

**PHY 711 Classical Mechanics and
Mathematical Methods**
9-9:50 AM MWF Olin 107
Plan for Lecture 26:

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Events

Meet the Speaker: Nov 1 at Noon
WFU Physics Meet the Speaker Event!
SPEAKER: Dr. Yves Chabal Department of Materials Science and Engineering, University of Texas at Dallas, Dallas, TX
TIME: Wed. Nov. 1, 2017, 12:00 - 1:00 ...

Colloquium: Nov. 1, 2017 at 4 PM
WFU Physics Colloquium TITLE: "The fascinating world of Metal Organic Framework materials" **SPEAKER:** Yves Chabal Department of Materials Science and Engineering, University of Texas at Dallas, Dallas, TX **TIME:** Wed.

19 Mon, 10/16/2017		Discuss exam questions and topics for presentations	Topic	10/18/2017
20 Wed, 10/18/2017	Chap. 7	Wave equation in one dimension	#11	10/20/2017
21 Fri, 10/20/2017	Chap. 7	Solutions of Sturm-Liouville equations	#12	10/27/2017
22 Mon, 10/23/2017	Chap. 7	Solutions of Sturm-Liouville equations		
23 Wed, 10/25/2017	Chap. 7	Solutions of Sturm-Liouville equations		
24 Fri, 10/27/2017	App. A	Laplace transforms and contour integrals	#13	11/01/2017
25 Mon, 10/30/2017	App. A	Contour integrals		
26 Wed, 11/01/2017	Chap. 8	Mechanics of Elastic Membranes	#14	11/06/2017
27 Fri, 11/03/2017				
28 Mon, 11/06/2017				
29 Wed, 11/08/2017				
30 Fri, 11/10/2017				
31 Mon, 11/13/2017				
32 Wed, 11/15/2017				
33 Fri, 11/17/2017				
34 Mon, 11/20/2017				
Wed, 11/22/2017		<i>Thanksgiving Holiday – No class</i>		
Fri, 11/24/2017		<i>Thanksgiving Holiday – No class</i>		
35 Mon, 11/27/2017				

Elastic media in two or more dimensions --

Review of wave equation in one-dimension – here $\mu(x,t)$ can describe either a longitudinal or transverse wave.

Traveling wave solutions --

$$\frac{\partial^2 \mu}{\partial t^2} - c^2 \frac{\partial^2 \mu}{\partial x^2} = 0$$

Note that for any function $f(q)$ or $g(q)$:

$$\mu(x,t) = f(x-ct) + g(x+ct)$$

satisfies the wave equation.

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Initial value problem : $\mu(x,0) = \phi(x)$ and $\frac{\partial \mu}{\partial t}(x,0) = \psi(x)$

then : $\mu(x,0) = \phi(x) = f(x) + g(x)$

$$\frac{\partial \mu}{\partial t}(x,0) = \psi(x) = -c \left(\frac{df(x)}{dx} - \frac{dg(x)}{dx} \right)$$

$$\Rightarrow f(x) - g(x) = -\frac{1}{c} \int_{x-ct}^x \psi(x') dx'$$

For each x , find $f(x)$ and $g(x)$:

$$f(x) = \frac{1}{2} \left(\phi(x) - \frac{1}{c} \int_{x-ct}^x \psi(x') dx' \right)$$

$$g(x) = \frac{1}{2} \left(\phi(x) + \frac{1}{c} \int_{x-ct}^x \psi(x') dx' \right)$$

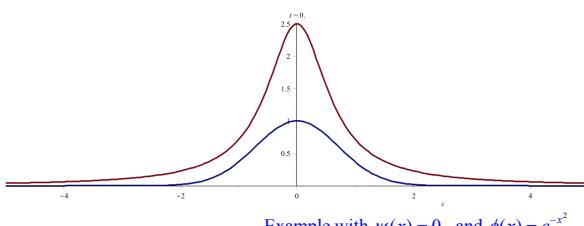
$$\Rightarrow \mu(x,t) = \frac{1}{2} (\phi(x-ct) + \phi(x+ct)) + \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(x') dx'$$

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Example with $\psi(x) = 0$ and $\phi(x) = \frac{1}{x^2 + 0.4}$



Example with $\psi(x) = 0$ and $\phi(x) = e^{-x^2}$

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Standing wave solutions of wave equation:

$$\frac{\partial^2 \mu}{\partial t^2} - c^2 \frac{\partial^2 \mu}{\partial x^2} = 0$$

with $\mu(0,t) = \mu(L,t) = 0$.

