

**PHY 711 Classical Mechanics and
Mathematical Methods**
11-11:50 AM MWF Olin 107

Plan for Lecture 30:
Chapter 11 in F & W:
Heat conduction

- 1. Basic equations**
- 2. Boundary value problems**

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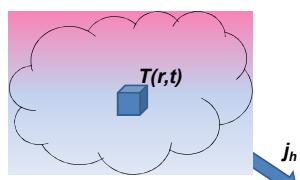
Mon	Tue	Chap.	Topic	Exam due	
20	Mon, 10/17/2016	Chap. 5	Mechanics of rigid bodies		
21	Wed, 10/19/2016	Chap. 5	Mechanics of rigid bodies	#16	10/24/2016
	Fri, 10/21/2016		Fall break - no class		
22	Mon, 10/24/2016	Chap. 8	Mechanics of Elastic Membranes	#17	10/28/2016
23	Wed, 10/26/2016	Chap. 9	Introduction to hydrodynamics	#18	10/31/2016
24	Fri, 10/28/2016	Chap. 9	Introduction to hydrodynamics	#18	10/31/2016
25	Mon, 10/31/2016	Chap. 9	Sound waves	#19	11/02/2016
26	Wed, 11/02/2016	Chap. 9	Sound waves	#20	11/04/2016
27	Fri, 11/04/2016	Chap. 9	Non-linear sound	#21	11/07/2016
28	Mon, 11/07/2016	Chap. 10	Surface waves in fluids		
29	Wed, 11/09/2016	Chap. 10	Surface waves in fluids	#22	11/11/2016
30	Fri, 11/11/2016	Chap. 11	Heat conductivity	#23	11/14/2016
31	Mon, 11/14/2016				
32	Wed, 11/16/2016				
33	Fri, 11/18/2016				
34	Mon, 11/21/2016				
	Wed, 11/23/2016		Thanksgiving Holiday -- no class		
	Fri, 11/25/2016		Thanksgiving Holiday -- no class		
	Wed, 12/07/2016		Presentations I		
	Fri, 12/09/2016		Presentations II		

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Conduction of heat



Enthalpy of a system at constant pressure p
non uniform temperature $T(\mathbf{r},t)$
mass density ρ and heat capacity c_p

$$H = \int_V \rho c_p (T(\mathbf{r},t) - T_0) d^3 r + H_0(T_0, p)$$

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Note that in this treatment we are considering a system at constant pressure p

Notation:	Heat added to system	$-dQ = TdS$
	External work done on system	$-dW = -pdV$
	Internal energy	$-dE = dQ + dW = TdS - pdV$
	Entropy	$-dS$
	Enthalpy	$-dH = d(E + pV) = TdS + Vdp$

Heat capacity at constant pressure:

$$C_p \equiv \left(\frac{\partial Q}{\partial T} \right)_p = \left(\frac{\partial H}{\partial T} \right)_p = T \left(\frac{\partial S}{\partial T} \right)_p$$

$$C_p = \rho c_p d^3 r$$

More generally, note that c_p can depend on T ; we are assume that dependence to be trivial.

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Conduction of heat -- continued

$$H = \int_V \rho c_p (T(\mathbf{r}, t) - T_0) d^3 r + H_0(T_0, p)$$

Time rate of change of enthalpy:

$$\frac{dH}{dt} = \int_V \rho c_p \frac{\partial T(\mathbf{r}, t)}{\partial t} d^3 r = - \int_A \mathbf{j}_h \cdot d\mathbf{A} + \int_V \rho \dot{q} d^3 r$$

heat flux heat source

$$\rho c_p \frac{\partial T(\mathbf{r}, t)}{\partial t} = -\nabla \cdot \mathbf{j}_h + \rho \dot{q}$$

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Conduction of heat -- continued

$$\rho c_p \frac{\partial T(\mathbf{r}, t)}{\partial t} = -\nabla \cdot \mathbf{j}_h + \rho \dot{q}$$

Empirically: $\mathbf{j}_h = -k_{th} \nabla T(\mathbf{r}, t)$

$$\Rightarrow \frac{\partial T(\mathbf{r}, t)}{\partial t} = \kappa \nabla^2 T(\mathbf{r}, t) + \frac{\dot{q}}{c_p}$$

$$\kappa \equiv \frac{k_{th}}{\rho c_p} \quad \text{thermal diffusivity}$$

Typical values (m²/s)

