

**PHY 711 Classical Mechanics and  
Mathematical Methods  
10-10:50 AM MWF Olin 103**

## Plan for Lecture 34:

## Chapter 10 in F & W: Surface waves

-- Non-linear contributions and soliton solutions

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1 Wed, 10/01/2014	Chap. 4	Normal modes of motion	Continue Take-Home
17 Fri, 10/03/2014	Chap. 4	Normal modes of motion	Take-Home due
18 Mon, 10/06/2014	Chap. 7	Wave motion	#12
19 Wed, 10/08/2014	Chap. 7	Sturm-Liouville Equations	#13
20 Fri, 10/10/2014	Chap. 7	Sturm-Liouville Equations	#13
21 Mon, 10/13/2014	Chap. 7	Sturm-Liouville Equations	#14
22 Wed, 10/15/2014	Appendix A	Contour integration methods	#15
Fri, 10/17/2014		Fall break -- no class	
23 Mon, 10/20/2014	Appendix A	Fourier transforms	#16
24 Wed, 10/22/2014	Chap. 5	Motion of Rigid Bodies	#17
25 Fri, 10/24/2014	Chap. 5	Motion of Rigid Bodies	#18
26 Mon, 10/27/2014	Chap. 5	Symmetric top in gravitational field	#18
27 Wed, 10/29/2014	Chap. 8	Vibrations of membranes	#19
28 Fri, 10/31/2014	Chap. 9	Physics of fluids	#20
29 Mon, 11/03/2014	Chap. 9	Physics of fluids	#21
30 Wed, 11/05/2014	Chap. 9	Sound waves	
31 Fri, 11/07/2014	Chap. 9	Sound waves	Begin Take-Home
32 Mon, 11/10/2014	Chap. 9	Non-linear effects	Continue Take-Home
33 Wed, 11/12/2014	Chap. 10	Surface waves in fluids	Continue Take-Home
34 Fri, 11/14/2014	Chap. 10	Surface waves in fluids	Continue Take-Home
35 Mon, 11/17/2014	Chap. 10	Surface waves in fluids	Take-Home due

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第十一章 聚合物的物理性质 11-21

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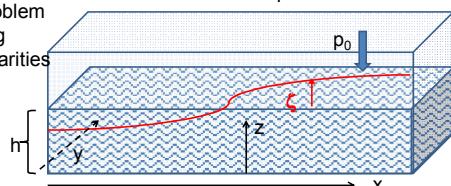
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Surface waves in an incompressible fluid



Within fluid:  $0 \leq z \leq h + \zeta$

$$-\frac{\partial \Phi}{\partial t} + \frac{1}{2}v^2 + g(z-h) = \text{constant} \quad \Phi = \Phi(x, y, z, t)$$

$$-\nabla^2 \Phi = 0 \quad \mathbf{v} = \mathbf{v}(x, y, z, t) = -\nabla \Phi(x, y, z, t)$$

$$\text{At surface: } z = h + \zeta \quad \text{with } \zeta = \zeta(x, y, t)$$

$$\frac{d\zeta}{dt} = \frac{\partial \zeta}{\partial t} + v_x \frac{\partial \zeta}{\partial x} + v_y \frac{\partial \zeta}{\partial y} \quad \text{where } v_{x,y} = v_{x,y}(x, y, h + \zeta, t)$$

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Further simplifications; assume trivial  $y$ -dependence  
 $\Phi = \Phi(x, z, t)$        $\zeta = \zeta(x, t)$   
 Within fluid:       $0 \leq z \leq h + \zeta$   
 At surface:       $v_z(x, z = h + \zeta, t) = -\frac{\partial \Phi}{\partial z} = \frac{d\zeta}{dt}$

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Non-linear effects in surface waves:

Dominant non-linear effects  $\Rightarrow$  soliton solutions  
 $\zeta(x, t) = \eta_0 \operatorname{sech}^2\left(\sqrt{\frac{3\eta_0}{h}} \frac{x - ct}{2h}\right)$        $\eta_0 = \text{constant}$   
 where  $c = \sqrt{\frac{gh}{1 - \eta_0/h}} \approx \sqrt{gh} \left(1 + \frac{\eta_0}{2h}\right)$

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**Detailed analysis of non-linear surface waves**  
 [Note that these derivations follow Alexander L. Fetter and John Dirk Walecka, *Theoretical Mechanics of Particles and Continua* (McGraw Hill, 1980), Chapt. 10.]

