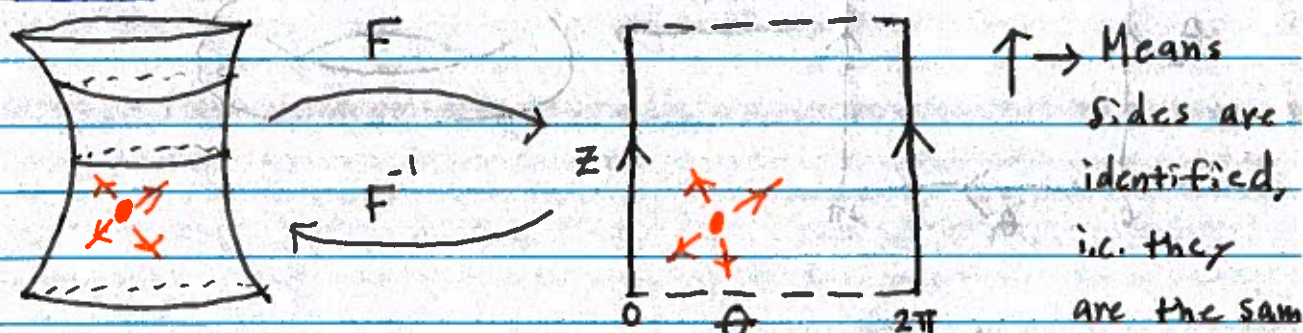


Lecture #1: Motion on Tori and cylinders

Manifolds

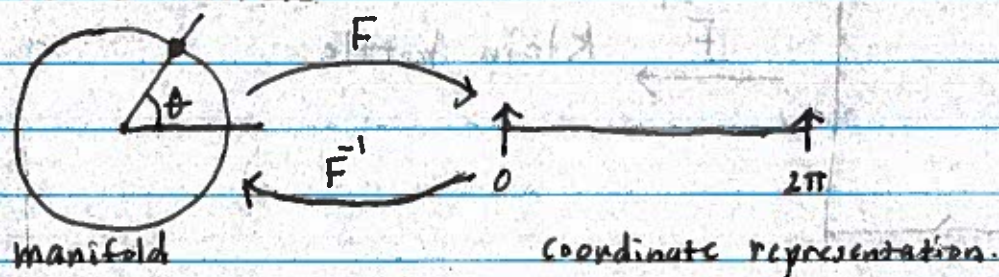


M = hyperboloid with a vector field.

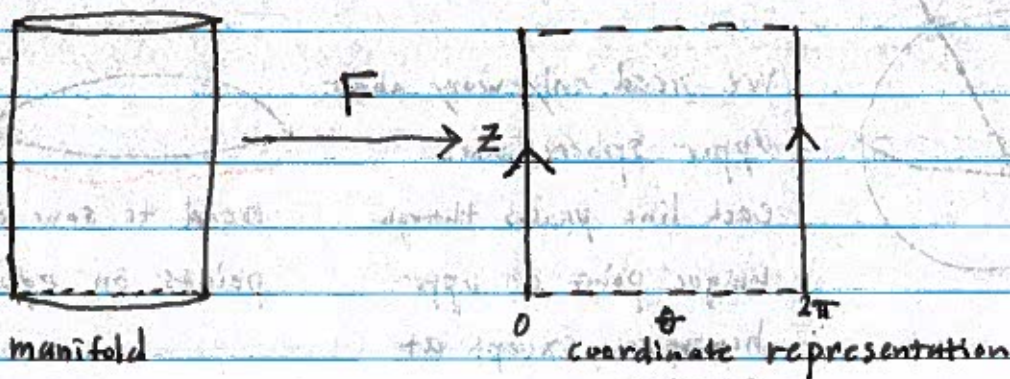
How do we describe flow on M ? Key idea is to create maps F which are diffeomorphisms into $\mathbb{R}^n$ on which calculus can be done. If this is possible, we call M an n -dimensional manifold and the collection of all maps is an atlas.

