

PHY 752 Solid State Physics
11-11:50 AM MWF Olin 103

Plan for Lecture 2:

Reading: Continue reading Chapter 1 in GGGPP; Electrons in one-dimensional periodic potentials

1. Review of Bloch's Theorem
2. Kronig-Penny model potential from the viewpoint of scattering or tunneling

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PHY 752 Solid State Physics

MWF 11 AM-11:50 PM | OPL 103 | <http://www.wfu.edu/~natalie/f15phy752/>

Instructor: [Natalie Holzwarth](#) **Phone:** 758-5510 **Office:** 300 OPL **e-mail:** natalie@wfu.edu

Course schedule

(Preliminary schedule – subject to frequent adjustment.)

Date	F&W Reading Topic	Assignment
1 Wed. 8/26/2015	Chap. 1 Electrons in a periodic one-dimensional potential	#1
2 Fri. 8/28/2015	Chap. 1 Electrons in a periodic one-dimensional potential	#2
3 Mon. 9/3/2015	Chap. 1	
4 Wed. 9/02/2015		
5 Fri. 9/04/2015		
6 Mon. 9/07/2015		
7 Wed. 9/09/2015		
8 Fri. 9/11/2015		
9 Mon. 9/14/2015		
10 Wed. 9/16/2015		

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1 Electrons in One-Dimensional Periodic Potentials

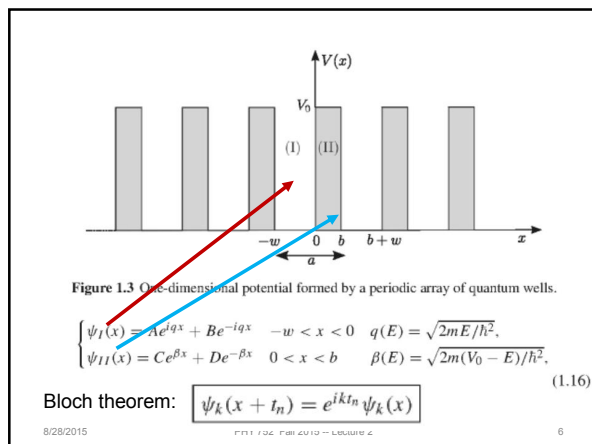
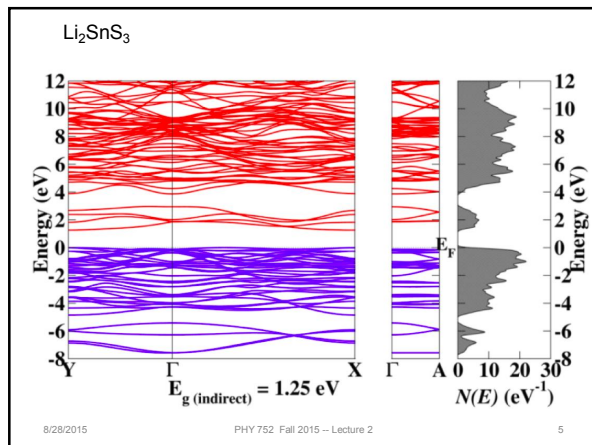
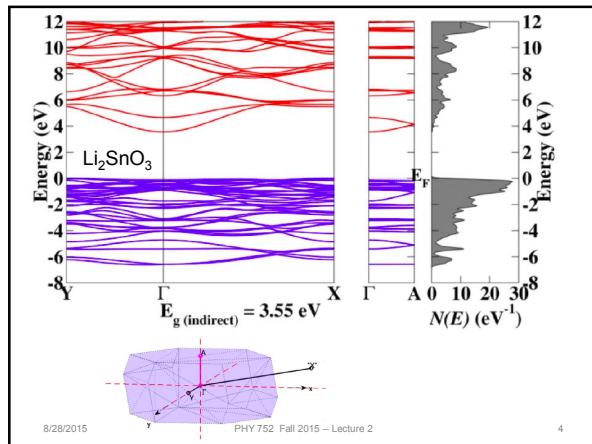
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Eigenvalue solutions to Kronig-Penny model

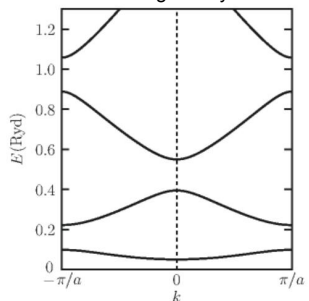


Figure 1.5 Energy bands of lowest energy in the Kronig-Penny model for $P = 3\pi/2$. The energy E (in Rydberg) is plotted as a function of k in the first Brillouin zone; the value of the lattice parameter considered is $a = 10a_B$ (a_B Bohr radius).

