

**PHY 711 Classical Mechanics and  
Mathematical Methods  
9-9:50 AM MWF Olin 107**

**Plan for Lecture 21:**

**Read Chapter 7 & Appendices A-D**

Generalization of the one dimensional wave equation → various mathematical problems and techniques including:

1. Sturm-Liouville equations
2. Orthogonal function expansions; Fourier analysis
3. Green's functions methods
4. Laplace transformation
5. Contour integration methods

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Mon, 10/02/2017		Take-home exam -- No class		
Wed, 10/04/2017		Take-home exam -- No class		
16 Fri, 10/06/2017	Chap. 4	Vibrational modes		
17 Mon, 10/09/2017	Chap. 5	Rigid body motion	#10	10/16/2017
18 Wed, 10/11/2017	Chap. 5	Rigid body motion		
Fri, 10/13/2017		Fall break -- No class		
19 Mon, 10/16/2017		Discuss exam questions and topics for presentations	Topic	10/18/2017
20 Wed, 10/18/2017	Chap. 7	Wave equation in one dimension	#11	10/20/2017
21 Fri, 10/20/2017	Chap. 7	Solutions of Sturm-Liouville equations	#12	10/25/2017
22 Mon, 10/23/2017				
23 Wed, 10/25/2017				
24 Fri, 10/27/2017				
25 Mon, 10/30/2017				
26 Wed, 11/01/2017				
27 Fri, 11/03/2017				
28 Mon, 11/06/2017				
29 Wed, 11/08/2017				
30 Fri, 11/10/2017				

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One dimensional wave equation for  $\mu(x,t)$ :

$$\frac{\partial^2 \mu}{\partial t^2} - c^2 \frac{\partial^2 \mu}{\partial x^2} = 0 \quad \text{where } c^2 = \frac{\tau}{\sigma}$$

Generalization for spacially dependent tension and mass density plus an extra potential energy density:

$$\sigma(x) \frac{\partial^2 \mu}{\partial t^2} - \frac{\partial}{\partial x} \left( \tau(x) \frac{\partial \mu}{\partial x} \right) + v(x) \mu = 0$$

Separating time and spatial variables:

$$\mu(x,t) = \rho(x) \cos(\omega t + \phi)$$

Sturm-Liouville equation for spatial function:

$$-\frac{d}{dx} \left( \tau(x) \frac{d\rho}{dx} \right) + v(x) \rho = \omega^2 \sigma(x) \rho$$

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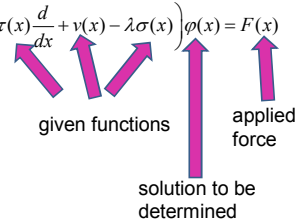
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Linear second-order ordinary differential equations  
Sturm-Liouville equations

Inhomogenous problem:  $\left(-\frac{d}{dx}\tau(x)\frac{d}{dx} + v(x) - \lambda\sigma(x)\right)\varphi(x) = F(x)$



Homogenous problem:  $F(x)=0$

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Examples of Sturm-Liouville eigenvalue equations --

$\left(-\frac{d}{dx}\tau(x)\frac{d}{dx} + v(x) - \lambda\sigma(x)\right)\varphi(x) = 0$

Bessel functions:  
 $\tau(x) = -x \quad v(x) = x \quad \sigma(x) = \frac{1}{x} \quad \lambda = \nu^2 \quad \varphi(x) = J_\nu(x)$

Legendre functions:  
 $\tau(x) = -(1-x^2) \quad v(x) = 0 \quad \sigma(x) = 1 \quad \lambda = l(l+1) \quad \varphi(x) = P_l(x)$

Fourier functions:  
 $\tau(x) = 1 \quad v(x) = 0 \quad \sigma(x) = 1 \quad \lambda = n^2\pi^2 \quad \varphi(x) = \sin(n\pi x)$

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Solution methods of Sturm-Liouville equations  
(assume all functions and constants are real):

Homogenous problem:  $\left(-\frac{d}{dx}\tau(x)\frac{d}{dx} + v(x) - \lambda\sigma(x)\right)\phi_0(x) = 0$

