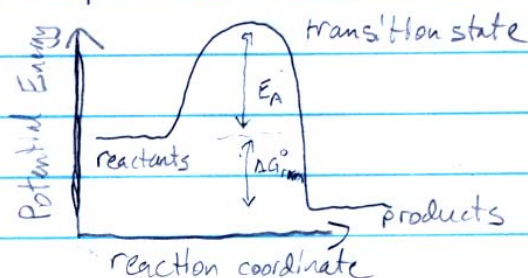


Kinetics Revisited

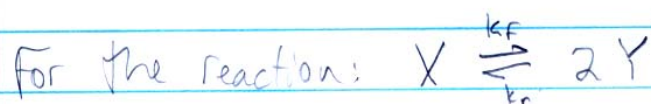
9.1

Last time we looked at kinetics we learned 2 important things:

- ① Order of reactions and the equations relating k , $t_{1/2}$, etc.
- ② how to interpret a 2-D or 3-D reaction diagram, eg:



Now we'll look at kinetics in more detail.

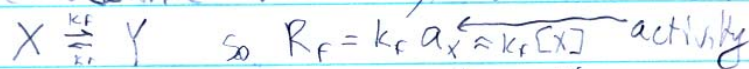


We have $\text{rate} = \frac{d\xi}{dt} = -\frac{1}{1} \frac{d[X]}{dt} = +\frac{1}{2} \frac{d[Y]}{dt}$ for 1 L

or, in general $\text{rate} = \text{reaction velocity } v = R_f - R_r = \frac{1}{V} \frac{d\xi}{dt} = \frac{1}{V} \frac{1}{\nu_i} \frac{d[Z_i]}{dt}$

where R_f and R_r are the forward and reverse rates (here they would be $k_f[X]$ and $k_r[Y]^2$) and V is volume while ν_i is the stoichiometric coefficient of component Z_i .

(Usually, we assume $V = 1$ L). Now take the reaction:



Since $[X] = [X]_0 - \xi$ $R_f = k_f [X] = k_f ([X]_0 - \xi)$

Same for $[Y]$, so we get:

$$\frac{d\xi}{dt} = R_f - R_r = k_f ([X]_0 - \xi) - k_r ([Y]_0 + \xi)$$

this is a differential eq. that we will solve...

9.2

To solve the differential eq., simplify: $\frac{d\xi}{dt} = 2k \left[\frac{X_0 - Y_0}{2} - \xi \right]$
 where $k_f = k_r = k$ and X_0 is short for $[X]_0$ this - sign is from changing $\xi \rightarrow z$
 if $K_{eq} = 1$

Now set $z = \frac{1}{2}(X_0 - Y_0) - \xi$ so we have $\frac{dz}{dt} = -2kz$

if $z(t) = e^{-2kt}$ then $\frac{dz(t)}{dt} = -2k e^{-2kt} = -2kz \checkmark$

so $\xi(t) = \frac{1}{2}(X_0 - Y_0) - z(t) = \frac{X_0 - Y_0}{2} - e^{-2kt}$ or $= \frac{X_0 - Y_0}{2} - \frac{X_0 - Y_0}{2} e^{-2kt}$

Using the second solution we can write:

$$R_f = k_f \left(X_0 - \xi \right) = k_f \left[\frac{2X_0}{2} - \left\{ \frac{(X_0 - Y_0)}{2} - \frac{(X_0 - Y_0)}{2} e^{-2kt} \right\} \right] = \frac{k}{2} (X_0 + Y_0) + \frac{k}{2} (X_0 - Y_0) e^{-2kt}$$

and:

$$R_r = k_r \left[\frac{2Y_0}{2} + \left\{ \frac{(X_0 - Y_0)}{2} - \frac{(X_0 - Y_0)}{2} e^{-2kt} \right\} \right] = \frac{k}{2} (X_0 + Y_0) - \frac{k}{2} (X_0 - Y_0) e^{-2kt}$$

still solves diff. eq. and $\xi(0) = 0$
these will cancel in $R_f - R_r$

(Internal) Entropy production per volume $= \frac{1}{V} \frac{d_i S}{dt} = \frac{1}{V} \frac{A d\xi}{T dt} \geq 0$ sum of chem. pot. $\mu_X - \mu_Y$
activity rate $= (R_f - R_r)$

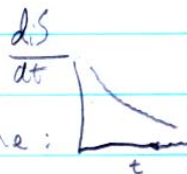
but $\mu_X = \mu_{X_0} + RT \ln a_X$, same for μ_Y

and $\mu_{X_0} - \mu_{Y_0} = -\Delta G_{rxn}^0 = RT \ln K_{eq}$, so $A = RT \ln K_{eq} + RT \ln \left(\frac{a_X}{a_Y} \right)$

but $K_{eq} = \frac{k_f}{k_r}$ so $A = RT \ln \frac{k_f a_X}{k_r a_Y} = RT \ln \frac{R_f}{R_r}$

