

PHY 752 Solid State Physics

11-11:50 AM MWF Olin 103

Plan for Lecture 4:

Reading: Chapter 1.6 & 2.1 in GGGPP – electron dynamics in one dimension & introduction to crystal structures

1.Brief discussion of electron dynamics in one dimension

2.Bravais lattice structures

3.Classification of crystal structures

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1

PHY 752 Solid State Physics

MWF 11 AM-11:50 PM | OPL 103 | <http://www.wfu.edu/~natalie/f15phy752/>

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Course schedule

(Preliminary schedule -- subject to frequent adjustment.)

Date	F&W	Reading Topic	Assignment
1 Wed, 8/26/2015	Chap. 1.1-1.2	Electrons in a periodic one-dimensional potential	#1
2 Fri, 8/28/2015	Chap. 1.3	Electrons in a periodic one-dimensional potential	#2
3 Mon, 8/31/2015	Chap. 1.4	Tight binding models	#3
4 Wed, 9/02/2015	Chap. 1.6, 2.1	Crystal structures	#4
5 Fri, 9/04/2015	Chap. 2	Group theory	
6 Mon, 9/07/2015	Chap. 2	Group theory	
7 Wed, 9/09/2015	Chap. 2	Group theory	
8 Fri, 9/11/2015			

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2

FOREST CITY

Department of Physics

News



Congratulations to Dr. Greg Smith, recent Ph.D. Recipient



Congratulations to Dr. Jie Liu, recent Ph.D. Recipient



Congratulations to Dr. Wei Li, recent Ph.D. Recipient

Events

Wed, Sept. 2, 2015

Laboratory tours for students -- Part I

Meet in Olin Lobby at 4 PM

Refreshments at 3:30 PM

Olin Lobby

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3

Brief comment on electron dynamics as described by electronic band structure --

Electron velocity:

$$v(k) = \langle \psi(k, x) | \frac{p}{m} | \psi(k, x) \rangle.$$

$$\langle \psi(k, x) | \frac{p^2}{2m} + V(x) | \psi(k, x) \rangle = E(k),$$

$$v(k) = \langle \psi(k, x) | \frac{p}{m} | \psi(k, x) \rangle = \frac{1}{\hbar} \frac{dE(k)}{dk};$$

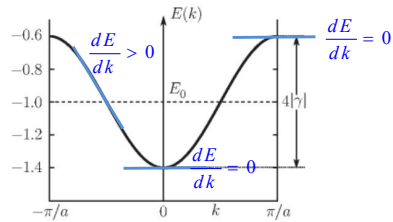
→ Band dispersion is related to electron velocity

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4

Example of electron band dispersion in one dimension



In practice, electron transport can be well approximated by the Boltzmann equation with the main contributions coming near the Fermi level.

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5

Electron effective mass from semi-classical analysis

$$\frac{d(\hbar k)}{dt} = -eF.$$

For applied force F

$$\frac{dv(k)}{dt} = \frac{d}{dt} \frac{1}{\hbar} \frac{dE(k)}{dk} = \frac{1}{\hbar} \frac{d^2E(k)}{dk^2} \frac{dk}{dt} = \frac{1}{\hbar^2} \frac{d^2E(k)}{dk^2} (-e)F.$$

$1/m^*$

$$\frac{1}{m^*} = \frac{1}{\hbar^2} \frac{d^2E(k)}{dk^2}.$$

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6

Classification of crystal structures

- 14 Bravais lattices
- 230 space groups



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7

2 Geometrical Description of Crystals: Direct and Reciprocal Lattices

Chapter Outline head

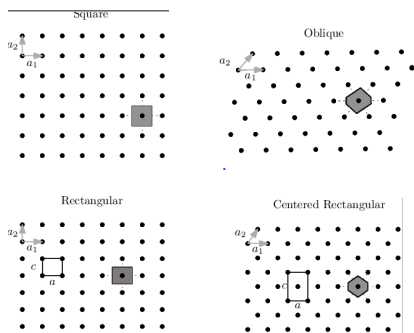
- ➔ 2.1 Simple Lattices and Composite Lattices 67
 - 2.1.1 Periodicity and Bravais Lattices 68
 - 2.1.2 Simple and Composite Crystal Structures 70
- 2.2 Geometrical Description of Some Crystal Structures 72
- 2.3 Wigner-Seitz Primitive Cells 83
- ➔ 2.4 Reciprocal Lattices 84
 - 2.4.1 Definitions and Basic Properties 84
 - 2.4.2 Planes and Directions in Bravais Lattices 85
- 2.5 Brillouin Zones 88