Examples:

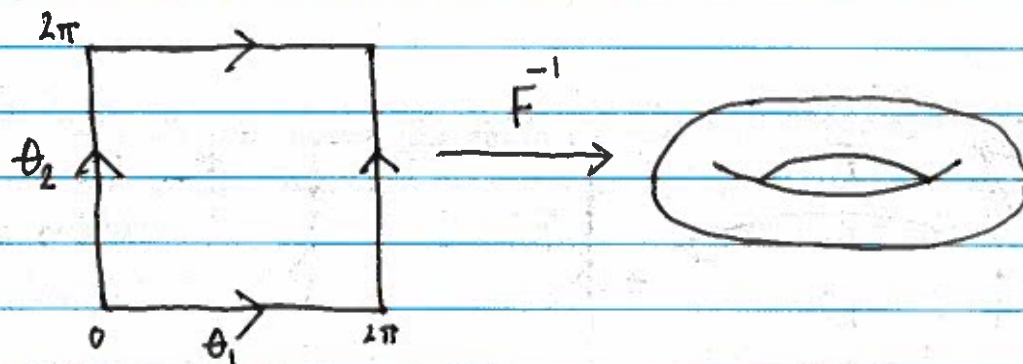
1. $M = S^1 = \text{circle}$



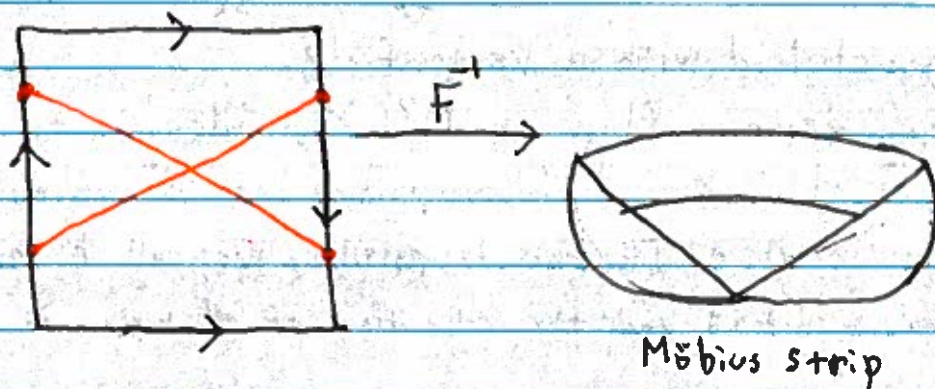
2. $M = S^1 \times \mathbb{R} = \text{cylinder}$



3. $M = S^1 \times S^1 = \text{torus}$.



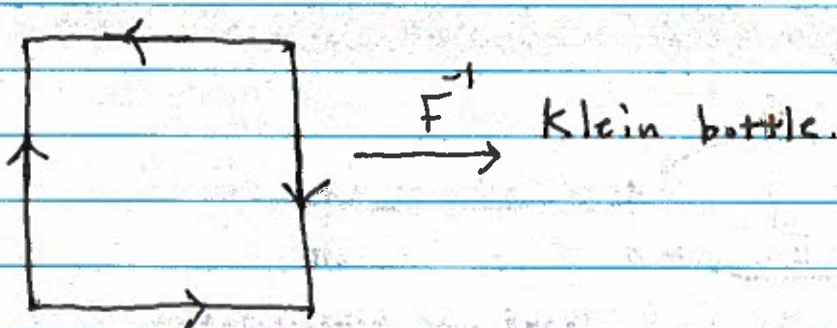
4.



Möbius strip

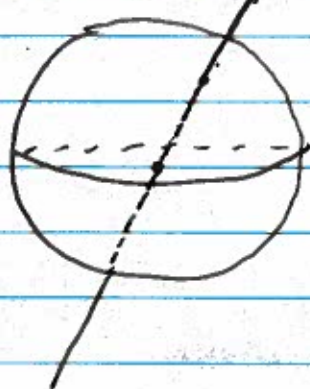
Dynamics difficult to interpret

5.

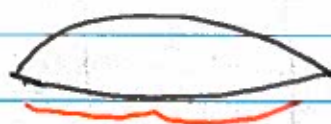


Klein bottle.

6. $\mathbb{R}P^2 = \{\text{set of lines through origin in } \mathbb{R}^3\}$.



