

**PHY 711 Classical Mechanics and  
Mathematical Methods  
10-10:50 AM MWF Olin 103**

**Plan for Lecture 17:**

**Continue reading Chapter 4**

- 1. Normal modes for extended one-dimensional systems**
- 2. Normal modes for 2 and 3 dimensional systems**

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**Course schedule**

(Preliminary schedule -- subject to frequent adjustment.)

Date	F&W Reading	Topic	Assignment
1 Wed. 8/29/2012	Chap. 1	Review of basic principles; Scattering theory	#1
2 Fri. 8/31/2012	Chap. 1	Scattering theory continued	#2
3 Mon. 9/03/2012	Chap. 1	Scattering theory continued	#3
4 Wed. 9/05/2012	Chap. 1 & 2	Scattering theory/Accelerated coordinate frame	#4
5 Fri. 9/07/2012	Chap. 2	Accelerated coordinate frame	#5
6 Mon. 9/10/2012	Chap. 3	Calculus of Variation	#6
7 Wed. 9/12/2012	Chap. 3	Calculus of Variation continued	
8 Fri. 9/14/2012	Chap. 3	Lagrangian	#7
9 Mon. 9/17/2012	Chap. 3 & 6	Lagrangian	#8
10 Wed. 9/19/2012	Chap. 3 & 6	Lagrangian	#9
11 Fri. 9/21/2012	Chap. 3 & 6	Lagrangian	#10
12 Mon. 9/24/2012	Chap. 3 & 6	Lagrangian and Hamiltonian	#11
13 Wed. 9/26/2012	Chap. 6	Lagrangian and Hamiltonian	#12
14 Fri. 9/28/2012	Chap. 6	Lagrangian and Hamiltonian	#13
15 Mon. 10/01/2012	Chap. 4	Small oscillations	#14
16 Wed. 10/03/2012	Chap. 4	Small oscillations	#15
17 Fri. 10/05/2012	Chap. 4	Small oscillations	
18 Mon. 10/08/2012	Chap. 7	Wave equation	Take Home Exam
19 Wed. 10/10/2012	Chap. 7	Wave equation	Take Home Exam
20 Fri. 10/12/2012	Chap. 7	Wave equation	Take Home Exam
21 Mon. 10/15/2012	Chap. 7	Wave equation	Exam due

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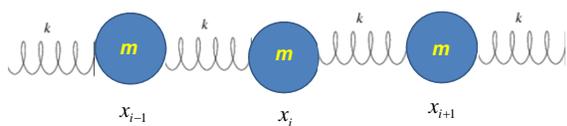
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Consider an infinite system of masses and springs:



Note: each mass coordinate is measured relative to its equilibrium position  $x_i^0$

$$L = T - V = \frac{1}{2} m \sum_{i=0}^{\infty} \dot{x}_i^2 - \frac{1}{2} k \sum_{i=0}^{\infty} (x_{i+1} - x_i)^2$$

Note: In this case we have an infinite number of identical masses and springs.

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In this case, the Euler-Lagrange equations all have the form :

$$m\ddot{x}_i = k(x_{i+1} - 2x_i + x_{i-1})$$

Again try:  $x_j(t) = Ae^{-i\omega t + iqa_j}$

$$-\omega^2 Ae^{-i\omega t + iqa_j} = \frac{k}{m}(e^{iqa} - 2 + e^{-iqa})Ae^{-i\omega t + iqa_j}$$