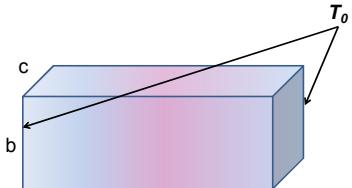
Air	2×10^{-5}
Water	1×10^{-7}
Copper	1×10^{-4}

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Boundary value problems for heat conduction



$$\frac{\partial T(\mathbf{r}, t)}{\partial t} - \kappa \nabla^2 T(\mathbf{r}, t) = \frac{\dot{q}}{c_p}$$

$$\text{Without source term: } \frac{\partial T(\mathbf{r}, t)}{\partial t} - \kappa \nabla^2 T(\mathbf{r}, t) = 0$$

$$\text{Example with boundary values: } T(0, y, z, t) = T(a, y, z, t) = T_0$$

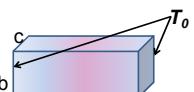
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Boundary value problems for heat conduction

$$\frac{\partial T(\mathbf{r}, t)}{\partial t} - \kappa \nabla^2 T(\mathbf{r}, t) = 0$$

$$T(0, y, z, t) = T(a, y, z, t) = T_0$$



$$\left. \begin{aligned} \frac{\partial T(x, 0, z, t)}{\partial y} &= \frac{\partial T(x, b, z, t)}{\partial y} = 0 \\ \frac{\partial T(x, y, 0, t)}{\partial z} &= \frac{\partial T(x, y, c, t)}{\partial z} = 0 \end{aligned} \right\} \text{Assuming thermally insulated boundaries}$$

$$\text{Separation of variables: } T(x, y, z, t) = T_0 + X(x)Y(y)Z(z)e^{-\lambda t}$$

$$\text{Let } \frac{d^2X}{dx^2} = -\alpha^2 X \quad \frac{d^2Y}{dy^2} = -\beta^2 Y \quad \frac{d^2Z}{dz^2} = -\gamma^2 Z$$

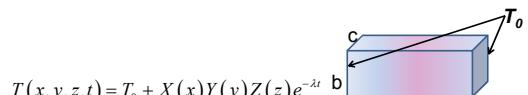
$$\Rightarrow -\lambda + \kappa(\alpha^2 + \beta^2 + \gamma^2) = 0$$

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Boundary value problems for heat conduction



$$T(x, y, z, t) = T_0 + X(x)Y(y)Z(z)e^{-\lambda t}$$

$$X(0) = X(a) = 0 \quad \Rightarrow X(x) = \sin\left(\frac{m\pi x}{a}\right)$$

$$\frac{dY(0)}{dy} = \frac{dY(b)}{dy} = 0 \quad \Rightarrow Y(y) = \cos\left(\frac{n\pi y}{b}\right)$$

$$\frac{dZ(0)}{dz} = \frac{dZ(c)}{dz} = 0 \quad \Rightarrow Z(z) = \cos\left(\frac{p\pi z}{c}\right)$$

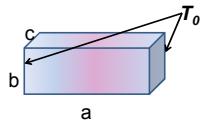
$$-\lambda_{nmp} + \kappa \left[\left(\frac{m\pi}{a}\right)^2 + \left(\frac{n\pi}{b}\right)^2 + \left(\frac{p\pi}{c}\right)^2 \right] = 0$$

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Boundary value problems for heat conduction



Full solution:

$$T(x, y, z, t) = T_0 + \sum_{nmp} C_{nmp} \sin\left(\frac{m\pi x}{a}\right) \cos\left(\frac{n\pi y}{b}\right) \cos\left(\frac{p\pi z}{c}\right) e^{-\lambda_{nmp} t}$$

$$\lambda_{nmp} = \kappa \left(\left(\frac{n\pi}{a}\right)^2 + \left(\frac{m\pi}{b}\right)^2 + \left(\frac{p\pi}{c}\right)^2 \right)$$

