Basic assumptions*

We assume that we have in incompressible fluid ($\rho = \text{constant}$) a velocity potential of the form $\Phi(x, z, t)$, where

$$\mathbf{v}(x,z,t) = -\nabla\Phi(x,z,t). \tag{1}$$

The surface of the fluid is described by $z = h + \zeta(x, t)$. It is assumed that the fluid is contained in a structure (lake, river, swimming pool, etc.) with a structureless bottom defined by the z = 0plane and filled to an equilibrium height of z = h. We assume that we have in incompressible fluid (ρ = constant) a velocity potential of the form $\Phi(x, z, t)$, where

$$\mathbf{v}(x,z,t) = -\nabla\Phi(x,z,t). \tag{2}$$

The surface of the fluid is described by $z = h + \zeta(x, t)$. It is assumed that the fluid is contained in a structure (lake, river, swimming pool, etc.) with a structureless bottom defined by the z = 0plane and filled to an equilibrium height of z = h.

^{*} Note that these derviations follow Alexander L. Fetter and John Dirk Walecka, **Theoretical Mechanics of Particles and Continua**, (McGraw Hill, 1980), Chapt. 10.



Defining equations for $\Phi(x, z, t)$ **and** $\zeta(x, t)$ $0 \le z \le h + \zeta(x, t)$

Continuity equation

$$\nabla \cdot \mathbf{v} = 0 \quad \Rightarrow \quad \frac{\partial^2 \Phi(x, z, t)}{\partial x^2} + \frac{\partial^2 \Phi(x, z, t)}{\partial z^2} = 0$$
 (3)

Bernoulli equation (assuming irrotational flow) and gravitation potential energy

$$-\frac{\partial\Phi(x,z,t)}{\partial t} + \frac{1}{2}\left[\left(\frac{\partial\Phi(x,z,t)}{\partial x}\right)^2 + \left(\frac{\partial\Phi(x,z,t)}{\partial z}\right)^2\right] + g(z-h) = 0.$$
(4)



Boundary conditions on functions

Zero vertical velocity at bottom of tank

$$\frac{\partial \Phi(x,0,t)}{\partial z} = 0. \tag{5}$$

Consistent vertical velocity at water surface

$$v_z(x,z,t)\rfloor_{z=h+\zeta} = \frac{d\zeta}{dt} = \mathbf{v} \cdot \nabla\zeta + \frac{\partial\zeta}{\partial t}.$$
(6)

$$-\frac{\partial\Phi(x,z,t)}{\partial z} + \frac{\partial\Phi(x,z,t)}{\partial x}\frac{\partial\zeta(x,t)}{\partial x} - \frac{\partial\zeta(x,t)}{\partial t}\Big]_{z=h+\zeta} = 0$$
(7)



Analysis assuming water height *z* is small relative to variations in the direction of wave motion (*x*)

Taylor's expansion about z = 0

$$\Phi(x,z,t) \approx \Phi(x,0,t) + z \frac{\partial \Phi}{\partial z}(x,0,t) + \frac{z^2}{2} \frac{\partial^2 \Phi}{\partial z^2}(x,0,t) + \frac{z^3}{3!} \frac{\partial^3 \Phi}{\partial z^3}(x,0,t) + \frac{z^4}{4!} \frac{\partial^4 \Phi}{\partial z^4}(x,0,t) \cdots$$
(8)

Note that the zero vertical velocity at the bottom ensures that all odd derivatives $\frac{\partial^n \Phi}{\partial z^n}(x, 0, t)$ vanish from the Taylor expansion. In addition, the Laplace equation allows us to convert all even derivatives with respect to z to derivatives with respect to x.

