

## **Announcements**

- 1. Physics lecture today 4 PM – Professor Brian Matthews from U. Oregon – “Tolerance and Intolerance in Protein Structure and Function”**
- 2. Third exam on Chaps. 15-20 – Tuesday, Nov. 25<sup>th</sup>**  
**Bring:**
  - Equation sheet (8.5" x 11")**
  - Calculator**
  - Clear head**
- 3. Extra problem solving session ?**
  - Sunday night 10-11 PM ?**
  - Monday evening 6-7 PM ?**
  - Other ?**

Advice for preparing for exam:

- Review lecture notes and text chapters
- Prepare equation sheet
- Work problems using equation sheet and calculator
  - From homework
  - From previous exams
  - From on line quizzes

Topics:

- Simple harmonic motion & resonance
- Wave motion, standing waves, sound waves
- The physics of fluids
- Thermodynamics & the ideal gas law

## Simple harmonic motion & resonance

Hooke's law:  $F = -kx = ma = m \frac{d^2 x}{dt^2}$       Solution :

$$\Rightarrow \frac{d^2 x}{dt^2} = -\frac{k}{m} x$$

$$x(t) = A \cos(\omega t + \varphi); \quad \omega = \sqrt{\frac{k}{m}}$$

Simple  
pendulum:

$$\tau = mgL \sin \Theta = -I\alpha = -I \frac{d^2 \Theta}{dt^2}$$

$$\Rightarrow \frac{d^2 \Theta}{dt^2} = -\frac{mgL}{I} \sin \Theta \approx -\frac{mgL}{I} \Theta$$

Solution (for small  $\Theta$ ):

$$\Theta(t) = A \cos(\omega t + \varphi); \quad \omega = \sqrt{\frac{mgL}{I}}$$

The notion of resonance:

$$\text{Suppose } F = -kx + F_0 \sin(\Omega t)$$

According to Newton's second law:

$$-kx + F_0 \sin(\Omega t) = m \frac{d^2 x}{dt^2}$$

Differential equation ("inhomogeneous"):

$$\frac{d^2 x}{dt^2} = -\frac{k}{m} x + \frac{F_0}{m} \sin(\Omega t)$$

A solution to this equation takes the form:

$$x(t) = \frac{F_0 / m}{k / m - \Omega^2} \sin(\Omega t) \equiv \frac{F_0 / m}{\omega^2 - \Omega^2} \sin(\Omega t)$$

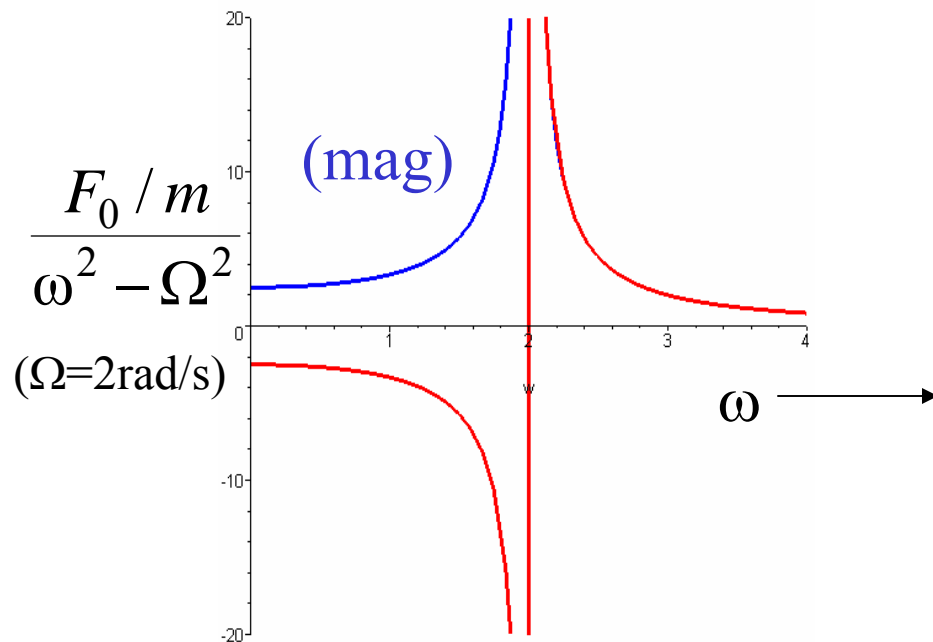
# Physics of a “driven” harmonic oscillator:

$$-kx + F_0 \sin(\Omega t) = m \frac{d^2 x}{dt^2}$$

“driving” frequency

$$x(t) = \frac{F_0 / m}{k / m - \Omega^2} \sin(\Omega t) \equiv \frac{F_0 / m}{\omega^2 - \Omega^2} \sin(\Omega t)$$

