

# **PHY 711 Classical Mechanics and Mathematical Methods**

## **10-10:50 AM MWF Olin 103**

**Plan for Lecture 31:**

**Wave equation for sound**

- 1. Standing waves**
- 2. Green's function for wave  
equation; wave scattering**

8	Fri, 9/14/2012	Chap. 3	Lagrangian	#7
9	Mon, 9/17/2012	Chap. 3 & 6	Lagrangian	#8
10	Wed, 9/19/2012	Chap. 3 & 6	Lagrangian	#9
11	Fri, 9/21/2012	Chap. 3 & 6	Lagrangian	#10
12	Mon, 9/24/2012	Chap. 3 & 6	Lagrangian and Hamiltonian	#11
13	Wed, 9/26/2012	Chap. 6	Lagrangian and Hamiltonian	#12
14	Fri, 9/28/2012	Chap. 6	Lagrangian and Hamiltonian	#13
15	Mon, 10/01/2012	Chap. 4	Small oscillations	#14
16	Wed, 10/03/2012	Chap. 4	Small oscillations	#15
17	Fri, 10/05/2012	Chap. 4	Small oscillations	
18	Mon, 10/08/2012	Chap. 7	Wave equation	Take Home Exam
19	Wed, 10/10/2012	Chap. 7	Wave equation	Take Home Exam
20	Fri, 10/12/2012	Chap. 7	Wave equation	Take Home Exam
21	Mon, 10/15/2012	Chap. 7	Wave equation	Exam due
22	Wed, 10/17/2012	Chap. 7, 5	Moment of inertia	
	Fri, 10/19/2012		Fall break	
23	Mon, 10/22/2012	Chap. 5	Rigid body rotation	#16
24	Wed, 10/24/2012	Chap. 5	Rigid body rotation	#17
25	Fri, 10/26/2012	Chap. 5	Rigid body rotation	#18
26	Mon, 10/29/2012	Chap. 8	Waves in elastic membranes	#19
27	Wed, 10/31/2012	Chap. 9	Introduction to hydrodynamics	
28	Fri, 11/01/2012	Chap. 9	Introduction to hydrodynamics	
29	Mon, 11/05/2012	Chap. 9	Introduction to hydrodynamics	#20
30	Wed, 11/07/2012	Chap. 9	Sound waves	
31	Fri, 11/09/2012	Chap. 9	Linear sound waves	#21



Linearization of the fluid dynamics relations:

Newton - Euler equation of motion :

$$\frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} = \mathbf{f}_{\text{applied}} - \frac{\nabla p}{\rho}$$

Continuity equation :  $\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0$

Near equilibrium :

$$\rho = \rho_0 + \delta\rho$$

$$p = p_0 + \delta p$$

$$\mathbf{v} = 0 + \delta\mathbf{v}$$

$$\mathbf{f}_{\text{applied}} = 0$$

Equations to lowest order in perturbation :

$$\frac{\partial \delta \mathbf{v}}{\partial t} = -\frac{\nabla \delta p}{\rho_0}$$

$$\frac{\partial \delta \rho}{\partial t} + \rho_0 \nabla \cdot \delta \mathbf{v} = 0$$

Velocity potential :  $\delta \mathbf{v} = -\nabla \Phi$

Pressure in terms of the density :

$$\delta p = \left( \frac{\partial p}{\partial \rho} \right)_{s, \rho_0, p_0} \delta \rho \equiv c^2 \delta \rho$$

$$\frac{\partial^2 \Phi}{\partial t^2} - c^2 \nabla^2 \Phi = 0$$

Wave equation for air :

$$\frac{\partial^2 \Phi}{\partial t^2} - c^2 \nabla^2 \Phi = 0$$

Here,  $c^2 = \left( \frac{\partial p}{\partial \rho} \right)_s = \frac{\gamma p_0}{\rho_0}$

$$\mathbf{v} = -\nabla \Phi$$

Boundary values :