We need only worry about upper sphere since each line passes through unique point on upper hemisphere, except at equator



need to show antipodal points on equator.

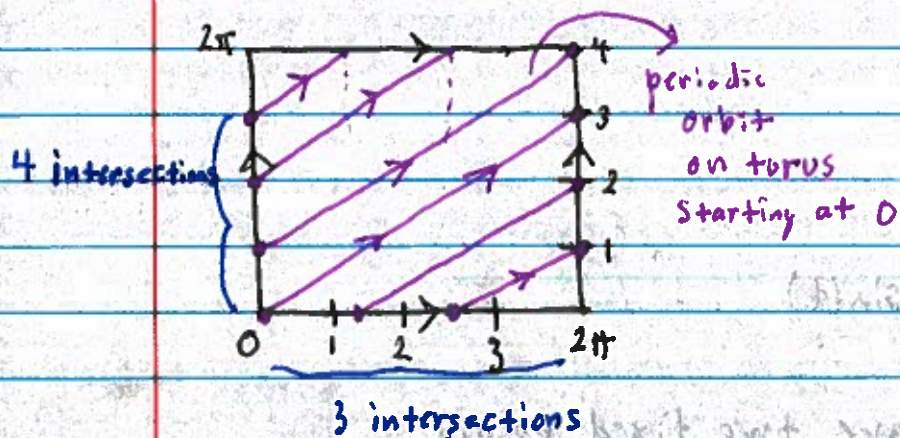
Motion on a torus

Example:

$$\dot{\theta}_1 = 4, \quad (\theta_1, \theta_2) \in S^1 \times S^1$$

$$\dot{\theta}_2 = 3$$

$$\frac{d\theta_2}{d\theta_1} = \frac{3}{4} \Rightarrow \theta_2(\theta_1) \text{ is a line of slope } \frac{3}{4}$$



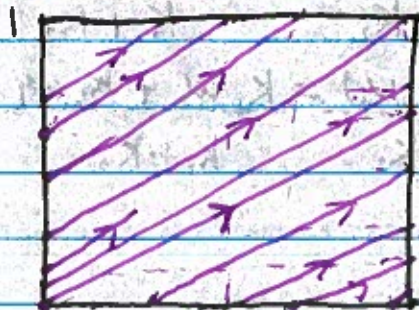
3 intersections

This is called a 3:4 resonant tori. The winding number is $\frac{3}{4}$.

Example:

$$\dot{\theta}_1 = \sqrt{2} \Rightarrow \frac{d\theta_2}{d\theta_1} = \frac{1}{\sqrt{2}}$$

$$\dot{\theta}_2 = 1$$



This is called a non-resonant tori. The orbit is called quasiperiodic.

Example:

$$\dot{\theta}_1 = \omega_1 + K_1 \sin(\theta_2 - \theta_1), \quad (\theta_1, \theta_2) \in S^1 \times S^1$$

$$\dot{\theta}_2 = \omega_2 + K_2 \sin(\theta_1 - \theta_2)$$

two joggers on a track
trying to match speed

$\omega_i \sim$ natural speed/frequency

$K_i \sim$ coupling strength

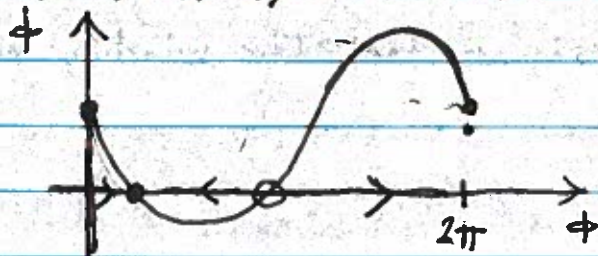
Let $\phi = \theta_2 - \theta_1$:

$$\dot{\phi} = \dot{\theta}_2 - \dot{\theta}_1$$

$$\Rightarrow \dot{\phi} = \omega_2 - \omega_1 - (K_1 + K_2) \sin(\phi)$$

$$\Rightarrow \dot{\phi} = \Omega - K \sin(\phi),$$

If $|\Omega/K| < 1$, we have two fixed points



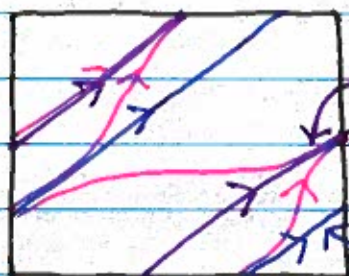
When $\sin(\phi) = \Omega/K$ we have:

$$\dot{\theta}_1 = \omega_1 + K_1 \left(\frac{\Omega}{K} \right) = \omega_1 + \frac{K_2 (\omega_2 - \omega_1)}{K_1 + K_2} = \frac{K_2 \omega_1 + K_1 \omega_2}{K_1 + K_2}$$

$$\dot{\theta}_2 = \omega_2 + K_2 \left(\frac{-\Omega}{K} \right) = \omega_2 - \frac{K_1 (\omega_2 - \omega_1)}{K_1 + K_2} = \frac{K_1 \omega_2 + K_2 \omega_1}{K_1 + K_2}$$

Therefore, when $\sin(\phi) = \Omega/K$

$$\frac{d\theta_2}{d\theta_1} = \frac{\dot{\theta}_2}{\dot{\theta}_1} = 1$$



attracting
limit cycle

unstable limit cycle.

Motion on a Cylinder

$$\dot{\phi} = y, \quad (\phi, y) \in S^1 \times \mathbb{R}, \quad I, \alpha > 0.$$
$$\dot{y} = I - \sin(\phi) - \alpha y$$

Nullclines satisfy:

$$\dot{\phi} = 0 \Rightarrow y = 0$$

$$\dot{y} = 0 \Rightarrow y = \frac{1}{\alpha}(I - \sin(\phi))$$

Equivalent terms

$$\ddot{\phi} = I - \sin(\phi) - \alpha \dot{\phi}$$

Damped oscillator.

Case 1 ($I < 1$):

Fixed point when $y = 0$ and $\sin(\phi) = I$

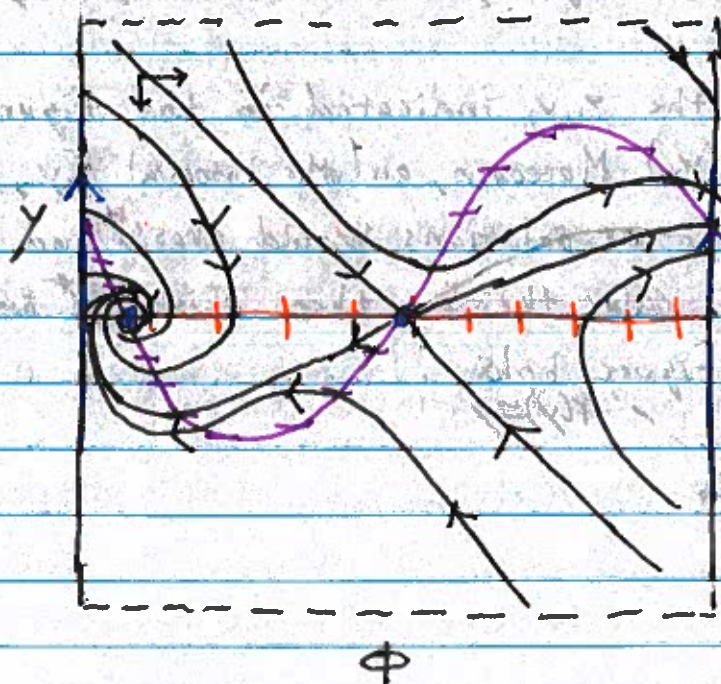
$$\Rightarrow (\phi^*, y^*) = (\sin^{-1}(I), 0), (\pi - \sin^{-1}(I), 0).$$

