

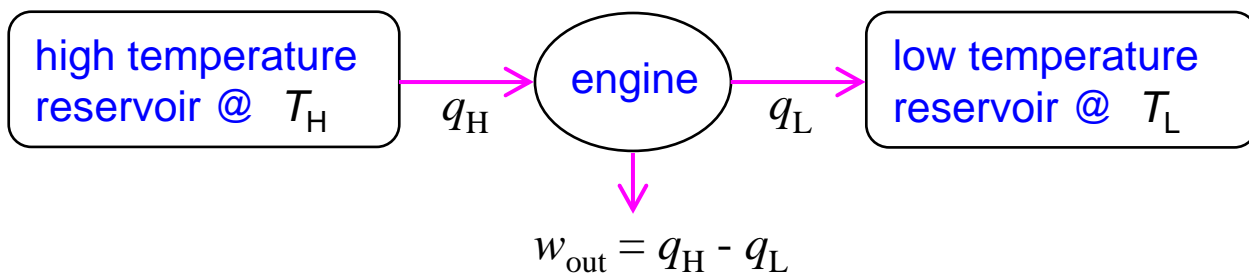
# Heat Engines

A **heat engine** is a system capable of transforming heat into \_\_\_\_\_ by some cyclic process.

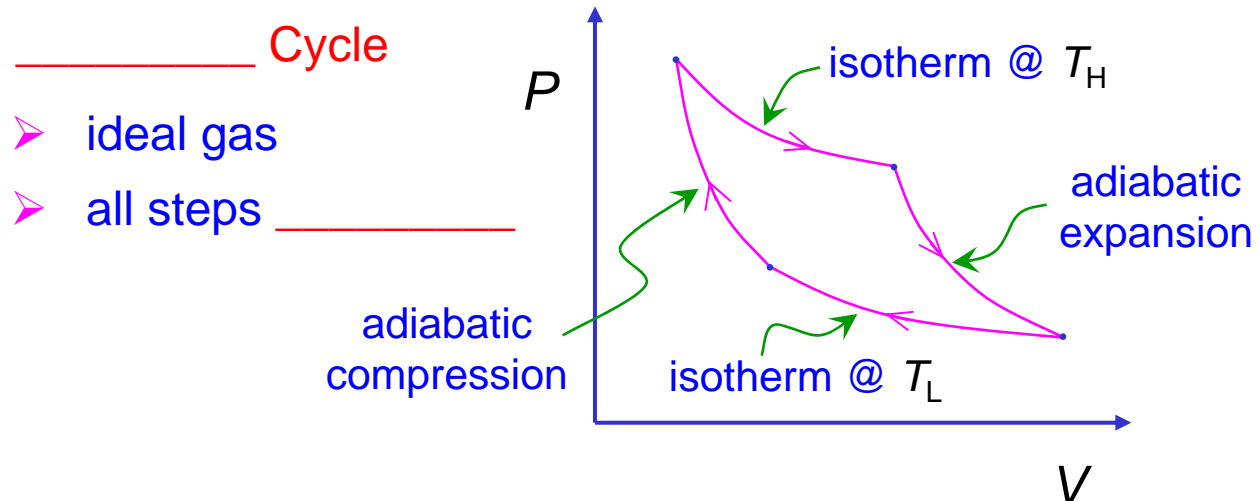
We will see that an \_\_\_\_\_ cyclic process can not produce net work. (2<sup>nd</sup> Law of Thermodynamics)

The \_\_\_\_\_ of a heat engine is defined as the ratio of the work produced to the heat \_\_\_\_\_ :

