

**PHY 711 Classical Mechanics and  
Mathematical Methods  
10-10:50 AM MWF Olin 103**

**Plan for Lecture 31:**

**Wave equation for sound**

- 1. Non-linear effects in traveling sound wave**
- 2. Shock wave**

11/15/2013

PHY 711 Fall 2013 – Lecture 31

1

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#	Day	Date	Topic	Notes
23	Mon	10/21/2013	Chap. 5 Rigid rotations	#20
24	Wed	10/23/2013	Chap. 5 Rigid rotations	#21
25	Fri	10/25/2013	Chap. 5 Rigid rotations	
	Mon	10/28/2013	No class	Take-home exam
	Wed	10/30/2013	No class	Take-home exam
	Fri	11/01/2013	No class	Take-home exam
26	Mon	11/04/2013	Chap. 8 Oscillations in two-dimensional membranes	Take-home exam due
27	Wed	11/06/2013	Chap. 9 Physics of fluids	#22
28	Fri	11/08/2013	Chap. 9 Physics of fluids	#23
29	Mon	11/11/2013	Chap. 9 Sound Waves	#24
30	Wed	11/13/2013	Chap. 9 Sound Waves	#25
31	Fri	11/15/2013	Chap. 9 Non linear effects in Sound	#26
32	Mon	11/18/2013	Chap. 10 Surface waves	
33	Wed	11/20/2013	Chap. 10 Surface waves	
34	Fri	11/22/2013		
35	Mon	11/25/2013		
	Wed	11/27/2013	Thanksgiving Holiday	
	Fri	11/29/2013	Thanksgiving Holiday	
36	Mon	12/02/2013	Student presentations I	
37	Wed	12/04/2013	Student presentations II	
38	Fri	12/06/2013	Student presentations III	
	Mon	12/09/2013	Begin Take-home final	

11/15/2013

PHY 711 Fall 2013 – Lecture 31

2

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**Effects of nonlinearities in fluid equations  
-- one dimensional case**

Newton - Euler equation of motion :

$$\frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} = \mathbf{f}_{\text{applied}} - \frac{\nabla p}{\rho}$$

Continuity equation :  $\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0$

Assume that spatial variation confined to x direction ;  
assume that  $\mathbf{v} = v \hat{\mathbf{x}}$  and  $\mathbf{f}_{\text{applied}} = 0$ .

$$\frac{\partial v}{\partial t} + v \frac{\partial v}{\partial x} = - \frac{1}{\rho} \frac{\partial p}{\partial x}$$

$$\frac{\partial \rho}{\partial t} + v \frac{\partial \rho}{\partial x} + \rho \frac{\partial v}{\partial x} = 0$$

11/15/2013

PHY 711 Fall 2013 – Lecture 31

3

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$$\frac{\partial v}{\partial t} + v \frac{\partial v}{\partial x} + \frac{1}{\rho} \frac{\partial p}{\partial x} = 0$$

$$\frac{\partial \rho}{\partial t} + v \frac{\partial \rho}{\partial x} + \rho \frac{\partial v}{\partial x} = 0$$

Expressing  $p$  in terms of  $\rho$ :  $p = p(\rho)$

$$\frac{\partial p}{\partial x} = \frac{\partial p}{\partial \rho} \frac{\partial \rho}{\partial x} \equiv c^2(\rho) \frac{\partial \rho}{\partial x} \quad \text{where} \quad \frac{\partial p}{\partial \rho} \equiv c^2(\rho)$$

For adiabatic ideal gas :

$$c^2(\rho) = \frac{\mathcal{P}}{\rho} = c_0^2 \left( \frac{\rho}{\rho_0} \right)^{\gamma-1} \quad \text{where} \quad c_0^2 \equiv \frac{\mathcal{P}_0}{\rho_0}$$

11/15/2013

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4

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$$\frac{\partial v}{\partial t} + v \frac{\partial v}{\partial x} + \frac{c^2(\rho)}{\rho} \frac{\partial \rho}{\partial x} = 0$$

$$\frac{\partial \rho}{\partial t} + v \frac{\partial \rho}{\partial x} + \rho \frac{\partial v}{\partial x} = 0$$

