

# Lecture #8: Weak Van der Pol Oscillator

$$\ddot{x} + x = \epsilon(1-x^2)\dot{x}$$

## Phase Portrait Analysis

$$\dot{x} = v$$

$$\dot{v} = -x + \epsilon(1-x^2)v$$

Nullclines:

$$v = 0 \quad (\dot{x} = 0)$$

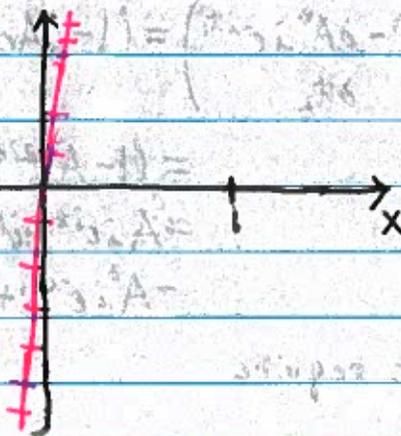
$$v = \frac{x}{\epsilon(1-x^2)} \quad (\dot{v} = 0)$$

Jacobian:

$$J(0,0) = \begin{bmatrix} 0 & 1 \\ -1 & \epsilon \end{bmatrix}$$

$$\lambda_{1,2} = \frac{\epsilon \pm \sqrt{\epsilon^2 - 4}}{2}$$

Thus, for  $\epsilon \ll 1$  the fixed point is an unstable spiral. Moreover, there is a Hopf bifurcation when  $\epsilon = 0$ .



The phase portrait is too hard to draw by hand for  $\epsilon \ll 1$ .

We can use method of multiple scales to determine if the Hopf bifurcation is sub or supercritical.

### Method of Multiple Scales

$$\ddot{x} + x = \varepsilon(1-x^2)\dot{x}$$

Assume two time scales

$$t_1 = \tau, \quad t_2 = \varepsilon t.$$

$$\rightarrow \frac{d}{dt} = \frac{\partial}{\partial t_1} + \varepsilon \frac{\partial}{\partial t_2}, \quad \frac{d^2}{dt^2} = \frac{\partial^2}{\partial t_1^2} + 2\varepsilon \frac{\partial^2}{\partial t_1 \partial t_2} + \varepsilon^2 \frac{\partial^2}{\partial t_2^2}$$

To  $\mathcal{O}(\varepsilon)$  we obtain the PDE:

$$\frac{\partial^2 x}{\partial t_1^2} + 2\varepsilon \frac{\partial^2 x}{\partial t_1 \partial t_2} + x = \varepsilon(1-x^2) \frac{\partial x}{\partial t_1}$$

We assume

$$x = x_0 + \varepsilon x_1 + \dots$$

Substituting into the PDE we obtain at  $\mathcal{O}(1)$ :

$$\frac{\partial^2 x_0}{\partial t_1^2} + x_0 = 0$$

$$\Rightarrow x_0 = A(t_2) e^{it_1} + A^*(t_2) e^{-it_1}$$

At  $\mathcal{O}(\varepsilon)$  we have:

$$\begin{aligned} \frac{\partial^2 x_1}{\partial t_1^2} + x_1 + 2 \left( \frac{\partial A}{\partial t_2} i e^{it_1} - \frac{\partial A^*}{\partial t_2} i e^{-it_1} \right) &= (1 - (A e^{it_1} + A^* e^{-it_1})) (A i e^{it_1} - A^* i e^{-it_1}) \\ &= (1 - A^2 e^{2it_1} - 2AA^* - A^{*2} e^{-2it_1}) (A i e^{it_1} - A^* i e^{-it_1}) \\ &= A i e^{it_1} - i A^3 e^{3it_1} - 2A^2 A^* e^{it_1} - i A^{*2} A e^{-it_1} \\ &\quad - A^* i e^{-it_1} + A^2 A^* i e^{it_1} + 2A A^{*3} e^{-it_1} + A^{*3} e^{-3it_1} \end{aligned}$$

To remove secular terms we require

$$2i \frac{\partial A}{\partial t_2} = A i - A^2 A^*$$

$$\Rightarrow \frac{\partial A}{\partial t_2} = \frac{1}{2} A - A^2 A^*$$

We express  $A$  in exponential form to yield:

$$A = r e^{i\theta}$$

$$\Rightarrow \frac{\partial A}{\partial t_2} = \frac{\partial r}{\partial t_2} e^{i\theta} + i r \frac{\partial \theta}{\partial t_2} e^{i\theta} = \frac{1}{2} (r e^{i\theta} - r^3 e^{i\theta})$$

$$\Rightarrow \frac{\partial r}{\partial t_2} = \frac{1}{2} (r - r^3), \quad \frac{\partial \theta}{\partial t_2} = 0$$

In the long run,  $r \rightarrow 1$  and  $\theta = \theta_0$ . Consequently, the limit cycle is approximated by:

$$\begin{aligned} x &= e^{i\theta_0} e^{it} + e^{-i\theta_0} e^{it} \\ &= 2 \cos(\theta_0 + t). \end{aligned}$$

The radius is 2 and the period is  $2\pi$ .

