

**PHY 711 Classical Mechanics and
Mathematical Methods
9-9:50 AM MWF Olin 107**

**Plan for Lecture 34:
Chapter 11 in F & W:
Heat conduction**

- 1. Basic equations**
- 2. Boundary value problems**

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23	Wed, 10/25/2017 Chap. 7	Solutions of Sturm-Liouville equations		
24	Fri, 10/27/2017 App. A	Laplace transforms and contour integrals	#13	11/01/2017
25	Mon, 10/30/2017 App. A	Contour integrals		
26	Wed, 11/01/2017 Chap. 8	Mechanics of Elastic Membranes	#14	11/08/2017
27	Fri, 11/03/2017 Chap. 9	Introduction to hydrodynamics		
28	Mon, 11/06/2017 Chap. 9	Introduction to hydrodynamics	#15	11/10/2017
29	Wed, 11/08/2017 Chap. 9	Sound waves		
30	Fri, 11/10/2017 Chap. 9	Sound waves	#16	11/17/2017
	Mon, 11/13/2017	Class cancelled		
31	Wed, 11/15/2017 Chap. 9	Sound waves – including non-linearities		
32	Fri, 11/17/2017 Chap. 10	Surface waves in fluids	#17	11/27/2017
33	Mon, 11/20/2017 Chap. 10	Surface waves in fluids		
	(Wed, 11/22/2017)	Thanksgiving Holiday -- No class		
	Fri, 11/24/2017	Thanksgiving Holiday -- No class		
34	Mon, 11/27/2017 Chap. 11	Heat conductivity		
35	Wed, 11/29/2017 Chap. 12	Viscous fluids		
36	Fri, 12/01/2017 Chap. 12	Viscous fluids		
	Mon, 12/04/2017	Presentations I		
	Wed, 12/06/2017	Presentations II		
	Fri, 12/08/2017	Presentations III		

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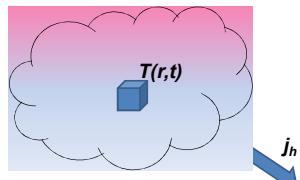
Schedule for PHY 711 presentations Monday Dec. 4, 2017		
Time	Presenter	Title
9:00-9:15 AM	Ali Daraei	Rheology of fibrin fibers
9:17-9:32 AM		
9:34-9:49 AM	Yan Li	Computational study of normal modes for one-dimensional mass-spring oscillating systems
Wednesday Dec. 6, 2017		
Time	Presenter	Title
9:00-9:15 AM	Matthew Waltrip	
9:17-9:32 AM	Nouf Aharbi	Vibrational Modes of Carbon Dioxide
9:34-9:49 AM	Kevin Roebuck	Equations of motion of the Foucault pendulum
Friday Dec. 8, 2017		
Time	Presenter	Title
9:00-9:15 AM	Haardik Pandey	Normal modes of vibrations for water molecule
9:17-9:32 AM	Ellie Alipour	
9:34-9:49 AM	TJ Colvin	Quasiharmonic analysis of biomolecules

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Conduction of heat

Enthalpy of a system at constant pressure p

non uniform temperature $T(\mathbf{r}, t)$
mass density ρ and heat capacity c_p

$$H = \int_V \rho c_p (T(\mathbf{r}, t) - T_0) d^3 r + H_0(T_0, p)$$

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Note that in this treatment we are considering a system at constant pressure p

Notation:	Heat added to system	-- $dQ = TdS$
	External work done on system	-- $dW = -pdV$
	Internal energy	-- $dE = dQ + dW = TdS - pdV$
	Entropy	-- dS
	Enthalpy	-- $dH = d(E + pV) = TdS + Vdp$
	Heat capacity at constant pressure:	

$$C_p \equiv \left(\frac{\partial Q}{\partial T} \right)_p = \left(\frac{\partial H}{\partial T} \right)_p = T \left(\frac{\partial S}{\partial T} \right)_p$$

$$C_p = \int \rho c_p d^3 r$$

More generally, note that c_p can depend on T ; we are assume that dependence to be trivial.

