{\mu\nu} (\Lambda^{-1})^\mu_2 (\Lambda^{-1})^\nu_0 = F_{10} (\Lambda^{-1})^1_2 (\Lambda^{-1})^0_0 + F_{20} (\Lambda^{-1})^2_2 (\Lambda^{-1})^0_0 = -E_1 \sin \theta + E_2 \cos \theta, \\ E'_3 = F'_{30} &= F_{\mu\nu} (\Lambda^{-1})^\mu_3 (\Lambda^{-1})^\nu_0 = F_{30} (\Lambda^{-1})^3_3 (\Lambda^{-1})^0_0 = E_3, \\ B'_1 = F'_{23} &= F_{\mu\nu} (\Lambda^{-1})^\mu_2 (\Lambda^{-1})^\nu_3 = F_{13} (\Lambda^{-1})^1_2 (\Lambda^{-1})^3_3 + F_{23} (\Lambda^{-1})^2_2 (\Lambda^{-1})^3_3 = B_2 \sin \theta + B_1 \cos \theta, \\ B'_2 = F'_{31} &= F_{\mu\nu} (\Lambda^{-1})^\mu_3 (\Lambda^{-1})^\nu_1 = F_{31} (\Lambda^{-1})^3_3 (\Lambda^{-1})^1_1 + F_{32} (\Lambda^{-1})^3_3 (\Lambda^{-1})^2_1 = B_2 \cos \theta - B_1 \sin \theta, \\ B'_3 = F'_{12} &= F_{\mu\nu} (\Lambda^{-1})^\mu_1 (\Lambda^{-1})^\nu_2 = F_{12} (\Lambda^{-1})^1_1 (\Lambda^{-1})^2_2 + F_{21} (\Lambda^{-1})^2_1 (\Lambda^{-1})^1_2 = B_3 \cos^2 \theta + B_3 \sin^2 \theta \\ &= B_3. \end{aligned}$$

For the boosts we have

$$\begin{aligned} E'_1 = F'_{10} &= F_{\mu\nu} (\Lambda^{-1})^\mu_1 (\Lambda^{-1})^\nu_0 = F_{10} (\Lambda^{-1})^1_1 (\Lambda^{-1})^0_0 + F_{01} (\Lambda^{-1})^0_1 (\Lambda^{-1})^1_0 = E_1 (\cosh^2 \phi - \sinh^2 \phi) \\ &= E_1, \\ E'_2 = F'_{20} &= F_{\mu\nu} (\Lambda^{-1})^\mu_2 (\Lambda^{-1})^\nu_0 = F_{20} (\Lambda^{-1})^2_2 (\Lambda^{-1})^0_0 + F_{21} (\Lambda^{-1})^2_2 (\Lambda^{-1})^1_0 = E_2 \cosh \phi - B_3 \sinh \phi, \\ E'_3 = F'_{30} &= F_{\mu\nu} (\Lambda^{-1})^\mu_3 (\Lambda^{-1})^\nu_0 = F_{30} (\Lambda^{-1})^3_3 (\Lambda^{-1})^0_0 + F_{31} (\Lambda^{-1})^3_3 (\Lambda^{-1})^1_0 = E_3 \cosh \phi + B_2 \sinh \phi, \\ B'_1 = F'_{23} &= F_{\mu\nu} (\Lambda^{-1})^\mu_2 (\Lambda^{-1})^\nu_3 = F_{23} (\Lambda^{-1})^2_2 (\Lambda^{-1})^3_3 = B_1, \\ B'_2 = F'_{31} &= F_{\mu\nu} (\Lambda^{-1})^\mu_3 (\Lambda^{-1})^\nu_1 = F_{30} (\Lambda^{-1})^3_3 (\Lambda^{-1})^0_1 + F_{31} (\Lambda^{-1})^3_3 (\Lambda^{-1})^1_1 = E_3 \sinh \phi + B_2 \cosh \phi, \\ B'_3 = F'_{12} &= F_{\mu\nu} (\Lambda^{-1})^\mu_1 (\Lambda^{-1})^\nu_2 = F_{02} (\Lambda^{-1})^0_1 (\Lambda^{-1})^2_2 + F_{12} (\Lambda^{-1})^1_1 (\Lambda^{-1})^2_2 = -E_2 \sinh \phi + B_3 \cosh \phi. \end{aligned}$$