$$-\omega^2 = \frac{k}{m}(2\cos(qa) - 2)$$

$$\Rightarrow \omega^2 = \frac{4k}{m}\sin^2\left(\frac{qa}{2}\right)$$

$$\omega = 2\sqrt{\frac{k}{m}}\sin\left(\frac{qa}{2}\right)$$

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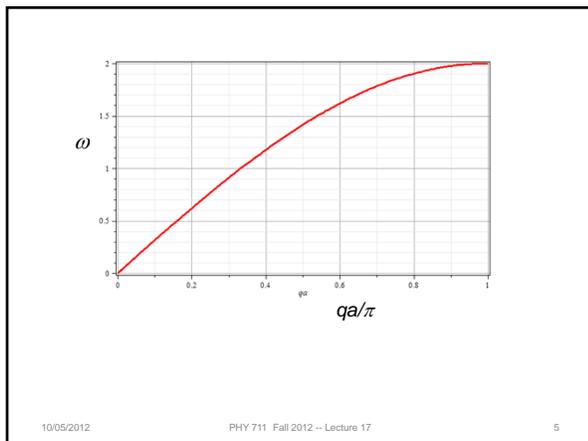
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Consider an infinite system of masses and springs now with two kinds of masses:



Note: each mass coordinate is measured relative to its equilibrium position  $x_i^0, y_i^0, \dots$

$$L = T - V$$

$$= \frac{1}{2}m\sum_{i=0}^{\infty} \dot{x}_i^2 + \frac{1}{2}M\sum_{i=0}^{\infty} \dot{y}_i^2 - \frac{1}{2}k\sum_{i=0}^{\infty} (x_{i+1} - y_i)^2 - \frac{1}{2}k\sum_{i=0}^{\infty} (y_i - x_i)^2$$

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$$L = T - V$$

$$= \frac{1}{2} m \sum_{i=0}^{\infty} \dot{x}_i^2 + \frac{1}{2} M \sum_{i=0}^{\infty} \dot{y}_i^2 - \frac{1}{2} k \sum_{i=0}^{\infty} (x_{i+1} - x_i)^2 - \frac{1}{2} k \sum_{i=0}^{\infty} (y_i - x_i)^2$$

Euler-Lagrange equations :

$$m\ddot{x}_j = k(y_{j-1} - 2x_j + y_j)$$

$$M\ddot{y}_j = k(x_j - 2y_j + x_{j+1})$$

Trial solution :

$$x_j(t) = A e^{-i\omega t + i2qa_j}$$

$$y_j(t) = B e^{-i\omega t + i2qa_j}$$

$$\begin{pmatrix} m\omega^2 - 2k & k(e^{-i2qa} + 1) \\ k(e^{i2qa} + 1) & M\omega^2 - 2k \end{pmatrix} \begin{pmatrix} A \\ B \end{pmatrix} = 0$$

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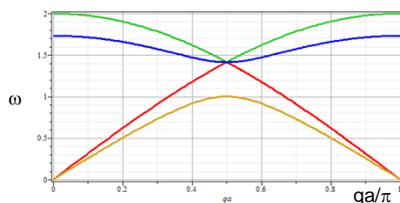
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$$\begin{pmatrix} m\omega^2 - 2k & k(e^{-i2qa} + 1) \\ k(e^{i2qa} + 1) & M\omega^2 - 2k \end{pmatrix} \begin{pmatrix} A \\ B \end{pmatrix} = 0$$

Solutions :

$$\omega_{\pm}^2 = \frac{k}{m} + \frac{k}{M} \pm k \sqrt{\frac{1}{m^2} + \frac{1}{M^2} + \frac{2\cos(2qa)}{mM}}$$



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Eigenvectors:

For  $qa = 0$ :

$$\omega_- = 0 \quad \omega_+ = \sqrt{\frac{2k}{m} + \frac{2k}{M}}$$

$$\begin{pmatrix} A \\ B \end{pmatrix}_- = N \begin{pmatrix} 1 \\ 1 \end{pmatrix} \quad \begin{pmatrix} A \\ B \end{pmatrix}_+ = N \begin{pmatrix} 1 \\ -1 \end{pmatrix}$$

For  $qa = \frac{\pi}{2}$ :

$$\omega_- = \sqrt{\frac{2k}{M}} \quad \omega_+ = \sqrt{\frac{2k}{m}}$$

$$\begin{pmatrix} A \\ B \end{pmatrix}_- = N \begin{pmatrix} 1 \\ 0 \end{pmatrix} \quad \begin{pmatrix} A \\ B \end{pmatrix}_+ = N \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