“natural” frequency

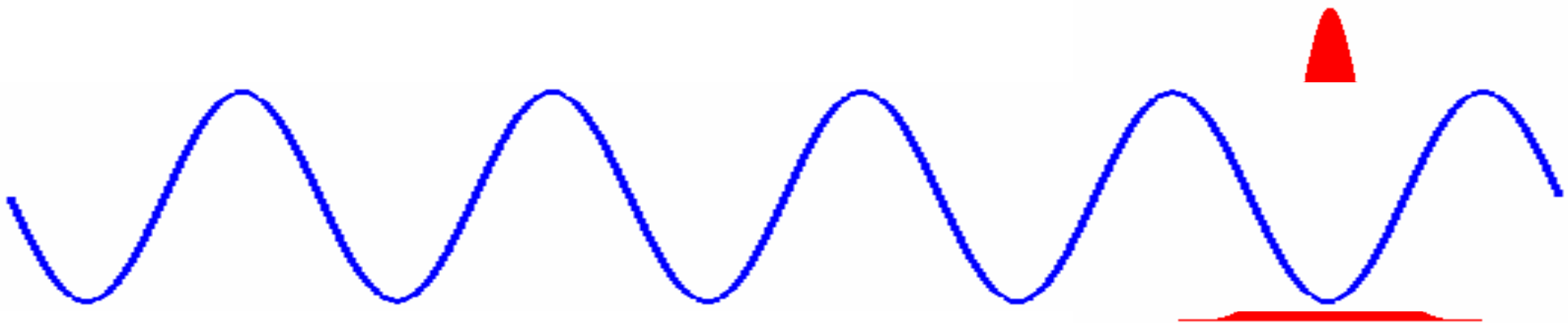


# The phenomenon of wave motion

The wave equation

$y(x, t)$   
position time

$$\frac{\partial^2 y}{\partial t^2} = v^2 \frac{\partial^2 y}{\partial x^2}$$



The wave equation:

$$\frac{\partial^2 y}{\partial t^2} = v^2 \frac{\partial^2 y}{\partial x^2} \quad \text{where } v \equiv \sqrt{\frac{T}{\mu}} \text{ (for a string)}$$

Solutions:  $y(x,t) = f(x \pm vt)$  ← function of *any* shape

Moving “pulse”:  $y(x,t) = y_0 e^{-(x-vt)^2}$

“wave vector”  
not spring constant!!!

Periodic wave:  $y(x,t) = y_0 \sin(k(x - vt) + \varphi)$

$$y(x,t) = y_0 \sin\left(2\pi\left(\frac{x}{\lambda} - \frac{t}{T}\right) + \varphi\right) \quad \frac{\lambda}{T} = v$$

Superposition of waves:

$$y_{right}(x, t) = y_0 \sin\left(2\pi\left(\frac{x}{\lambda} - \frac{t}{T}\right) + \varphi\right) \quad y_{left}(x, t) = y_0 \sin\left(2\pi\left(\frac{x}{\lambda} + \frac{t}{T}\right) + \varphi\right)$$

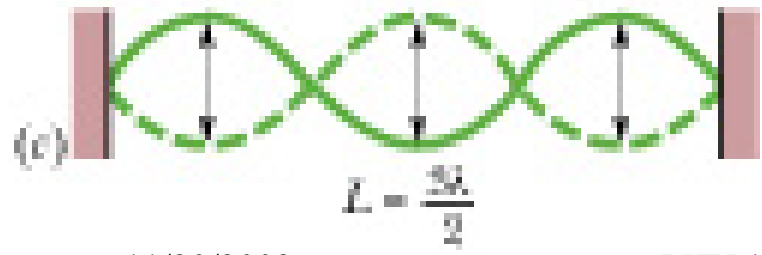
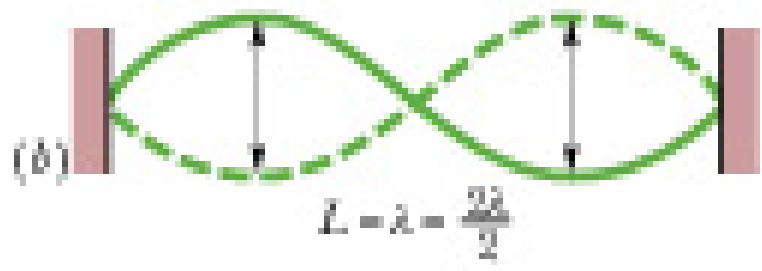
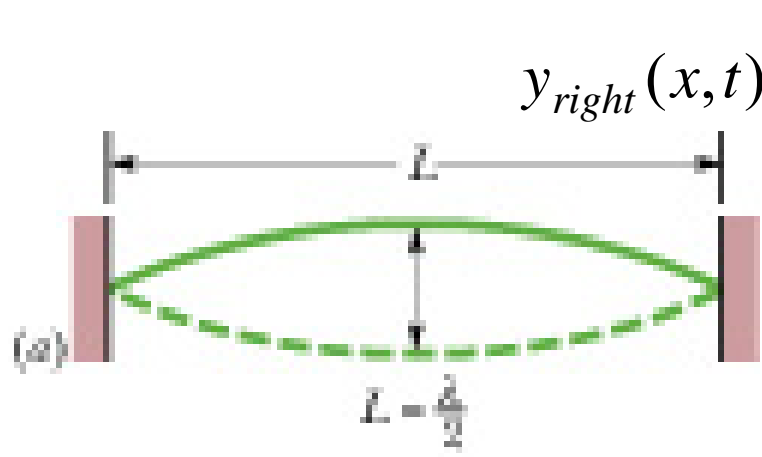
$$\sin A \pm \sin B = 2 \sin\left[\frac{1}{2}(A \pm B)\right] \cos\left[\frac{1}{2}(A \mp B)\right]$$

“Standing” wave:

$$y_{right}(x, t) + y_{left}(x, t) = 2y_0 \sin\left(\frac{2\pi x}{\lambda} + \varphi\right) \cos\left(\frac{2\pi t}{T}\right)$$