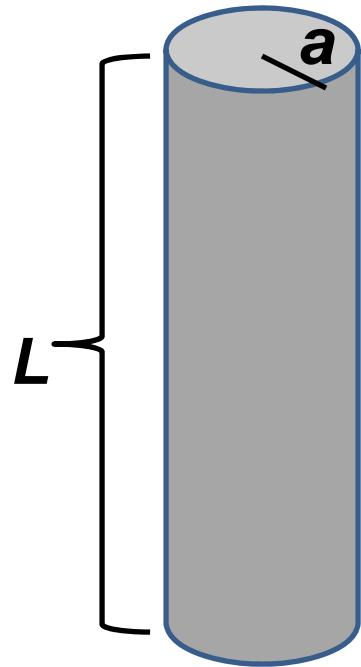
Impenetrable surface with normal  $\hat{\mathbf{n}}$  moving at velocity  $\mathbf{V}$  :

$$\hat{\mathbf{n}} \cdot \mathbf{V} = \hat{\mathbf{n}} \cdot \delta \mathbf{v} = -\hat{\mathbf{n}} \cdot \nabla \Phi$$

Free surface :

$$\delta p = 0 \quad \Rightarrow \rho_0 \frac{\partial \Phi}{\partial t} = 0$$

## Time harmonic standing waves in a pipe



$$\frac{\partial^2 \Phi}{\partial t^2} - c^2 \nabla^2 \Phi = 0$$

Boundary values :

At fixed surface :  $\hat{\mathbf{n}} \cdot \nabla \Phi = 0$

At free surface :  $\frac{\partial \Phi}{\partial t} = 0$

$$\nabla^2 \Phi - \frac{1}{c^2} \frac{\partial^2 \Phi}{\partial t^2} = 0 \quad \text{Define : } k \equiv \frac{\omega}{c}$$

In cylindrical coordinates :

$$\Phi(r, \varphi, z, t) = R(r)F(\varphi)Z(z)e^{-i\omega t}$$

$$\nabla^2 = \frac{1}{r} \frac{\partial}{\partial r} r \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \varphi^2} + \frac{\partial^2}{\partial z^2}$$

$$\left( \frac{1}{r} \frac{\partial}{\partial r} r \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \varphi^2} + \frac{\partial^2}{\partial z^2} + k^2 \right) \Phi(r, \varphi, z, t) = 0$$

$$\left( \frac{1}{r} \frac{\partial}{\partial r} r \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \varphi^2} + \frac{\partial^2}{\partial z^2} + k^2 \right) \Phi(r, \varphi, z, t) = 0$$

$$\Phi(r, \varphi, z, t) = R(r) F(\varphi) Z(z) e^{-i\omega t}$$

$$F(\varphi) = e^{im\varphi}; \quad F(\varphi) = F(\varphi + 2\pi N) \Rightarrow m = \text{integer}$$

$$Z(z) = e^{i\alpha z}; \quad \alpha = \text{real plus other restrictions}$$

$$\left( \frac{d^2}{dr^2} + \frac{1}{r} \frac{d}{dr} - \frac{m^2}{r^2} - \alpha^2 + k^2 \right) R(r) = 0$$

$$\left( \frac{d^2}{dr^2} + \frac{1}{r} \frac{d}{dr} - \frac{m^2}{r^2} - \alpha^2 + k^2 \right) R(r) = 0$$