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Conduction of heat -- continued

$$H = \int_V \rho c_p (T(\mathbf{r}, t) - T_0) d^3 r + H_0(T_0, p)$$

Time rate of change of enthalpy:

$$\frac{dH}{dt} = \int_V \rho c_p \frac{\partial T(\mathbf{r}, t)}{\partial t} d^3 r = - \int_A \mathbf{j}_h \cdot d\mathbf{A} + \int_V \rho \dot{q} d^3 r$$

heat flux heat source

$$\rho c_p \frac{\partial T(\mathbf{r}, t)}{\partial t} = -\nabla \cdot \mathbf{j}_h + \rho \dot{q}$$

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Conduction of heat -- continued

$$\rho c_p \frac{\partial T(\mathbf{r}, t)}{\partial t} = -\nabla \cdot \mathbf{j}_h + \rho \dot{q}$$

Empirically: $\mathbf{j}_h = -k_{th} \nabla T(\mathbf{r}, t)$

$$\Rightarrow \frac{\partial T(\mathbf{r},t)}{\partial t} = \kappa \nabla^2 T(\mathbf{r},t) + \frac{\dot{q}}{c_p}$$

$$\kappa \equiv \frac{k_{th}}{\rho c_p} \quad \text{thermal diffusivity}$$

https://www.engineersedge.com/heat_transfer/thermal_diffusivity_table_13953.htm

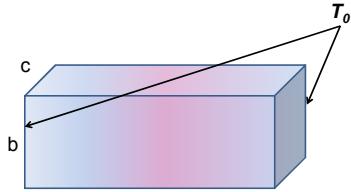
Typical values (m²/s)

Air	2×10^{-5}
Water	1×10^{-7}
Copper	1×10^{-4}

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Boundary value problems for heat conduction



$$\frac{\partial T(\mathbf{r},t)}{\partial t} - \kappa \nabla^2 T(\mathbf{r},t) = \frac{\mathbf{a}}{c_p}$$

$$\text{Without source term: } \frac{\partial T(\mathbf{r}, t)}{\partial t} - \kappa \nabla^2 T(\mathbf{r}, t) = 0$$

Example with boundary values: $T(0, y, z, t) = T(a, y, z, t) = T_0$
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Journal of Oral Rehabilitation 2013; 40(12): 938-945

Boundary value problems for heat conduction

$$\frac{\partial T(\mathbf{r},t)}{\partial t} - \kappa \nabla^2 T(\mathbf{r},t) = 0$$

$$T(0, y, z, t) = T(a, y, z, t) = T_0$$

$$\frac{\partial T(x, 0, z, t)}{\partial v} = \frac{\partial T(x, b, z, t)}{\partial v} = 0 \quad \boxed{a}$$

$$\left. \frac{\partial T(x,y,0,t)}{\partial y} = \frac{\partial T(x,y,c,t)}{\partial y} = 0 \right\} \text{ Assuming thermally insulated boundaries}$$

$$\text{Separation of variables: } T(x, y, z, t) = T_0 + X(x)Y(y)Z(z)e^{-\lambda t}$$

$$\text{Let } \frac{d^2X}{t^2} = -\alpha^2 X \quad \frac{d^2Y}{t^2} = -\beta^2 Y \quad \frac{d^2Z}{t^2} = -\gamma^2 Z$$

$$\Rightarrow -\lambda + \kappa(\alpha^2 + \beta^2 + \gamma^2) \equiv 0$$

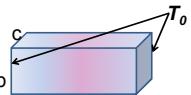
$$\rightarrow \pi(\alpha + \beta + \gamma) = 0.$$

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Boundary value problems for heat conduction



$$T(x,y,z,t) = T_0 + X(x)Y(y)Z(z)e^{-\lambda t} \quad \text{b) a}$$

$$X(0) = X(a) = 0 \quad \Rightarrow X(x) = \sin\left(\frac{m\pi x}{a}\right)$$

$$\frac{dY(0)}{dy} = \frac{dY(b)}{dy} = 0 \quad \Rightarrow Y(y) = \cos\left(\frac{n\pi y}{b}\right)$$

$$\frac{dZ(0)}{dz} = \frac{dZ(c)}{dz} = 0 \quad \Rightarrow Z(z) = \cos\left(\frac{p\pi z}{c}\right)$$

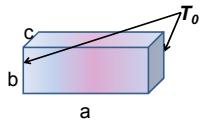
$$-\lambda_{nmp} + \kappa \left(\left(\frac{m\pi}{a} \right)^2 + \left(\frac{n\pi}{b} \right)^2 + \left(\frac{p\pi}{c} \right)^2 \right) = 0$$

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Boundary value problems for heat conduction



Full solution:

$$T(x, y, z, t) = T_0 + \sum_{nmp} C_{nmp} \sin\left(\frac{m\pi x}{a}\right) \cos\left(\frac{n\pi y}{b}\right) \cos\left(\frac{p\pi z}{c}\right) e^{-\lambda_{nmp} t}$$

$$\lambda_{nmp} = \kappa \left(\left(\frac{n\pi}{a} \right)^2 + \left(\frac{m\pi}{b} \right)^2 + \left(\frac{p\pi}{c} \right)^2 \right)$$