Constraints of standing waves: ( $\phi = 0$ )



$$y_{right}(x,t) + y_{left}(x,t) = 2y_0 \sin\left(\frac{2\pi x}{\lambda}\right) \cos\left(\frac{2\pi t}{T}\right)$$

for string:

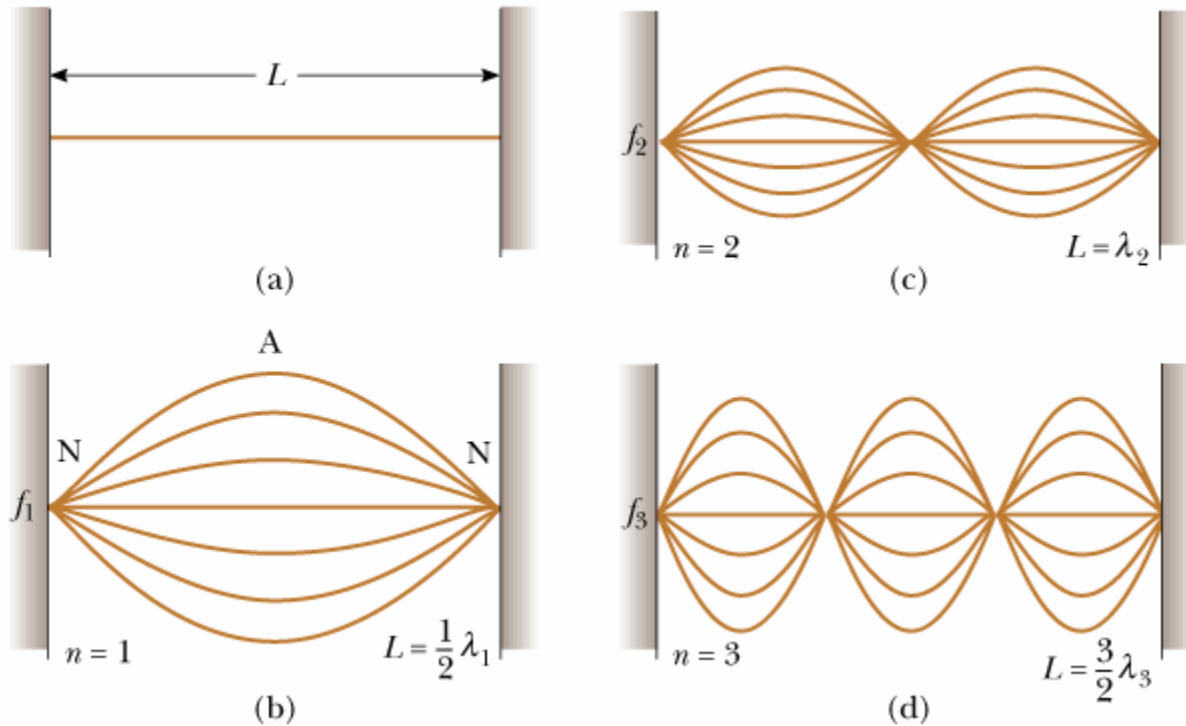
$$\lambda = \frac{2L}{n}; n = 1, 2, 3, \dots$$

$$\frac{\lambda}{T} = v \Rightarrow T = \frac{2L}{nv}$$

$$f = \frac{1}{T} = \frac{nv}{2L}$$

# The sound of music

String instruments (Guitar, violin, etc.)



$$\lambda_n = \frac{2L}{n}$$

$$f_n = \frac{v}{\lambda_n} = \frac{nv}{2L}$$

$$v = \sqrt{\frac{T}{\mu}}$$

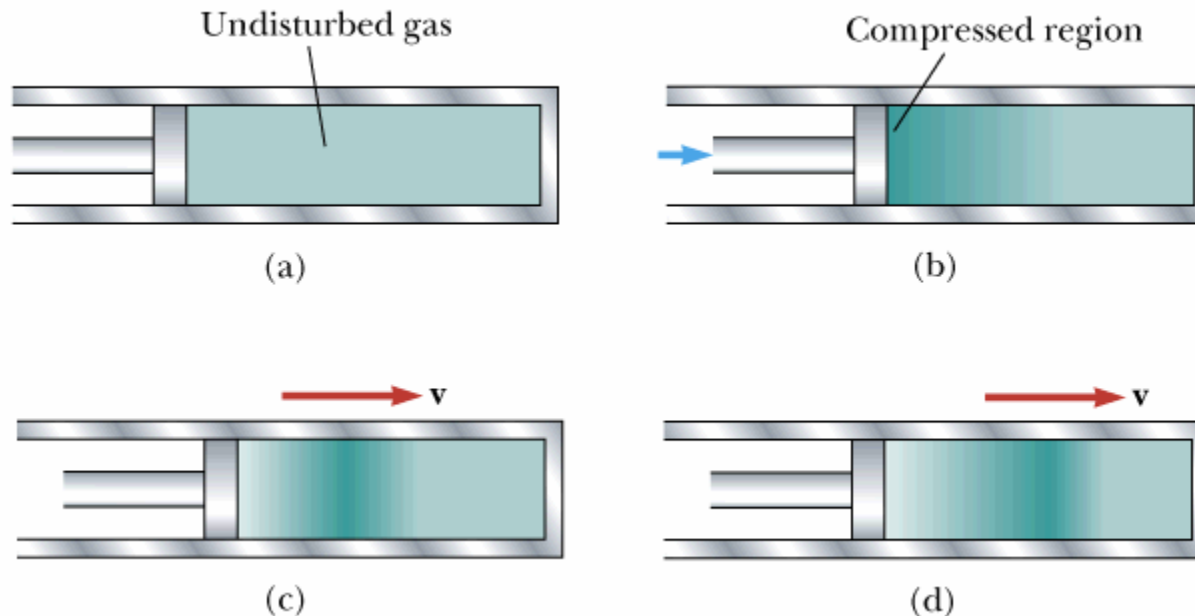
$$y_{\text{standing}}(x, t) = 2y_0 \sin\left(\frac{2\pi x}{\lambda}\right) \cos\left(\frac{2\pi t}{T}\right) \rightarrow \text{Couples to sound in the air}$$

