

**PHY 711 Classical Mechanics and
Mathematical Methods**
9-9:50 AM MWF Olin 107

Plan for Lecture 29: Chap. 9 of F&W

Wave equation for sound in linear approximation

1. Wave equations for sound
2. Speed of sound
3. Standing wave solutions

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Events

Colloquium: Nov. 8, 2017 at 4 PM
WFU Physics Colloquium TITLE: "Materials and Measurements for Physical and Plastic Electronics" SPEAKER: Dean M. DeLongchamp, Leader, Polymers Processing Group, National Institute of Standards and Technology, Gaithersburg, MD 20899 USA

Colloquium: Nov. 15, 2017 at 4 PM
WFU Physics Colloquium TITLE: "What is fossilization? Casting new light (and electrical) on an old question" SPEAKER: Professor Elizabeth Spathan, Department of Engineering, Wake Forest University, Winston-Salem, NC TIME: Wed. ...

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18	Wed, 10/11/2017	Chap. 5	Rigid body motion		
	Fri, 10/13/2017		Fail break -- No class		
19	Mon, 10/16/2017		Discuss exam questions and topics for presentations	Topic	10/18/2017
20	Wed, 10/18/2017	Chap. 7	Wave equation in one dimension	#11	10/20/2017
21	Fri, 10/20/2017	Chap. 7	Solutions of Sturm-Liouville equations	#12	10/27/2017
22	Mon, 10/23/2017	Chap. 7	Solutions of Sturm-Liouville equations		
23	Wed, 10/25/2017	Chap. 7	Solutions of Sturm-Liouville equations		
24	Fri, 10/27/2017	App. A	Laplace transforms and contour integrals	#13	11/01/2017
25	Mon, 10/30/2017	App. A	Contour integrals		
26	Wed, 11/01/2017	Chap. 8	Mechanics of Elastic Membranes	#14	11/06/2017
27	Fri, 11/03/2017	Chap. 9	Introduction to hydrodynamics		
28	Mon, 11/06/2017	Chap. 9	Introduction to hydrodynamics	#15	11/10/2017
29	Wed, 11/08/2017	Chap. 9	Sound waves		
30	Fri, 11/10/2017				
31	Mon, 11/13/2017				
32	Wed, 11/15/2017				
33	Fri, 11/17/2017				
34	Mon, 11/20/2017				
	Wed, 11/22/2017		Thanksgiving Holiday -- No class		
	Fri, 11/24/2017		Thanksgiving Holiday -- No class		
35	Mon, 11/27/2017				

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Application of fluid equations to the case of air in equilibrium plus small perturbation

Newton - Euler equation of motion :

$$\frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} = \mathbf{f}_{\text{applied}} - \frac{\nabla p}{\rho}$$

Continuity equation : $\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0$

Near equilibrium :

$$\rho = \rho_0 + \delta \rho$$

$$p = p_0 + \delta p$$

$$\mathbf{v} = 0 + \delta \mathbf{v}$$

$$\mathbf{f}_{\text{applied}} = 0$$

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Equations to lowest order in perturbation :

$$\frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} = \mathbf{f}_{\text{applied}} - \frac{\nabla p}{\rho} \quad \Rightarrow \quad \frac{\partial \delta \mathbf{v}}{\partial t} = - \frac{\nabla \delta p}{\rho_0}$$

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0 \quad \Rightarrow \quad \frac{\partial \delta \rho}{\partial t} + \rho_0 \nabla \cdot \delta \mathbf{v} = 0$$

In terms of the velocity potential :

$$\delta \mathbf{v} = -\nabla \Phi$$

$$\frac{\partial \delta \mathbf{v}}{\partial t} = - \frac{\nabla \delta p}{\rho_0} \quad \Rightarrow \quad \nabla \left(- \frac{\partial \Phi}{\partial t} + \frac{\delta p}{\rho_0} \right) = 0$$

$$\frac{\partial \delta \rho}{\partial t} + \rho_0 \nabla \cdot \delta \mathbf{v} = 0 \quad \Rightarrow \quad \frac{\partial \delta \rho}{\partial t} - \rho_0 \nabla^2 \Phi = 0$$

