Physics 745 - Group Theory Solution Set 27

- 1. [5] This problem has to do with demonstrating that according the isospin symmetry, the three $|\Sigma\rangle$'s all have the same mass.
 - (a) [2] Work out the effects of the isospin operators \mathcal{I}_{\pm} on all three of the $|\Sigma\rangle$ states.

These are in an isospin 1 state, so we have

$$\begin{split} \mathcal{I}_{_{+}}\left|\Sigma^{_{-}}\right\rangle &= \sqrt{2}\left|\Sigma^{_{0}}\right\rangle, \quad \mathcal{I}_{_{+}}\left|\Sigma^{_{0}}\right\rangle = \sqrt{2}\left|\Sigma^{_{+}}\right\rangle, \quad \mathcal{I}_{_{+}}\left|\Sigma^{_{+}}\right\rangle = 0, \\ \mathcal{I}_{_{-}}\left|\Sigma^{_{+}}\right\rangle &= \sqrt{2}\left|\Sigma^{_{0}}\right\rangle, \quad \mathcal{I}_{_{-}}\left|\Sigma^{_{0}}\right\rangle = \sqrt{2}\left|\Sigma^{_{-}}\right\rangle, \quad \mathcal{I}_{_{-}}\left|\Sigma^{_{-}}\right\rangle = 0. \end{split}$$

(b) [3] Assuming that isospin commutes with the mass portion of the Hamiltonian, show that all three of the $|\Sigma\rangle$'s have the same mass.

$$\begin{split} \left\langle \Sigma^{+} \left| H \right| \Sigma^{+} \right\rangle &= \frac{1}{\sqrt{2}} \left\langle \Sigma^{+} \left| H \mathcal{I}_{+} \right| \Sigma^{0} \right\rangle = \frac{1}{\sqrt{2}} \left\langle \Sigma^{+} \left| \mathcal{I}_{+} H \right| \Sigma^{0} \right\rangle = \left\langle \Sigma^{0} \left| H \right| \Sigma^{0} \right\rangle \\ &= \frac{1}{\sqrt{2}} \left\langle \Sigma^{0} \left| H \mathcal{I}_{+} \right| \Sigma^{-} \right\rangle = \frac{1}{\sqrt{2}} \left\langle \Sigma^{0} \left| \mathcal{I}_{+} H \right| \Sigma^{-} \right\rangle = \left\langle \Sigma^{-} \left| H \right| \Sigma^{-} \right\rangle, \\ &m_{\Sigma^{+}} = m_{\Sigma^{0}} = m_{\Sigma^{-}}. \end{split}$$

- 2. [10] There is an excited set of two baryons $|N^+\rangle$ and $|N^0\rangle$ that have the same isospin properties of the neutron and proton. They can decay, in principle, into a $|\Delta \pi\rangle$ combination of two particles.
 - (a) [4] Suppose that the Hamiltonian that performs this transition takes the form

$$H\left|N^{+}\right\rangle = a\left|\Delta^{++};\pi^{-}\right\rangle + b\left|\Delta^{+};\pi^{0}\right\rangle + c\left|\Delta^{0};\pi^{+}\right\rangle$$

Find the relative sizes of the factors a, b, and c. I recommend doing this by letting \mathcal{I}_+ act on both sides.

Following the suggestion, we have

$$\begin{split} \mathcal{I}_{+}H\left|N^{+}\right\rangle &= a\mathcal{I}_{+}\left|\Delta^{++};\pi^{-}\right\rangle + b\mathcal{I}_{+}\left|\Delta^{+};\pi^{0}\right\rangle + c\mathcal{I}_{+}\left|\Delta^{0};\pi^{+}\right\rangle,\\ 0 &= \sqrt{2}a\left|\Delta^{++};\pi^{0}\right\rangle + \sqrt{3}b\left|\Delta^{++};\pi^{0}\right\rangle + \sqrt{2}b\left|\Delta^{+};\pi^{+}\right\rangle + 2c\left|\Delta^{+};\pi^{+}\right\rangle,\\ 0 &= \left(\sqrt{2}a + \sqrt{3}b\right)\left|\Delta^{++};\pi^{0}\right\rangle + \left(\sqrt{2}b + 2c\right)\left|\Delta^{+};\pi^{+}\right\rangle \end{split}$$