Inhomogenous problem:  $\left(-\frac{d}{dx}\tau(x)\frac{d}{dx} + v(x) - \lambda\sigma(x)\right)\phi(x) = F(x)$

Eigenfunctions:

$\left(-\frac{d}{dx}\tau(x)\frac{d}{dx} + v(x)\right)f_n(x) = \lambda_n\sigma(x)f_n(x)$

Orthogonality of eigenfunctions:  $\int_a^b \sigma(x)f_n(x)f_m(x)dx = \delta_{nm}N_n$ ,  
where  $N_n \equiv \int_a^b \sigma(x)(f_n(x))^2 dx$ .

Completeness of eigenfunctions:  
 $\sigma(x)\sum \frac{f_n(x)f_n(x')}{N_n} = \delta(x-x')$

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Comment on "completeness"

It can be shown that for any reasonable function  $h(x)$ , defined within the interval  $a < x < b$ , we can expand that function as a linear combination of the eigenfunctions  $f_n(x)$

$$h(x) \approx \sum_n C_n f_n(x),$$

where  $C_n = \frac{1}{N_n} \int_a^b \sigma(x') h(x') f_n(x') dx'$ .

These ideas lead to the notion that the set of eigenfunctions  $f_n(x)$  form a "complete" set in the sense of "spanning" the space of all functions in the interval  $a < x < b$ , as summarized by the statement:

$$\sigma(x) \sum_n \frac{f_n(x) f_n(x')}{N_n} = \delta(x - x').$$

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Variation approximation to lowest eigenvalue

In general, there are several techniques to determine the eigenvalues  $\lambda_n$  and eigenfunctions  $f_n(x)$ . When it is not possible to find the "exact" functions, there are several powerful approximation techniques. For example, the lowest eigenvalue can be approximated by minimizing the function

$$\lambda_0 \leq \frac{\langle \tilde{h} | S | \tilde{h} \rangle}{\langle \tilde{h} | \sigma | \tilde{h} \rangle}, \quad S(x) \equiv -\frac{d}{dx} \tau(x) \frac{d}{dx} + v(x)$$

where  $\tilde{h}(x)$  is a variable function which satisfies the correct boundary values. The "proof" of this inequality is based on the notion that  $\tilde{h}(x)$  can in principle be expanded in terms of the (unknown) exact eigenfunctions  $f_n(x)$ :

$$\tilde{h}(x) = \sum_n C_n f_n(x), \quad \text{where the coefficients } C_n \text{ can be}$$

assumed to be real.

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Estimation of the lowest eigenvalue – continued:

From the eigenfunction equation, we know that

$$S(x) \tilde{h}(x) = S(x) \sum_n C_n f_n(x) = \sum_n C_n \lambda_n \sigma(x) f_n(x).$$

It follows that:

$$\langle \tilde{h} | S | \tilde{h} \rangle = \int_a^b \tilde{h}(x) S(x) \tilde{h}(x) dx = \sum_n |C_n|^2 N_n \lambda_n.$$

It also follows that:

$$\langle \tilde{h} | \sigma | \tilde{h} \rangle = \int_a^b \tilde{h}(x) \sigma(x) \tilde{h}(x) dx = \sum_n |C_n|^2 N_n,$$

Therefore  $\frac{\langle \tilde{h} | S | \tilde{h} \rangle}{\langle \tilde{h} | \sigma | \tilde{h} \rangle} = \frac{\sum_n |C_n|^2 N_n \lambda_n}{\sum_n |C_n|^2 N_n} \geq \lambda_0.$

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Rayleigh-Ritz method of estimating the lowest eigenvalue

$$\lambda_0 \leq \frac{\langle \tilde{h} | S | \tilde{h} \rangle}{\langle \tilde{h} | \sigma | \tilde{h} \rangle},$$

Example:  $-\frac{d^2}{dx^2} f_n(x) = \lambda_n f_n(x)$  with  $f_n(0) = f_n(a) = 0$   
 trial function  $f_{\text{trial}}(x) = x(x-a)$