We assume that we have an incompressible fluid:  $\rho = \text{constant}$   
 Velocity potential:  $\Phi(x, z, t)$ ;  $v(x, z, t) = -\nabla \Phi(x, z, t)$

The surface of the fluid is described by  $z = h + \zeta(x, t)$ . It is assumed that the fluid is contained in a structure (lake, river, swimming pool, etc.) with a structureless bottom defined by the  $z = 0$  plane and filled to an equilibrium height of  $z = h$ .

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## Defining equations for $\Phi(x,z,t)$ and $\zeta(x,t)$

where  $0 \leq z \leq h + \zeta(x, t)$

Continuity equation:

$$\nabla \cdot \mathbf{v} = 0 \quad \Rightarrow \quad \frac{\partial^2 \Phi(x, z, t)}{\partial x^2} + \frac{\partial^2 \Phi(x, z, t)}{\partial z^2} = 0$$

Bernoulli equation (assuming irrotational flow) and gravitational potential energy

$$-\frac{\partial \Phi(x,z,t)}{\partial t} + \frac{1}{2} \left[ \left( \frac{\partial \Phi(x,z,t)}{\partial x} \right)^2 + \left( \frac{\partial \Phi(x,z,t)}{\partial z} \right)^2 \right] + g(z-h) = 0.$$

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## Boundary conditions on functions –

Zero velocity at bottom of tank:

$$\frac{\partial \Phi(x, 0, t)}{\partial z} = 0.$$

Consistent vertical velocity at water surface

$$v_z(x, z, t)|_{z=h+\zeta} = \frac{d\zeta}{dt} = \mathbf{v} \cdot \nabla \zeta + \frac{\partial \zeta}{\partial t}.$$

$$-\frac{\partial \Phi(x,z,t)}{\partial z} + \frac{\partial \Phi(x,z,t)}{\partial x} \frac{\partial \zeta(x,t)}{\partial x} - \frac{\partial \zeta(x,t)}{\partial t} \Big|_{z=h+\zeta} = 0$$

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**Analysis assuming water height  $z$  is small relative to variations in the direction of wave motion ( $x$ )**  
**Taylor's expansion about  $z = 0$ :**

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Note that the zero vertical velocity at the bottom ensures

Note that the zero vertical velocity at the bottom ensures that all odd derivatives  $\frac{\partial^n \Phi}{\partial z^n}(x, 0, t)$  vanish from the

Taylor expansion . In addition, the Laplace equation allows us to convert all even derivatives with respect to  $z$  to derivatives with respect to  $x$ .

Modified Taylor's expansion:  $\Phi(x, z, t) \approx \Phi(x, 0, t) - \frac{z^2}{2} \frac{\partial^2 \Phi}{\partial x^2}(x, 0, t) + \frac{z^4}{4!} \frac{\partial^4 \Phi}{\partial x^4}(x, 0, t) \dots$

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Check linearized equations and their solutions:  
Bernoulli equations --

Bernoulli equation evaluated at  $z = h + \zeta(x, t)$

$$-\frac{\partial \Phi(x, h, t)}{\partial t} + g\zeta(x, t) = 0$$

Consistent vertical velocity at  $z = h + \zeta(x, t)$

$$-\frac{\partial \Phi(x, z, t)}{\partial z} - \frac{\partial \zeta(x, t)}{\partial t} \Big|_{z=h+\zeta} = 0$$

