

Physics 745 - Group Theory
Solution Set 25

1. [15] In the electric dipole approximation, the rate at which an atom decaying from one state to another by the emission of a photon is given by

$$\Gamma(I \rightarrow F) = \frac{4}{3} \alpha \omega_{IF}^3 |\mathbf{r}_{FI}|^2 / c^2 \quad \text{where}$$

$$\mathbf{r}_{FI} = \langle F | \mathbf{r} | I \rangle$$

The absolute value symbol means that $|\mathbf{r}_{FI}|^2$ contains not only a sum of the three components of \mathbf{r} , but also the real and imaginary parts.

- (a) [4] Demonstrate first that

$$|\mathbf{r}_{FI}|^2 = \sum_{q=-1}^1 |\langle F | r_q^{(1)} | I \rangle|^2$$

where $r_q^{(1)}$ are the three components of the spherical tensor corresponding to the vector operator \mathbf{r} .

This is straightforward. We have

$$\begin{aligned} \sum_{q=-1}^1 |\langle F | r_q^{(1)} | I \rangle|^2 &= |\langle F | r_{-1}^{(1)} | I \rangle|^2 + |\langle F | r_0^{(1)} | I \rangle|^2 + |\langle F | r_{+1}^{(1)} | I \rangle|^2 \\ &= \frac{1}{2} |\langle F | r_1 | I \rangle - i \langle F | r_2 | I \rangle|^2 + |\langle F | r_3 | I \rangle|^2 + \frac{1}{2} |-\langle F | r_1 | I \rangle - i \langle F | r_2 | I \rangle|^2 \\ &= \frac{1}{2} |\langle F | r_1 | I \rangle|^2 + \frac{1}{2} |\langle F | r_2 | I \rangle|^2 + |\langle F | r_3 | I \rangle|^2 + \frac{1}{2} |\langle F | r_1 | I \rangle|^2 + \frac{1}{2} |\langle F | r_2 | I \rangle|^2 \\ &= |\langle F | \mathbf{r} | I \rangle|^2 = |\mathbf{r}_{FI}|^2 \end{aligned}$$

- (b) [4] An atom in the state $|njm\rangle$ with $j = 3/2$ is about to decay via dipole radiation. What possible j' values might be allowed for the final state $|n'j'm'\rangle$?

According to the Wigner Eckart theorem, the matrix elements we need will be of the form

$$\langle n'j'm' | r_q^{(1)} | n \frac{3}{2} m \rangle = \frac{\langle 1 \frac{3}{2}; qm | j'm' \rangle}{\sqrt{2j'+1}} \langle n'j' || r || nj \rangle$$

For this rate to be non-vanishing, j' must be in the range from $\frac{3}{2} + 1, \dots, \left| \frac{3}{2} - 1 \right| = \frac{5}{2}, \frac{3}{2}, \frac{1}{2}$, so these are the three possibilities.

(c) [7] The atom is actually going to decay to a state with $j' = 1/2$. Using the Wigner Eckart Theorem, find the *relative* rate of decay

$$\Gamma(njm \rightarrow n'j'm')$$

for all non-vanishing possible values of m and m' .

We need the Clebsch-Gordan coefficients $\langle 1\frac{3}{2}; qm | \frac{1}{2}m' \rangle$. The Clebsch routine from the web is happy to help us out. The only non-zero ones will be when $m' = q + m$, so these work out to six possibilities total. We calculate them using commands like this:

```
> for q from -1 to 1 do clebsch(1,3/2,q,1/2-q,1/2,1/2);
  clebsch(1,3/2,q,-1/2-q,1/2,-1/2); end do;
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$$\langle 1\frac{3}{2}; -1, \frac{3}{2} | \frac{1}{2}, \frac{1}{2} \rangle = \langle 1\frac{3}{2}; 1, -\frac{3}{2} | \frac{1}{2}, -\frac{1}{2} \rangle = \sqrt{\frac{1}{2}},$$

$$\langle 1\frac{3}{2}; 1, -\frac{1}{2} | \frac{1}{2}, \frac{1}{2} \rangle = \langle 1\frac{3}{2}; -1, \frac{1}{2} | \frac{1}{2}, -\frac{1}{2} \rangle = \sqrt{\frac{1}{6}}$$

$$\langle 1\frac{3}{2}; 0, \frac{1}{2} | \frac{1}{2}, \frac{1}{2} \rangle = \langle 1\frac{3}{2}; 0, -\frac{1}{2} | \frac{1}{2}, -\frac{1}{2} \rangle = \sqrt{\frac{1}{3}}$$

The probabilities for these various rates is therefore proportional to

$$\Gamma\left(\frac{3}{2}\frac{3}{2} \rightarrow \frac{1}{2}\frac{1}{2}\right), \Gamma\left(\frac{3}{2}, -\frac{3}{2} \rightarrow \frac{1}{2}, -\frac{1}{2}\right) \propto \frac{1}{2},$$

$$\Gamma\left(\frac{3}{2}\frac{1}{2} \rightarrow \frac{1}{2}\frac{1}{2}\right), \Gamma\left(\frac{3}{2}, -\frac{1}{2} \rightarrow \frac{1}{2}, -\frac{1}{2}\right) \propto \frac{1}{3},$$

$$\Gamma\left(\frac{3}{2}, -\frac{1}{2} \rightarrow \frac{1}{2}, \frac{1}{2}\right), \Gamma\left(\frac{3}{2}\frac{1}{2} \rightarrow \frac{1}{2}, -\frac{1}{2}\right) \propto \frac{1}{6}.$$

Or, doubling everything, another way to put this would be in a little probability table like the one at right:

$ \psi_I\rangle =$	$ \frac{3}{2}, \frac{3}{2}\rangle$	$ \frac{3}{2}, \frac{1}{2}\rangle$	$ \frac{3}{2}, -\frac{1}{2}\rangle$	$ \frac{3}{2}, -\frac{3}{2}\rangle$
$ \psi_F\rangle = \frac{1}{2}, \frac{1}{2}\rangle$	1	$\frac{2}{3}$	$\frac{1}{3}$	0
$ \psi_F\rangle = \frac{1}{2}, -\frac{1}{2}\rangle$	0	$\frac{1}{3}$	$\frac{2}{3}$	1