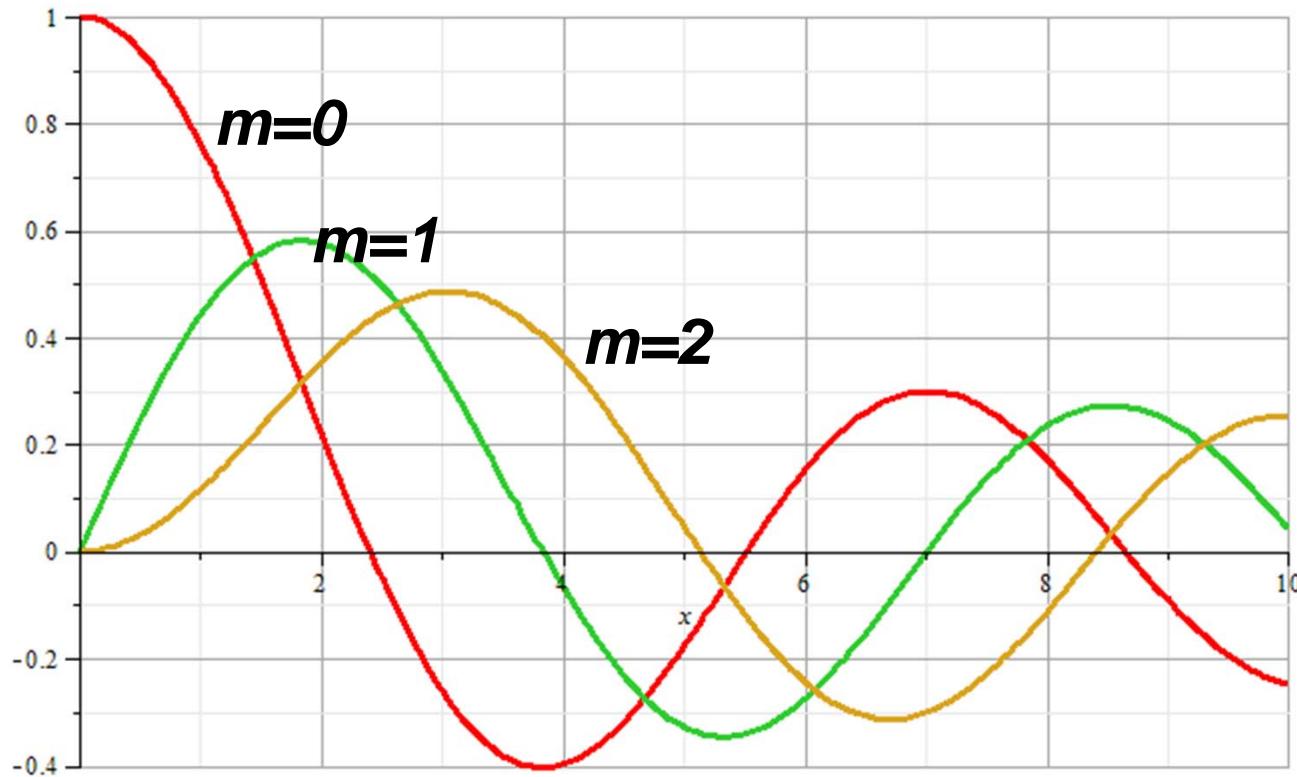
For  $k^2 \geq \alpha^2$  define  $\kappa^2 \equiv k^2 - \alpha^2$

$$\left( \frac{d^2}{dr^2} + \frac{1}{r} \frac{d}{dr} - \frac{m^2}{r^2} + \kappa^2 \right) R(r) = 0$$

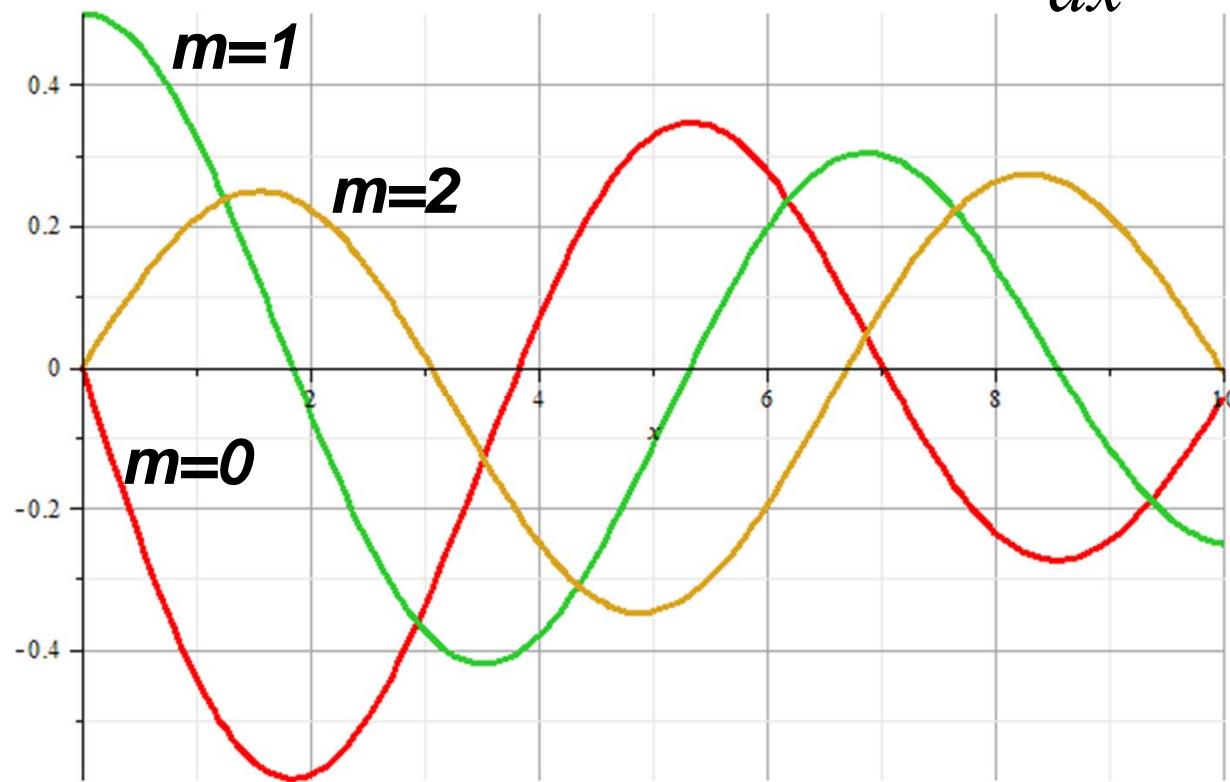
Cylinder surface boundary conditions :  $\left. \frac{dR}{dr} \right|_{r=a} = 0$