Expressing variation  $v$  in terms of  $v(\rho)$ :

$$\frac{\partial v}{\partial \rho} \frac{\partial \rho}{\partial t} + v \frac{\partial v}{\partial \rho} \frac{\partial \rho}{\partial x} + \frac{c^2(\rho)}{\rho} \frac{\partial \rho}{\partial x} = 0$$

$$\frac{\partial \rho}{\partial t} + v \frac{\partial \rho}{\partial x} + \rho \frac{\partial v}{\partial \rho} \frac{\partial \rho}{\partial x} = 0$$

11/15/2013

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5

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Some more algebra :

$$\text{From Euler equation : } \frac{\partial v}{\partial \rho} \left( \frac{\partial \rho}{\partial t} + v \frac{\partial \rho}{\partial x} \right) + \frac{c^2(\rho)}{\rho} \frac{\partial \rho}{\partial x} = 0$$

$$\text{From continuity equation : } \frac{\partial \rho}{\partial t} + v \frac{\partial \rho}{\partial x} = -\rho \frac{\partial v}{\partial \rho} \frac{\partial \rho}{\partial x}$$

$$\text{Combined equation : } \frac{\partial v}{\partial \rho} \left( -\rho \frac{\partial v}{\partial \rho} \frac{\partial \rho}{\partial x} \right) + \frac{c^2(\rho)}{\rho} \frac{\partial \rho}{\partial x} = 0$$

$$\Rightarrow \left( \frac{\partial v}{\partial \rho} \right)^2 = \frac{c^2(\rho)}{\rho^2} \quad \frac{\partial v}{\partial \rho} = \pm \frac{c}{\rho}$$

$$\Rightarrow \frac{\partial \rho}{\partial t} + (v \pm c) \frac{\partial \rho}{\partial x} = 0$$

11/15/2013

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6

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Assuming adiabatic process :  $c^2 = c_0^2 \left( \frac{\rho}{\rho_0} \right)^{\gamma-1}$   $c_0^2 = \frac{\gamma P_0}{\rho_0}$

$$\frac{\partial v}{\partial \rho} = \frac{dv}{d\rho} = \pm \frac{c}{\rho} \Rightarrow v = \pm c_0 \int \left( \frac{\rho'}{\rho_0} \right)^{(\gamma-1)/2} \frac{d\rho'}{\rho'}$$

$$\Rightarrow v = \pm \frac{2c_0}{\gamma-1} \left( \left( \frac{\rho}{\rho_0} \right)^{(\gamma-1)/2} - 1 \right)$$

$$\Rightarrow c = c_0 \left( \frac{\rho}{\rho_0} \right)^{(\gamma-1)/2}$$

11/15/2013

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7

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Summary :

$$\frac{dv}{d\rho} = \pm \frac{c}{\rho}$$

$$\frac{\partial \rho}{\partial t} + (v \pm c) \frac{\partial \rho}{\partial x} = 0$$

Assuming adiabatic process :  $c^2 = c_0^2 \left( \frac{\rho}{\rho_0} \right)^{\gamma-1}$   $c_0^2 = \frac{\gamma P_0}{\rho_0}$

$$c = c_0 \left( \frac{\rho}{\rho_0} \right)^{(\gamma-1)/2} \quad v = \pm \frac{2c_0}{\gamma-1} \left( \left( \frac{\rho}{\rho_0} \right)^{(\gamma-1)/2} - 1 \right)$$

11/15/2013

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8

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Traveling wave solution:

Assume :  $\rho = \rho_0 + f(x - u(\rho)t)$

Need to find self-consistent equations for propagation velocity  $u(\rho)$  using equations

From previous derivations :  $\frac{\partial \rho}{\partial t} + (v \pm c) \frac{\partial \rho}{\partial x} = 0$