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8

Two-dimensional crystals

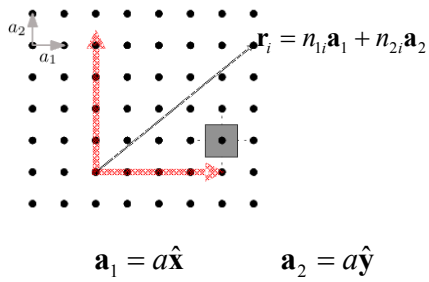


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9

Some details –
square lattice example

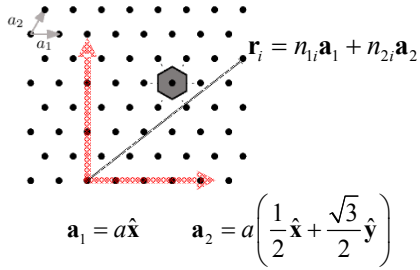


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10

Two-dimensional crystals -- continued
Hexagonal

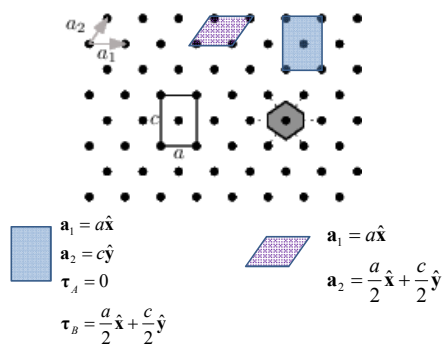


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11

“Conventional” versus “Primitive” unit cell
Centered Rectangular



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12

Primitive translation vectors

$$\mathbf{t}_n = n_1 \mathbf{t}_1 + n_2 \mathbf{t}_2 + n_3 \mathbf{t}_3;$$

Volume of primitive unit cell

$$\Omega = |\mathbf{t}_1 \cdot (\mathbf{t}_2 \times \mathbf{t}_3)|$$

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13

Crystal system	Bravais lattices			
	primitive	base-centered	body-centered	face-centered
Triclinic a≠b≠c α≠β≠γ				
Monoclinic a≠b≠c α=β=γ≠90°				
Orthorhombic a≠b≠c α=β=γ=90°				
Trigonal a=b=c α=β=γ=120°				

Tetragonal $a=b \neq c$ $\alpha=\beta=\gamma=90^\circ$		
Hexagonal a=b≠c $\alpha=\beta=90^\circ \neq \gamma=120^\circ$		
Cubic a=b=c $\alpha=\beta=\gamma=90^\circ$		

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14

Common examples

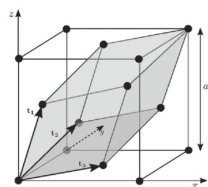
2.2 Geometrical Description of Some Crystal Structures

Crystal Structure of Rare-Gas Solids (Face-Centered Cubic Lattice)

We begin our brief description of some crystals of interest with the case of rare-gas solids (Ne, Ar, Kr, Xe). This crystal structure (illustrated in Figure 2.4) is obtained repeating periodically in space a *face-centered cube* (fcc), i.e. a conventional cubic cell (of edge a) with atoms at the corners and at the center of the faces. The primitive translation vectors of the fcc Bravais lattice are

$$\mathbf{t}_1 = \frac{a}{2}(0, 1, 1), \quad \mathbf{t}_2 = \frac{a}{2}(1, 0, 1), \quad \mathbf{t}_3 = \frac{a}{2}(1, 1, 0). \quad (2.6)$$

Nearest neighbor distance between
fcc lattice points: $a / \sqrt{2}$ (12)



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15

https://en.wikipedia.org/wiki/Periodic_table_(crystal_structure)

