

**PHY 711 Classical Mechanics and
Mathematical Methods
10-10:50 AM MWF Olin 103**

Plan for Lecture 33:

Effects of viscosity in fluid motion – Chap.12 in Fetter & Walecka

1. Navier-Stokes equation
 2. Terminal velocity of a sphere moving with constant applied force in a viscous medium
 3. Stokes' viscosity relation

11/16/2015

PHY 711 Fall 2015 -- Lecture 33

1

22 Wed, 10/21/2015 Chap. 5	Motion of Rigid Bodies	#21
23 Fri, 10/23/2015 Chap. 8	Motion of Elastic membranes	#22
24 Mon, 10/26/2015 Chap. 9	Hydrodynamics	#23
25 Wed, 10/28/2015 Chap. 9	Hydrodynamics	#24
26 Fri, 10/30/2015 Chap. 9	Sound waves	#25
27 Mon, 11/02/2015 Chap. 9	Sound waves	#26
28 Wed, 11/04/2015 Chap. 9	Sound waves	
29 Fri, 11/06/2015 Chap. 10	Surface waves on fluids	#27
30 Mon, 11/09/2015 Chap. 10	Surface waves on fluids	#28
31 Wed, 11/11/2015 Chap. 11	Heat Conduction	#29
32 Fri, 11/13/2015 Chap. 12	Viscosity	#30
33 Mon, 11/16/2015 Chap. 12	Viscosity	Prepare presentation.
34 Wed, 11/18/2015 Chap. 12	Viscosity	Prepare presentation.
35 Fri, 11/20/2015 Chap. 12	More viscosity	Prepare presentation.
36 Mon, 11/23/2015 Chap. 13	Elastic Continua	Prepare presentation.
Wed, 11/25/2015	Thanksgiving Holiday	
Fri, 11/27/2015	Thanksgiving Holiday	
37 Mon, 11/30/2015 Chap. 13	Elastic Continua	Prepare presentation.
Wed, 12/02/2015	Student presentations I	
Fri, 12/04/2015	Student presentations II	
Mon, 12/07/2015	Open Topic, home final	

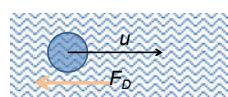
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PHY 711 Fall 2015 – Lecture 38

2

Brief introduction to viscous effects in incompressible fluids
 Stokes' analysis of viscous drag on a sphere of radius R
 moving at speed u in medium with viscosity η :

$$F_D = -\eta(6\pi Ru)$$



Plan:

1. Consider the general effects of viscosity on fluid equations
 2. Consider the solution to the linearized equations for the case of steady-state flow of a sphere of radius R
 3. Infer the drag force needed to maintain the steady-state flow

11/16/2015

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3

Newton-Euler equation for incompressible fluid,
modified by viscous contribution (Navier-Stokes equation):

$$\frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} = \mathbf{f}_{\text{applied}} - \frac{\nabla p}{\rho} + \frac{\eta}{\nu} \nabla^2 \mathbf{v}$$

ν Kinematic viscosity

Typical kinematic viscosities at 20° C and 1 atm:

Fluid	ν (m ² /s)
Water	1.00×10^{-6}
Air	14.9×10^{-6}
Ethyl alcohol	1.52×10^{-6}
Glycerine	1183×10^{-6}

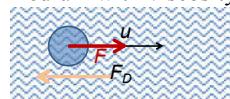
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4

Stokes' analysis of viscous drag on a sphere of radius R moving at speed u in medium with viscosity η :

$$F_D = -\eta(6\pi R u)$$



Effects of drag force on motion of

particle of mass m with constant force F :

$$F - 6\pi R \eta u = m \frac{du}{dt} \quad \text{with } u(0) = 0$$

$$\Rightarrow u(t) = \frac{F}{6\pi R \eta} \left(1 - e^{-\frac{6\pi R \eta}{m} t} \right)$$

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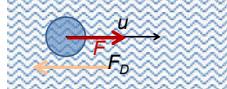
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Effects of drag force on motion of

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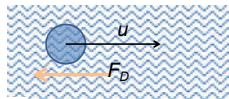
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6

Effects of drag force on motion of particle of mass m
with an initial velocity with $u(0) = U_0$ and no external force

$$\begin{aligned} -6\pi R\eta u &= m \frac{du}{dt} \\ \Rightarrow u(t) &= U_0 e^{-\frac{6\pi R\eta}{m}t} \end{aligned}$$



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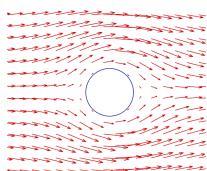
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Recall: PHY 711 -- Assignment #24 Oct. 28, 2015

Determine the form of the velocity potential for an incompressible fluid representing uniform velocity in the \mathbf{z} direction at large distances from a spherical obstruction of radius a . Find the form of the velocity potential and the velocity field for all $r > a$. Assume that the velocity in the radial direction is 0 for $r = a$ and assume that the velocity is uniform in the azimuthal direction.

$$\nabla^2 \Phi = 0$$

$$\Phi(r, \theta) = \left(-v_0 r + \frac{v_0 R^3}{2r^2} \right) \cos \theta$$



11/16/2015

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8

Newton-Euler equation for incompressible fluid,
modified by viscous contribution (Navier-Stokes equation):

$$\frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} = \mathbf{f}_{applied} - \frac{\nabla p}{\rho} + \frac{\eta}{\rho} \nabla^2 \mathbf{v}$$