Modified Taylor's expansion

$$\Phi(x,z,t) \approx \Phi(x,0,t) - \frac{z^2}{2} \frac{\partial^2 \Phi}{\partial x^2}(x,0,t) + \frac{z^4}{4!} \frac{\partial^4 \Phi}{\partial x^4}(x,0,t) \cdots$$
(9)



Linearized equations and their solution

Bernoulli equation evaluated at $z = h + \zeta(x, t)$

$$-\frac{\partial\Phi(x,h,t)}{\partial t} + g\zeta(x,t) = 0$$
(10)

Consistent vertical velocity at $z = h + \zeta(x, t)$

$$-\frac{\partial\Phi(x,z,t)}{\partial z} - \frac{\partial\zeta(x,t)}{\partial t}\bigg|_{z=h+\zeta} = 0$$
(11)

Using Taylor's expansion results to lowest order

$$-\frac{\partial \Phi(x,z,t)}{\partial z} \approx h \frac{\partial^2 \Phi(x,0,t)}{\partial x^2} - \frac{\partial \Phi(x,h,t)}{\partial t} \approx -\frac{\partial \Phi(x,0,t)}{\partial t}$$
(12)
Decoupled equations

$$\frac{\partial^2 \Phi(x,0,t)}{\partial t^2} = gh \frac{\partial^2 \Phi(x,0,t)}{\partial x^2}.$$
(13)



Nonlinear equations and solutions – let $\phi(x, t) \equiv \Phi(x, 0, t)$

Approximate forms of Bernoulli equation evaluated at $z = h + \zeta$ surface

$$-\frac{\partial\phi}{\partial t} + \frac{(h+\zeta)^2}{2} \frac{\partial^3\phi}{\partial t\partial x^2} + \frac{1}{2} \left[\left(\frac{\partial\phi}{\partial x} \right)^2 + \left((h+\zeta) \frac{\partial^2\phi}{\partial x^2} \right)^2 \right] + g\zeta = 0 \quad (14)$$
$$-\frac{\partial\phi}{\partial t} + \frac{h^2}{2} \frac{\partial^3\phi}{\partial t\partial x^2} + \frac{1}{2} \left(\frac{\partial\phi}{\partial x} \right)^2 + g\zeta = 0. \quad (15)$$

Approximate form of surface velocity expression

$$\frac{\partial}{\partial x}\left((h+\zeta(x,t))\frac{\partial\phi}{\partial x}\right) - \frac{h^3}{3!}\frac{\partial^4\phi}{\partial x^4} - \frac{\partial\zeta}{\partial t} = 0.$$
 (16)

The expressions keep the lowest order nonlinear terms and include up to 4th order derivatives in the linear terms.



Traveling wave solutions with u = x - ct: $\phi(x,t) = \chi(u)$ and $\zeta(x,t) = \eta(u)$

Note that the wave "speed" c will be consistently determined Modified Bernoulli equation

$$c\frac{d\chi(u)}{du} - \frac{ch^2}{2}\frac{d^3\chi(u)}{du^3} + \frac{1}{2}\left(\frac{d\chi(u)}{du}\right)^2 + g\eta(u) = 0.$$
 (17)

Modified surface velocity equation

$$\frac{d}{du}\left((h+\eta(u))\frac{d\chi(u)}{du}\right) - \frac{h^3}{6}\frac{d^4\chi(u)}{du^4} + c\frac{d\eta(u)}{du} = 0.$$
 (18)



Integrating and rearranging coupled equations

Modified surface velocity equation

$$\frac{d}{du}\left((h+\eta(u))\frac{d\chi(u)}{du}\right) - \frac{h^3}{6}\frac{d^4\chi(u)}{du^4} + c\frac{d\eta(u)}{du} = 0.$$
 (19)

$$\Rightarrow (h+\eta)\chi' - \frac{h^3}{6}\chi''' + c\eta = 0$$
⁽²⁰⁾

Modified Bernoulli equation

$$\chi' - \frac{h^2}{2}\chi''' + \frac{1}{2c}(\chi')^2 + \frac{g}{c}\eta = 0.$$
 (21)

$$\Rightarrow \chi' = -\frac{g}{c}\eta + \frac{h^2}{2}\chi''' - \frac{1}{2c}(\chi')^2 \approx -\frac{g}{c}\eta - \frac{h^2g}{2c}\eta'' - \frac{g^2}{2c^3}\eta^2.$$
(22)