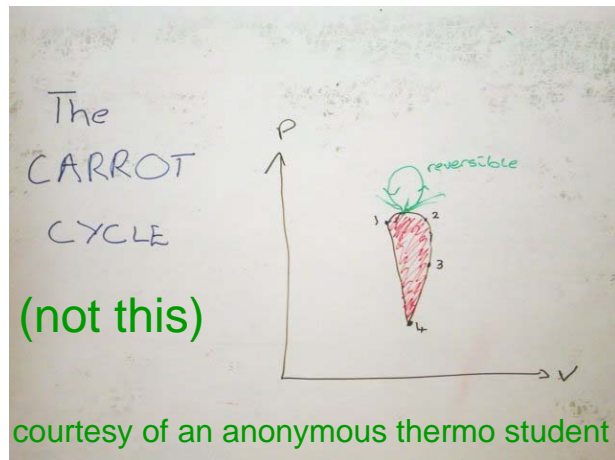
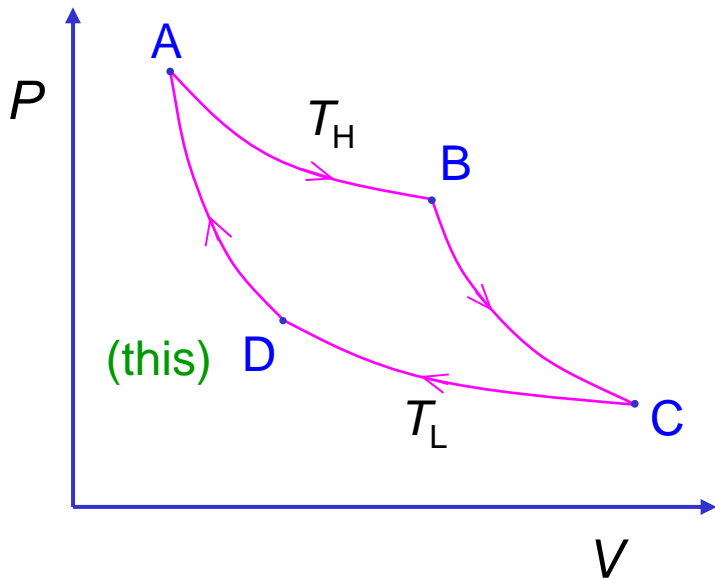
$$\varepsilon = \frac{w}{q_H} = \frac{q_H - q_L}{q_H} = \boxed{\phantom{0000}}$$



A heat \_\_\_\_\_ is a heat engine in reverse. Work is needed to transfer heat from a lower to a higher temperature reservoir.



# The Carnot Cycle



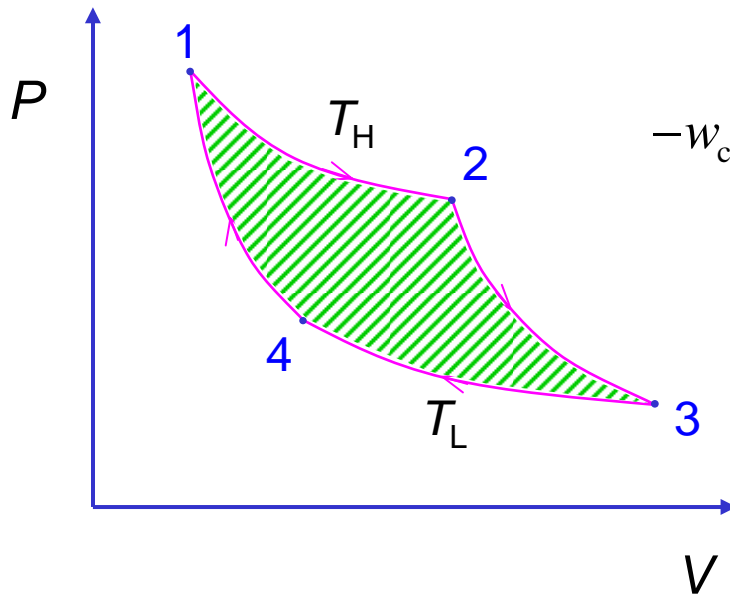
step	$w$	$q$	$\Delta U$
$A \rightarrow B$	$-nRT_H \ln(\text{_____})$	$-w_{AB}$	0
$B \rightarrow C$	$\Delta U_{BC}$	0	$-C_V(T_H - T_L)$
$C \rightarrow D$	$-nRT_L \ln(\text{_____})$	$-w_{CD}$	0
$D \rightarrow A$	$\Delta U_{DA}$	0	$C_V(T_H - T_L)$
<b>Total</b>	$nR(T_H - T_L) \ln(\text{_____})$	$-w_{\text{cyc}}$	0

