

**PHY 711 Classical Mechanics and
Mathematical Methods**
9-9:50 AM MWF Olin 107

Plan for Lecture 3:
Text reference: Chapter 1

1. Scattering experiments
2. Center of mass and laboratory reference frames
3. Calculation of scattering cross sections for given interaction models

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PHY 711 Classical Mechanics and Mathematical Methods

MWF 9 AM-9:50 AM OPL 107 <http://www.wfu.edu/~natalie/f17phy711/>

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Course schedule

(Preliminary schedule – subject to frequent adjustment.)

Date	F&W Reading	Topic	Assignment	Due
1 Mon, 8/28/2017	Chap. 1	Introduction	#1	9/6/2017
2 Wed, 8/30/2017	Chap. 1	Scattering theory	#2	9/6/2017
3 Fri, 9/01/2017	Chap. 1	Scattering theory		
4 Mon, 9/04/2017				
5 Wed, 9/06/2017				
6 Fri, 9/08/2017				
7 Mon, 9/11/2017				

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Scattering theory:

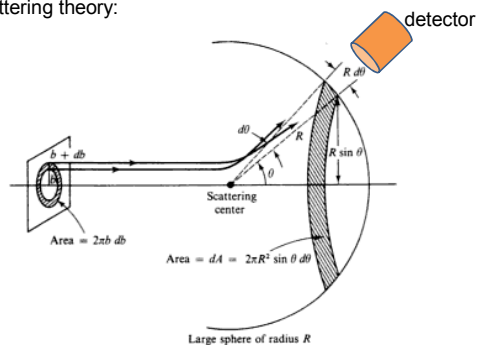


Figure 5.5 The scattering problem and relation of cross section to impact parameter.

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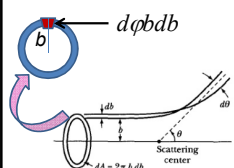
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Differential cross section

$$\left(\frac{d\sigma}{d\Omega}\right) = \frac{\text{Number of detected particles at } \theta \text{ per target particle}}{\text{Number of incident particles per unit area}}$$

= Area of incident beam that is scattered into detector at angle θ



$$\left(\frac{d\sigma}{d\Omega}\right) = \frac{d\phi b db}{d\phi \sin\theta d\theta} = \frac{b}{\sin\theta} \left|\frac{db}{d\theta}\right|$$

Figure from Marion & Thornton, Classical Dynamics

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Note: Notion of cross section is common to many areas of physics including classical mechanics, quantum mechanics, optics, etc. Only in the classical mechanics can we calculate it using geometric considerations

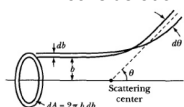


Figure from Marion & Thornton, Classical Dynamics

$$\left(\frac{d\sigma}{d\Omega}\right) = \frac{d\phi b db}{d\phi \sin\theta d\theta} = \frac{b}{\sin\theta} \left|\frac{db}{d\theta}\right|$$

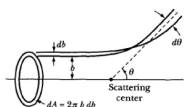
Note: We are assuming that the process is isotropic in ϕ

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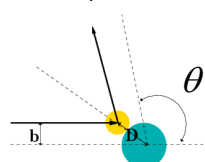
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Simple example – collision of hard spheres



$$\left(\frac{d\sigma}{d\Omega}\right) = \frac{b}{\sin\theta} \left|\frac{db}{d\theta}\right|$$

Microscopic view:



$$b(\theta) = ?$$

$$b(\theta) = D \sin\left(\frac{\pi}{2} - \frac{\theta}{2}\right)$$

$$\left(\frac{d\sigma}{d\Omega}\right) = \frac{D^2}{4}$$

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Simple example – collision of hard spheres -- continued

Total scattering cross section:

$$\sigma = \int \left(\frac{d\sigma}{d\Omega} \right) d\Omega$$

Hard sphere:

$$\left(\frac{d\sigma}{d\Omega} \right) = \frac{D^2}{4}$$

$$\sigma = \pi D^2$$

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Relationship of scattering cross-section to particle interactions --
Classical mechanics of a conservative 2-particle system.

1 $\frac{d\mathbf{p}_1}{dt} = \mathbf{F}_{12}$ $\frac{d\mathbf{p}_2}{dt} = \mathbf{F}_{21}$ 2

$\mathbf{F}_{12} = -\nabla_1 V(\mathbf{r}_1 - \mathbf{r}_2) \Rightarrow E = \frac{1}{2} m_1 v_1^2 + \frac{1}{2} m_2 v_2^2 + V(\mathbf{r}_1 - \mathbf{r}_2)$

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Typical two-particle interactions –

Central potential: $V(\mathbf{r}_1 - \mathbf{r}_2) = V(|\mathbf{r}_1 - \mathbf{r}_2|) \equiv V(r)$

Hard sphere: $V(r) = \begin{cases} \infty & r \leq a \\ 0 & r > a \end{cases}$

Coulomb or gravitational: $V(r) = \frac{K}{r}$

Lennard-Jones: $V(r) = \frac{A}{r^{12}} - \frac{B}{r^6}$

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Relationship between center of mass and laboratory frames of reference