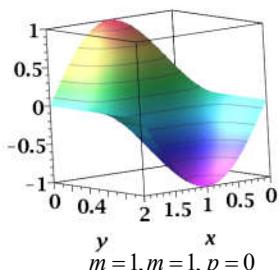
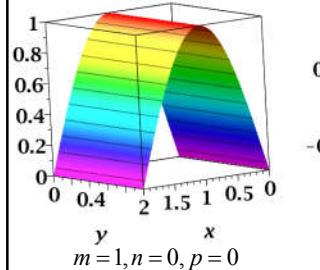
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Full solution:

$$T(x, y, z, t) = T_0 + \sum_{nmp} C_{nmp} \sin\left(\frac{m\pi x}{a}\right) \cos\left(\frac{n\pi y}{b}\right) \cos\left(\frac{p\pi z}{c}\right) e^{-\lambda_{nmp} t}$$



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Oscillatory thermal behavior

$$T(z=0,t) = \Re(T_0 e^{-i\omega t})$$

$\frac{\partial T}{\partial t} = \kappa \frac{\partial^2 T}{\partial z^2}$

Assume: $T(z,t) = \Re(f(z)e^{-i\omega t})$

$$(-i\omega)f = \kappa \frac{d^2 f}{dz^2}$$

Let $f(z) = Ae^{\alpha z}$

$$\alpha^2 = -\frac{i\omega}{\kappa} = e^{3i\pi/2} \frac{\omega}{\kappa}$$

$$\alpha = \pm(1-i)\sqrt{\frac{\omega}{2\kappa}}$$

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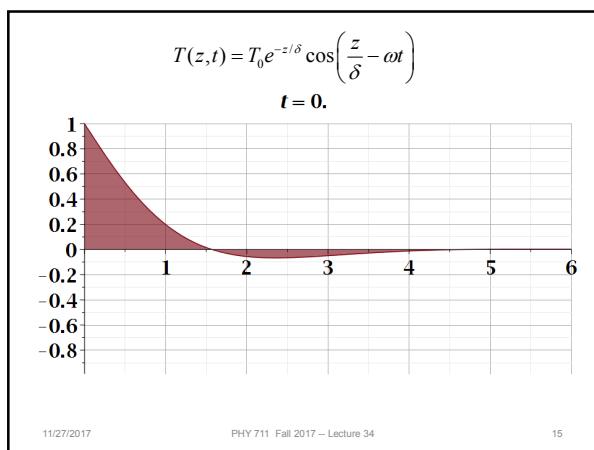
Oscillatory thermal behavior -- continued

$$T(z=0,t) = \Re(T_0 e^{-i\omega t})$$

$$T(z,t) = \Re(A e^{\pm(1-i)z/\delta} e^{-i\omega t})$$

where $\delta \equiv \sqrt{\frac{2\kappa}{\omega}}$

Physical solution: $T(z,t) = T_0 e^{-z/\delta} \cos\left(\frac{z}{\delta} - \omega t\right)$



Initial value problem in an infinite domain; Fourier transform

$$\frac{\partial T(\mathbf{r},t)}{\partial t} - \kappa \nabla^2 T(\mathbf{r},t) = 0$$

$$T(\mathbf{r},0) = f(\mathbf{r})$$

$$\text{Let : } \tilde{T}(\mathbf{q}, t) = \int d^3r e^{-i\mathbf{q}\cdot\mathbf{r}} T(\mathbf{r}, t)$$

$$\tilde{f}(\mathbf{q}) = \int d^3r e^{-i\mathbf{q}\cdot\mathbf{r}} f(\mathbf{r})$$

$$\Rightarrow \tilde{T}(\mathbf{q},0) = \tilde{f}(\mathbf{q})$$

$$\Rightarrow \frac{\partial \tilde{T}(\mathbf{q},t)}{\partial t} = -\kappa q^2 \tilde{T}(\mathbf{q},t)$$

$$\tilde{T}(\mathbf{q}, t) = \tilde{T}(\mathbf{q}, 0)e^{-\kappa q^2 t}$$

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Initial value problem in an infinite domain; Fourier transform

$$\tilde{T}(\mathbf{q}, t) = \int d^3r e^{-i\mathbf{q}\cdot\mathbf{r}} T(\mathbf{r}, t) \quad \Rightarrow T(\mathbf{r}, t) = \frac{1}{(2\pi)^3} \int d^3q e^{i\mathbf{q}\cdot\mathbf{r}} \tilde{T}(\mathbf{q}, t)$$

