

**PHY 711 Classical Mechanics and
Mathematical Methods
10-10:50 AM MWF Olin 103**

Plan for Lecture 14:

Finish reading Chapter 6

- 1. Liouville's theorem**
- 2. Hamilton-Jacobi formalism**

Course schedule

(Preliminary schedule -- subject to frequent adjustment.)

	Date	F&W Reading	Topic	Assignment
1	Wed, 8/29/2012	Chap. 1	Review of basic principles; Scattering theory	#1
2	Fri, 8/31/2012	Chap. 1	Scattering theory continued	#2
3	Mon, 9/03/2012	Chap. 1	Scattering theory continued	#3
4	Wed, 9/05/2012	Chap. 1 & 2	Scattering theory/Accelerated coordinate frame	#4
5	Fri, 9/07/2012	Chap. 2	Accelerated coordinate frame	#5
6	Mon, 9/10/2012	Chap. 3	Calculus of Variation	#6
7	Wed, 9/12/2012	Chap. 3	Calculus of Variation continued	
8	Fri, 9/14/2012	Chap. 3	Lagrangian	#7
9	Mon, 9/17/2012	Chap. 3 & 6	Lagrangian	#8
10	Wed, 9/19/2012	Chap. 3 & 6	Lagrangian	#9
11	Fri, 9/21/2012	Chap. 3 & 6	Lagrangian	#10
12	Mon, 9/24/2012	Chap. 3 & 6	Lagrangian and Hamiltonian	#11
13	Wed, 9/26/2012	Chap. 6	Lagrangian and Hamiltonian	#12
14	Fri, 9/28/2012	Chap. 6	Lagrangian and Hamiltonian	#13



Liouville's theorem:

Imagine a collection of particles obeying the Canonical equations of motion in phase space.

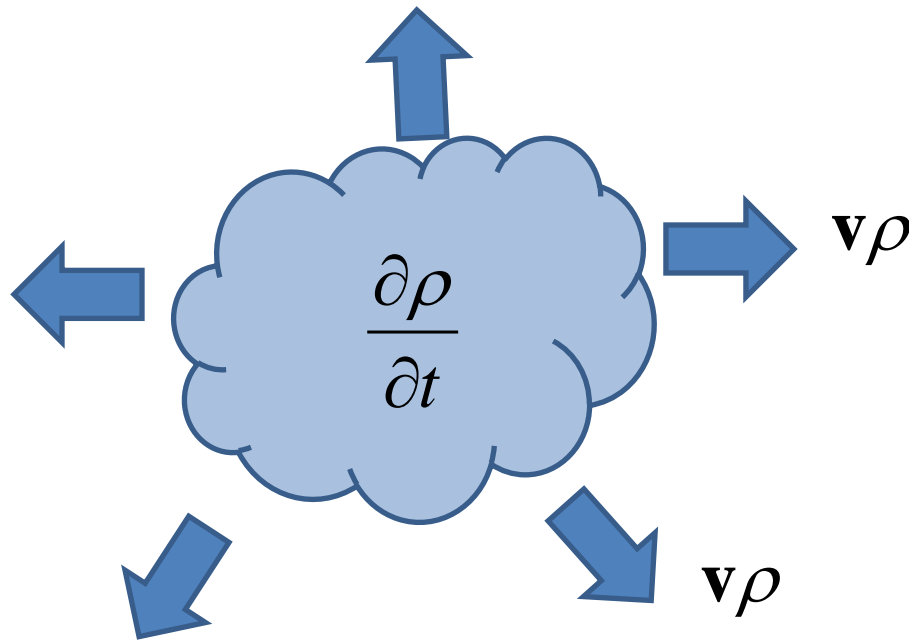
Let ρ denote the "distribution" of particles in phase space :

$$\rho = \rho(\{q_1 \cdots q_{3N}\}, \{p_1 \cdots p_{3N}\}, t)$$

Liouville's theorem shows that :

$$\frac{d\rho}{dt} = 0 \quad \Rightarrow \rho \text{ is constant in time}$$

Proof of Liouville's theorem:



Continuity equation :

$$\frac{\partial \rho}{\partial t} = -\nabla \cdot (\mathbf{v}\rho)$$

Note : in this case, the velocity is the $6N$ dimensional vector :

$$\mathbf{v} = (\dot{\mathbf{r}}_1, \dot{\mathbf{r}}_2, \dots, \dot{\mathbf{r}}_N, \dot{\mathbf{p}}_1, \dot{\mathbf{p}}_2, \dots, \dot{\mathbf{p}}_N)$$