$$\varepsilon = \frac{q_{\text{in}}}{q_{\text{out}}} = \frac{nR(T_H - T_L) \ln(V_A/V_B)}{nR(T_H - T_L)} = 1 - \frac{T_L}{T_H}$$

$$\varepsilon = \frac{(T_H - T_L)}{T_H}$$

for best efficiency,  
\_\_\_\_\_  $T_H$   
\_\_\_\_\_  $T_L$

## Changes in the Carnot Cycle



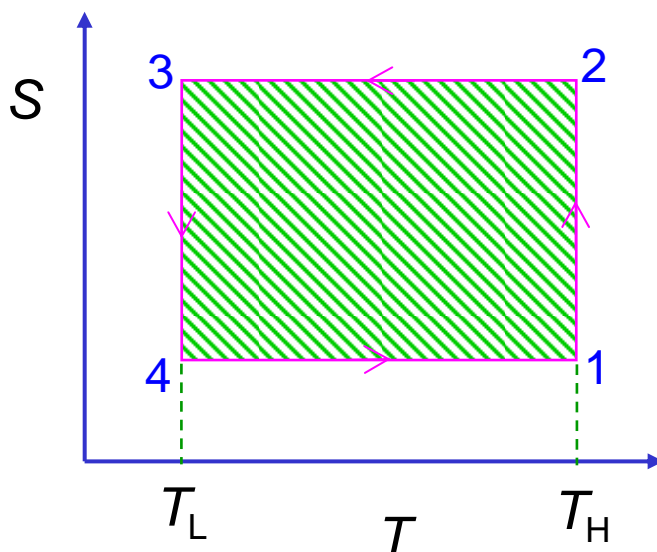
$$\begin{aligned}
 -w_{\text{cycle}} &= \oint P dV \\
 &= \text{area of } PV \text{ plot} \\
 &= q_{\text{cycle}}
 \end{aligned}$$

expansion (1  $\rightarrow$  2):

$$\Delta U = 0 \Rightarrow -w = q > 0 \quad \text{Define Entropy } \Delta S = \frac{q}{T} \therefore \Delta S > 0$$

compression (4  $\leftarrow$  3):  $\Delta S < 0$

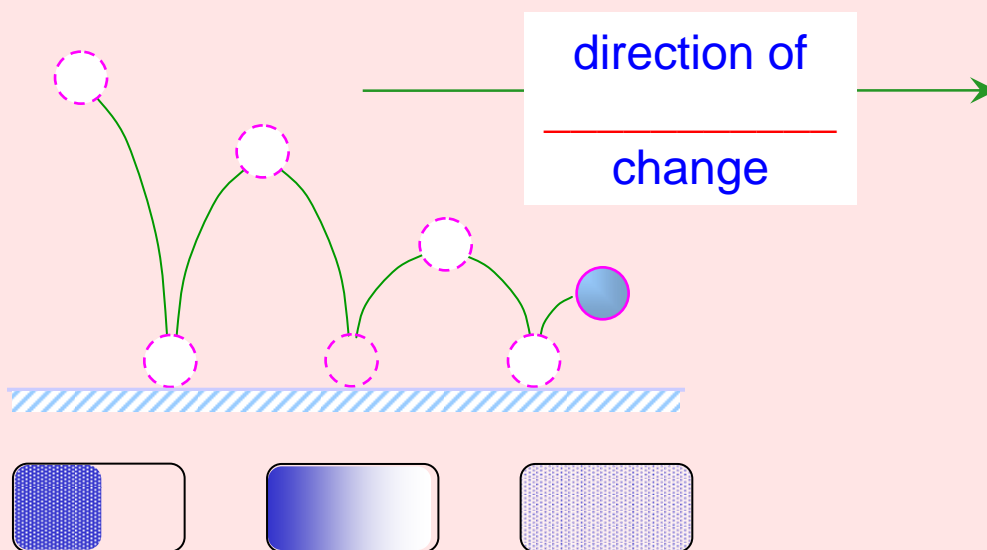
steps (2  $\rightarrow$  3 and 1  $\leftarrow$  4):  $q = 0 \Rightarrow \Delta S = 0$



$$\begin{aligned}
 q_{\text{tot}} &= T_H (S_2 - S_1) \\
 &\quad + T_L (S_4 - S_3) \\
 &= (T_H - T_L)(S_2 - S_1) \\
 &= \text{area of } ST \text{ plot} \\
 &= -w_{\text{cycle}}
 \end{aligned}$$

# Spontaneous Change

(So, why do we need entropy, anyway?)