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Expressing pressure in terms of the density :

$$p = p(s, \rho) = p_0 + \delta p \quad \text{where } s \text{ denotes the (constant) entropy}$$

$$p_0 = p(s, \rho_0)$$

$$\delta p = \left(\frac{\partial p}{\partial \rho} \right)_s \delta \rho \equiv c^2 \delta \rho$$

$$\nabla \left(- \frac{\partial \Phi}{\partial t} + \frac{\delta p}{\rho_0} \right) = 0 \quad \Rightarrow \quad - \frac{\partial \Phi}{\partial t} + c^2 \frac{\delta \rho}{\rho_0} = (\text{constant})$$

$$\Rightarrow - \frac{\partial^2 \Phi}{\partial t^2} + \frac{c^2}{\rho_0} \frac{\partial \delta \rho}{\partial t} = 0$$

$$\frac{\partial \delta \rho}{\partial t} - \rho_0 \nabla^2 \Phi = 0 \quad \Rightarrow \quad \frac{\partial^2 \Phi}{\partial t^2} - c^2 \nabla^2 \Phi = 0$$

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Wave equation for air :

$$\frac{\partial^2 \Phi}{\partial t^2} - c^2 \nabla^2 \Phi = 0$$

Here, $c^2 = \left(\frac{\partial p}{\partial \rho} \right)_s$

$\mathbf{v} = -\nabla \Phi$

Note that, we also have :

$$\frac{\partial^2 \delta p}{\partial t^2} - c^2 \nabla^2 \delta p = 0$$

$$\frac{\partial^2 \delta p}{\partial t^2} - c^2 \nabla^2 \delta p = 0$$

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Analysis of wave velocity in an ideal gas:

$$c^2 = \left(\frac{\partial p}{\partial \rho} \right)_s \leftarrow \text{adiabatic conditions}$$

Equation of state for ideal gas:

$$pV = NkT \quad N = \frac{M}{M_0}$$

$$p = \frac{M}{V} \frac{k}{M_0} T = \rho \frac{k}{M_0} T$$

$$k = 1.38 \times 10^{-23} \text{ J / K}$$

$M_0 =$ average mass of each molecule

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Internal energy for ideal gas :

$$E = \frac{f}{2} NkT = M\varepsilon \quad \varepsilon = \frac{f}{2} \frac{k}{M_0} T = \frac{f}{2} \frac{p}{\rho}$$

In terms of specific heat ratio : $\gamma \equiv \frac{C_p}{C_v}$

$$dE = dQ - dW$$

$$C_v = \left(\frac{dQ}{dT} \right)_v = \left(\frac{\partial E}{\partial T} \right)_v = \frac{f}{2} \frac{Mk}{M_0}$$

$$C_p = \left(\frac{dQ}{dT} \right)_p = \left(\frac{\partial E}{\partial T} \right)_p + p \left(\frac{\partial V}{\partial T} \right)_p = \frac{f}{2} \frac{Mk}{M_0} + \frac{Mk}{M_0}$$

$$\gamma = \frac{C_p}{C_v} = \frac{\frac{f}{2} + 1}{\frac{f}{2}} \Rightarrow \frac{f}{2} = \frac{1}{\gamma - 1}$$

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Internal energy for ideal gas :

$$E = \frac{1}{\gamma-1} NkT = M\varepsilon \quad \varepsilon = \frac{1}{\gamma-1} \frac{k}{M_0} T = \frac{1}{\gamma-1} \frac{p}{\rho}$$

Internal energy for ideal gas under isentropic conditions :

$$d\varepsilon = -\frac{p}{M} dV = \frac{p}{\rho^2} d\rho$$

$$\left(\frac{\partial \varepsilon}{\partial \rho}\right)_s = \frac{p}{\rho^2} = \frac{\partial}{\partial \rho} \left(\frac{1}{\gamma-1} \frac{p}{\rho}\right)_s = \left(\frac{\partial p}{\partial \rho}\right)_s \frac{1}{(\gamma-1)\rho} - \frac{p}{(\gamma-1)\rho^2}$$

$$\Rightarrow \left(\frac{\partial p}{\partial \rho}\right)_s = \frac{p\gamma}{\rho}$$

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Alternative derivation:

Isentropic or adiabatic equation of state :

$$\frac{dp}{p} = \gamma \frac{d\rho}{\rho} \quad \Rightarrow \frac{p}{p_0} = \left(\frac{\rho}{\rho_0}\right)^\gamma$$

$$\left(\frac{\partial p}{\partial \rho}\right)_s = \frac{p\gamma}{\rho}$$

Linearized speed of sound

$$c_0^2 = \left(\frac{\partial p}{\partial \rho}\right)_{s, p_0, \rho_0} = \frac{p_0 \gamma}{\rho_0}$$

$$c_0^2 \approx \frac{1.5 \cdot 1.0 \times 10^5 \text{ Pa}}{1.3 \text{ kg / m}^3} \quad c_0 \approx 340 \text{ m/s}$$