%29

Crystal structure of elements in the periodic table

Crystal structure of elements in the periodic table																	
H HEX 2 He																	
3 4 5 6 7 8 9 10 Li Be B C N O F Ne																	
BCC HCP RHO HEX 13 14 15 16 17 18 Na Mg Al Si P S Cl Ar																	
BCC HCP FCC CUB ORTH ORTH ORTH FCC																	
19 20 21 22 23 24 25 26 27 28 29 30 31 32 33 34 35 36 K Ca Sc Ti V Cr Mn Fe Co Ni Cu Zn Ga Ge As Se Br Kr																	
BCC FCC HCP HCP BCC BCC BCC BCC HCP FCC FCC HCP ORTH CUB RHO HEX ORTH FCC																	
37 38 39 40 41 42 43 44 45 46 47 48 49 50 51 52 53 54 Rb Sr Y Zr Nb Mo Tc Ru Rh Pd Ag Cd In Sn Sb Te I Xe																	
BCC FCC HCP BCC BCC BCC HCP FCC FCC FCC HCP TETR TETR RHO HEX ORTH FCC																	
55 56 72 73 74 75 76 77 78 79 80 81 82 83 84 85 86 Cs Ba * Hf Ta W Re Os Ir Pt Au Hg Tl Pb Bi Po At Rn																	
BCC BCC BCC BCC BCC BCC BCC FCC FCC FCC FCC RHO HCP RHO CUB FCC FCC																	
87 88 104 105 106 107 108 109 110 111 112 113 114 115 116 117 118 Fr Ra ** Rf Db Sg Bh Hs Mt Ds Rg Cn Uut Uup Uuq Uuh Uus Uuo																	
[BCC] BCC [HCP] [BCC] [BCC] [BCC] [HCP] [HCP] [FCC] [BCC] [BCC] [HCP] [FCC] [FCC] [FCC] [FCC] [FCC]																	

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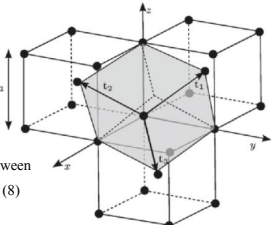
16

Crystal Structure of Alkali Metals (Body-Centered Cubic Lattice)

The crystal structure of alkali metals (see [Figure 2.5](#)) is obtained repeating periodically in space a body-centered cube (bcc), i.e. a conventional cubic cell with atoms at the corners and at the center of a cube of edge a . The primitive translation vectors of the bcc Bravais lattice are

$$\mathbf{t}_1 = \frac{a}{2}(-1, 1, 1), \quad \mathbf{t}_2 = \frac{a}{2}(1, -1, 1), \quad \mathbf{t}_3 = \frac{a}{2}(1, 1, -1) \quad (2.7)$$

Notice that $\mathbf{t}_1 + \mathbf{t}_2 + \mathbf{t}_3 = (a/2)(1, 1, 1)$, and $\mathbf{t}_2 + \mathbf{t}_3 + \mathbf{t}_1 = a(1, 0, 0)$, $\mathbf{t}_1 + \mathbf{t}_3 + \mathbf{t}_2 = a(0, 1, 0)$, and $\mathbf{t}_1 + \mathbf{t}_2 = a(0, 0, 1)$. The atom at the origin has 8 nearest neighbors in the positions



Nearest neighbor distance between bcc lattice points: $a\sqrt{3}/2$ (8)

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17

2.4 Reciprocal Lattices

2.4.1 Definitions and Basic Properties

For the study of crystals, besides the *direct lattice in the ordinary space*, it is important to consider also the *reciprocal lattice in the dual (or reciprocal) space*.

Given a crystal with primitive translation vectors $\mathbf{t}_1, \mathbf{t}_2, \mathbf{t}_3$ in the direct space, we consider the three primitive vectors $\mathbf{g}_1, \mathbf{g}_2, \mathbf{g}_3$ in the reciprocal space, defined by the relations

$$\mathbf{t}_i \cdot \mathbf{g}_j = 2\pi \delta_{ij} \quad (2.16a)$$

$$\mathbf{g}_1 = \frac{2\pi}{\Omega} \mathbf{t}_2 \times \mathbf{t}_3, \quad \mathbf{g}_2 = \frac{2\pi}{\Omega} \mathbf{t}_3 \times \mathbf{t}_1, \quad \mathbf{g}_3 = \frac{2\pi}{\Omega} \mathbf{t}_1 \times \mathbf{t}_2, \quad \text{and} \quad \Omega = \mathbf{t}_1 \cdot (\mathbf{t}_2 \times \mathbf{t}_3), \quad (2.17)$$