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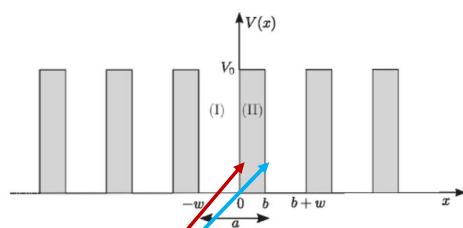


Figure 1.3 One-dimensional potential formed by a periodic array of quantum wells.

$$\begin{aligned} \psi_I(0) &= \psi_{II}(0), & \left(\frac{d\psi_I}{dx}\right)_{x=0} &= \left(\frac{d\psi_{II}}{dx}\right)_{x=0}, \\ \psi_{II}(b) &= e^{ika}\psi_I(-w), & \left(\frac{d\psi_{II}}{dx}\right)_{x=b} &= e^{ika}\left(\frac{d\psi_I}{dx}\right)_{x=-w}. \end{aligned}$$

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From wavefunction matching conditions:

$$\begin{cases} A + B = C + D \\ Aiq - Biq = C\beta - D\beta \\ Ce^{\beta b} + De^{-\beta b} = e^{ika}[Ae^{-iqw} + Be^{+iqw}] \\ C\beta e^{\beta b} - D\beta e^{-\beta b} = e^{ika}[Aiqe^{-iqw} - Biqe^{+iqw}]. \end{cases}$$

Condition for non-trivial solution:

$$\begin{vmatrix} 1 & 1 & -1 & -1 \\ iq & -iq & -\beta & \beta \\ -e^{ika-iqw} & -e^{ika+iqw} & e^{\beta b} & e^{-\beta b} \\ -iqe^{ika-iqw} & iqe^{ika+iqw} & \beta e^{\beta b} & -\beta e^{-\beta b} \end{vmatrix} = 0.$$

Simplified result:

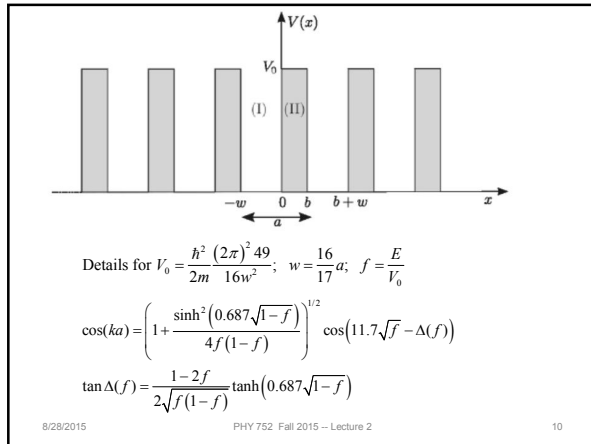
$$\frac{\beta^2 - q^2}{2q\beta} \sinh \beta b \sin qw + \cosh \beta b \cos qw = \cos ka.$$

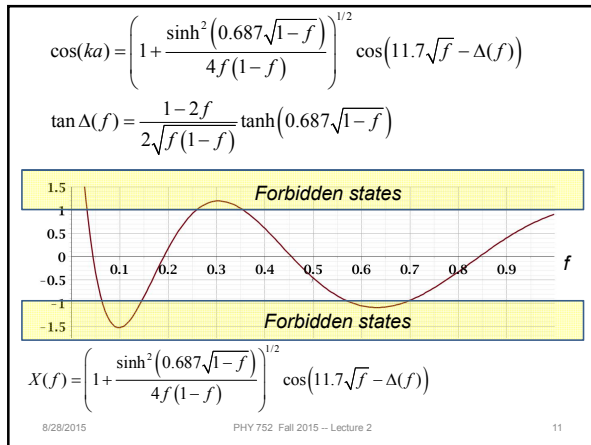
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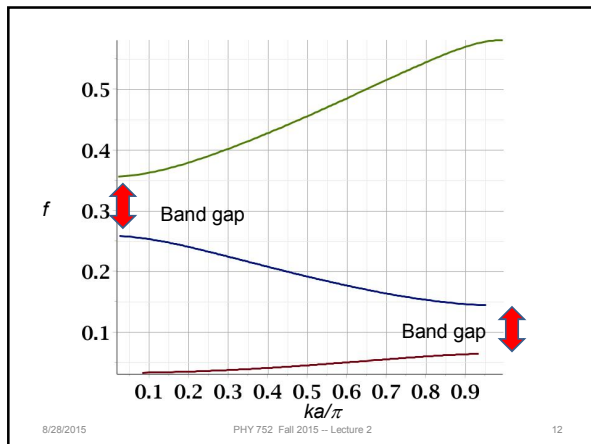
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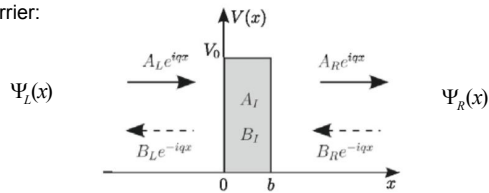








Treatment of same problem in terms of transmission and reflection from a periodic barrier. First consider a single barrier:



$$\begin{cases} \psi_L(x) = A_L e^{iqx} + B_L e^{-iqx} & x \leq x_L \\ \psi_R(x) = A_R e^{iqx} + B_R e^{-iqx} & x \geq x_R \end{cases} \quad q(E) = \sqrt{2mE/\hbar^2} \quad E > 0.$$

