

# The Third Law of Thermodynamics

- "If the entropy of every element in its stable state at  $T = 0$  is taken as zero, every substance has a \_\_\_\_\_ entropy which at  $T = 0$  may become zero, and does become zero for all perfect \_\_\_\_\_ substances, including compounds."

\_\_\_\_\_ Heat Theorem "The entropy change accompanying transformation between condensed phases in equilibrium, including chemical \_\_\_\_\_, approaches zero as  $T \rightarrow 0$ ."

$$\lim_{T \rightarrow 0} \Delta S = 0$$

Practical consequence: Set  $S(0) = 0$  for \_\_\_\_\_ by convention. Apply Nernst to determine  $S(0)$  for all else.

- "It is impossible to reach absolute zero in a \_\_\_\_\_ number of steps."

The 1<sup>st</sup> Law says  $U$  cannot be \_\_\_\_\_ or \_\_\_\_\_.

The 2<sup>nd</sup> Law says  $S$  cannot \_\_\_\_\_.

The 3<sup>rd</sup> Law says zero \_\_\_\_\_ cannot be reached.

# The \_\_\_\_\_ Inequality

Given  $dS = \frac{\delta q_{\text{rev}}}{T}$  and  $\delta w_{\text{rev}} = -PdV$

Substitute into the 1<sup>st</sup> Law:

$$dU = TdS - PdV \quad \text{Equation of Thermodynamics}$$

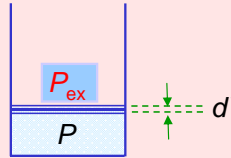
All \_\_\_\_\_ differentials, so path independent.

But  $dU = \delta q + \delta w = \delta q - P_{\text{ex}} dV$  In general, i.e. any path

$$\Rightarrow dS = \frac{\delta q}{T} + \frac{(\text{---} - \text{---})}{T} dV$$

If  $P > P_{\text{ex}}$ ,  $dV > 0$ ; if  $P < P_{\text{ex}}$ ,  $dV < 0$

$$\therefore (P - P_{\text{ex}})dV \geq 0$$



Clausius Inequality  $dS \geq \frac{\delta q}{T}$  Equal for change

Even more generally,

$$dS_{\text{univ}} = \left( \frac{T - T_{\text{surr}}}{T} \right) dS_{\text{surr}} + \frac{(P - P_{\text{surr}})}{T} dV$$

Conditions for: \_\_\_\_\_ equilibrium \_\_\_\_\_ equilibrium

# The Fundamental \_\_\_\_\_ of Thermodynamics

Combine  $dU = \delta q - PdV$  with  $dS = \frac{\delta q_{\text{rev}}}{T}$  Reversible change but true for all paths since  $dV$  exact

$$\Rightarrow dU = TdS - PdV$$

or  $dS = \left( \frac{1}{T} \right) dU + \left( \frac{P}{T} \right) dV$   $\delta w_{\text{rev}} = -PdV$

This fundamental equation generates many more \_\_\_\_\_.

Example 1: Comparison with  $dU = \left( \frac{\partial U}{\partial S} \right)_V dS + \left( \frac{\partial U}{\partial V} \right)_S dV$   
 $\Rightarrow \left( \frac{\partial U}{\partial S} \right)_V = T$  and  $\left( \frac{\partial U}{\partial V} \right)_S = -P$

Example 2: Consider that  $dU$  is exact and cross \_\_\_\_\_.

use this relation here  $\left( \frac{\partial T}{\partial V} \right)_S = - \left( \frac{\partial P}{\partial S} \right)_V$  which is a \_\_\_\_\_ relation

Example 3:  $\left( \frac{\partial S}{\partial V} \right)_T = \left( \frac{\partial S}{\partial T} \right)_V \left( \frac{\partial V}{\partial T} \right)_T$  cyclic rule...