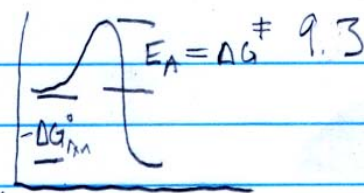
so $\frac{1}{V} \frac{d_i S}{dt} = \frac{1}{V} \frac{d\xi}{dt} \cdot \frac{RT \ln \frac{R_f}{R_r}}{T} = (R_f - R_r) R \ln \frac{R_f}{R_r}$ substituting:

$$= 2 \frac{k}{2} (X_0 - Y_0) e^{-2kt} \cdot R \ln \frac{\frac{k}{2} (X_0 + Y_0 + [X_0 - Y_0] e^{-2kt})}{\frac{k}{2} (X_0 + Y_0 - [X_0 - Y_0] e^{-2kt})}$$



This allows us to plot entropy production vs. time:

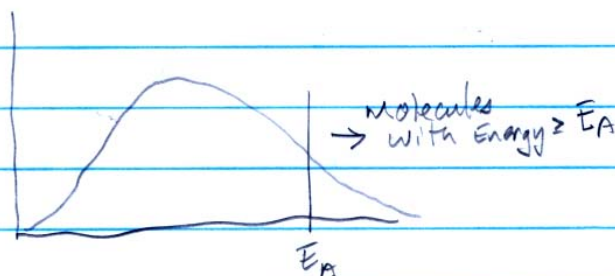
Arrhenius equation, again.



$$k_F = k_0 e^{-E_A/RT} \quad k_r = k_0 e^{-(E_A - \Delta G_{rxn}^{\circ})/RT}$$

$$K_{eq} = \frac{k_F}{k_r} = \frac{k_0 e^{-E_A/RT}}{k_0 e^{-E_A/RT} e^{\Delta G_{rxn}^{\circ}/RT}} = e^{-\Delta G_{rxn}^{\circ}/RT} \quad \checkmark$$

Not only does this eq flow from $\Delta G_{rxn}^{\circ} = -RT \ln K_{eq}$ but it also relates to the Boltzmann distribution:



What about k_0 ? Collision Theory says $k_0 \propto \left(\frac{1}{T}\right)^{1/2}$

Transition state theory says $k_0 \propto \frac{k_B T}{h} e^{-\frac{\Delta G^{\ddagger}}{RT}}$ (liquids)

TST $k_0 \propto \frac{(k_B T)^{3/2}}{h} e^{-\frac{\Delta G^{\ddagger}}{RT}}$ (gases)

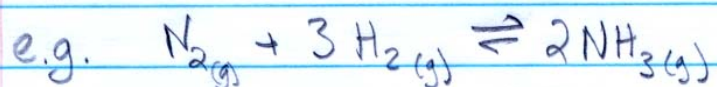
$\frac{1}{K_{eq}}$

The key point of Arrhenius is that rate increases as E_A decreases, which happens due to catalysis.

TST adds the insight that catalysts (enzymes) reduce the energy of the transition state

Thermodynamic's most important eq.

How do we use $\Delta G_{\text{rxn}}^{\circ} = -RT \ln K_{\text{eq}}$?

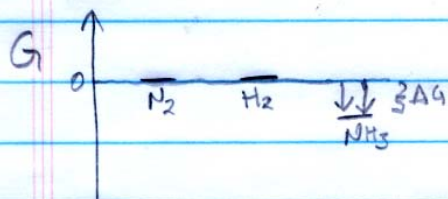


$$\Delta G_{\text{rxn}}^{\circ} = 2\Delta G_{\text{f, NH}_3(\text{g})}^{\circ} - \cancel{\Delta G_{\text{f, N}_2(\text{g})}^{\circ}} - 3\cancel{\Delta G_{\text{f, H}_2(\text{g})}^{\circ}} = -32.90 \text{ kJ/mol}$$

-16.45 kJ/mol

$$\text{@ } 25^{\circ}\text{C} \quad K_{\text{eq}} = e^{\frac{32,900}{8.34(298.15)}} = 5.80 \times 10^5 = K^{\circ} = K_p (\text{atm}^{-2})$$

The process of finding $\Delta G_{\text{rxn}}^{\circ}$ is like finding $\Delta H_{\text{rxn}}^{\circ}$:



but $\Delta G_{\text{rxn}}^{\circ}$ gives K_{eq} so it's more useful than $\Delta H_{\text{rxn}}^{\circ}$