$$\text{Assume: } \mu(x,t) = \Re(e^{-i\omega t} \rho(x))$$

where $\frac{d^2 \rho(x)}{dx^2} + k^2 \rho(x) = 0$

$$k = \frac{\omega}{c}$$

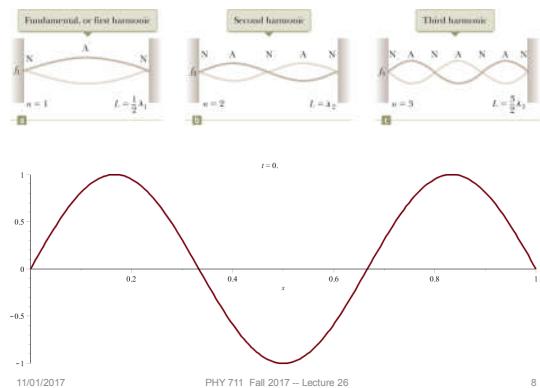
$$\rho_v(x) = A \sin\left(\frac{v\pi x}{L}\right)$$

$$k_\nu = \frac{\nu\pi}{L} \quad \omega_\nu = ck_\nu$$

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Wave motion on a two-dimensional surface – elastic membrane (transverse wave; linear regime).

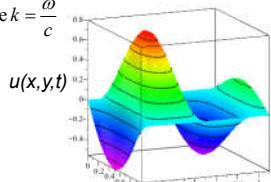
Two-dimensional wave equation:

$$\frac{\partial^2 u}{\partial t^2} - c^2 \nabla^2 u = 0 \quad \text{where } c^2 = \frac{\tau}{\sigma}$$

Standing wave solutions:

$$u(x, y, t) = \Re(e^{-i\omega t} \rho(x, y))$$

$$(\nabla^2 + k^2)\rho(x, y) = 0 \quad \text{where } k =$$



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Lagrangian density: $\mathcal{L}\left(u, \frac{\partial u}{\partial x}, \frac{\partial u}{\partial y}, \frac{\partial u}{\partial t}; x, y, t\right)$

$$L = \int \mathcal{L}\left(u, \frac{\partial u}{\partial x}, \frac{\partial u}{\partial y}, \frac{\partial u}{\partial t}; x, y, t\right) dx dy$$

Hamilton's principle:

$$\delta \int_{t_1}^{t_2} L dt = 0$$

$$\frac{\partial \mathcal{L}}{\partial u} - \frac{\partial}{\partial t} \frac{\partial \mathcal{L}}{\partial (\partial u / \partial t)} - \frac{\partial}{\partial x} \frac{\partial \mathcal{L}}{\partial (\partial u / \partial x)} - \frac{\partial}{\partial y} \frac{\partial \mathcal{L}}{\partial (\partial u / \partial y)} = 0$$

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Lagrangian density for elastic membrane with constant σ and τ :

$$\mathcal{L}\left(u, \frac{\partial u}{\partial x}, \frac{\partial u}{\partial y}, \frac{\partial u}{\partial t}; x, y, t\right) = \frac{1}{2} \sigma \left(\frac{\partial u}{\partial t}\right)^2 - \frac{1}{2} \tau (\nabla u)^2$$

$$\frac{\partial \mathcal{L}}{\partial u} - \frac{\partial}{\partial t} \frac{\partial \mathcal{L}}{\partial (\partial u / \partial t)} - \frac{\partial}{\partial x} \frac{\partial \mathcal{L}}{\partial (\partial u / \partial x)} - \frac{\partial}{\partial y} \frac{\partial \mathcal{L}}{\partial (\partial u / \partial y)} = 0$$

$$\frac{\partial^2 u}{\partial t^2} - c^2 \nabla^2 u = 0 \quad \text{where } c^2 = \frac{\tau}{\sigma}$$

Two-dimensional wave equation:

$$\frac{\partial^2 u}{\partial t^2} - c^2 \nabla^2 u = 0 \quad \text{where } c^2 = \frac{\tau}{\sigma}$$

Standing wave solutions:

$$u(x, y, t) = \Re(e^{-i\omega t} \rho(x, y))$$

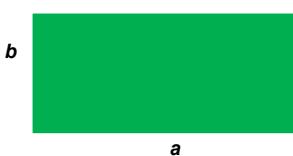
$$(\nabla^2 + k^2) \rho(x, y) = 0 \quad \text{where } k = \frac{\omega}{c}$$