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Density dependence of speed of sound for ideal gas:

$$c^2 = \left(\frac{\partial p}{\partial \rho}\right)_s = \frac{p\gamma}{\rho}$$

$$\frac{p}{p_0} = \left(\frac{\rho}{\rho_0}\right)^\gamma$$

$$c^2 = \frac{p_0 \gamma}{\rho_0} \frac{p / p_0}{\rho / \rho_0} = c_0^2 \left(\frac{\rho}{\rho_0}\right)^{\gamma-1}$$

where $c_0^2 = \frac{p_0 \gamma}{\rho_0}$

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Wave equation for air :

$$\frac{\partial^2 \Phi}{\partial t^2} - c^2 \nabla^2 \Phi = 0$$

Note that, we also have :

$$\frac{\partial^2 \delta p}{\partial t^2} - c^2 \nabla^2 \delta p = 0$$

Here, $c^2 = \left(\frac{\partial p}{\partial \rho} \right)_s$

$$\frac{\partial^2 \delta p}{\partial t^2} - c^2 \nabla^2 \delta p = 0$$

$\mathbf{v} = -\nabla \Phi$

Boundary values:

Impenetrable surface with normal $\hat{\mathbf{n}}$ moving at velocity \mathbf{V} :

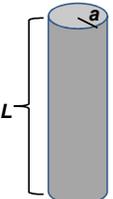
$$\hat{\mathbf{n}} \cdot \mathbf{V} = \hat{\mathbf{n}} \cdot \delta \mathbf{v} = -\hat{\mathbf{n}} \cdot \nabla \Phi$$

Free surface:

$$\delta p = 0 \quad \Rightarrow \quad \rho_0 \frac{\partial \Phi}{\partial t} = 0$$

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Time harmonic standing waves in a pipe



$\frac{\partial^2 \Phi}{\partial t^2} - c^2 \nabla^2 \Phi = 0$

Boundary values :

At fixed surface : $\hat{\mathbf{n}} \cdot \nabla \Phi = 0$

At free surface : $\frac{\partial \Phi}{\partial t} = 0$

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$$\nabla^2 \Phi - \frac{1}{c^2} \frac{\partial^2 \Phi}{\partial t^2} = 0 \quad \text{Define : } k \equiv \frac{\omega}{c}$$

In cylindrical coordinates :

$$\Phi(r, \varphi, z, t) = R(r)F(\varphi)Z(z)e^{-i\omega t} \equiv R(r)F(\varphi)Z(z)e^{-ikct}$$

$$\nabla^2 = \frac{1}{r} \frac{\partial}{\partial r} r \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \varphi^2} + \frac{\partial^2}{\partial z^2}$$

$$\left(\frac{1}{r} \frac{\partial}{\partial r} r \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \varphi^2} + \frac{\partial^2}{\partial z^2} + k^2 \right) \Phi(r, \varphi, z, t) = 0$$

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$$\left(\frac{1}{r} \frac{\partial}{\partial r} r \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \varphi^2} + \frac{\partial^2}{\partial z^2} + k^2 \right) \Phi(r, \varphi, z, t) = 0$$

$\Phi(r, \varphi, z, t) = R(r)F(\varphi)Z(z)e^{-i\alpha t}$

$F(\varphi) = e^{im\varphi}; F(\varphi) = F(\varphi + 2\pi N) \Rightarrow m = \text{integer}$

$Z(z) = e^{i\alpha z}; \alpha = \text{real plus other restrictions}$

$$\left(\frac{d^2}{dr^2} + \frac{1}{r} \frac{d}{dr} - \frac{m^2}{r^2} - \alpha^2 + k^2 \right) R(r) = 0$$

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$$\left(\frac{d^2}{dr^2} + \frac{1}{r} \frac{d}{dr} - \frac{m^2}{r^2} - \alpha^2 + k^2 \right) R(r) = 0$$