Using Taylor's expansion results to lowest order

$$-\frac{\partial \Phi(x, z, t)}{\partial z} \approx h \frac{\partial^2 \Phi(x, 0, t)}{\partial x^2} \quad -\frac{\partial \Phi(x, h, t)}{\partial t} \approx -\frac{\partial \Phi(x, 0, t)}{\partial t}$$

$$\text{Decoupled equations: } \frac{\partial^2 \Phi(x, 0, t)}{\partial t^2} = gh \frac{\partial^2 \Phi(x, 0, t)}{\partial x^2}.$$

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Analysis of non-linear equations -- keeping the lowest order nonlinear terms and include up to 4th order derivatives in the linear terms. Let  $\phi(x, t) \equiv \Phi(x, 0, t)$

Approximate form of Bernoulli equation evaluated at surface:  $z = h + \zeta$

$$-\frac{\partial \phi}{\partial t} + \frac{(h + \zeta)^2}{2} \frac{\partial^3 \phi}{\partial t \partial x^2} + \frac{1}{2} \left[ \left( \frac{\partial \phi}{\partial x} \right)^2 + \left( (h + \zeta) \frac{\partial^2 \phi}{\partial x^2} \right)^2 \right] + g\zeta = 0$$

$$\Rightarrow -\frac{\partial \phi}{\partial t} + \frac{h^2}{2} \frac{\partial^3 \phi}{\partial t \partial x^2} + \frac{1}{2} \left( \frac{\partial \phi}{\partial x} \right)^2 + g\zeta = 0.$$

Approximate form of surface velocity expression :

$$\frac{\partial}{\partial x} \left( (h + \zeta(x, t)) \frac{\partial \phi}{\partial x} \right) - \frac{h^3}{3!} \frac{\partial^4 \phi}{\partial x^4} - \frac{\partial \zeta}{\partial t} = 0.$$

The expressions keep the lowest order nonlinear terms and include up to 4th order derivatives in the linear terms.

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Coupled equations:  $-\frac{\partial \phi}{\partial t} + \frac{h^2}{2} \frac{\partial^3 \phi}{\partial t \partial x^2} + \frac{1}{2} \left( \frac{\partial \phi}{\partial x} \right)^2 + g\zeta = 0.$

$$\frac{\partial}{\partial x} \left( (h + \zeta(x, t)) \frac{\partial \phi}{\partial x} \right) - \frac{h^3}{3!} \frac{\partial^4 \phi}{\partial x^4} - \frac{\partial \zeta}{\partial t} = 0.$$

Traveling wave solutions with new notation:

$$u \equiv x - ct \quad \phi(x, t) \equiv \chi(u) \quad \text{and} \quad \zeta(x, t) \equiv \eta(u)$$

Note that the wave "speed"  $c$  will be consistently determined

$$c \frac{d\chi(u)}{du} - \frac{ch^2}{2} \frac{d^3\chi(u)}{du^3} + \frac{1}{2} \left( \frac{d\chi(u)}{du} \right)^2 + g\eta(u) = 0.$$

$$\frac{d}{du} \left( (h + \eta(u)) \frac{d\chi(u)}{du} \right) - \frac{h^3}{6} \frac{d^4\chi(u)}{du^4} + c \frac{d\eta(u)}{du} = 0.$$

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## Integrating and re-arranging coupled equations

$$\begin{aligned} & c \frac{d\chi(u)}{du} - \frac{ch^2}{2} \frac{d^3\chi(u)}{du^3} + \frac{1}{2} \left( \frac{d\chi(u)}{du} \right)^2 + g\eta(u) = 0, \\ \Rightarrow & \chi' = -\frac{g}{c}\eta + \frac{h^2}{2}\chi''' - \frac{1}{2c}(\chi')^2 \approx -\frac{g}{c}\eta - \frac{h^2\eta}{2c} - \frac{g^2}{2c^3}\eta^2 \\ & \frac{d}{du} \left( (h+\eta(u)) \frac{d\chi(u)}{du} \right) - \frac{h^3}{6} \frac{d^4\chi(u)}{du^4} + c \frac{d\eta(u)}{du} = 0, \\ \Rightarrow & (h+\eta)\chi' - \frac{h^3}{6}\chi''' + c\eta = 0 \end{aligned}$$