The direction of spontaneous change is that which

- ☐ leads to \_\_\_\_\_ dispersal of the total energy
- ☐ moves from a state of low intrinsic probability towards one of \_\_\_\_\_ probability.

Work is needed to reverse a spontaneous process.

We need a quantity – \_\_\_\_\_ – to describe energy dispersal, i.e. the probability of a state.

Spontaneous processes are \_\_\_\_\_ .

They “\_\_\_\_\_” entropy

Reversible processes do not generate entropy – but they may \_\_\_\_\_ it from one part of the universe to another.

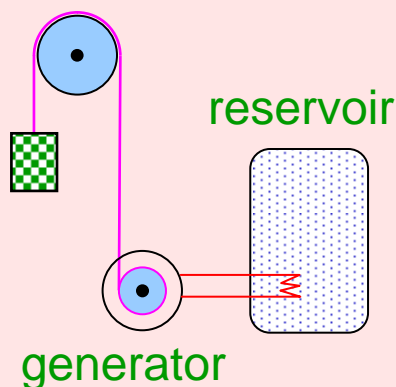
# Entropy 1

- ❖ Entropy is a \_\_\_\_\_ variable (property) which determines if a state is accessible from another by a \_\_\_\_\_ change.
- ❖ Entropy is a measure of chaotic dispersal of energy.
- ❖ The natural tendency of spontaneous change is towards states of higher entropy.
- ❖ There are both thermodynamic (how much \_\_\_\_\_ is produced?) and statistical definitions (how \_\_\_\_\_ is a state?). They both become equivalent when statistics is applied to a \_\_\_\_\_ number of molecules.

Consider a falling weight which drives a generator and thus results in heat  $q$  being added to the reservoir (the surroundings).

Define a \_\_\_\_\_ variable  $S$

$$dS(\text{surr}) = -\delta q / T$$



Then use stored energy to restore the weight to its original height. The reservoir gives up  $\delta q_{\text{rev}}$  to the system, and there is no overall change in the \_\_\_\_\_ .

this would only work for infinitesimal changes

$$dS(\text{sys}) = \text{_____} dS(\text{surr}) = \frac{\delta q_{\text{rev}}}{T}$$

# Entropy 2

In general,  $dS(\text{sys}) + dS(\text{surr}) \geq 0$

$$dS(\text{sys}) \geq -dS(\text{surr})$$

Equality for reversible processes only

or, for the \_\_\_\_\_,  $dS \geq \frac{\delta q}{T}$  \_\_\_\_\_ inequality

For an \_\_\_\_\_ system,  $q = 0$  hence  $\Delta S \geq 0$

Isothermal Processes

$$\Delta S = q_{\text{rev}} / T$$

e.g.  $\Delta S(\text{fusion}) = \frac{\Delta H_{\text{fus}}}{T_m}$        $\Delta S(\text{vap}) = \frac{\Delta H_{\text{vap}}}{T_b}$

\_\_\_\_\_ Rule:  $\Delta S(\text{vap}) \approx$  \_\_\_\_\_ J K<sup>-1</sup> mol<sup>-1</sup>

Can be used to estimate \_\_\_\_\_ if  $T_b$  is known. Not good for \_\_\_\_\_ liquids.