$$= \left( \frac{\partial P}{\partial T} \right)_V \left( \frac{\partial V}{\partial T} \right)_T = \left( \frac{\partial P}{\partial T} \right)_V \dots \text{again} \dots$$

$$\dots \text{and again!} = \left( \frac{\partial P}{\partial T} \right)_V \left( \frac{\partial V}{\partial T} \right)_T = - \left( \frac{\partial V}{\partial T} \right)_P \left( \frac{\partial P}{\partial V} \right)_T = \alpha / \kappa$$

# How Entropy Depends on $T$ and $V$

$$dS = \left( \frac{1}{T} \right) dU + \left( \frac{P}{T} \right) dV$$

$$\Rightarrow dS = \frac{C_V}{T} dT + \frac{1}{T} \left[ P + \left( \frac{\partial U}{\partial V} \right)_T \right] dV$$

Compare with  $dS = \left( \frac{\partial S}{\partial T} \right)_V dT + \left( \frac{\partial S}{\partial V} \right)_T dV$

$$\left( \frac{\partial S}{\partial T} \right)_V = \frac{C_V}{T}$$

$$\left( \frac{\partial S}{\partial V} \right)_T = \frac{1}{T} \left[ P + \left( \frac{\partial U}{\partial V} \right)_T \right]$$

$$\Delta S = \int \frac{C_V}{T} dT$$

0 for ideal gases

$$\Delta S_T = \int \frac{P}{T} dV = \int \frac{1}{V} dV = nR \ln(V_2/V_1)$$

For \_\_\_\_\_ substance,  $dS = \frac{C_V}{T} dT + \frac{\alpha}{\kappa} dV$

For \_\_\_\_\_ gases,  $\Delta S = C_V \ln \frac{T_2}{T_1} + nR \ln \frac{V_2}{V_1}$  assuming  $C_V$  is  $T$  \_\_\_\_\_  
 (eqs 3.7.4, 6.1.6)

Entropy Depends on  $T$  and  $P$  (not in text)

$$dS = \left(\frac{1}{T}\right)dU + \left(\frac{P}{T}\right)dV$$

First problem: replace \_\_\_\_\_ ; second problem: replace \_\_\_\_\_.  
Use  $d$  \_\_\_\_\_ =  $dH - Pd$  \_\_\_\_\_ -  $VdP$  and both are solved!

$$\Rightarrow dS = \left(\frac{1}{T}\right)dH - \left(\frac{V}{T}\right)dP$$

But  $dH = \left(\frac{\partial H}{\partial T}\right)_P dT + \left(\frac{\partial H}{\partial P}\right)_T dP =$  \_\_\_\_\_  $dT + \left(\frac{\partial H}{\partial P}\right)_T dP$

$$\Rightarrow dS = \frac{C_P}{T} dT + \frac{1}{T} \left[ \left(\frac{\partial H}{\partial P}\right)_T - V \right] dP$$

Compare with  $dS = \left(\frac{\partial S}{\partial T}\right)_P dT + \left(\frac{\partial S}{\partial P}\right)_T dP$  \_\_\_\_\_ variables

$$\left(\frac{\partial S}{\partial T}\right)_P = \frac{C_P}{T}$$
  
$$\left(\frac{\partial S}{\partial P}\right)_T = \frac{1}{T} \left[ \left(\frac{\partial H}{\partial P}\right)_T - V \right]$$
  
$$\Delta S = \int \frac{C_P}{T} dT$$
  
$$\Delta S_T = - \int \frac{V}{T} dP = - \int \frac{1}{P} dP$$
  
$$\Delta S_T = -nR \ln(P_2/P_1) = nR \ln(V_2/V_1)$$
  
$$\Delta S_T = -nR \ln(P_2/P_1) = nR \ln(V_2/V_1)$$

Entropy Changes in \_\_\_\_\_ Processes

Entropy is a \_\_\_\_\_ function, so  
 $\Delta S(\text{sys}) = S_2 - S_1$  \_\_\_\_\_ of path  
This can be used to calculate  $\Delta S$  for an irreversible process.