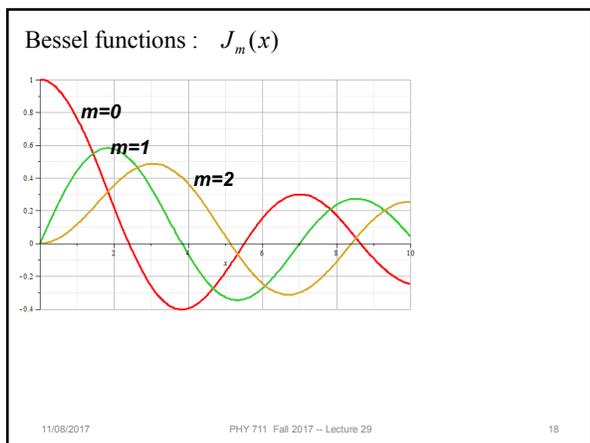
For $k^2 \geq \alpha^2$ define $\kappa^2 \equiv k^2 - \alpha^2$

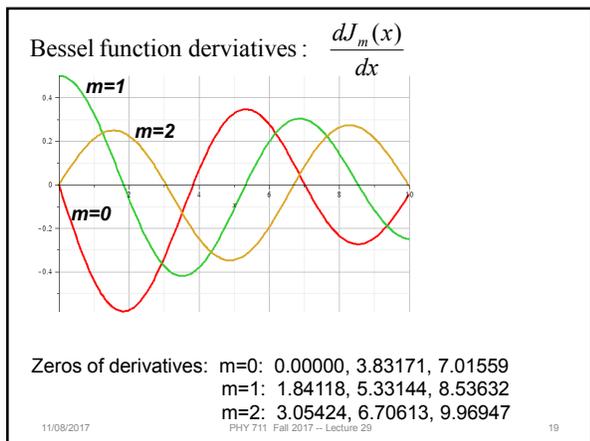
$$\left(\frac{d^2}{dr^2} + \frac{1}{r} \frac{d}{dr} - \frac{m^2}{r^2} + \kappa^2 \right) R(r) = 0$$

Cylinder surface boundary conditions: $\left. \frac{dR}{dr} \right|_{r=a} = 0$

$\Rightarrow R(r) = J_m(\kappa r)$ where for $\frac{dJ_m(x'_{mn})}{dx} = 0, \kappa_{mn} = \frac{x'_{mn}}{a}$

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Boundary condition for $z=0, z=L$:
 For open - open pipe:
 $Z(0) = Z(L) = 0 \Rightarrow Z(z) = \sin\left(\frac{p\pi z}{L}\right)$
 $\Rightarrow \alpha_p = \frac{p\pi}{L}, \quad p = 1, 2, 3, \dots$
 Resonant frequencies:
 $\frac{\omega^2}{c^2} = k^2 = \kappa_{mn}^2 + \alpha_p^2$
 $k_{mnp}^2 = \left(\frac{x'_{mn}}{a}\right)^2 + \left(\frac{\pi p}{L}\right)^2$

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Example

$$k_{mnp}^2 = \left(\frac{x'_{mn}}{a}\right)^2 + \left(\frac{\pi p}{L}\right)^2 = \left(\frac{\pi p}{L}\right)^2 \left(1 + \left(\frac{L}{a}\right)^2 \left(\frac{x'_{mn}}{\pi p}\right)^2\right)$$

$\pi p = 3.14, 6.28, 9.42, \dots$
 $x'_{mn} = 0.00, 1.84, 3.05$

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Alternate boundary condition for $z=0, z=L$:

For open - closed pipe :

$$\frac{dZ(0)}{dz} = Z(L) = 0 \quad \Rightarrow \quad Z(z) = \cos\left(\frac{(2p+1)\pi z}{2L}\right)$$

$$\Rightarrow \alpha_p = \frac{(2p+1)\pi}{2L}, \quad p = 0, 1, 2, 3, \dots$$

$$k_{mnp}^2 = \left(\frac{x'_{mn}}{a}\right)^2 + \left(\frac{\pi(2p+1)}{2L}\right)^2$$

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Other solutions to wave equation:

$$\nabla^2 \Phi - \frac{1}{c^2} \frac{\partial^2 \Phi}{\partial t^2} = 0$$

Plane wave solution :

$$\Phi(\mathbf{r}, t) = A e^{i\mathbf{k}\cdot\mathbf{r} - i\omega t} \quad \text{where} \quad k^2 = \left(\frac{\omega}{c}\right)^2$$

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