Apparently :  $u(\rho) \Leftrightarrow v \pm c$

For adiabatic ideal gas and + signs :

$$u = v + c = c_0 \left( \frac{\gamma+1}{\gamma-1} \left( \frac{\rho}{\rho_0} \right)^{(\gamma-1)/2} - \frac{2}{\gamma-1} \right)$$

11/15/2013

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9

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Traveling wave solution -- continued:

$$\frac{\partial \rho}{\partial t} + (v \pm c) \frac{\partial \rho}{\partial x} = 0$$

Assume:  $\rho = \rho_0 + f(x - u(\rho)t) = \rho_0 + f(x - (v \pm c)t)$

For adiabatic ideal gas and + signs:

$$u = v + c = c_0 \left( \frac{\gamma + 1}{\gamma - 1} \left( \frac{\rho}{\rho_0} \right)^{(\gamma-1)/2} - \frac{2}{\gamma - 1} \right)$$

Solution in linear approximation:

$$u = v + c \approx v_0 + c_0 = c_0 \left( \frac{\gamma + 1}{\gamma - 1} - \frac{2}{\gamma - 1} \right) = c_0$$

$$\Rightarrow \rho = \rho_0 + f(x - c_0 t)$$

11/15/2013 PHY 711 Fall 2013 - Lecture 31

10

Traveling wave solution -- full non-linear case:

Visualization for particular waveform:  $\rho = \rho_0 + f(x - u(\rho)t)$

Assume:  $f(w) \equiv f_0 s(w)$

$$\frac{\rho}{\rho_0} = 1 + \frac{f_0}{\rho_0} s(x - ut)$$

For adiabatic ideal gas:

$$u = c_0 \left( \frac{\gamma + 1}{\gamma - 1} \left( \frac{\rho}{\rho_0} \right)^{(\gamma-1)/2} - \frac{2}{\gamma - 1} \right)$$

$$u = c_0 \left( \frac{\gamma + 1}{\gamma - 1} \left( 1 + \frac{f_0}{\rho_0} s(x - ut) \right)^{(\gamma-1)/2} - \frac{2}{\gamma - 1} \right)$$

11/15/2013

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11

Visualization continued:

$$u = c_0 \left( \frac{\gamma + 1}{\gamma - 1} \left( 1 + \frac{f_0}{\rho_0} s(x - ut) \right)^{(\gamma-1)/2} - \frac{2}{\gamma - 1} \right)$$

Plot  $s(x - ut)$  for fixed  $t$ , as a function of  $x$ :

Let  $w = x - ut$

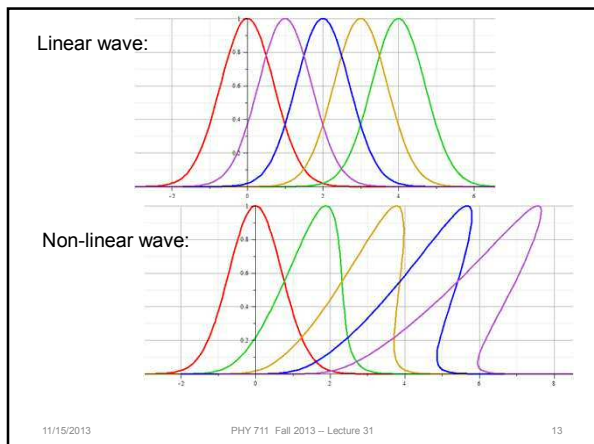
$$x = w + ut = w + u(w)t$$

$$u(w) = c_0 \left( \frac{\gamma + 1}{\gamma - 1} \left( 1 + \frac{f_0}{\rho_0} s(w) \right)^{(\gamma-1)/2} - \frac{2}{\gamma - 1} \right)$$