Example: fcc bravais lattice

$$\mathbf{t}_1 = \frac{a}{2}(0, 1, 1), \quad \mathbf{t}_2 = \frac{a}{2}(1, 0, 1), \quad \mathbf{t}_3 = \frac{a}{2}(1, 1, 0).$$

$$\mathbf{g}_1 = \frac{2\pi}{a}(-1, 1, 1), \quad \mathbf{g}_2 = \frac{2\pi}{a}(1, -1, 1), \quad \mathbf{g}_3 = \frac{2\pi}{a}(1, 1, -1).$$

9/2/2015
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18

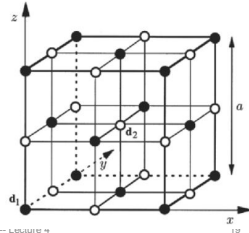
Crystal basis vectors

$$\text{Crystal structure} \Rightarrow \left\{ \begin{array}{l} \mathbf{t}_1, \mathbf{t}_2, \mathbf{t}_3 \quad (\text{primitive translation vectors}) \\ \mathbf{d}_1, \mathbf{d}_2, \dots, \mathbf{d}_v \quad (\text{basis of equal or different atoms}) \end{array} \right.$$

Example – NaCl having fcc lattice with a basis

$$\mathbf{t}_1 = \frac{a}{2}(0, 1, 1), \quad \mathbf{t}_2 = \frac{a}{2}(1, 0, 1), \quad \mathbf{t}_3 = \frac{a}{2}(1, 1, 0)$$

$$\mathbf{d}_1 = 0, \quad \mathbf{d}_2 = \frac{a}{2}(1, 1, 1)$$



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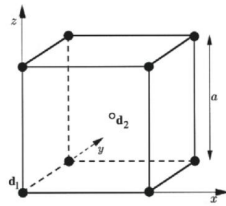
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19

Example – CsCl having sc lattice with a basis

$$\mathbf{t}_1 = a(1, 0, 0), \quad \mathbf{t}_2 = a(0, 1, 0), \quad \mathbf{t}_3 = a(0, 0, 1)$$

$$\mathbf{d}_1 = 0, \quad \mathbf{d}_2 = \frac{a}{2}(1, 1, 1)$$



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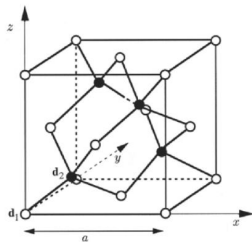
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20

Example – diamond having fcc lattice with a basis

$$\mathbf{t}_1 = \frac{a}{2}(0, 1, 1), \quad \mathbf{t}_2 = \frac{a}{2}(1, 0, 1), \quad \mathbf{t}_3 = \frac{a}{2}(1, 1, 0)$$

$$\mathbf{d}_1 = 0, \quad \mathbf{d}_2 = \frac{a}{4}(1, 1, 1)$$



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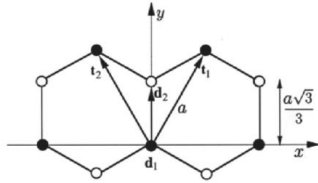
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21

Two-dimensional graphene

$$\mathbf{t}_1 = a \left(\frac{1}{2}, \frac{\sqrt{3}}{2}, 0 \right), \quad \mathbf{t}_2 = a \left(-\frac{1}{2}, \frac{\sqrt{3}}{2}, 0 \right)$$

$$\mathbf{d}_1 = 0, \quad \mathbf{d}_2 = a \left(0, \frac{\sqrt{3}}{3}, 0 \right)$$



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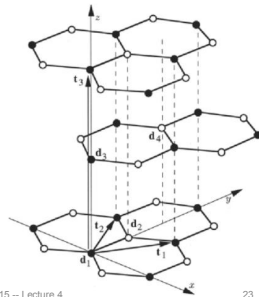
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22

Graphite lattice

$$\mathbf{t}_1 = a \left(\frac{1}{2}, \frac{\sqrt{3}}{2}, 0 \right), \quad \mathbf{t}_2 = a \left(-\frac{1}{2}, \frac{\sqrt{3}}{2}, 0 \right), \quad \mathbf{t}_3 = c(0, 0, 1)$$

$$\mathbf{d}_1 = 0, \quad \mathbf{d}_2 = a \left(0, \frac{\sqrt{3}}{3}, 0 \right), \quad \mathbf{d}_3 = c \left(0, 0, \frac{1}{2} \right), \quad \mathbf{d}_4 = \left(0, a \frac{2\sqrt{3}}{3}, \frac{c}{2} \right)$$



9/2/2015

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23