# Sound waves

Longitudinal waves propagating in a fluid or solid

$$\frac{\partial^2 y}{\partial t^2} = v^2 \frac{\partial^2 y}{\partial x^2} \quad y(x,t) \text{ describes density or pressure variations}$$

$$v = \sqrt{\frac{T}{\mu}} \Rightarrow \sqrt{\frac{B}{\rho}} \left( \sqrt{\frac{\text{compressibility}}{\text{density}}} \right) = \sqrt{\frac{\gamma P}{\rho}} \approx 343 \text{ m/s}$$



## Periodic sound wave

In terms of pressure:

$$P(x, t) = P_0 + \Delta P_{\max} \sin\left(\frac{2\pi x}{\lambda} \pm \frac{2\pi t}{T}\right)$$

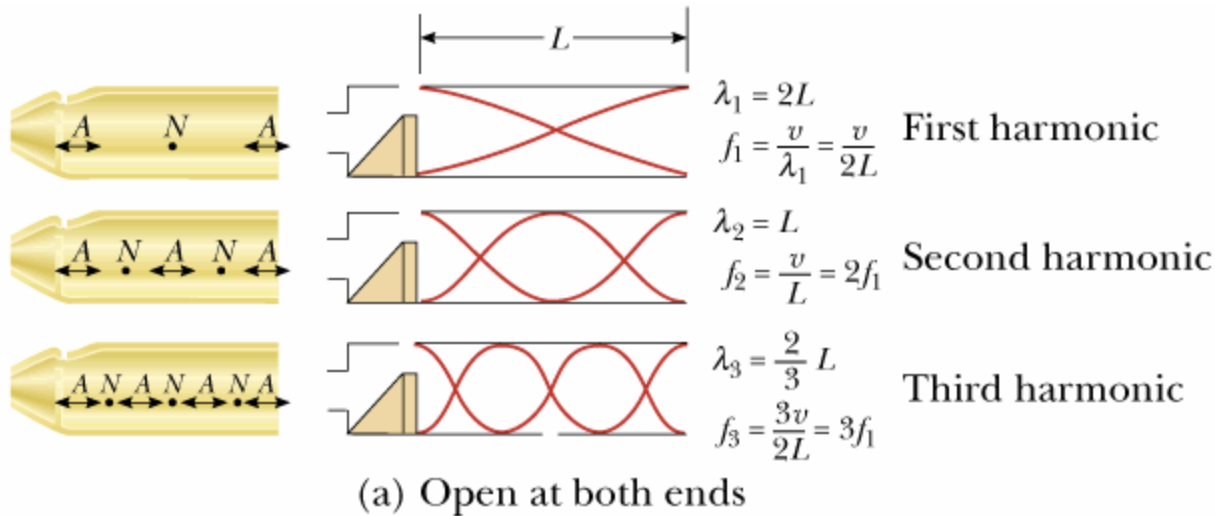
Sound intensity: (energy/(unit time · unit area))

$$I \equiv \frac{(\Delta P_{\max})^2}{2\rho v}$$

Decibel scale:

$$\beta \equiv 10 \log\left(\frac{I}{I_0}\right) \quad I_0 = 10^{-12} \text{ W/m}^2$$

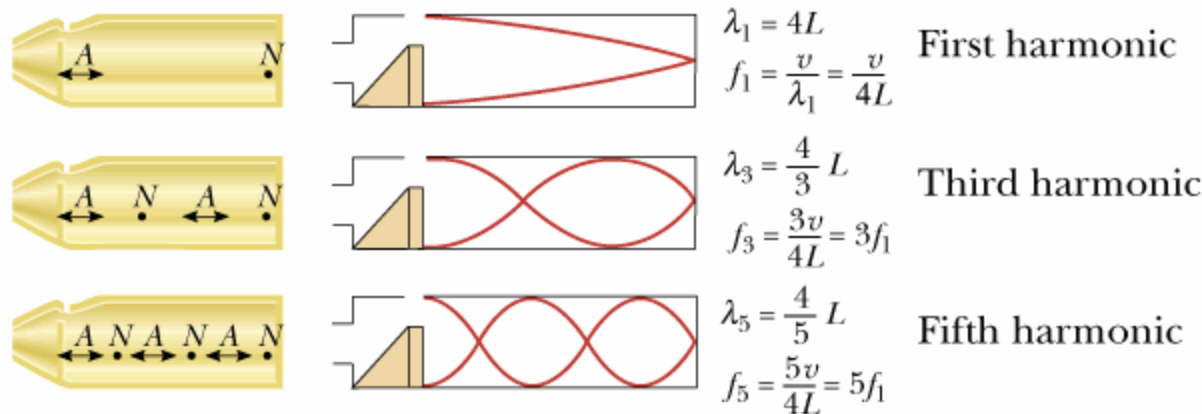
# “Wind” instruments (standing waves in air)