$\Rightarrow R(r) = J_m(kr)$  where for  $\frac{dJ_m(x')}{dx} = 0$ ,  $\kappa_{mn} = \frac{x'}{a}$

# Bessel functions : $J_m(x)$



Bessel function derivatives :  $\frac{dJ_m(x)}{dx}$



Zeros of derivatives:   
  $m=0$ : 0.00000, 3.83171, 7.01559  
  $m=1$ : 1.84118, 5.33144, 8.53632  
  $m=2$ : 3.05424, 6.70613, 9.96947

Boundary condition for  $z=0$ ,  $z=L$ :

For open - open pipe :

$$Z(0) = Z(L) = 0 \quad \Rightarrow \quad Z(z) = \sin\left(\frac{p\pi z}{L}\right)$$

$$\Rightarrow \alpha_p = \frac{p\pi}{L}, \quad p = 1, 2, 3, \dots$$

Resonant frequencies :

$$\frac{\omega^2}{c^2} = k^2 = \kappa_{mn}^2 + \alpha_p^2$$

$$k_{mnp}^2 = \left(\frac{x'_{mn}}{a}\right)^2 + \left(\frac{\pi p}{L}\right)^2$$

## Example

$$k_{mnp}^2 = \left( \frac{x'_{mn}}{a} \right)^2 + \left( \frac{\pi p}{L} \right)^2 = \left( \frac{\pi p}{L} \right)^2 \left( 1 + \left( \frac{L}{a} \right)^2 \left( \frac{x'_{mn}}{\pi p} \right)^2 \right)$$

$$\pi p = 3.14, 6.28, 9.42\dots$$

$$x'_{mn} = 0.00, 1.84, 3.05$$

Alternate boundary condition for  $z=0$ ,  $z=L$ :

For open - closed pipe :

$$\frac{dZ(0)}{dz} = Z(L) = 0 \quad \Rightarrow \quad Z(z) = \cos\left(\frac{(2p+1)\pi z}{2L}\right)$$

$$\Rightarrow \alpha_p = \frac{(2p+1)\pi}{2L}, \quad p = 0, 1, 2, 3, \dots$$

$$k_{mnp}^2 = \left(\frac{x'_{mn}}{a}\right)^2 + \left(\frac{\pi(2p+1)}{2L}\right)^2$$

Other solutions to wave equation:

$$\nabla^2 \Phi - \frac{1}{c^2} \frac{\partial^2 \Phi}{\partial t^2} = 0$$

Plane wave solution :

$$\Phi(\mathbf{r}, t) = A e^{i \mathbf{k} \cdot \mathbf{r} - i \omega t} \quad \text{where} \quad k^2 = \left( \frac{\omega}{c} \right)^2$$

Wave equation with source:

$$\nabla^2 \Phi - \frac{1}{c^2} \frac{\partial^2 \Phi}{\partial t^2} = -f(\mathbf{r}, t)$$

Solution in terms of Green's function :

$$\Phi(\mathbf{r}, t) = \int d^3 r' \int dt' G(\mathbf{r} - \mathbf{r}', t - t') f(\mathbf{r}', t')$$

where

$$\left( \nabla^2 - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} \right) G(\mathbf{r} - \mathbf{r}', t - t') = -\delta(\mathbf{r} - \mathbf{r}') \delta(t - t')$$

Wave equation with source -- continued:

We can show that :

$$G(\mathbf{r} - \mathbf{r}', t - t') = \frac{\delta\left(t' - \left(t \mp \frac{|\mathbf{r} - \mathbf{r}'|}{c}\right)\right)}{4\pi|\mathbf{r} - \mathbf{r}'|}$$