Jacobian is

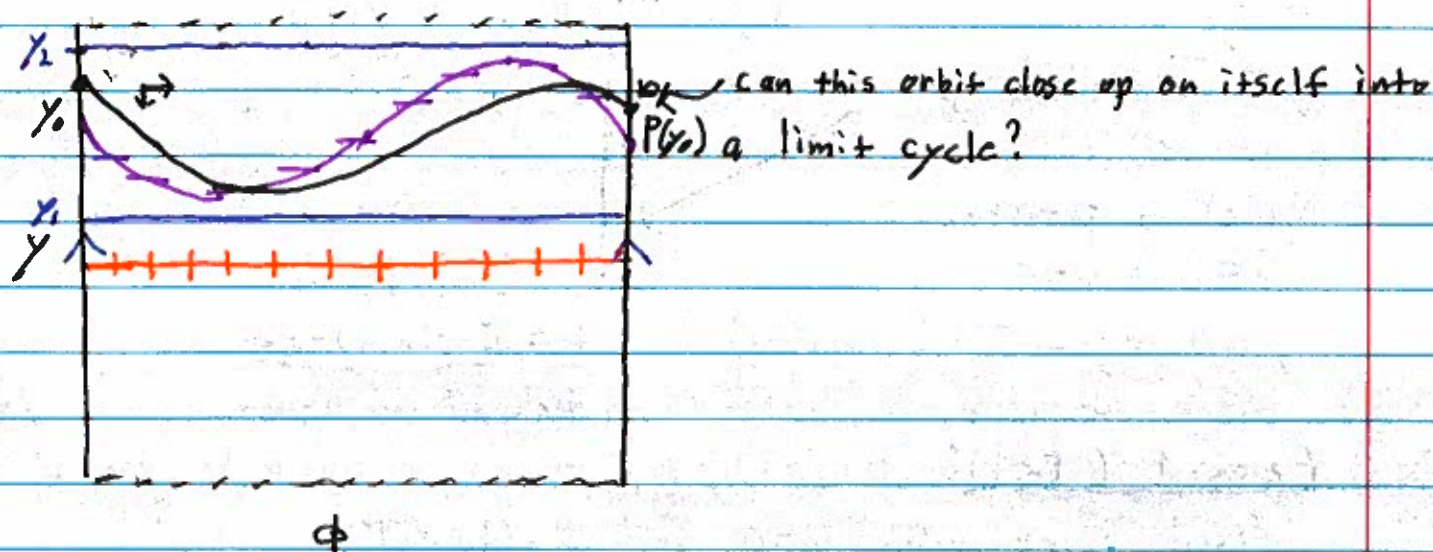
$$J = \begin{bmatrix} 0 & 1 \\ -\cos(\phi) & -\alpha \end{bmatrix}$$

$$J(\sin^{-1}(I), 0) = \begin{bmatrix} 0 & 1 \\ -\sqrt{1-I^2} & -\alpha \end{bmatrix} \Rightarrow \text{stable node/stable spiral}$$

$$J(\pi - \sin^{-1}(I), 0) = \begin{bmatrix} 0 & 1 \\ \sqrt{1-I^2} & -\alpha \end{bmatrix} \rightarrow \text{saddle}$$



Case II' ($I > 1$):



Define $P: \mathbb{R} \rightarrow \mathbb{R}$, called a Poincaré map, to be the value $P(y)$ after one loop around the cylinder for a trajectory starting at $(0, y)$.

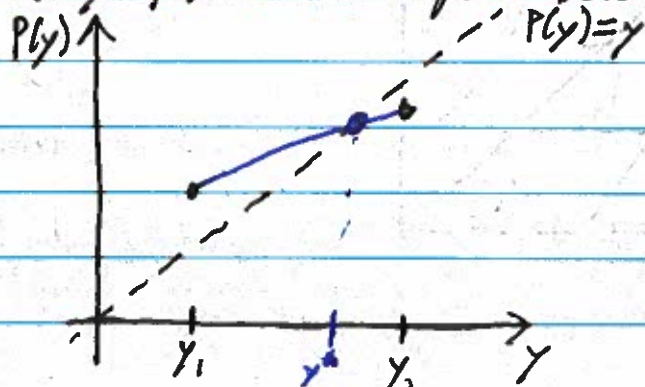
Key idea:

To show that there exists a limit cycle we need to find y^* such that $P(y^*) = y^*$. Such a y^* is called a fixed point.