Exact value of  $\lambda_0 = \frac{\pi^2}{a^2} = \frac{9.869604404}{a^2}$

Rayleigh-Ritz estimate:  $\frac{\langle x(a-x) | -\frac{d^2}{dx^2} | x(a-x) \rangle}{\langle x(a-x) | x(a-x) \rangle} = \frac{10}{a^2}$

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Green's function solution methods

Suppose that we can find a Green's function defined as follows:

$$\left( -\frac{d}{dx} \tau(x) \frac{d}{dx} + v(x) - \lambda \sigma(x) \right) G_\lambda(x, x') = \delta(x - x')$$

Completeness of eigenfunctions:

Recall:  $\sigma(x) \sum_n \frac{f_n(x) f_n(x')}{N_n} = \delta(x - x')$

In terms of eigenfunctions:

$$\left( -\frac{d}{dx} \tau(x) \frac{d}{dx} + v(x) - \lambda \sigma(x) \right) G_\lambda(x, x') = \sigma(x) \sum_n \frac{f_n(x) f_n(x')}{N_n}$$

$$\Rightarrow G_\lambda(x, x') = \sum_n \frac{f_n(x) f_n(x') / N_n}{\lambda_n - \lambda}$$

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Solution to inhomogeneous problem by using Green's functions

Inhomogeneous problem:

$$\left( -\frac{d}{dx} \tau(x) \frac{d}{dx} + v(x) - \lambda \sigma(x) \right) \varphi(x) = F(x)$$

Green's function :

$$\left( -\frac{d}{dx} \tau(x) \frac{d}{dx} + v(x) - \lambda \sigma(x) \right) G_\lambda(x, x') = \delta(x - x')$$

Formal solution:

$$\varphi_\lambda(x) = \varphi_{\lambda 0}(x) + \int_0^L G_\lambda(x, x') F(x') dx'$$

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Solution to homogeneous problem

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Example Sturm-Liouville problem:

Example:  $\tau(x)=1; \sigma(x)=1; v(x)=0; a=0$  and  $b=L$

$$\lambda=1; F(x)=F_0 \sin\left(\frac{\pi x}{L}\right)$$

Inhomogenous equation :

$$\left(-\frac{d^2}{dx^2}-1\right)\phi(x)=F_0 \sin\left(\frac{\pi x}{L}\right)$$

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Eigenvalue equation :

$$\left(-\frac{d^2}{dx^2}\right)f_n(x)=\lambda_n f_n(x)$$

Eigenfunctions

$$f_n(x)=\sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right)$$

Eigenvalues :

$$\lambda_n=\left(\frac{n\pi}{L}\right)^2$$

Completeness of eigenfunctions :

$$\sigma(x)\sum_n \frac{f_n(x)f_n(x')}{N_n}=\delta(x-x')$$

In this example :  $\frac{2}{L}\sum_n \sin\left(\frac{n\pi x}{L}\right)\sin\left(\frac{n\pi x'}{L}\right)=\delta(x-x')$

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Green's function :

$$\left(-\frac{d}{dx}\tau(x)\frac{d}{dx}+v(x)-\lambda\sigma(x)\right)G_\lambda(x,x')=\delta(x-x')$$

Green's function for the example :

$$G(x,x')=\sum_n \frac{f_n(x)f_n(x')/N_n}{\lambda_n-\lambda}=\frac{2}{L}\sum_n \frac{\sin\left(\frac{n\pi x}{L}\right)\sin\left(\frac{n\pi x'}{L}\right)}{\left(\frac{n\pi}{L}\right)^2-1}$$

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Using Green's function to solve inhomogenous equation :

$$\left(-\frac{d^2}{dx^2} - 1\right)\phi(x) = F_0 \sin\left(\frac{\pi x}{L}\right)$$

$$\phi(x) = \phi_0(x) + \int_0^L G(x, x') F_0 \sin\left(\frac{\pi x'}{L}\right) dx'$$

$$= \phi_0(x) + \frac{2}{L} \sum_n \left[ \frac{\sin\left(\frac{n\pi x}{L}\right)}{\left(\frac{n\pi}{L}\right)^2 - 1} \int_0^L \sin\left(\frac{n\pi x'}{L}\right) F_0 \sin\left(\frac{\pi x'}{L}\right) dx' \right]$$

$$= \phi_0(x) + \frac{F_0}{\left(\frac{\pi}{L}\right)^2 - 1} \sin\left(\frac{\pi x}{L}\right)$$