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Lattice vibrations for 3-dimensional lattice

Example: diamond lattice

Ref: [http://phycomp.technion.ac.il/~nika/diamond\\_structure.html](http://phycomp.technion.ac.il/~nika/diamond_structure.html)

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Atoms located at the positions :

$$\mathbf{R}^a = \mathbf{R}_0^a + \mathbf{u}^a$$

Potential energy function near equilibrium :

$$U(\{\mathbf{R}^a\}) \approx U(\{\mathbf{R}_0^a\}) + \frac{1}{2} \sum_{a,b} (\mathbf{R}^a - \mathbf{R}_0^a) \cdot \left. \frac{\partial^2 U}{\partial \mathbf{R}^a \partial \mathbf{R}^b} \right|_{\{\mathbf{R}_0^a\}} (\mathbf{R}^b - \mathbf{R}_0^b)$$

Define :

$$D_{jk}^{ab} \equiv \left. \frac{\partial^2 U}{\partial \mathbf{R}_j^a \partial \mathbf{R}_k^b} \right|_{\{\mathbf{R}_0^a\}}$$

so that

$$U(\{\mathbf{R}^a\}) \approx U_0 + \frac{1}{2} \sum_{a,b,j,k} u_j^a D_{jk}^{ab} u_k^b$$

$$L(\{u_j^a, \dot{u}_j^a\}) = \frac{1}{2} \sum_{a,j} m_a (\dot{u}_j^a)^2 - U_0 - \frac{1}{2} \sum_{a,b,j,k} u_j^a D_{jk}^{ab} u_k^b$$

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$$L(\{u_j^a, \dot{u}_j^a\}) = \frac{1}{2} \sum_{a,j} m_a (\dot{u}_j^a)^2 - U_0 - \frac{1}{2} \sum_{a,b,j,k} u_j^a D_{jk}^{ab} u_k^b$$

Equations of motion :

$$m_a \ddot{u}_j^a = - \sum_{b,k} D_{jk}^{ab} u_k^b$$

Solution form :

$$u_j^a(t) = \frac{1}{\sqrt{m_a}} A_j^a e^{-i\omega t + i\mathbf{q} \cdot \mathbf{R}_0^a}$$

Details:  $\mathbf{R}_0^a = \boldsymbol{\tau}^a + \mathbf{T}$  where  $\boldsymbol{\tau}^a$  denotes unique sites and  $\mathbf{T}$  denotes replicas

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Define :

$$W_{jk}^{ab}(\mathbf{q}) = \sum_{\mathbf{r}} \frac{D_{jk}^{ab} e^{i\mathbf{q} \cdot (\mathbf{r} - \mathbf{r}')} }{\sqrt{m_a m_b}} e^{i\mathbf{q} \cdot \mathbf{r}}$$

Eigenvalue equations :

$$\omega^2 A_j^a = \sum_{b,k} W(\mathbf{q})_{jk}^{ab} A_k^b$$

In this equation the summation is only over unique atomic sites.

⇒ Find "dispersion curves"  $\omega(\mathbf{q})$

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B. P. Pandy and B. Dayal, J. Phys. C. Solid State Phys. 6 2943 (1973)

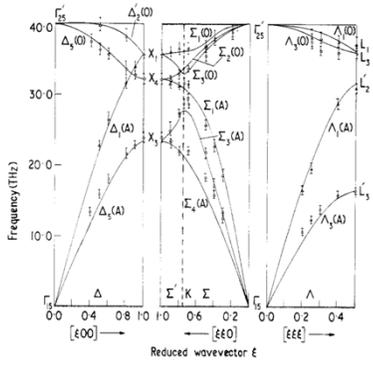


Figure 2. Phonon dispersion curves of diamond. Experimental points *et al* (1965, 1967).  $\Delta$  and  $\circ$  represent the longitudinal and transverse modes.

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