Integrating and rearranging coupled equations – continued

Modified surface velocity equation in terms of η

$$(h+\eta)\left(-\frac{g}{c}\eta - \frac{h^2g}{2c}\eta'' - \frac{g^2}{2c^3}\eta^2\right) + \frac{h^3g}{6c}\eta'' + c\eta = 0$$
(23)

$$\Rightarrow \left(1 - \frac{gh}{c^2}\right)\eta - \frac{gh^3}{3c^2}\eta'' - \frac{g}{c^2}\left(1 + \frac{gh}{2c^2}\right)\eta^2 = 0.$$
 (24)

$$\Rightarrow \left(1 - \frac{hg}{c^2}\right)\eta(u) - \frac{h^2}{3}\eta''(u) - \frac{3}{2h}\left[\eta(u)\right]^2 = 0.$$
 (25)



Solution of the famous Korteweg-de Vries equation

Modified surface velocity equation in terms of η

$$\Rightarrow \left(1 - \frac{hg}{c^2}\right)\eta(u) - \frac{h^2}{3}\eta''(u) - \frac{3}{2h}\left[\eta(u)\right]^2 = 0.$$
 (26)

Soliton solution

$$\zeta(x,t) = \eta(x-ct) = \eta_0 \operatorname{sech}^2\left(\sqrt{\frac{3\eta_0}{h}} \,\frac{x-ct}{2h}\right)$$
(27)

$$c = \sqrt{\frac{gh}{1 - \eta_0/h}} \approx \sqrt{gh} \left(1 + \frac{\eta_0}{2h}\right).$$
(28)

Here η_0 is a constant to be determined.



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Relationship with "standard" form of Korteweg-de Vries equation

New variables

$$\beta = 2\eta_0, \quad \bar{x} = \sqrt{\frac{3}{2h}} \frac{x}{h}, \quad \text{and} \quad \bar{t} = \sqrt{\frac{3}{2h}} \frac{ct}{2\eta_0 h}.$$
 (29)

Standard Korteweg-de Vries equation

$$\frac{\partial \eta}{\partial \bar{t}} + 6\eta \frac{\partial \eta}{\partial \bar{x}} + \frac{\partial^3 \eta}{\partial \bar{x}^3} = 0.$$
(30)
Soliton solution

$$\eta(\bar{x},\bar{t}) = \frac{\beta}{2} \operatorname{sech}^2 \left[\frac{\sqrt{\beta}}{2} (\bar{x} - \beta \bar{t}) \right].$$
(31)



More details

Modified surface velocity equation in terms of η

$$\left(1 - \frac{hg}{c^2}\right)\eta(u) - \frac{h^2}{3}\eta''(u) - \frac{3}{2h}\left[\eta(u)\right]^2 = 0.$$
(32)

Some identities

$$\frac{\eta_0}{h} = 1 - \frac{gh}{c^2}; \quad \frac{\partial\eta}{\partial t} = -c\frac{d\eta}{du}; \quad \frac{\partial\eta}{\partial x} = \frac{d\eta}{du}.$$
(33)

Derivative of surface velocity equation

$$\frac{\eta_0}{h}\eta' - \frac{h^2}{3}\eta''' - \frac{3}{h}\eta\eta' = 0.$$
(34)

In terms of \boldsymbol{x} and \boldsymbol{t}

$$-\frac{\eta_0}{ch}\frac{\partial\eta}{\partial t} - \frac{h^2}{3}\frac{\partial^3\eta}{\partial x^3} - \frac{3}{h}\eta\frac{\partial\eta}{\partial x} = 0.$$
 (35)

In terms of \bar{x} and \bar{t}

$$\frac{\partial \eta}{\partial \bar{t}} + 6\eta \frac{\partial \eta}{\partial \bar{x}} + \frac{\partial^3 \eta}{\partial \bar{x}^3} = 0.$$
(36)



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