We also have a $6N$ dimensional gradient :

$$\nabla = (\nabla_{\mathbf{r}_1}, \nabla_{\mathbf{r}_2}, \dots, \nabla_{\mathbf{r}_N}, \nabla_{\mathbf{p}_1}, \nabla_{\mathbf{p}_2}, \dots, \nabla_{\mathbf{p}_N})$$

$$\frac{\partial \rho}{\partial t} = -\nabla \cdot (\mathbf{v}\rho)$$

$$= -\sum_{j=1}^{3N} \left[\frac{\partial}{\partial q_j} (\dot{q}_j \rho) + \frac{\partial}{\partial p_j} (\dot{p}_j \rho) \right]$$

$$= -\sum_{j=1}^{3N} \left[\frac{\partial \rho}{\partial q_j} \dot{q}_j + \frac{\partial \rho}{\partial p_j} \dot{p}_j \right] - \rho \sum_{j=1}^{3N} \left[\frac{\partial \dot{q}_j}{\partial q_j} + \frac{\partial \dot{p}_j}{\partial p_j} \right]$$

$$\frac{\partial \dot{q}_j}{\partial q_j} + \frac{\partial \dot{p}_j}{\partial p_j} = \frac{\partial^2 H}{\partial q_j \partial p_j} + \left(-\frac{\partial^2 H}{\partial p_j \partial q_j} \right) = 0$$

$$\frac{\partial \rho}{\partial t} = - \sum_{j=1}^{3N} \left[\frac{\partial \rho}{\partial q_j} \dot{q}_j + \frac{\partial \rho}{\partial p_j} \dot{p}_j \right] - \cancel{\rho \sum_{j=1}^{3N} \left[\frac{\partial \dot{q}_j}{\partial q_j} + \frac{\partial \dot{p}_j}{\partial p_j} \right]} \quad 0$$

$$\frac{\partial \rho}{\partial t} = - \sum_{j=1}^{3N} \left[\frac{\partial \rho}{\partial q_j} \dot{q}_j + \frac{\partial \rho}{\partial p_j} \dot{p}_j \right]$$

$$\Rightarrow \frac{\partial \rho}{\partial t} + \sum_{j=1}^{3N} \left[\frac{\partial \rho}{\partial q_j} \dot{q}_j + \frac{\partial \rho}{\partial p_j} \dot{p}_j \right] = \frac{d\rho}{dt} = 0$$

Notion of “Canonical” distributions

$$q_\sigma = q_\sigma(\{Q_1 \cdots Q_n\}, \{P_1 \cdots P_n\}, t) \quad \text{for each } \sigma$$

$$p_\sigma = p_\sigma(\{Q_1 \cdots Q_n\}, \{P_1 \cdots P_n\}, t) \quad \text{for each } \sigma$$

$$\sum_{\sigma} p_\sigma \dot{q}_\sigma - H(\{q_\sigma\}, \{p_\sigma\}, t) =$$

$$\sum_{\sigma} P_\sigma \dot{Q}_\sigma - \tilde{H}(\{Q_\sigma\}, \{P_\sigma\}, t) + \frac{d}{dt} F(\{q_\sigma\}, \{Q_\sigma\}, t)$$

Apply Hamilton's principle :

$$\delta \int_{t_i}^{t_f} \left[\sum_{\sigma} P_\sigma \dot{Q}_\sigma - \tilde{H}(\{Q_\sigma\}, \{P_\sigma\}, t) + \frac{d}{dt} F(\{q_\sigma\}, \{Q_\sigma\}, t) \right] dt = 0$$

$$\dot{Q}_\sigma = \frac{\partial H}{\partial P_\sigma} \quad \dot{P}_\sigma = -\frac{\partial H}{\partial Q_\sigma}$$

Note that it is conceivable that if we were extraordinarily clever, we could find all of the constants of the motion!