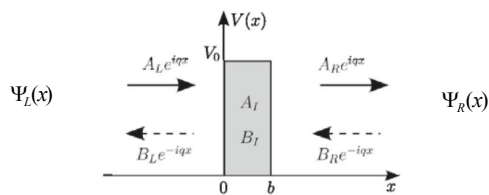
Matrix for determining coefficients:

$$\begin{pmatrix} A_R \\ B_R \end{pmatrix} = M(E) \begin{pmatrix} A_L \\ B_L \end{pmatrix} = \begin{bmatrix} m_{11}(E) & m_{12}(E) \\ m_{21}(E) & m_{22}(E) \end{bmatrix} \begin{pmatrix} A_L \\ B_L \end{pmatrix}$$

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$$\begin{cases} \psi_L(x) = A_L e^{iqx} + B_L e^{-iqx} & x < 0, \\ \psi_I(x) = A_I e^{\beta x} + B_I e^{-\beta x} & 0 < x < b, \\ \psi_R(x) = A_R e^{iqx} + B_R e^{-iqx} & x > b, \end{cases}$$

where $q^2(E) = 2mE/\hbar^2$ and $\beta^2(E) = 2m(V_0 - E)/\hbar^2$.

Matching conditions for wavefunction coefficients:

$$\begin{pmatrix} A_I \\ B_I \end{pmatrix} = \frac{1}{2\beta} \begin{bmatrix} iq + \beta & -iq + \beta \\ -iq + \beta & iq + \beta \end{bmatrix} \begin{pmatrix} A_L \\ B_L \end{pmatrix}, \quad \begin{pmatrix} A_R \\ B_R \end{pmatrix} = \frac{1}{2iq} \begin{bmatrix} (iq + \beta)e^{(-iq+\beta)b} & (iq - \beta)e^{(-iq-\beta)b} \\ (iq - \beta)e^{(+iq+\beta)b} & (iq + \beta)e^{(+iq-\beta)b} \end{bmatrix} \begin{pmatrix} A_I \\ B_I \end{pmatrix}$$

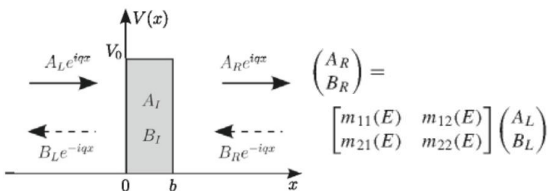
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Solving for transfer matrix elements:

$$\begin{cases} m_{11} = e^{-iqb} \left[\cosh \beta b + i \frac{q^2 - \beta^2}{2q\beta} \sinh \beta b \right] & m_{22} = m_{11}^* \\ m_{12} = e^{-iqb} (-i) \frac{q^2 + \beta^2}{2q\beta} \sinh \beta b & m_{21} = m_{12}^* \end{cases}$$



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Transmittance

$$t = \frac{A_R}{A_L} \quad T = |t|^2 = \left| \frac{1}{m_{22}} \right|^2$$

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After some algebra:

$$T(E) = \frac{1}{1 + \frac{V_0^2}{4E(V_0 - E)} \sinh^2 \sqrt{\frac{2m(V_0 - E)b^2}{\hbar^2}}} \quad 0 \leq E \leq V_0.$$

$$T(E) = \frac{1}{1 + \frac{V_0^2}{4E(E - V_0)} \sin^2 \sqrt{\frac{2m(E - V_0)b^2}{\hbar^2}}} \quad E \geq V_0.$$

Figure 1.8 Transmission coefficient through a barrier of height $V_0 = 1$ eV and length $b = 5$ Å (dashed line), and through a barrier of height $V_0 = 1$ eV and length $b = 20$ Å (solid line).

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Electron transmission through a one-dimensional periodic barrier

Because of Bloch theorem:

$$\psi_R(a) = e^{ika} \psi_L(0) \quad \text{and} \quad \left(\frac{d\psi_R}{dx} \right)_{x=a} = e^{ika} \left(\frac{d\psi_L}{dx} \right)_{x=0}.$$

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Bloch theorem & transmission coefficients:

$$\begin{pmatrix} A_R \\ B_R \end{pmatrix} = \begin{pmatrix} m_{11}e^{iqa} & m_{12}e^{iqa} \\ m_{21}e^{-iqa} & m_{22}e^{-iqa} \end{pmatrix} \begin{pmatrix} A_L \\ B_L \end{pmatrix} = e^{ika} \begin{pmatrix} A_L \\ B_L \end{pmatrix}$$

Condition for solution:

$$\text{Re} \left[m_{11}e^{iqa} \right] = \cos ka.$$

$$\cosh \beta b \cos qw - \frac{q^2 - \beta^2}{2q\beta} \sinh \beta b \sin qw = \cos ka.$$

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