❖ Consider isothermal expansion of a gas from  $V_1$  to  $V_2$ :  
 $\Delta S(\text{sys}) = nR \ln(V_2/V_1)$  reversible and irreversible cases  
For the reversible case  $\Delta S(\text{surr}) = -\Delta S(\text{sys})$   
For the irreversible case  $\Delta S(\text{surr})$  \_\_\_\_\_  $-\Delta S(\text{sys})$   
$$\Delta S(\text{surr}) = -\frac{q}{T} = \frac{w}{T}$$
  
e.g. for \_\_\_\_\_ expansion,  $w = 0$   
$$\Rightarrow \Delta S(\text{surr}) = 0, \Delta S(\text{univ}) = \Delta S(\text{sys}) > 0$$

❖ Consider freezing of \_\_\_\_\_ water at  $T < 273 \text{ K}$

water,  $0^\circ\text{C}$   $\xrightarrow{\text{reversible}}$  ice,  $0^\circ\text{C}$   
rev.  $\uparrow$   $\downarrow$  rev.  
water,  $T$   $\xrightarrow{\text{irreversible}}$  ice,  $T$

$\Delta S(\text{sys}) =$  \_\_\_\_\_  
 $\Delta S(\text{surr}) =$  \_\_\_\_\_

Entropy of \_\_\_\_\_

❖ Consider the mixing of two ideal gases :

$P, T, V_1, n_1$   $\rightarrow$   $P, T, V_2, n_2$   $\rightarrow$   $P, T, V_1+V_2, n_1+n_2$

$$\Delta S_1 = -n_1 R \ln \frac{V_1}{V_1+V_2} = -n_1 R \ln \frac{n_1}{n_1+n_2} = -n_1 R \ln \frac{1}{1+\frac{n_2}{n_1}}$$
  
$$\Delta S_2 = -n_2 R \ln \frac{V_2}{V_1+V_2} = -n_2 R \ln \frac{n_2}{n_1+n_2} = -n_2 R \ln \frac{1}{1+\frac{n_1}{n_2}}$$
  
$$\Delta S_{\text{mix}} = -n_1 R \ln \chi_1 - n_2 R \ln \chi_2 = -(n_1+n_2) R \ln \left( \frac{1}{\chi_1^{\frac{n_1}{n_1+n_2}} \chi_2^{\frac{n_2}{n_1+n_2}}} \right)$$
  
In general 
$$\Delta S_{\text{mix}} = -R \sum_i \chi_i \ln \chi_i$$

This expression applies to the arrangement of objects (\_\_\_\_\_ ) just as well as \_\_\_\_\_ (gases and liquids).  
For example, arrange  $N$  identical atoms in  $N$  sites in a crystal:  
$$\Omega = N! / N! = 1 \quad S = k \ln \Omega = 0$$
  
Compare with the arrangement of two types of atoms, A and B.  
$$\Omega = \frac{N!}{N_A! N_B!} \quad \Delta S = k \ln \Omega$$
  
Application of \_\_\_\_\_ approximation:  $\ln(z!) = z \ln z - z$   
leads to 
$$\Delta S_{\text{config}} = -kN(\chi_A \ln \chi_A + \chi_B \ln \chi_B)$$
 do the math

Using Entropy to Achieve Low  $T$

$$\Delta S_P = \int \frac{C_P}{T} dT \approx C_P \ln(T_2/T_1)$$
 if  $C_P$  is  $\approx$  constant

To achieve lower temperatures,  $S$  must be reduced.  
Choose some property  $X$  that varies with  $S$ , i.e.  $S = f(X, T)$ .  
This could be the pressure of a gas or, for example, the magnetic moment of a paramagnetic salt (whose energy varies with magnetic field).

1. Alter  $X$  isothermally. Entropy changes.  $\Delta S_T \approx \left(\frac{\partial S}{\partial X}\right)_T \Delta X, \quad q = T \Delta S$
2. Restore  $X$  by a reversible adiabatic process.  $q = 0, \quad \Delta S = 0$
3. Repeat cycle.

curves coincide at  $0 \text{ K}$ , a consequence of the Nernst Heat Theorem