$$\begin{aligned}
 \sum_{\sigma} p_{\sigma} \dot{q}_{\sigma} - H(\{q_{\sigma}\}, \{p_{\sigma}\}, t) &= \\
 \sum_{\sigma} P_{\sigma} \dot{Q}_{\sigma} - \tilde{H}(\{Q_{\sigma}\}, \{P_{\sigma}\}, t) + \frac{d}{dt} F(\{q_{\sigma}\}, \{Q_{\sigma}\}, t) & \\
 \frac{d}{dt} F(\{q_{\sigma}\}, \{Q_{\sigma}\}, t) = \sum_{\sigma} \left(\left(\frac{\partial F}{\partial q_{\sigma}} \right) \dot{q}_{\sigma} + \left(\frac{\partial F}{\partial Q_{\sigma}} \right) \dot{Q}_{\sigma} \right) + \frac{\partial F}{\partial t} & \\
 \sum_{\sigma} \left(p_{\sigma} - \left(\frac{\partial F}{\partial q_{\sigma}} \right) \right) \dot{q}_{\sigma} - H(\{q_{\sigma}\}, \{p_{\sigma}\}, t) = & \\
 \sum_{\sigma} \left(P_{\sigma} + \left(\frac{\partial F}{\partial Q_{\sigma}} \right) \right) \dot{Q}_{\sigma} - \tilde{H}(\{Q_{\sigma}\}, \{P_{\sigma}\}, t) + \frac{\partial F}{\partial t} &
 \end{aligned}$$

$$\sum_{\sigma} \left(p_{\sigma} - \left(\frac{\partial F}{\partial q_{\sigma}} \right) \right) \dot{q}_{\sigma} - H(\{q_{\sigma}\}, \{p_{\sigma}\}, t) =$$

$$\sum_{\sigma} \left(P_{\sigma} + \left(\frac{\partial F}{\partial Q_{\sigma}} \right) \right) \dot{Q}_{\sigma} - \tilde{H}(\{Q_{\sigma}\}, \{P_{\sigma}\}, t) + \frac{\partial F}{\partial t}$$

$$\Rightarrow p_{\sigma} = \left(\frac{\partial F}{\partial q_{\sigma}} \right) \quad P_{\sigma} = - \left(\frac{\partial F}{\partial Q_{\sigma}} \right)$$

$$\Rightarrow \tilde{H}(\{Q_{\sigma}\}, \{P_{\sigma}\}, t) = H(\{q_{\sigma}\}, \{p_{\sigma}\}, t) + \frac{\partial F}{\partial t}$$

Note that it is conceivable that if we were extraordinarily clever, we could find all of the constants of the motion!

$$\dot{Q}_\sigma = \frac{\partial H}{\partial P_\sigma} \quad \dot{P}_\sigma = -\frac{\partial H}{\partial Q_\sigma}$$

$$\text{Suppose : } \dot{Q}_\sigma = \frac{\partial H}{\partial P_\sigma} = 0 \quad \text{and} \quad \dot{P}_\sigma = -\frac{\partial H}{\partial Q_\sigma} = 0$$

$\Rightarrow Q_\sigma, P_\sigma$ are constants of the motion

Possible solution – Hamilton-Jacobi theory:

$$\text{Suppose : } F(\{q_\sigma\}, \{Q_\sigma\}, t) \Rightarrow -\sum_\sigma P_\sigma Q_\sigma + S(\{q_\sigma\}, \{P_\sigma\}, t)$$

$$\begin{aligned}
& \sum_{\sigma} p_{\sigma} \dot{q}_{\sigma} - H(\{q_{\sigma}\}, \{p_{\sigma}\}, t) = \\
& \sum_{\sigma} P_{\sigma} \dot{Q}_{\sigma} - \tilde{H}(\{Q_{\sigma}\}, \{P_{\sigma}\}, t) + \frac{d}{dt} \left(- \sum_{\sigma} P_{\sigma} Q_{\sigma} + S(\{q_{\sigma}\}, \{P_{\sigma}\}, t) \right) \\
& = -\tilde{H}(\{Q_{\sigma}\}, \{P_{\sigma}\}, t) - \sum_{\sigma} \dot{P}_{\sigma} Q_{\sigma} + \sum_{\sigma} \left(\frac{\partial S}{\partial q_{\sigma}} \dot{q}_{\sigma} + \frac{\partial S}{\partial P_{\sigma}} \dot{P}_{\sigma} \right) + \frac{\partial S}{\partial t}
\end{aligned}$$

Solution :

$$p_{\sigma} = \frac{\partial S}{\partial q_{\sigma}} \qquad Q_{\sigma} = \frac{\partial S}{\partial P_{\sigma}}$$

$$\tilde{H}(\{Q_{\sigma}\}, \{P_{\sigma}\}, t) = H(\{q_{\sigma}\}, \{p_{\sigma}\}, t) + \frac{\partial S}{\partial t}$$

When the dust clears :

Assume $\{Q_\sigma\}, \{P_\sigma\}, \tilde{H}$ are constants; choose $\tilde{H} = 0$