$$\tilde{T}(\mathbf{q}, t) = \tilde{T}(\mathbf{q}, 0) e^{-\kappa q^2 t}$$

$$T(\mathbf{r}, t) = \frac{1}{(2\pi)^3} \int d^3 q e^{i\mathbf{q} \cdot \mathbf{r}} \tilde{T}(\mathbf{q}, 0) e^{-\kappa q^2 t}$$

$$\tilde{T}(\mathbf{q},0) = \tilde{f}(\mathbf{q}) = \int d^3r e^{-i\mathbf{q}\cdot\mathbf{r}} f(\mathbf{r})$$

$$T(\mathbf{r}, t) = \int d^3 r' G(\mathbf{r} - \mathbf{r}', t) T(\mathbf{r}', 0)$$

$$\text{with } G(\mathbf{r} - \mathbf{r}', t) \equiv \frac{1}{(2\pi)^3} \int d^3 q e^{i\mathbf{q} \cdot (\mathbf{r} - \mathbf{r}')} e^{-\kappa q^2 t}$$

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Initial value problem in an infinite domain; Fourier transform

$$T(\mathbf{r}, t) = \int d^3 r' G(\mathbf{r} - \mathbf{r}', t) T(\mathbf{r}', 0)$$

$$\text{with } G(\mathbf{r} - \mathbf{r}', t) \equiv \frac{1}{(2\pi)^3} \int d^3 q e^{i\mathbf{q} \cdot (\mathbf{r} - \mathbf{r}')} e^{-\kappa q^2 t}$$

$$G(\mathbf{r} - \mathbf{r}', t) = \frac{1}{(4\pi k t)^{3/2}} e^{-|\mathbf{r} - \mathbf{r}'|^2 / (4k t)}$$

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Heat equation in half-space

$$\frac{\partial T(\mathbf{r}, t)}{\partial t} - \kappa \nabla^2 T(\mathbf{r}, t) = 0$$

$T(\mathbf{r},t) \Rightarrow T(z,t)$ with initial and boundary values:

$$T(z,t) \equiv 0 \quad \text{for } z < 0$$

$$T(z,0) = 0 \quad \text{for } z > 0$$

$$T(0,t) = T_0 \quad \text{for } t \geq 0$$

$$\text{Solution : } T = T_0 \operatorname{erfc} \left(\frac{z}{2\sqrt{kt}} \right)$$

$$\text{where } \operatorname{erfc}(x) \equiv \frac{2}{\sqrt{\pi}} \int_x^{\infty} e^{-u^2} du$$

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Heat equation in half-space -- continued

$$\frac{\partial T(z,t)}{\partial t} - \kappa \frac{\partial^2 T(z,t)}{\partial z^2} = 0$$

$$\text{Solution : } T = T_0 \operatorname{erfc} \left(\frac{z}{2\sqrt{kt}} \right)$$

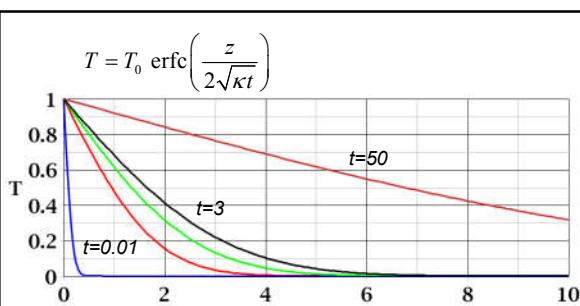
$$\text{where } \operatorname{erfc}(x) \equiv \frac{2}{\sqrt{\pi}} \int_x^{\infty} e^{-u^2} du$$

$$\text{Note that } \frac{d}{dx} \operatorname{erfc}(x) = \frac{d}{dx} \frac{2}{\sqrt{\pi}} \int_x^{\infty} e^{-u^2} du = -\frac{2}{\sqrt{\pi}} e^{-x^2}$$

$$\frac{\partial}{\partial t} \operatorname{erfc}\left(\frac{z}{2\sqrt{kt}}\right) = \frac{2}{\sqrt{\pi}} e^{-(z^2/(4kt))} \left(\frac{z}{4\sqrt{kt^3}} \right)$$

$$\frac{\partial^2}{\partial z^2} \operatorname{erfc}\left(\frac{z}{2\sqrt{\kappa t}}\right) = \frac{2}{\sqrt{\pi}} e^{-(z^2/(4\kappa t))} \left(\frac{z}{4\kappa\sqrt{\kappa t^3}} \right)$$

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