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Alternate Green's function method :

$$G(x, x') = \frac{1}{W} g_a(x_<) g_b(x_>)$$

$$\left(-\frac{d^2}{dx^2} - 1\right)g_i(x) = 0 \quad \Rightarrow g_a(x) = \sin(x); \quad g_b(x) = \sin(L-x);$$

$$W = g_a(x) \frac{dg_b(x)}{dx} - g_b(x) \frac{dg_a(x)}{dx} = \sin(L-x)\cos(x) + \sin(x)\cos(L-x)$$

$$= \sin(L)$$

$$\phi(x) = \phi_0(x) + \frac{\sin(L-x)}{\sin(L)} \int_0^x \sin(x') F_0 \sin\left(\frac{\pi x'}{L}\right) dx'$$

$$+ \frac{\sin(x)}{\sin(L)} \int_x^L \sin(L-x') F_0 \sin\left(\frac{\pi x'}{L}\right) dx'$$

$$\phi(x) = \phi_0(x) + \frac{F_0}{\left(\frac{\pi}{L}\right)^2 - 1} \sin\left(\frac{\pi x}{L}\right)$$

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General method of constructing Green's functions using homogeneous solution

Green's function :

$$\left(-\frac{d}{dx} \tau(x) \frac{d}{dx} + v(x) - \lambda \sigma(x)\right) G_i(x, x') = \delta(x - x')$$

Two homogeneous solutions

$$\left(-\frac{d}{dx} \tau(x) \frac{d}{dx} + v(x) - \lambda \sigma(x)\right) g_i(x) = 0 \quad \text{for } i = a, b$$

Let

$$G_i(x, x') = \frac{1}{W} g_a(x_<) g_b(x_>)$$

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For  $\epsilon \rightarrow 0$ :

$$\int_{x'-\epsilon}^{x'+\epsilon} dx \left( -\frac{d}{dx} \tau(x) \frac{d}{dx} + v(x) - \lambda \sigma(x) \right) G_{\lambda}(x, x') = \int_{x'-\epsilon}^{x'+\epsilon} dx \delta(x - x')$$

$$\int_{x'-\epsilon}^{x'+\epsilon} dx \left( -\frac{d}{dx} \tau(x) \frac{d}{dx} \right) \frac{1}{W} g_a(x_{<}) g_b(x_{>}) = 1$$

$$-\frac{\tau(x)}{W} \left( \frac{d}{dx} g_a(x_{<}) g_b(x_{>}) \right) \Big|_{x'-\epsilon}^{x'+\epsilon} = \frac{\tau(x')}{W} \left( g_a(x') \frac{d}{dx} g_b(x') - g_b(x') \frac{d}{dx} g_a(x') \right)$$

$$\Rightarrow W = \tau(x') \left( g_a(x') \frac{d}{dx} g_b(x') - g_b(x') \frac{d}{dx} g_a(x') \right)$$

Note --  $W$  (Wronskian) is constant, since  $\frac{dW}{dx'} = 0$ .

$\Rightarrow$  Useful Green's function construction in one dimension:

$$G_{\lambda}(x, x') = \frac{1}{W} g_a(x_{<}) g_b(x_{>})$$

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$$\left( -\frac{d}{dx} \tau(x) \frac{d}{dx} + v(x) - \lambda \sigma(x) \right) \phi(x) = F(x)$$

Green's function solution:

$$\phi_{\lambda}(x) = \phi_{\lambda 0}(x) + \int_{x_i}^{x_f} G_{\lambda}(x, x') F(x') dx'$$

$$= \phi_{\lambda 0}(x) + \frac{g_b(x)}{W} \int_{x_i}^x g_a(x') F(x') dx' + \frac{g_a(x)}{W} \int_x^{x_f} g_b(x') F(x') dx'$$

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