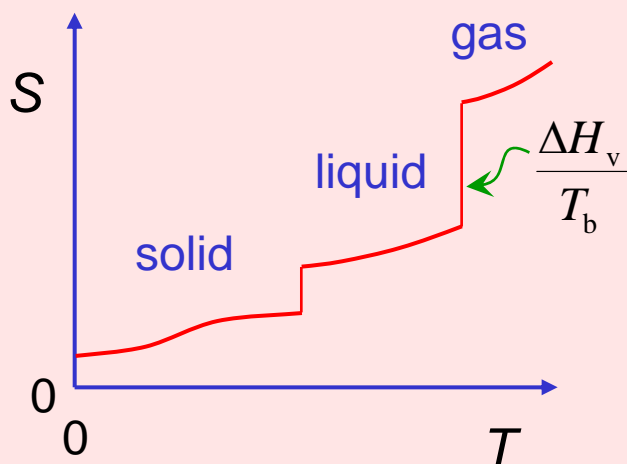
Temperature Variation  $\delta q_{\text{rev}} =$  \_\_\_\_\_  $dT$

$$(\Delta S)_V = \int_{T_1}^{T_2} \frac{C_V}{T} dT$$

and  $(\Delta S)_P = \int_{T_1}^{T_2} \frac{C_P}{T} dT$

\_\_\_\_\_ Entropy

$$S(T) = S(0) + \int_0^T \frac{C_P}{T} dT$$

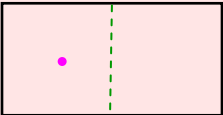
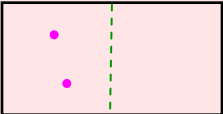
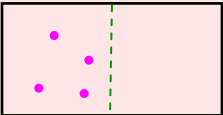
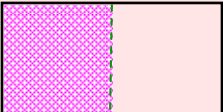


# Entropy 3 (Statistical Mechanics)

Entropy depends on Probability.

Consider the number of ways  $\Omega$  of arranging  $n$  molecules between two sides (A and B) of a container.

The probability  $\mathcal{P}_A$  that all molecules are on side A depends on the \_\_\_\_\_ of  $\Omega$  \_\_\_\_\_ to the total number of arrangements.

A	B			
		$\Omega_A = 1$	$\Omega_{\text{tot}} = 2$	$\mathcal{P}_A = \frac{1}{2}$
		$\Omega_A = 1$	$\Omega_{\text{tot}} = 4$	$\mathcal{P}_A = \frac{1}{4}$
		$\Omega_A = 1$	$\Omega_{\text{tot}} = 16$	$\mathcal{P}_A = \frac{1}{16}$
		$\Omega_A = 1$	$\Omega_{\text{tot}} = \underline{\hspace{1cm}}$	$\mathcal{P}_A = \underline{\hspace{1cm}}$

State A becomes less and less probable as  $n$  increases. Conversely, the probability of the less ordered, roughly evenly distributed states, increases.

Since entropy is a measure of \_\_\_\_\_, it follows that  $S$  depends on  $\Omega$ .

**Boltzmann equation**  $S = k \underline{\hspace{1cm}} \Omega$

Since  $\mathcal{P}(x \text{ AND } y) = \mathcal{P}_x \underline{\hspace{1cm}} \mathcal{P}_y$ ,  $\ln \mathcal{P}_{x+y} = \ln \mathcal{P}_x + \ln \mathcal{P}_y$

# The \_\_\_\_\_ Law of Thermodynamics

- “An \_\_\_\_\_ cyclic process in which there is a net conversion of \_\_\_\_\_ into work is impossible.”
- “No process is possible in which the sole result is the absorption of heat from a reservoir and its conversion into work.” It is possible to convert \_\_\_\_\_ work into heat!
- “It is impossible for heat to be transformed from a body at a lower temperature to one at a higher temperature unless \_\_\_\_\_ is done.”
- “The entropy of an isolated system \_\_\_\_\_ during any natural process.” The universe is an isolated system.  
 $\Delta S(\text{sys}) < 0$  is allowed provided  $\Delta S(\text{sys}) + \Delta S(\text{surr}) > 0$
- “All reversible \_\_\_\_\_ cycles operating between the same two temperatures have the same thermodynamic efficiency.”
- “There is a state function called entropy  $S$  that can be calculated from  $S = \delta q_{\text{rev}}/T$ . The change in entropy in any process is given by  $dS \geq \delta q/T$ , where the inequality refers to a spontaneous (irreversible) process.”

The 1<sup>st</sup> Law uses  $U$  to identify \_\_\_\_\_ changes of state.

The 2<sup>nd</sup> Law uses  $S$  to identify \_\_\_\_\_ changes among the permissible ones.