11/15/2013

PHY 711 Fall 2013 - Lecture 31

12




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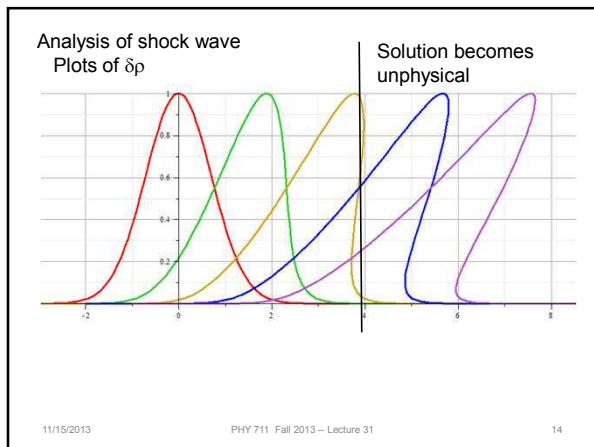
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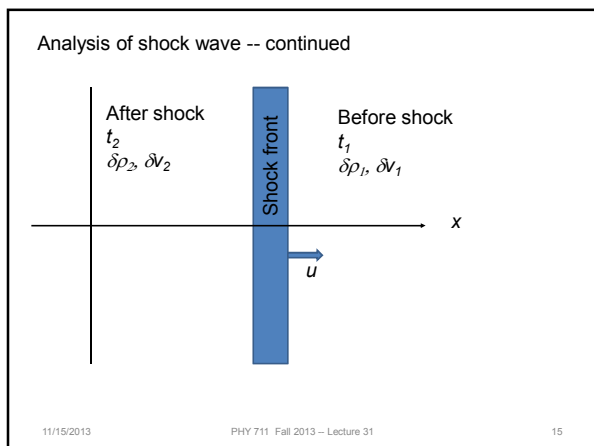
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## Analysis of shock wave – continued

While analysis in the shock region is complicated, we can use conservation laws to analyze regions 1 and 2

Assume  $\rho(x,t) = \rho(x-ut)$

$$p(x,t) = p(x-ut)$$

$$v(x,t) = v(x-ut)$$

Continuity equation :

$$\frac{\partial \rho}{\partial t} + \frac{\partial(\rho v)}{\partial x} = 0 = \frac{\partial(\rho v - \rho u)}{\partial x} \Rightarrow (v_2 - u)\rho_2 = (v_1 - u)\rho_1$$

Conservation of momentum :

$$\Rightarrow p_2 + \rho_2(v_2 - u)^2 = p_1 + \rho_1(v_1 - u)^2$$

11/15/2013

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16

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## Analysis of shock wave – continued

For adiabatic ideal gas, also considering energy conservation:

$$\frac{\rho_2}{\rho_1} = \frac{\frac{\gamma+1}{\gamma-1} p_2 + 1}{\frac{\gamma+1}{\gamma-1} p_1 + \frac{p_2}{p_1}} \leq \frac{\gamma+1}{\gamma-1}$$

Velocity relationships:

$$\frac{(v_1 - u)^2}{c_1^2} = \frac{1}{2\gamma} \left( \gamma - 1 + (\gamma + 1) \frac{p_2}{p_1} \right) \quad \frac{(v_2 - u)^2}{c_2^2} = \frac{1}{2\gamma} \left( \gamma - 1 + (\gamma + 1) \frac{p_1}{p_2} \right)$$

11/15/2013

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17

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## Analysis of shock wave – continued

For adiabatic ideal gas, entropy considerations:

$$\text{Internal energy density: } \varepsilon = \frac{p}{(\gamma-1)\rho} = C_v T$$

$$\text{First law of thermo: } d\varepsilon = Tds - pd\left(\frac{1}{\rho}\right)$$

$$ds = \frac{1}{T} \left( d\left(\frac{p}{(\gamma-1)\rho}\right) - pd\left(\frac{1}{\rho}\right) \right) = C_v d \ln\left(\frac{p}{\rho^\gamma}\right)$$

$$s = C_v \ln\left(\frac{p}{\rho^\gamma}\right) + (\text{constant})$$

$$s_2 - s_1 = C_v \ln\left(\frac{p_2}{p_1} \left(\frac{\rho_1}{\rho_2}\right)^\gamma\right)$$

11/15/2013

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18

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