$$\lambda_n = \frac{2L}{n}; n = 1, 2, 3, \dots$$

$$f_n = \frac{v}{\lambda_n} = \frac{nv}{2L}$$

$$v = \sqrt{\frac{\gamma P}{\rho}} \approx 343 \text{ m/s}$$



$$\lambda_n = \frac{4L}{n}; n = 1, 3, 5, \dots$$

$$f_n = \frac{v}{\lambda_n} = \frac{nv}{4L}$$

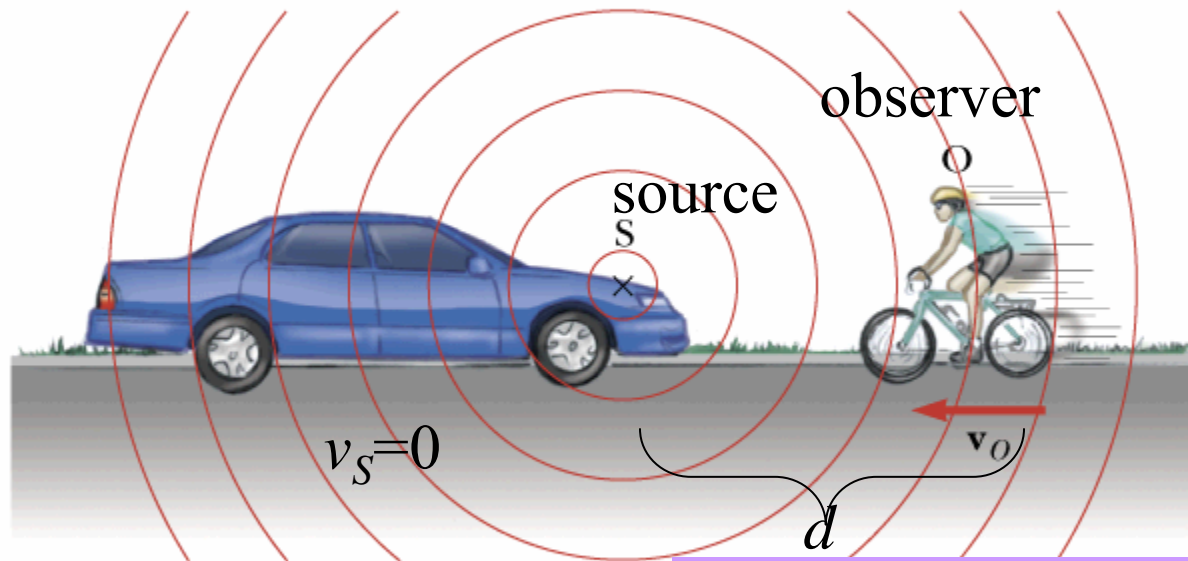
$$v = \sqrt{\frac{\gamma P}{\rho}} \approx 343 \text{ m/s}$$

# The “Doppler” effect

$v$ =sound velocity

observer moving, source stationary

Serway, Physics for Scientists and Engineers, 5/e  
Figure 17.10



$$vt_1 = d - v_o t_1$$

$$v(t_2 - T) = d - v_o t_2$$

$$t_2 - t_1 = \frac{1}{f_o} = T \frac{v}{v + v_o}$$

$$f_o = f_s \frac{v + v_o}{v}$$

Summary of sound Doppler effect :

$$f_o = f_s \frac{v \pm v_o}{v \mp v_s}$$

toward

away

The physics of fluids.

- Fluids include liquids (usually “incompressible”) and gases (highly “compressible”).

- **Fluids obey Newton’s equations of motion**, but because they move within their containers, the application of Newton’s laws to fluids introduces some new forms.

  - Pressure:  $P = \text{force/area}$        $1 \text{ (N/m}^2\text{)} = 1 \text{ Pascal}$

  - Density:  $\rho = \text{mass/volume}$        $1 \text{ kg/m}^3 = 0.001 \text{ gm/ml}$

**Note: In this chapter  $P \equiv$  pressure (NOT MOMENTUM)**

Density:  $\rho = \text{mass/volume}$

Effects of the weight of a fluid:

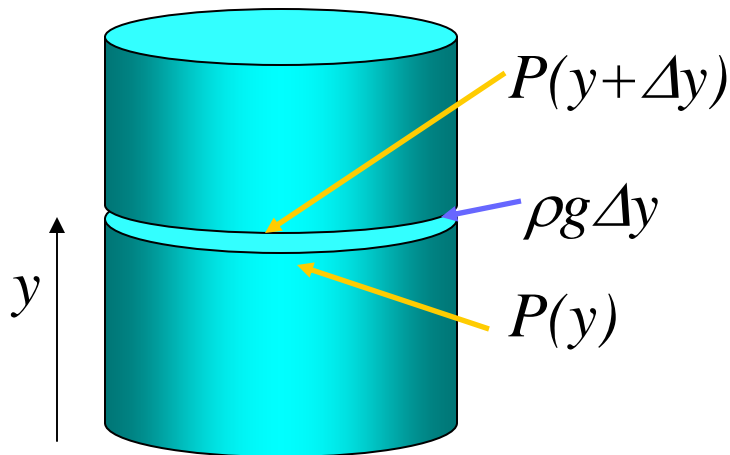
$$F(y) = F(y + \Delta y) + mg$$

$$\frac{F(y)}{A} = \frac{F(y + \Delta y)}{A} + \frac{mg}{A}$$

$$P(y) = P(y + \Delta y) + \rho g \Delta y$$

$$\lim_{\Delta y \rightarrow 0} \frac{P(y + \Delta y) - P(y)}{\Delta y} = \frac{dP}{dy}$$

$$\Rightarrow \frac{dP}{dy} = -\rho g$$



Note: In this formulation **+y** is defined to be in the **up** direction.