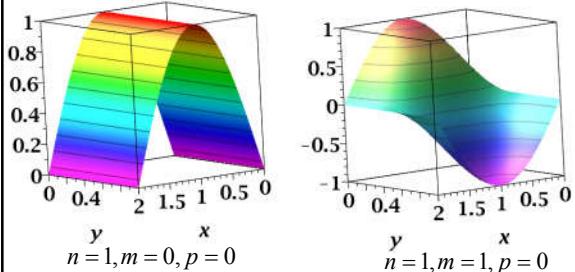
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Full solution:

$$T(x, y, z, t) = T_0 + \sum_{nmp} C_{nmp} \sin\left(\frac{m\pi x}{a}\right) \cos\left(\frac{n\pi y}{b}\right) \cos\left(\frac{p\pi z}{c}\right) e^{-\lambda_{nmp} t}$$



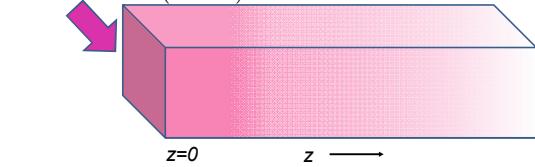
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Oscillatory thermal behavior

$$T(z = 0, t) = \Re(T_0 e^{-i\omega t})$$



$$\frac{\partial T}{\partial t} = \kappa \frac{\partial^2 T}{\partial z^2}$$

$$\text{Assume: } T(z, t) = \Re(f(z) e^{-i\omega t})$$

$$(-i\omega) f = \kappa \frac{d^2 f}{dz^2}$$

$$\text{Let } f(z) = A e^{i\alpha z}$$

$$\alpha^2 = -\frac{i\omega}{\kappa} = e^{3i\pi/2} \frac{\omega}{\kappa}$$

$$\alpha = \pm(1-i)\sqrt{\frac{\omega}{2\kappa}}$$

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Oscillatory thermal behavior -- continued

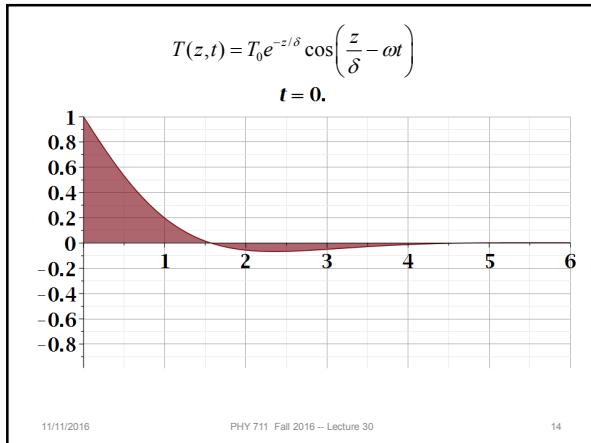
$$T(z=0, t) = \Re(T_0 e^{-i\omega t})$$

$$T(z, t) = \Re(A e^{\pm(1-i)z/\delta} e^{-i\omega t})$$

$$\text{where } \delta \equiv \sqrt{\frac{2\kappa}{\omega}}$$

Physical solution: $T(z, t) = T_0 e^{-z/\delta} \cos\left(\frac{z}{\delta} - \omega t\right)$

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Initial value problem in an infinite domain; Fourier transform

$$\frac{\partial T(\mathbf{r}, t)}{\partial t} - \kappa \nabla^2 T(\mathbf{r}, t) = 0$$

$$T(\mathbf{r}, 0) = f(\mathbf{r})$$

Let: $\tilde{T}(\mathbf{q}, t) = \int d^3 r e^{-i\mathbf{q} \cdot \mathbf{r}} T(\mathbf{r}, t)$

$$\tilde{f}(\mathbf{q}) = \int d^3 r e^{-i\mathbf{q} \cdot \mathbf{r}} f(\mathbf{r})$$

$$\Rightarrow \tilde{T}(\mathbf{q}, 0) = \tilde{f}(\mathbf{q})$$

$$\Rightarrow \frac{\partial \tilde{T}(\mathbf{q}, t)}{\partial t} = -\kappa q^2 \tilde{T}(\mathbf{q}, t)$$

$$\tilde{T}(\mathbf{q}, t) = \tilde{T}(\mathbf{q}, 0) e^{-\kappa q^2 t}$$

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Initial value problem in an infinite domain; Fourier transform