Continuity equation: $\nabla \cdot \mathbf{v} = 0$

Assume steady state: $\Rightarrow \frac{\partial \mathbf{v}}{\partial t} = 0$

Assume non-linear effects small

Initially set $\mathbf{f}_{applied} = 0$;

$$\Rightarrow \nabla p = \eta \nabla^2 \mathbf{v}$$



11/16/2015

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6

$$\nabla p = \eta \nabla^2 \mathbf{v}$$

Take curl of both sides of equation:

$$\nabla \times (\nabla p) = 0 = \eta \nabla^2 (\nabla \times \mathbf{v})$$

Assume (with a little insight from Landau):

$$\mathbf{v} = \nabla \times (\nabla \times f(r) \mathbf{u}) + \mathbf{u}$$

where $f(r) \xrightarrow[r \rightarrow \infty]{} 0$

Note that:

$$\nabla \times (\nabla \times \mathbf{A}) = \nabla(\nabla \cdot \mathbf{A}) - \nabla^2 \mathbf{A}$$

11/16/2015

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10

Digression

Some comment on assumption: $\mathbf{v} = \nabla \times (\nabla \times f(r)\mathbf{u}) + \mathbf{u}$

$$\nabla \times (\nabla \times \mathbf{A}) = \nabla(\nabla \cdot \mathbf{A}) - \nabla^2 \mathbf{A}$$

Here $\mathbf{A} = \nabla \times f(r)\mathbf{u}$

$$\nabla \times \mathbf{v} = 0 = \nabla \times (\nabla \times (\nabla \times \mathbf{A})) = -\nabla \times (\nabla^2 \mathbf{A})$$

Also note : $\nabla p = \eta \nabla^2 \mathbf{v}$

$$\Rightarrow \nabla \times \nabla p = \nabla \times \eta \nabla^2 \mathbf{v} \quad \text{or} \quad \nabla^2(\nabla \times \mathbf{v}) = 0$$

11/16/2015

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11

$$\mathbf{v} = \nabla \times (\nabla \times f(r) \mathbf{u}) + \mathbf{u}$$

$$\mathbf{u} = u\hat{\mathbf{z}}$$

$$\nabla \times (\nabla \times f(r) \hat{\mathbf{z}}) = \nabla(\nabla \cdot f(r) \hat{\mathbf{z}}) - \nabla^2 f(r) \hat{\mathbf{z}}$$

$$\nabla \times \mathbf{v} = 0 \quad \Rightarrow \nabla^2(\nabla \times \mathbf{v}) = 0$$

$$\nabla^4(\nabla \times f(r)\hat{\mathbf{z}}) = 0 \quad \Rightarrow \nabla^4(\nabla f(r) \times \hat{\mathbf{z}}) = 0 \quad \Rightarrow \nabla^4 f(r) = 0$$

$$f(r) = C_1 r^2 + C_2 r + C_3 + \frac{C_4}{r}$$

$$v_r = u \cos \theta \left(1 - \frac{2}{r} \frac{df}{dr} \right) = u \cos \theta \left(1 - 4C_1 - \frac{2C_2}{r} - \frac{2C_4}{r^3} \right)$$

$$v_\theta = -u \sin \theta \left(1 - \frac{d^2 f}{dr^2} - \frac{1}{r} \frac{df}{dr} \right) = -u \sin \theta \left(1 - 4C_1 - \frac{C_2}{r} + \frac{C_4}{r^3} \right)$$

11/16/2015

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12

$$v_r = u \cos \theta \left(1 - \frac{2}{r} \frac{df}{dr} \right) = u \cos \theta \left(1 - 4C_1 - \frac{2C_2}{r} - \frac{2C_4}{r^3} \right)$$

$$v_\theta = -u \sin \theta \left(1 - \frac{d^2 f}{dr^2} - \frac{1}{r} \frac{df}{dr} \right) = -u \sin \theta \left(1 - 4C_1 - \frac{C_2}{r} + \frac{C_4}{r^3} \right)$$

To satisfy $\mathbf{v}(r \rightarrow \infty) = \mathbf{u}$: $\Rightarrow C_1 = 0$

To satisfy $v(R) = 0$ solve for C_2, C_4

$$v_r = u \cos \theta \left(1 - \frac{3R}{2r} + \frac{R^3}{2r^3} \right)$$

$$v_\theta = -u \sin \theta \left(1 - \frac{3R}{4r} - \frac{R^3}{4r^3} \right)$$

11/16/2015

PHY 711 Fall 2015 -- Lecture 33

13

$$v_r = u \cos \theta \left(1 - \frac{3R}{2r} + \frac{R^3}{2r^3} \right)$$

$$v_\theta = -u \sin \theta \left(1 - \frac{3R}{4r} - \frac{R^3}{4r^3} \right)$$

Determining pressure:

$$\nabla p = \eta \nabla^2 \mathbf{v} = -\eta \nabla \left(u \cos \theta \left(\frac{3R}{2r^2} \right) \right)$$

$$\Rightarrow p = p_0 - \eta u \cos \theta \left(\frac{3R}{2r^2} \right)$$

11/16/2015

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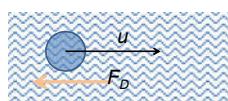
14

$$p = p_0 - \eta u \cos \theta \left(\frac{3R}{2r^2} \right)$$

Corresponds to:

$$F_D \cos \theta = (p(R) - p_0) 4\pi R^2$$

$$\Rightarrow F_D = -\eta u(6\pi R)$$



11/16/2015

PHY 711 Fall 2015 – Lecture 33

15