Theorem- There exists a unique y^* such that $P(y^*) = y^*$.

proof:

By construction, for the y_1, y_2 indicated in the figure above $y_1 < P(y_1)$ and $y_2 > P(y_2)$. Moreover, on the interval $[y_1, y_2]$ $P(y)$ must be monotone or solutions would cross and thus by the intermediate value theorem there exists y^* such that $P(y^*) = y^*$; see the figure below:



Now, suppose there are two limit cycles $y_+(t) > y_-(t)$. Define

$$E[y(\phi)] = \frac{1}{2}y^2 - \cos(\phi)$$

Along any periodic orbit we have

$$0 = E[y_+(2\pi)] - E[y_+(0)] = E[y_-(2\pi)] - E[y_-(0)]$$

$$\Rightarrow 0 = \int_0^{2\pi} \frac{dE}{d\phi} \bigg|_{y_+} d\phi = \int_0^{2\pi} \frac{dE}{d\phi} \bigg|_{y_-} d\phi$$

$$\text{But, } \frac{dE}{d\phi} = y \frac{dy}{d\phi} + \sin(\phi),$$

$$= y \left(\frac{1 - \sin(\phi) - \alpha y}{y} \right) + \sin(\phi),$$

$$= 1 - \sin(\phi) - \alpha y + \sin(\phi),$$

$$= 1 - \alpha y.$$

Therefore,

$$\int_0^{2\pi} (1 - \alpha y_+(\phi)) d\phi = \int_0^{2\pi} (1 - \alpha y_-(\phi)) d\phi,$$

$$\Rightarrow \int_0^{2\pi} y_+(\phi) d\phi = \int_0^{2\pi} y_-(\phi) d\phi,$$

which is a contradiction.

Stability when $I > 1$:

A Poincaré map allows us to think of the dynamical system as an iterated map

$$y_{n+1} = P(y_n), \quad n \in \{0, 1, 2, \dots\}$$

A value y^* is a fixed point of the Poincaré map if

$$P(y^*) = y^*.$$

A fixed point y^* is attracting if there exists $\delta > 0$ such that $|y_0 - y^*| < \delta$.

Theorem - If y^* is a fixed point of P and $P'(y^*) = \lambda$ satisfies $|\lambda| < 1$ then y^* is attracting. The value λ is called the Floquet multiplier.

proof:

Since $|\lambda| < 1$, it follows from continuity of P' that there exists $\delta > 0$ such that for all $y \in \mathcal{N}_\delta(y^*) = \{y : |y - y^*| < \delta\}$, $|P'(y)| \leq M < 1$.

Now, if we let $y_0 = y^* + \Delta y \in \mathcal{N}_\delta(y^*)$ it follows that $|\Delta y| < \delta$ and

$$\begin{aligned} |y_1 - y^*| &= |P(y_0) - y^*|, \\ &= |P(y^* + \Delta y) - y^*|, \\ &= |P(y^*) + P'(c_1)\Delta y - y^*|, \quad (\text{Taylor's Theorem with Remainder}) \\ &= |y^* + P'(c_1)\Delta y - y^*|, \\ &\leq M|\Delta y| < \Delta y, \end{aligned}$$

where $c_1 \in \mathcal{N}_\delta(y^*)$. Thus, $y_1 \in \mathcal{N}_\delta(y^*)$ and

$$\begin{aligned} |y_2 - y^*| &= |P(y_1) - y^*|, \\ &= |P(P(y_0)) - y^*|, \\ &= |P(P(y^* + \Delta y)) - y^*|, \\ &= |P(P(y^*)) + P'(P(c_2))P'(c_1)\Delta y - y^*|, \\ &= |y^* + P'(P(c_2))P'(c_1)\Delta y - y^*|, \\ &\leq M^2|\Delta y| < \Delta y, \end{aligned}$$

where $c_2 \in \mathcal{N}_\delta(y^*)$ and thus $y_2 \in \mathcal{N}_\delta(y^*)$. Continuing inductively we have

$$\begin{aligned} \bullet |y_n - y^*| &< M^n \Delta y \\ \bullet y_n &\in \mathcal{N}_\delta(y^*). \end{aligned}$$

Consequently,

$$\lim_{n \rightarrow \infty} |y_n - y^*| = 0 \Rightarrow \lim_{n \rightarrow \infty} y_n = y^*.$$