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Consider a rectangular boundary:



Clamped boundary conditions:

$$\rho(0, y) = \rho(a, y) = \rho(x, 0) = \rho(x, b) = 0 \quad (\nabla^2 + k^2) \rho(x, y) = 0$$

$$\Rightarrow \rho_{mn}(x, y) = A \sin\left(\frac{m\pi x}{a}\right) \sin\left(\frac{n\pi y}{b}\right) \quad \text{where } k = \frac{\omega}{c}$$

$$k_{mn}^2 = \left(\frac{m\pi}{a}\right)^2 + \left(\frac{n\pi}{b}\right)^2 \quad \omega_{mn} = ck_{mn}$$

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Series expansions of Bessel and Neumann functions

$$J_v(z) = \left(\frac{1}{2}z\right)^v \sum_{k=0}^{\infty} (-1)^k \frac{\left(\frac{1}{4}z^2\right)^k}{k!\Gamma(v+k+1)}.$$

$$Y_n(z) = -\frac{\left(\frac{1}{2}z\right)^{-n}}{\pi} \sum_{k=0}^{n-1} \frac{(n-k-1)!}{k!} \left(\frac{1}{4}z^2\right)^k + \frac{2}{\pi} \ln\left(\frac{1}{2}z\right) J_n(z)$$

$$-\frac{\left(\frac{1}{2}z\right)^n}{\pi} \sum_{k=0}^{\infty} (\psi(k+1) + \psi(n+k+1)) \frac{\left(-\frac{1}{4}z^2\right)^k}{k!(n+k)!}$$

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Some properties of Bessel functions -- continued

Note : It is possible to prove the following

identity for the functions $J_m\left(\frac{z_{mn}}{R}r\right)$:

$$\int_0^R J_m\left(\frac{z_{mn}}{R}r\right) J_m\left(\frac{z_{mn'}}{R}r\right) r dr = \frac{R^2}{2} (J_{m+1}(z_{mn}))^2 \delta_{mn'}$$

Returning to differential equation for radial function :

$$\left(\frac{d^2}{dr^2} + \frac{1}{r} \frac{d}{dr} - \frac{m^2}{r^2} + k^2 \right) f(r) = 0$$

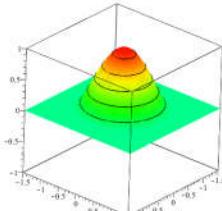
$$\Rightarrow f_{mn}(r) = AJ_m\left(\frac{z_{mn}}{R}r\right); \quad k_{mn} = \frac{z_{mn}}{R}$$

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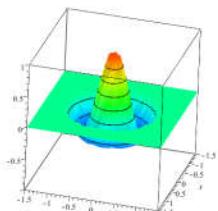
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$$\rho_{01}(r, \varphi) = f_{01}(r) = AJ_0\left(\frac{z_{01}}{R}r\right) \quad \rho_{02}(r, \varphi) = f_{02}(r) = AJ_0\left(\frac{z_{02}}{R}r\right)$$



$$k_{01} = \frac{2.406}{R}$$



$$k_{02} = \frac{5.520}{R}$$

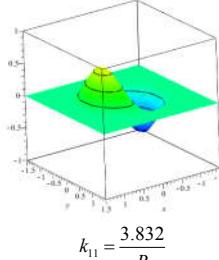
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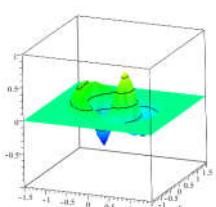
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$$\begin{aligned}\rho_{11}(r, \varphi) &= f_{11}(r) \cos(\varphi) \\ &= AJ_1\left(\frac{z_{11}}{R} r\right) \cos(\varphi)\end{aligned}$$

$$\begin{aligned}\rho_{12}(r, \varphi) &= f_{12}(r) \cos(\varphi) \\ &= AJ_1\left(\frac{z_{12}}{R} r\right) \cos(\varphi)\end{aligned}$$



$$k_{11} = \frac{3.852}{R}$$



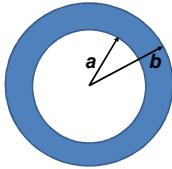
$$r_{12} = \frac{7.016}{R}$$

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More complicated geometry – annular membrane