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## Solution of the famous Korteweg-de Vries equation

Modified surface amplitude equation in terms of  $\eta$

$$\Rightarrow \left(1 - \frac{hg}{c^2}\right)\eta(u) - \frac{h^2}{3}\eta''(u) - \frac{3}{2h}[\eta(u)]^2 = 0.$$

### Soliton solution

$$\zeta(x,t) = \eta(x-ct) = \eta_0 \operatorname{sech}^2\left(\sqrt{\frac{3\eta_0}{h}} \frac{x-ct}{2h}\right)$$

$$c = \sqrt{\frac{gh}{1 - \eta_0/h}} \approx \sqrt{gh} \left(1 + \frac{\eta_0}{2h}\right) \quad \text{where } \eta_0 \text{ is a constant}$$

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Relationship to “standard” form of Korteweg-de Vries equation

New variables:

$$\beta = 2\eta_0, \quad \bar{x} = \sqrt{\frac{3}{2h}} \frac{x}{h}, \quad \text{and} \quad \bar{t} = \sqrt{\frac{3}{2h}} \frac{ct}{2\eta_0 h}.$$

## Standard Korteweg-de Vries equation

$$\frac{\partial \eta}{\partial \bar{t}} + 6\eta \frac{\partial \eta}{\partial \bar{x}} + \frac{\partial^3 \eta}{\partial \bar{x}^3} = 0.$$

Soliton solution:

$$\eta(\bar{x}, \bar{t}) = \frac{\beta}{2} \operatorname{sech}^2 \left[ \frac{\sqrt{\beta}}{2} (\bar{x} - \beta \bar{t}) \right].$$

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## More details

Modified surface amplitude equation in terms of  $\eta$ :

$$\left(1 - \frac{hg}{c^2}\right)\eta(u) - \frac{h^2}{3}\eta''(u) - \frac{3}{2h}[\eta(u)]^2 = 0.$$

Some identities:  $\frac{\eta_0}{h} = 1 - \frac{gh}{c^2}$ ;  $\frac{\partial \eta}{\partial t} = -c \frac{d\eta}{du}$ ;  $\frac{\partial \eta}{\partial x} = \frac{d\eta}{du}$ .

Derivative of surface amplitude equation:

$$\frac{\eta_0}{h}\eta' - \frac{h^2}{3}\eta''' - \frac{3}{h}\eta\eta' = 0.$$

Expression in terms of  $x$  and  $t$ :

$$-\frac{\eta_0}{ch} \frac{\partial \eta}{\partial t} - \frac{h^2}{3} \frac{\partial^3 \eta}{\partial x^3} - \frac{3}{h} \eta \frac{\partial \eta}{\partial x} = 0.$$

Expression in terms of  $\bar{x}$  and  $\bar{t}$ :

$$\frac{\partial \eta}{\partial \bar{t}} + 6\eta \frac{\partial \eta}{\partial \bar{x}} + \frac{\partial^3 \eta}{\partial \bar{x}^3} = 0.$$

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## Summary

## Soliton solution

$$\zeta(x,t) = \eta(x-ct) = \eta_0 \operatorname{sech}^2\left(\sqrt{\frac{3\eta_0}{h}} \frac{x-ct}{2h}\right)$$

$$c = \sqrt{\frac{gh}{1 - \eta_0/h}} \approx \sqrt{gh} \left(1 + \frac{\eta_0}{2h}\right) \quad \text{where } \eta_0 \text{ is a constant}$$

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Photo of canal soliton <http://www.ma.hw.ac.uk/solitons/>



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