$$\tilde{T}(\mathbf{q}, t) = \int d^3r e^{-i\mathbf{q} \cdot \mathbf{r}} T(\mathbf{r}, t) \Rightarrow T(\mathbf{r}, t) = \frac{1}{(2\pi)^3} \int d^3q e^{i\mathbf{q} \cdot \mathbf{r}} \tilde{T}(\mathbf{q}, t)$$

$$\tilde{T}(\mathbf{q}, t) = \tilde{T}(\mathbf{q}, 0) e^{-\kappa q^2 t}$$

$$T(\mathbf{r}, t) = \frac{1}{(2\pi)^3} \int d^3q e^{i\mathbf{q} \cdot \mathbf{r}} \tilde{T}(\mathbf{q}, 0) e^{-\kappa q^2 t}$$

$$\tilde{T}(\mathbf{q}, 0) = \tilde{f}(\mathbf{q}) = \int d^3r e^{-i\mathbf{q} \cdot \mathbf{r}} f(\mathbf{r})$$

$$T(\mathbf{r}, t) = \int d^3r' G(\mathbf{r} - \mathbf{r}', t) T(\mathbf{r}', 0)$$

$$\text{with } G(\mathbf{r} - \mathbf{r}', t) \equiv \frac{1}{(2\pi)^3} \int d^3q e^{i\mathbf{q} \cdot (\mathbf{r} - \mathbf{r}')} e^{-\kappa q^2 t}$$

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Initial value problem in an infinite domain; Fourier transform

$$T(\mathbf{r}, t) = \int d^3r' G(\mathbf{r} - \mathbf{r}', t) T(\mathbf{r}', 0)$$

$$\text{with } G(\mathbf{r} - \mathbf{r}', t) \equiv \frac{1}{(2\pi)^3} \int d^3q e^{i\mathbf{q} \cdot (\mathbf{r} - \mathbf{r}')} e^{-\kappa q^2 t}$$

$$G(\mathbf{r} - \mathbf{r}', t) = \frac{1}{(4\pi\kappa t)^{3/2}} e^{-|\mathbf{r} - \mathbf{r}'|^2 / (4\kappa t)}$$

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Heat equation in half-space

$$\frac{\partial T(\mathbf{r}, t)}{\partial t} - \kappa \nabla^2 T(\mathbf{r}, t) = 0$$

$T(\mathbf{r}, t) \Rightarrow T(z, t)$ with initial and boundary values :

$$T(z, t) \equiv 0 \text{ for } z < 0$$

$$T(z, 0) = 0 \text{ for } z > 0$$

$$T(0, t) = T_0 \text{ for } t \geq 0$$

$$\text{Solution : } T = T_0 \operatorname{erfc}\left(\frac{z}{2\sqrt{\kappa t}}\right)$$

$$\text{where } \operatorname{erfc}(x) \equiv \frac{2}{\sqrt{\pi}} \int_x^\infty e^{-u^2} du$$

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Heat equation in half-space -- continued

$$\frac{\partial T(z,t)}{\partial t} - \kappa \frac{\partial^2 T(z,t)}{\partial z^2} = 0$$

Solution : $T = T_0 \operatorname{erfc}\left(\frac{z}{2\sqrt{\kappa t}}\right)$

where $\operatorname{erfc}(x) = \frac{2}{\sqrt{\pi}} \int_x^\infty e^{-u^2} du$

Note that $\frac{d \operatorname{erfc}(x)}{dx} = \frac{d}{dx} \frac{2}{\sqrt{\pi}} \int_x^\infty e^{-u^2} du = -\frac{2}{\sqrt{\pi}} e^{-x^2}$

$$\frac{\partial}{\partial t} \operatorname{erfc}\left(\frac{z}{2\sqrt{\kappa t}}\right) = \frac{2}{\sqrt{\pi}} e^{-(z^2/(4\kappa t))} \left(\frac{z}{4\sqrt{\kappa t^3}} \right)$$

$$\frac{\partial^2}{\partial z^2} \operatorname{erfc}\left(\frac{z}{2\sqrt{\kappa t}}\right) = \frac{2}{\sqrt{\pi}} e^{-(z^2/(4\kappa t))} \left(\frac{z}{4\kappa \sqrt{\kappa t^3}} \right)$$

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