

## Lecture #7: Solving PDEs using Fourier Transforms

Example:

$$u_t + cu_x = 0$$

$$u(x, 0) = u_0(x)$$

Take Fourier transform of both sides

$$\hat{u}_t = ikc\hat{u}$$

$$\hat{u}(k, 0) = \hat{u}_0(k)$$

$$\Rightarrow \hat{u}(k, t) = \hat{u}_0(k)e^{ikct}$$

$$\begin{aligned}\Rightarrow u(x, t) &= \frac{1}{2\pi} \int_{-\infty}^{\infty} \hat{u}_0(k) e^{ikct} e^{-ikx} dk \\ &= \frac{1}{2\pi} \int_{-\infty}^{\infty} \hat{u}_0(k) e^{-ik(x-ct)} dk \\ &= u_0(x-ct).\end{aligned}$$

Example:

$$u_t = D u_{xx}$$

$$u(x, 0) = f(x)$$

Take Fourier transform of both sides

$$\hat{u}_t = -k^2 \hat{u}$$

$$\hat{u}(k, 0) = \hat{f}(k)$$

$$\Rightarrow \hat{u}(k, t) = \hat{f}(k) e^{-k^2 t}$$

$$\Rightarrow u(x, t) = f(x) * \mathcal{F}^{-1}[e^{-k^2 t}](x)$$

Now,

$$\begin{aligned} \mathcal{F}^{-1}[e^{-k^2 t}] &= \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{-k^2 t} e^{-ikx} dk \\ &= \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{-u^2 t} e^{iux} du \\ &= \frac{1}{2\pi} \sqrt{\frac{\pi}{t}} e^{-x^2/4t} \\ &= \frac{1}{\sqrt{4\pi t}} e^{-x^2/4t} \end{aligned}$$

Consequently, a solution is given by:

$$v(x, t) = f(x) * \frac{e^{-x^2/4t}}{\sqrt{4\pi t}}$$

$$\Rightarrow v(x, t) = \int_{-\infty}^{\infty} f(y) \frac{1}{\sqrt{4\pi t}} e^{-(x-y)^2/4t} dy$$

The function

$$G(x, t) = \frac{1}{\sqrt{4\pi t}} e^{-x^2/4t}$$

is called

1. The heat kernel
2. Green's function
3. Fundamental solution.

### Properties of Heat Equation:

$$u_t = D u_{xx}, \quad u \in \mathcal{X}$$

$$u(x, 0) = f(x)$$

1.  $\int_{-\infty}^{\infty} u(x, t) dx = \int_{-\infty}^{\infty} f(x) dx$

proof:

$$\frac{d}{dt} \int_{-\infty}^{\infty} u(x, t) dx = \int_{-\infty}^{\infty} u_t dx = \int_{-\infty}^{\infty} u_{xx} dx = u_x(\infty) - u_x(-\infty) = 0.$$

Therefore,  $\int_{-\infty}^{\infty} u(x, t) dx$  is constant and thus

$$\int_{-\infty}^{\infty} u(x, t) dx = \int_{-\infty}^{\infty} u(x, 0) dx = \int_{-\infty}^{\infty} f(x) dx.$$

2. Let  $E(t) = \int_{-\infty}^{\infty} u(x, t)^2 dx$ . Then  $\frac{d}{dt} E(t) \leq 0$ .

proof:

$$\frac{dE}{dt} = \int_{-\infty}^{\infty} 2u u_t dx = \int_{-\infty}^{\infty} 2u u_{xx} dx = 2u u_x \Big|_{-\infty}^{\infty} - \int_{-\infty}^{\infty} 2u_x^2 dx$$

$$\Rightarrow \frac{dE}{dt} = -2 \int_{-\infty}^{\infty} u_x^2 dx \leq 0.$$

3. Solutions are unique

proof:

Let  $u, v$  be solutions. By linearity, if we let  $w = u - v$  then  $w$  satisfies

$$w_t = D w_{xx}$$

$$w(x, 0) = 0$$

If we let  $E = \int_{-\infty}^{\infty} w^2 dx$ , then

1.  $E(0) = 0$

2.  $E(t) \geq 0$

3.  $\frac{dE}{dt} \leq 0$

Consequently,  $E = 0$  which implies

$$0 = \int_{-\infty}^{\infty} w^2 dx \Rightarrow w = 0.$$

4. If  $|f(x)| \leq M$  then  $|v(x,t)| \leq M$ .

proof:

$$|v(x,t)| = \left| \int_{-\infty}^{\infty} f(y) e^{-(x-y)^2/4\pi Dt} \cdot \frac{1}{\sqrt{4\pi Dt}} dy \right|$$

$$\leq \int_{-\infty}^{\infty} |f(y)| e^{-(x-y)^2/4\pi Dt} \cdot \frac{1}{\sqrt{4\pi Dt}} dy$$

$$\leq M \int_{-\infty}^{\infty} \frac{1}{\sqrt{4\pi Dt}} e^{-(x-y)^2/4\pi Dt} dy$$

$$= M.$$

Example:

Solve

$$v_t + cv_x = Dv_{xx}, \quad v(x,0) = f(x)$$

where  $c \in \mathbb{R}$ ,  $D > 0$ . Take Fourier transform of both sides

$$\hat{v}_t - ick\hat{v} = -k^2 D\hat{v}$$

$$\hat{v}(k,0) = \hat{f}(k)$$

$$\Rightarrow \hat{v}_t = (ick - k^2 D)\hat{v}$$

$$\Rightarrow \hat{v} = \hat{f}(k) e^{(ick - k^2 D)t}$$

Therefore,

$$v = \mathcal{F}^{-1}[\hat{f}(k) e^{ick t}] * \mathcal{F}^{-1}[e^{-k^2 D t}]$$

$$= f(x - ct) * e^{-k^2 D t}$$

Thus

$$v(x,t) = \frac{1}{\sqrt{4\pi Dt}} \int_{-\infty}^{\infty} f(y - ct) e^{-(x-y)^2/4Dt} dy.$$