In cylindrical coordinate system

Assume : $\rho(r, \varphi) = f(r)\Phi(\varphi)$

$$\text{Let : } \Phi(\varphi) = e^{im\varphi}$$

Note : $\Phi(\varphi) = \Phi(\varphi + 2\pi)$
 $\Rightarrow m = \text{integer}$

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Consider circular boundary -- continued

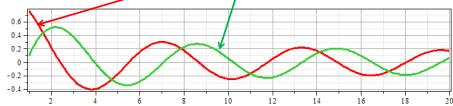
Differential equation for radial function :

$$\left(\frac{d^2}{dr^2} + \frac{1}{r} \frac{d}{dr} - \frac{m^2}{r^2} + k^2 \right) f(r) = 0$$

⇒ Bessel equation of integer order with transcendental solutions

Cylindrical Bessel function $J_m(z)$

Cylindrical Neumann function $N_m(z)$

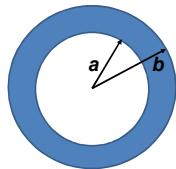


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Normal modes of an annular membrane -- continued



Differential equation for radial function:

$$\left(\frac{d^2}{dr^2} + \frac{1}{r} \frac{d}{dr} - \frac{m^2}{r^2} + k^2 \right) f(r) = 0$$

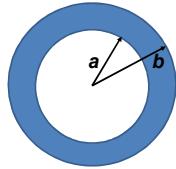
General form of radial function: $f(r) = AJ_m(kr) + BN_m(kr)$

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Normal modes of an annular membrane -- continued

Boundary conditions:
 $f(a) = 0$ $f(b) = 0$

$$AJ_m(ka) + BN_m(ka) = 0$$

$$AJ_m(kb) + BN_m(kb) = 0$$

\Rightarrow 2 equations and 2 unknowns -- k and $\frac{B}{A}$

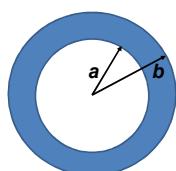
$$\frac{B}{A} = \frac{-J_m(ka)}{N_m(ka)} = \frac{-J_m(kb)}{N_m(kb)} \quad (\text{transcendental equation for } k)$$

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Normal modes of an annular membrane -- continued

Boundary conditions:
 $f(a) = 0$ $f(b) = 0$

$$\frac{B}{A} = \frac{-J_m(ka)}{N_m(ka)} = \frac{-J_m(kb)}{N_m(kb)} \quad \text{-- in terms of solution } k_{mn} :$$

$$f(r) = A \left(J_m(k_{mn}r) - \frac{J_m(k_{mn}a)}{N_m(k_{mn}a)} N_m(k_{mn}r) \right)$$

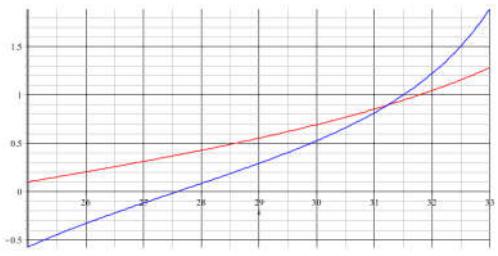
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Analysis for $m=0$ and $a=0.1$, $b=0.2$:

```
> plot( [ -BesselJ(0, 0.1·k) / BesselY(0, 0.1·k), -BesselJ(0, 0.2·k) / BesselY(0, 0.2·k)], k = 25 .. 33, color = [red, blue]);
```



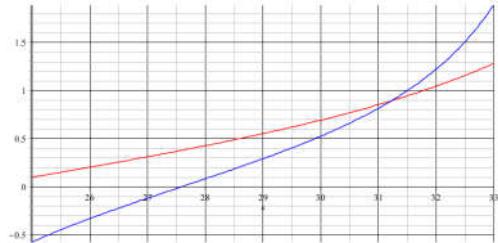
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```
> solve( -BesselJ(0, 0.1·k) / BesselY(0, 0.1·k) = -BesselJ(0, 0.2·k) / BesselY(0, 0.2·k), k, 30 .. 33);
```

31.23030920

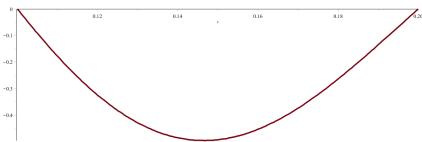


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$$f(r) = A \left(J_m(k_{mn}r) - \frac{J_m(k_{mn}a)}{N_m(k_{mn}a)} N_m(k_{mn}r) \right) \quad k_{01} = 31.23030920$$



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