Also, the 2nd most important eq, $\Delta G = \Delta H - T\Delta S$, tells us that $\Delta G_{\text{rxn}}^{\circ}$ and $\Delta H_{\text{rxn}}^{\circ}$ are related to $\Delta S_{\text{rxn}}^{\circ}$:

$$\Delta G_{\text{rxn}}^{\circ} = \Delta H_{\text{rxn}}^{\circ} - T\Delta S_{\text{rxn}}^{\circ}$$

total entropy of universe
entropy released to surroundings
entropy of system
depends on T

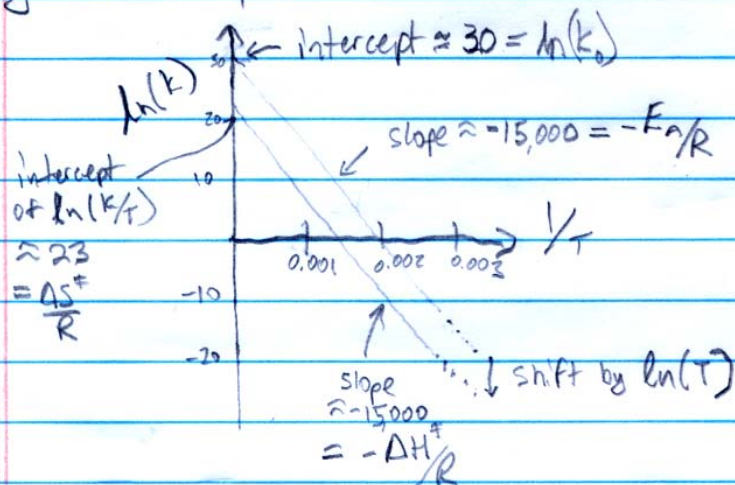
} All 3 terms are T-independent

Take home message: The reason these 2 eqns are so important is because they tell us whether products are favored in any reaction based on constants that can be found in a table, and they tell us why products are favored (heat released, internal entropy or both).

Van't Hoff and Arrhenius plots

Problem 9.3 (due next week) illustrates how these two plots are used:

given temperatures and rate constants plot:



this plot shows an Arrhenius plot gives E_a and k_0 , but a Van't Hoff plot gives ΔH and ΔS , or in this case $\Delta H^\ddagger$ and $\Delta S^\ddagger$

(* we plotted $\ln(\frac{k}{T})$ instead of $\ln(k_{eq})$, but the result of TST is $k \propto (\frac{k_B T}{h}) e^{-\Delta G^\ddagger / RT}$ or $k/T \propto (\frac{k_B}{h}) K_{eq}^\ddagger$ so we were really plotting $\ln(\frac{k_B}{h} K_{eq}^\ddagger)$, which has the form of a van't Hoff plot)

9.6

final note about activity, fugacity, etc

the definition: $a_i = \gamma_i X_i$ would indicate that the activity of an ideal gas or solution ($\gamma_i = 1$) would range from $0 \rightarrow 1$.

However, there are 2 other "definitions":

$$a_i = \gamma_i P_i / P^\circ \leftarrow \begin{array}{l} \text{partial pressure} \\ P^\circ \leftarrow 1 \text{ bar} \end{array}$$

and

$$a_i = \gamma_i C_i / C^\circ \leftarrow \begin{array}{l} \text{concentration} \\ C^\circ \leftarrow 1 \text{ M or } 1 \text{ m} \end{array}$$

In these cases, a_i can range from $0 \rightarrow \infty$ (or $0 \rightarrow 55 \text{ M}$, e.g.)

$\leftarrow$ fugacity coeff. ($\Phi_i = 1$ for $P_i \rightarrow 0$)

Fugacity, $F_i = \gamma_i \Phi_i P_i \leftarrow$ partial pressure

So, fugacity has units of pressure and corrects for deviations due to mixing (γ_i) and pressure (Φ_i)

We can separate these two corrections, e.g.,

$$\mu_i(T, P, \text{mix}) = \mu_i(T, P^\circ, \text{pure}) + RT \ln F_{i(\text{pure})} + RT \ln a_i$$

in general, we correct for pressure using $F_{i(\text{pure})} = \Phi_i P_i$
and mixing using $a_i = \gamma_i X_i$

We correct for T using
$$\frac{\mu_i(T)}{T} = \frac{\mu_i(T_0)}{T_0} + \int_{T_0}^T \frac{-\bar{H}_i(T')}{T'^2} dT'$$

where T' is a place-holder for the integral from $T_0 \rightarrow T$