The coefficients of each of these states must vanish, so we find

$$a = -\sqrt{\frac{3}{2}}b$$
 and $c = -\sqrt{\frac{1}{2}}b$

so that we have

$$H\left|N^{+}\right\rangle = -\sqrt{\frac{3}{2}}b\left|\Delta^{++};\pi^{-}\right\rangle + b\left|\Delta^{+};\pi^{0}\right\rangle - b\sqrt{\frac{1}{2}}\left|\Delta^{0};\pi^{+}\right\rangle$$

(b) [2] Calculate the relative rate for the decay rates $\Gamma(N^+ \to \Delta^{++}\pi^-)$,

$$\Gamma(N^+ \to \Delta^+ \pi^0)$$
, and $\Gamma(N^+ \to \Delta^0 \pi^+)$.

As usual, we simply have

$$\Gamma\left(N^{+} \to \Delta^{++} \pi^{-}\right) \propto \frac{3}{2} b^{2}, \quad \Gamma\left(N^{+} \to \Delta^{+} \pi^{0}\right) \propto b^{2}, \quad \Gamma\left(N^{+} \to \Delta^{0} \pi^{+}\right) \propto \frac{1}{2} b^{2}$$

So they are in the ratio 3:2:1.

(c) [4] Now I want you to figure out how the N^0 decays. Using isospin symmetry, determine that the same interaction discussed in part (a) also leads to three decay processes for the N^0 . Find the relative probability amplitudes for these processes, and predict the corresponding decay rates, and show how they relate to those found in part (b).

We start by letting \mathcal{I}_{-} act on both sides, to yield

$$\begin{split} \mathcal{I}_{-}H\left|N^{+}\right\rangle &= -\sqrt{\frac{3}{2}}b\mathcal{I}_{-}\left|\Delta^{++};\pi^{-}\right\rangle + b\mathcal{I}_{-}\left|\Delta^{+};\pi^{0}\right\rangle - b\sqrt{\frac{1}{2}}\mathcal{I}_{-}\left|\Delta^{0};\pi^{+}\right\rangle,\\ H\left|N^{0}\right\rangle &= -\sqrt{\frac{3}{2}}\sqrt{3}b\left|\Delta^{+};\pi^{-}\right\rangle + \sqrt{2}b\left|\Delta^{+};\pi^{-}\right\rangle + 2b\left|\Delta^{0};\pi^{0}\right\rangle - b\sqrt{\frac{1}{2}}\sqrt{2}\left|\Delta^{0};\pi^{0}\right\rangle\\ &- b\sqrt{\frac{1}{2}}\sqrt{3}\left|\Delta^{-};\pi^{+}\right\rangle\\ &= -\sqrt{\frac{1}{2}}b\left|\Delta^{+};\pi^{-}\right\rangle + b\left|\Delta^{0};\pi^{0}\right\rangle - \sqrt{\frac{3}{2}}b\left|\Delta^{-};\pi^{+}\right\rangle \end{split}$$

and the rates come out to

$$\Gamma(N^0 \to \Delta^+ \pi^-) \propto \frac{1}{2} b^2, \quad \Gamma(N^0 \to \Delta^0 \pi^0) \propto b^2, \quad \Gamma(N^0 \to \Delta^- \pi^+) \propto \frac{3}{2} b^2$$

The ratios are 1:2:3 this time, or putting it all together, we have

$$\frac{\Gamma\left(N^{+} \to \Delta^{++} \pi^{-}\right)}{3} = \frac{\Gamma\left(N^{+} \to \Delta^{+} \pi^{0}\right)}{2} = \Gamma\left(N^{+} \to \Delta^{0} \pi^{+}\right)$$
$$= \Gamma\left(N^{0} \to \Delta^{+} \pi^{-}\right) = \frac{\Gamma\left(N^{0} \to \Delta^{0} \pi^{0}\right)}{2} = \frac{\Gamma\left(N^{0} \to \Delta^{-} \pi^{+}\right)}{3}$$