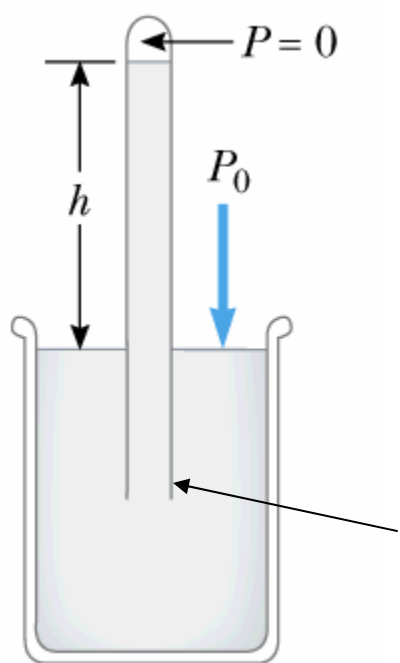


For an “incompressible” fluid (such as mercury):

$$\rho = 13.585 \times 10^3 \text{ kg/m}^3 \text{ (constant)}$$

$$\frac{dP}{dy} = -\rho g \quad \Rightarrow \quad P = P_0 - \rho g(y - y_0)$$

Example:


$$h = y - y_0 = \frac{P_0}{\rho g}$$
$$= \frac{1.013 \times 10^5 \text{ Pa}}{13.595 \times 10^3 \text{ kg/m}^3 \cdot 9.8 \text{ m/s}^2}$$
$$= 0.76 \text{ m}$$
$$\rho = 13.595 \times 10^3 \text{ kg/m}^3$$

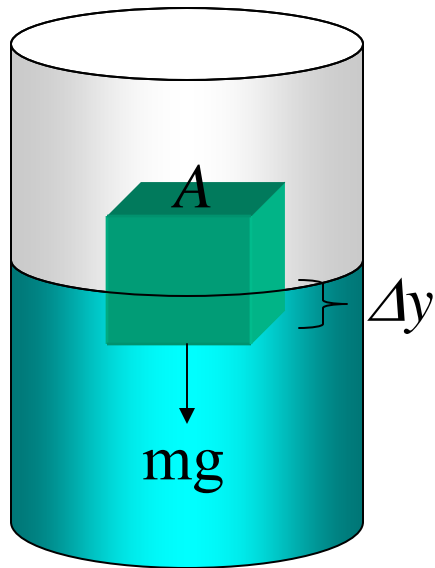
## Buoyant force for fluid acting on a solid:

$$F_B = \rho_{\text{fluid}} V_{\text{displaced}} g$$

$$P(y) = P(y + \Delta y) + \rho_{\text{fluid}} g \Delta y$$

$$\text{Buoyant force: } F_B = F_{\text{bottom}} - F_{\text{top}}$$

$$F_B = \{P(y) - P(y + \Delta y)\}A = \rho_{\text{fluid}} g \Delta y A = \rho_{\text{fluid}} g V_{\text{submerged}}$$



$$F_B - mg = 0$$

$$\rho_{\text{fluid}} V_{\text{submerged}} g - \rho_{\text{solid}} V_{\text{solid}} g = 0$$

$$\frac{V_{\text{submerged}}}{V_{\text{solid}}} = \frac{\rho_{\text{solid}}}{\rho_{\text{fluid}}}$$