Need to find $S(\{q_\sigma\}, \{P_\sigma\}, t)$

$$p_\sigma = \frac{\partial S}{\partial q_\sigma} \quad Q_\sigma = \frac{\partial S}{\partial P_\sigma}$$

$$\Rightarrow H\left(\{q_\sigma\}, \left\{\frac{\partial S}{\partial q_\sigma}\right\}, t\right) + \frac{\partial S}{\partial t} = 0$$

Note: S is the "action":

$$\sum_{\sigma} p_{\sigma} \dot{q}_{\sigma} - H(\{q_{\sigma}\}, \{p_{\sigma}\}, t) =$$

$$\sum_{\sigma} P_{\sigma} \dot{Q}_{\sigma} - \tilde{H}(\{Q_{\sigma}\}, \{P_{\sigma}\}, t) + \frac{d}{dt} \left(- \sum_{\sigma} P_{\sigma} Q_{\sigma} + S(\{q_{\sigma}\}, \{P_{\sigma}\}, t) \right)$$

$$\sum_{\sigma} p_{\sigma} \dot{q}_{\sigma} - H(\{q_{\sigma}\}, \{p_{\sigma}\}, t) =$$

$$\sum_{\sigma} P_{\sigma} \dot{Q}_{\sigma} - \tilde{H}(\{Q_{\sigma}\}, \{P_{\sigma}\}, t) + \frac{d}{dt} \left(- \sum_{\sigma} P_{\sigma} Q_{\sigma} + S(\{q_{\sigma}\}, \{P_{\sigma}\}, t) \right)$$

$$\int_{t_i}^{t_f} \left(\sum_{\sigma} p_{\sigma} \dot{q}_{\sigma} - H(\{q_{\sigma}\}, \{p_{\sigma}\}, t) \right) dt = \int_{t_i}^{t_f} \left(\frac{d}{dt} (S(\{q_{\sigma}\}, \{P_{\sigma}\}, t)) \right) dt$$

$$= S(\{q_{\sigma}\}, \{P_{\sigma}\}, t) \Big|_{t_i}^{t_f}$$

Differential equation for **S**:

$$H\left(\{q_\sigma\}, \left\{\frac{\partial S}{\partial q_\sigma}\right\}, t\right) + \frac{\partial S}{\partial t} = 0$$

Example: $H(\{q\}, \{p\}, t) = \frac{p^2}{2m} + \frac{1}{2}m\omega^2 q^2$

Hamilton - Jacobi Eq: $H\left(\{q\}, \left\{\frac{\partial S}{\partial q}\right\}, t\right) + \frac{\partial S}{\partial t} = 0$

$$\frac{1}{2m}\left(\frac{\partial S}{\partial q}\right)^2 + \frac{1}{2}m\omega^2 q^2 + \frac{\partial S}{\partial t} = 0$$

Assume: $S(q, t) \equiv W(q) - Et$ (E constant)

Continued:

$$\frac{1}{2m} \left(\frac{\partial S}{\partial q} \right)^2 + \frac{1}{2} m \omega^2 q^2 + \frac{\partial S}{\partial t} = 0$$

Assume: $S(q, t) \equiv W(q) - Et$ (E constant)

$$\frac{1}{2m} \left(\frac{dW}{dq} \right)^2 + \frac{1}{2} m \omega^2 q^2 = E$$

$$\frac{dW}{dq} = \sqrt{2mE - (m\omega)^2 q^2}$$

$$W(q) = \int \sqrt{2mE - (m\omega)^2 q^2} dq$$

Continued:

$$W(q) = \int \sqrt{2mE - (m\omega)^2 q^2} dq$$

$$= \frac{1}{2} q \sqrt{2mE - (m\omega)^2 q^2} + \frac{E}{\omega} \sin^{-1} \left(\frac{m\omega q}{\sqrt{2mE}} \right) + C$$

$$S(q, E, t) = \frac{1}{2} q \sqrt{2mE - (m\omega)^2 q^2} + \frac{E}{\omega} \sin^{-1} \left(\frac{m\omega q}{\sqrt{2mE}} \right) - Et$$

$$\frac{\partial S}{\partial E} = Q = \frac{1}{\omega} \sin^{-1} \left(\frac{m\omega q}{\sqrt{2mE}} \right) - t$$

$$\Rightarrow q(t) = \frac{\sqrt{2mE}}{m\omega} \sin(\omega(t + Q))$$