

PHY 113 A General Physics I
9-9:50 AM MWF Olin 101

Plan for Lecture 31:

Chapter 21: Ideal gas equations

- 1. Molecular view of ideal gas**
- 2. Internal energy of ideal gas**
- 3. Distribution of molecular speeds in ideal gas**

11/16/2012

PHY 113 A Fall 2012 – Lecture 31

1

22	10/29/2012	Kepler's laws and satellite motion	13.4-13.6	13.28, 13.34	10/31/2012
	10/31/2012	Review	10-13, 15		
	11/02/2012	Exam	10-13, 15		
23	11/05/2012	Fluid mechanics	14.1-14.4	14.8, 14.24	11/07/2012
24	11/07/2012	Fluid mechanics	14.6-14.7	14.39, 14.51	11/09/2012
25	11/09/2012	Temperature	18.1-18.6	18.1, 18.20	11/12/2012
26	11/12/2012	Heat	20.1-20.4	20.3, 20.14, 20.24	11/14/2012
27	11/14/2012	First law of thermodynamics	20.5-20.7	20.28, 20.35	11/16/2012
28	11/16/2012	Ideal gases	21.1-21.5	21.10, 21.19	11/19/2012
29	11/19/2012	Engines	22.1-22.8	22.3, 22.62	11/26/2012
	11/21/2012	Thanksgiving Holiday			
	11/23/2012	Thanksgiving Holiday			
	11/26/2012	Review	14, 19, 22		
	11/28/2012	Exam	14, 19, 22		
30	11/30/2012	Wave motion	16.1-16.6		12/03/2012
31	12/03/2012	Sound & standing waves	17.1-17.8		12/05/2012
	12/05/2012	Review	1, 22		
	12/13/2012	Final Exam – 9 AM			

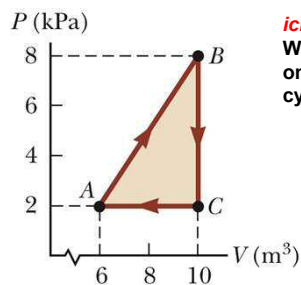
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2

Review:

Consider the process described by $A \rightarrow B \rightarrow C \rightarrow A$



iclicker exercise:
What is the net work done on the system in this cycle?

- A. -12000 J
 B. 12000 J
 C. 0

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3

Equation of "state" for ideal gas
(from experiment)

$$PV = nRT$$

pressure in Pascals → P
 volume in m^3 → V
 # of moles → n
 temperature in K → T
 $R = 8.314 \text{ J}/(\text{mol K})$

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4

Ideal gas -- continued

Equation of state: $PV = nRT$

Internal energy: $E_{\text{int}} = \frac{1}{\gamma - 1} nRT = \frac{1}{\gamma - 1} PV$

γ = parameter depending on type of ideal gas

$$= \begin{cases} \frac{5}{3} & \text{for monoatomic} \\ \frac{7}{5} & \text{for diatomic} \\ \dots & \dots \end{cases}$$

Note that at this point, the above equation for E_{int} is completely unjustified...

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5

From The New Yorker Magazine, November 2003



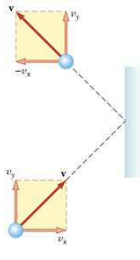
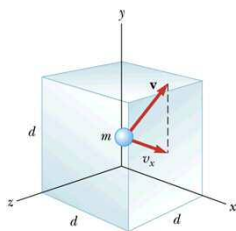
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6

Microscopic model of ideal gas:

Each atom is represented as a tiny hard sphere of mass m with velocity \mathbf{v} . Collisions and forces between atoms are neglected. Collisions with the walls of the container are assumed to be elastic.



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7

What we can show is the pressure exerted by the atoms by their collisions with the walls of the container is given by:

$$P = \frac{2}{3} \frac{N}{V} \frac{1}{2} m \langle v^2 \rangle_{\text{avg}} = \frac{2}{3} \frac{N}{V} \langle K \rangle_{\text{avg}}$$

Proof:

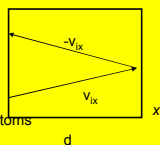
Force exerted on wall perpendicular to x-axis by an atom which collides with it:

$$|F_{ix}| = \left| \frac{\Delta p_{ix}}{\Delta t} \right| = \frac{2m_i |v_{ix}|}{\Delta t} \quad \Delta t \approx 2d / v_{ix}$$

$$\Rightarrow F_{ix} \approx \frac{2m_i v_{ix}}{2d / v_{ix}} = \frac{m_i v_{ix}^2}{d}$$

$$P = \sum_i \frac{F_{ix}}{A} = \sum_i \frac{m_i v_{ix}^2}{dA} = \frac{N}{V} \langle m_i v_{ix}^2 \rangle$$

number of atoms
average over atoms



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8



$$F_{ix} \approx \frac{2m_i v_{ix}}{2d / v_{ix}} = \frac{m_i v_{ix}^2}{d}$$

$$P = \sum_i \frac{F_{ix}}{A} = \sum_i \frac{m_i v_{ix}^2}{dA} = \frac{N}{V} \langle m_i v_{ix}^2 \rangle$$