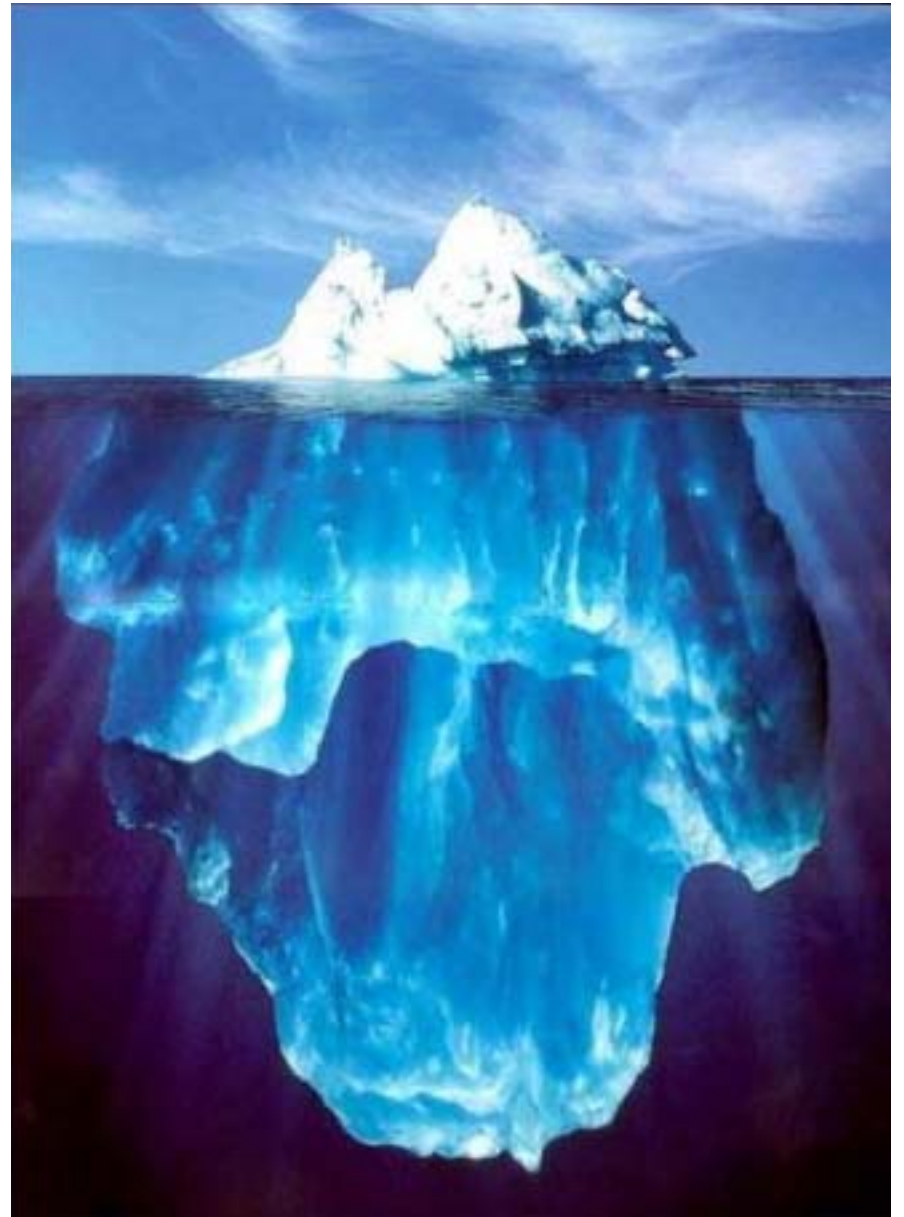
Bouyant forces:

the tip of the iceberg

$$F_B = m_{ice} g$$

$$\rho_{water} g V_{submerged} = \rho_{ice} g V_{total}$$

$$\frac{V_{submerged}}{V_{total}} = \frac{\rho_{ice}}{\rho_{water}} \approx \frac{0.917}{1.024} = 93.9\%$$



Source: <http://bb-bird.com/iceburg.html>

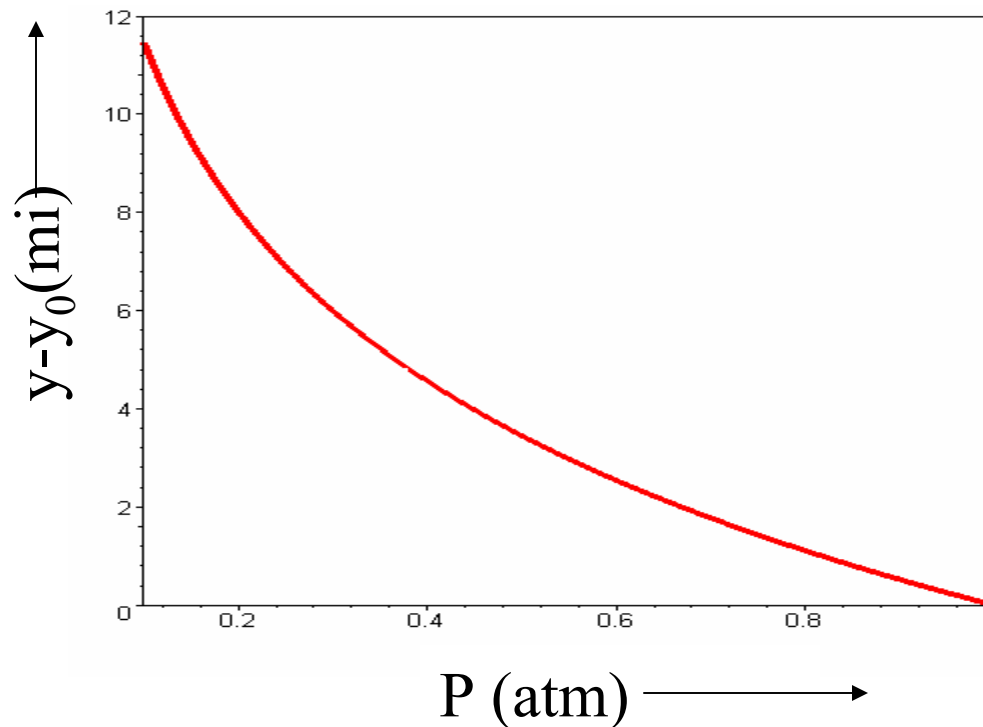
Effects of the weight of a “compressible” fluid on pressure.

$$\frac{dP}{dy} = -\rho g$$

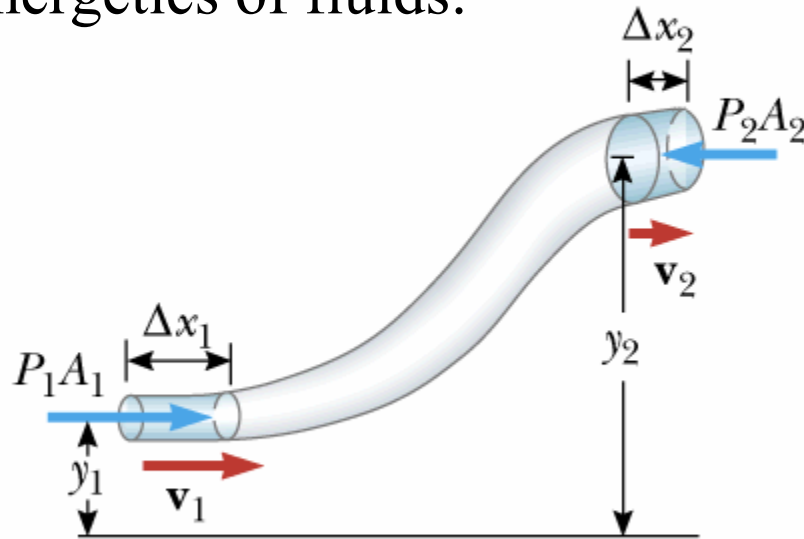
$$\text{For a gas, } \rho = P \frac{\rho_0}{P_0}$$

$$\frac{dP}{dy} = -P \left( \frac{\rho_0 g}{P_0} \right)$$

$$\text{Solution: } P(y) = P_0 e^{-\frac{\rho_0 g}{P_0}(y-y_0)} \approx P_0 e^{-\frac{y-y_0}{8000m}} \approx P_0 e^{-\frac{y-y_0}{5mi}}$$



## Energetics of fluids:



$$\Delta x_1 = v_1 \Delta t$$

$$\Delta x_2 = v_2 \Delta t$$

$$A_1 \Delta x_1 = A_2 \Delta x_2$$

$$m = \rho A_1 \Delta x_1$$

Harcourt, Inc.

$$K_2 + U_2 = K_1 + U_1 + W_{12}$$

$$\frac{1}{2} m v_2^2 + m g h_2 = \frac{1}{2} m v_1^2 + m g h_1 + (P_1 A_1 \Delta x_1 - P_2 A_2 \Delta x_2)$$

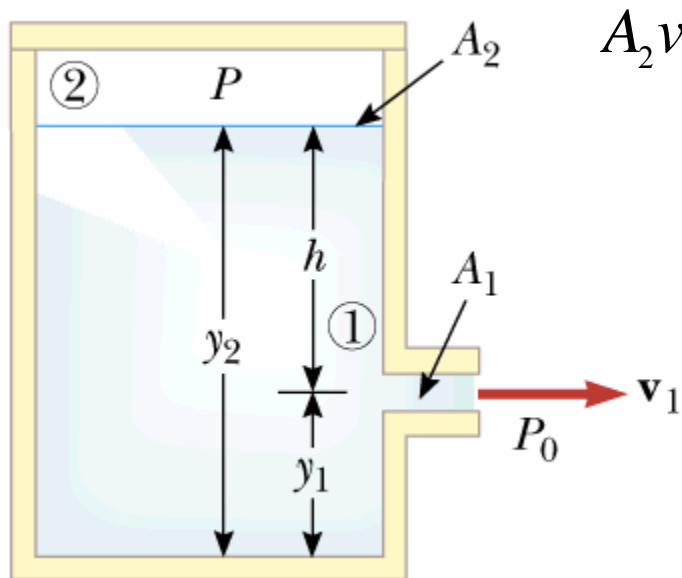
$$\rightarrow P_2 + \frac{1}{2} \rho v_2^2 + \rho g h_2 = P_1 + \frac{1}{2} \rho v_1^2 + \rho g h_1$$

Bernoulli's equation:

$$P_2 + \frac{1}{2} \rho v_2^2 + \rho g h_2 = P_1 + \frac{1}{2} \rho v_1^2 + \rho g h_1$$

Example: Suppose we know  $A_1, A_2, \rho, P, y_1, y_2$

$$P + \frac{1}{2} \rho v_2^2 + \rho g y_2 = P_0 + \frac{1}{2} \rho v_1^2 + \rho g y_1$$



$$A_2 v_2 = A_1 v_1$$

$$v_1 = \sqrt{\frac{2(gh - (P_0 - P)/\rho)}{1 - (A_1/A_2)^2}}$$

Thermodynamic statement of conservation of energy –

### First Law of Thermodynamics

$$\Delta E_{\text{int}} = Q - W$$

Work done by system

Heat added to system

“Internal” energy of system

How is temperature related to  $E_{\text{int}}$ ?

Consider an ideal gas

- Analytic expressions for physical variables
- Approximates several real situations

Ideal Gas Law:  $P V = n R T$

The diagram shows the equation  $P V = n R T$  with the variables  $P$ ,  $V$ ,  $n$ ,  $R$ , and  $T$  in orange. Pink arrows point from each variable to its corresponding unit or value:  $P$  to pressure (Pa),  $V$  to volume ( $\text{m}^3$ ),  $n$  to number of moles,  $R$  to gas constant ( $8.31 \text{ J}/(\text{mole} \cdot \text{K})$ ), and  $T$  to temperature (K).

pressure (Pa)

volume ( $\text{m}^3$ )

number of moles

gas constant ( $8.31 \text{ J}/(\text{mole} \cdot \text{K})$ )

temperature (K)



Review of results from ideal gas analysis in terms of the specific heat ratio  $\gamma \equiv C_p/C_v$ :

$$\Delta E_{\text{int}} = \frac{n}{\gamma-1} R\Delta T = nC_v\Delta T$$

$$C_v = \frac{R}{\gamma-1} \qquad C_p = \frac{\gamma R}{\gamma-1}$$

$$\text{Note: } \gamma = \frac{C_p}{C_v}$$