Since molecules are equally likely to move along x, y, z :

$$\langle m_i v_{ix}^2 \rangle = \langle m_i v_{iy}^2 \rangle = \langle m_i v_{iz}^2 \rangle$$

Also note that $v_i^2 = v_{ix}^2 + v_{iy}^2 + v_{iz}^2$

$$\langle v_i^2 \rangle = \langle v_{ix}^2 \rangle + \langle v_{iy}^2 \rangle + \langle v_{iz}^2 \rangle = 3 \langle v_{ix}^2 \rangle$$

$$\Rightarrow P = \frac{N}{V} \langle m_i v_{ix}^2 \rangle = \frac{N}{3V} \langle m_i v_i^2 \rangle = \frac{2N}{3V} \langle \frac{1}{2} m_i v_i^2 \rangle$$

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9

iclicker question:

What should we call $\left\langle \frac{1}{2} m_i v_i^2 \right\rangle$?

- A. Average kinetic energy of atom.
- B. We cannot use our macroscopic equations at the atomic scale -- so this quantity will go unnamed.
- C. We made too many approximations, so it is not worth naming/discussion.
- D. Very boring.

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10

$$P = \frac{2N}{3V} \left\langle \frac{1}{2} m_i v_i^2 \right\rangle$$

Note: N = number of atoms

m_i = mass of atom (6.6×10^{-27} kg for He atom)

$Nm_i = nM$ where M denotes the molar mass (0.004 kg for He atom)

n = number of moles of atoms

Connection to ideal gas law :

$$PV = n \left(\frac{2}{3} \left\langle \frac{1}{2} M v_i^2 \right\rangle \right) = nRT$$

$$\Rightarrow \left(\frac{2}{3} \left\langle \frac{1}{2} M v_i^2 \right\rangle \right) = RT \quad \text{or} \quad \left\langle \frac{1}{2} M v_i^2 \right\rangle = \frac{3}{2} RT$$

$\left\langle \frac{1}{2} M v_i^2 \right\rangle$ = average kinetic energy of mole of ideal gas atoms

$$E_{\text{int}} = \frac{3}{2} nRT$$

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11

Average atomic velocities:
(note $\langle v_i \rangle = 0$)

$$\left\langle \frac{1}{2} M v_i^2 \right\rangle = \frac{3}{2} RT$$

$$\left\langle v_i^2 \right\rangle = \sqrt{\frac{3RT}{M}}$$

Relationship between
average atomic velocities
with T

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12

For monoatomic ideal gas:

$$E_{\text{int}} = \frac{3}{2} nRT$$

General form for ideal gas (including mono-, di-, poly-atomic ideal gases):

$$E_{\text{int}} = \frac{1}{\gamma - 1} nRT$$

$$\gamma = \begin{cases} \frac{5}{3} & \text{for monoatomic} \\ \frac{7}{5} & \text{for diatomic} \\ \dots & \dots \end{cases}$$

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13

Internal energy of an ideal gas:

$$E_{\text{int}} = N \left\langle \frac{1}{2} m \langle v^2 \rangle \right\rangle = N \frac{3}{2} k_B T \Rightarrow \frac{N}{\gamma - 1} k_B T = \frac{n}{\gamma - 1} RT$$

derived for
monoatomic ideal gas

more general relation
for polyatomic ideal gas

Gas	γ (theory)	γ (exp)
He	5/3	1.67
N ₂	7/5	1.41
H ₂ O	4/3	1.30

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14

Determination of Q for various processes in an ideal gas:

$$E_{\text{int}} = \frac{n}{\gamma - 1} RT$$

$$\Delta E_{\text{int}} = \frac{n}{\gamma - 1} R \Delta T = Q + W$$

Example: Isovolumetric process – ($V = \text{constant} \rightarrow W = 0$)

$$\Delta E_{\text{int } i \rightarrow f} = \frac{n}{\gamma - 1} R \Delta T_{i \rightarrow f} = Q_{i \rightarrow f}$$

In terms of "heat capacity": $Q_{i \rightarrow f} = \frac{n}{\gamma - 1} R \Delta T_{i \rightarrow f} \equiv n C_V \Delta T_{i \rightarrow f}$

$$C_V = \frac{R}{\gamma - 1}$$

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15

Example: Isobaric process ($P=\text{constant}$):

$$\Delta E_{\text{int } i \rightarrow f} = \frac{n}{\gamma-1} R \Delta T_{i \rightarrow f} = Q_{i \rightarrow f} + W_{i \rightarrow f}$$

In terms of "heat capacity":

$$Q_{i \rightarrow f} = \frac{n}{\gamma-1} R \Delta T_{i \rightarrow f} + P_i (V_f - V_i) = \frac{n}{\gamma-1} R \Delta T_{i \rightarrow f} + n R \Delta T_{i \rightarrow f} \equiv n C_p \Delta T_{i \rightarrow f}$$

$$\Rightarrow C_p = \frac{R}{\gamma-1} + R = \frac{\gamma R}{\gamma-1}$$

Note: $\gamma = C_p/C_v$

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16

More examples:

Isothermal process ($T=\text{const}$)

$$E_{\text{int}} = \frac{n}{\gamma-1} R T$$

$$\Delta E_{\text{int}} = \frac{n}{\gamma-1} R \Delta T = Q + W$$

$$\Rightarrow \Delta T = 0 \Rightarrow \Delta E_{\text{int}} = 0 \Rightarrow Q = -W$$

$$-W = \int_{V_i}^{V_f} P dV = n R T \int_{V_i}^{V_f} \frac{dV}{V} = n R T \ln \left(\frac{V_f}{V_i} \right)$$

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17

Even more examples:

Adiabatic process ($Q=0$)

$$\Delta E_{\text{int}} = W$$

$$\frac{n}{\gamma-1} R \Delta T = -P \Delta V$$

$$P V = n R T$$

$$\Delta P V + P \Delta V = n R \Delta T$$

$$n R \Delta T = -(\gamma-1) P \Delta V = \Delta P V + P \Delta V$$

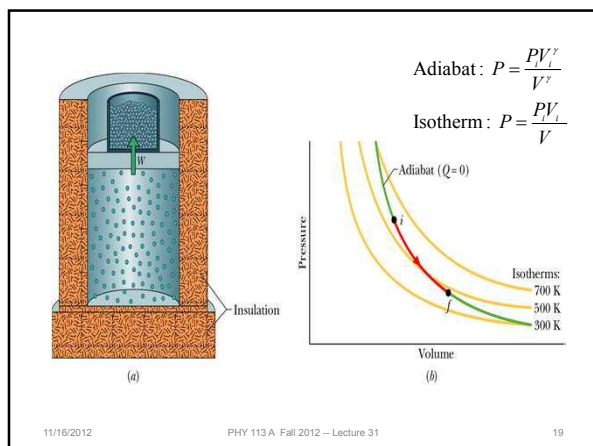
$$-\gamma \frac{\Delta V}{V} = \frac{\Delta P}{P}$$

$$\Rightarrow -\ln \left(\frac{V_f}{V_i} \right) = \ln \left(\frac{P_f}{P_i} \right) \Rightarrow P_i V_i^\gamma = P_f V_f^\gamma$$

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18



iclicker question:

Suppose that an ideal gas expands adiabatically. Does the temperature

- (A) Increase (B) Decrease (C) Remain the same

$$P_i V_i^\gamma = P_f V_f^\gamma$$

$$P_i V_i = nRT_i \Rightarrow P_i = nR \frac{T_i}{V_i}$$

$$T_i V_i^{\gamma-1} = T_f V_f^{\gamma-1}$$

$$T_f = T_i \left(\frac{V_i}{V_f} \right)^{\gamma-1}$$

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20

Review of results from ideal gas analysis in terms of the specific heat ratio $\gamma \equiv C_P/C_V$:

$$\Delta E_{\text{int}} = \frac{n}{\gamma-1} R \Delta T = n C_V \Delta T \quad ; \quad C_V = \frac{R}{\gamma-1}$$

$$C_P = \frac{\gamma R}{\gamma-1}$$

For an isothermal process, $\Delta E_{\text{int}} = 0 \rightarrow Q = -W$

$$-W = \int_{V_i}^{V_f} P dV = nRT \ln \left(\frac{V_f}{V_i} \right) = P_i V_i \ln \left(\frac{V_f}{V_i} \right)$$

For an adiabatic process, $Q = 0$

$$P_i V_i^\gamma = P_f V_f^\gamma$$

$$T_i V_i^{\gamma-1} = T_f V_f^{\gamma-1}$$

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21

Note:

It can be shown that the work done by an ideal gas which has an initial pressure P_i and initial volume V_i when it expands *adiabatically* to a volume V_f is given by:

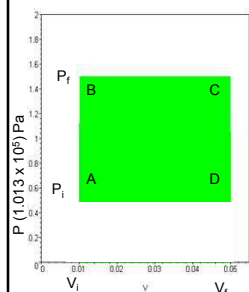
$$W = - \int_{V_i}^{V_f} P dV = - \frac{P_i V_i}{\gamma - 1} \left(1 - \left(\frac{V_i}{V_f} \right)^{\gamma-1} \right)$$

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22

Examples process by an ideal gas:



	A→B	B→C	C→D	D→A
Q	$\frac{V_f(P_f - P_i)}{\gamma - 1}$	$\frac{P_f(V_f - V_i)}{\gamma - 1}$	$\frac{-V_f(P_f - P_i)}{\gamma - 1}$	$\frac{-P_f(V_f - V_i)}{\gamma - 1}$
W	0	$-P_f(V_f - V_i)$	0	$P_i(V_f - V_i)$
ΔE_{int}	$\frac{V_f(P_f - P_i)}{\gamma - 1}$	$\frac{P_f(V_f - V_i)}{\gamma - 1}$	$\frac{-V_f(P_f - P_i)}{\gamma - 1}$	$\frac{-P_f(V_f - V_i)}{\gamma - 1}$

Efficiency as an engine:

$$e = W_{net} / Q_{input}$$

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23