Definition of center of mass $\mathbf{R}_{CM}$

$$m_1 \mathbf{r}_1 + m_2 \mathbf{r}_2 = (m_1 + m_2) \mathbf{R}_{CM}$$

$$m_1 \dot{\mathbf{r}}_1 + m_2 \dot{\mathbf{r}}_2 = (m_1 + m_2) \dot{\mathbf{R}}_{CM} = (m_1 + m_2) \mathbf{V}_{CM}$$

$$E = \frac{1}{2} m_1 v_1^2 + \frac{1}{2} m_2 v_2^2 + V(\mathbf{r}_1 - \mathbf{r}_2)$$

$$= \frac{1}{2} (m_1 + m_2) V_{CM}^2 + \frac{1}{2} \mu |\mathbf{v}_1 - \mathbf{v}_2|^2 + V(\mathbf{r}_1 - \mathbf{r}_2)$$

where: $\mu \equiv \frac{m_1 m_2}{m_1 + m_2}$

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Classical mechanics of a conservative 2-particle system -- continued

$$E = \frac{1}{2} (m_1 + m_2) V_{CM}^2 + \frac{1}{2} \mu |\mathbf{v}_1 - \mathbf{v}_2|^2 + V(\mathbf{r}_1 - \mathbf{r}_2)$$

For central potentials: $V(\mathbf{r}_1 - \mathbf{r}_2) = V(|\mathbf{r}_1 - \mathbf{r}_2|) \equiv V(r_{12})$

Relative angular momentum is also conserved:

$$E = \frac{1}{2} (m_1 + m_2) V_{CM}^2 + \frac{1}{2} \mu v_{12}^2 + \frac{L_{12}^2}{2\mu r_{12}^2} + V(r_{12})$$

Simpler notation:

$$E = \underbrace{\frac{1}{2} (m_1 + m_2) V_{CM}^2}_{E_{CM}(\text{constant})} + \underbrace{\frac{1}{2} \mu \dot{r}^2 + \frac{\ell^2}{2\mu r^2}}_{E_{rel}} + V(r)$$

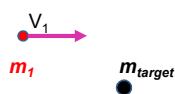
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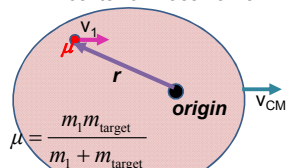
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Separating the collective (center of mass) and relative motion:

In laboratory frame:



In center-of-mass frame:



We will assume that the interaction between particle and target $V(r)$ conserves energy and angular momentum.

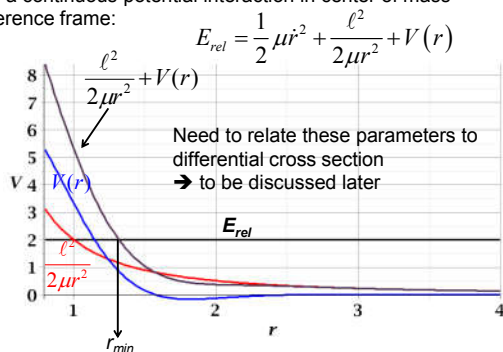
$$\ell = |\mathbf{r} \times \mu \mathbf{v}_1|$$

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For a continuous potential interaction in center of mass reference frame:



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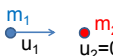
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It is often convenient to analyze the scattering cross section in the center of mass reference frame.

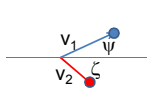
Relationship between normal laboratory reference and center of mass:

Laboratory reference frame:

Before

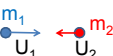


After

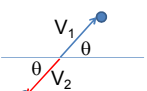


Center of mass reference frame:

Before



After



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Relationship between center of mass and laboratory frames of reference -- continued

Since m_2 is initially at rest :

$$\mathbf{V}_{CM} = \frac{m_1}{m_1 + m_2} \mathbf{u}_1 \quad \mathbf{u}_1 = \mathbf{U}_1 + \mathbf{V}_{CM} \Rightarrow \mathbf{U}_1 = \frac{m_2}{m_1 + m_2} \mathbf{u}_1 = \frac{m_2}{m_1} \mathbf{V}_{CM}$$

$$\mathbf{u}_2 = \mathbf{U}_2 + \mathbf{V}_{CM} \Rightarrow \mathbf{U}_2 = -\frac{m_1}{m_1 + m_2} \mathbf{u}_1 = -\mathbf{V}_{CM}$$

$$\mathbf{v}_1 = \mathbf{V}_1 + \mathbf{V}_{CM}$$

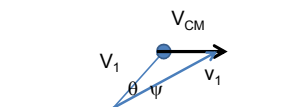
$$\mathbf{v}_2 = \mathbf{V}_2 + \mathbf{V}_{CM}$$

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Relationship between center of mass and laboratory frames of reference



$$\mathbf{v}_1 = \mathbf{V}_1 + \mathbf{V}_{CM}$$

$$v_1 \sin \psi = V_1 \sin \theta$$

$$v_1 \cos \psi = V_1 \cos \theta + V_{CM}$$

$$\tan \psi = \frac{\sin \theta}{\cos \theta + V_{CM} / V_1} = \frac{\sin \theta}{\cos \theta + m_1 / m_2} \quad \text{For elastic scattering}$$

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Digression – elastic scattering

$$\begin{aligned} \frac{1}{2} m_1 U_1^2 + \frac{1}{2} m_2 U_2^2 + \frac{1}{2} (m_1 + m_2) V_{CM}^2 \\ = \frac{1}{2} m_1 V_1^2 + \frac{1}{2} m_2 V_2^2 + \frac{1}{2} (m_1 + m_2) V_{CM}^2 \end{aligned}$$

Also note:

$$m_1 \mathbf{U}_1 + m_2 \mathbf{U}_2 = 0 \quad m_1 \mathbf{V}_1 + m_2 \mathbf{V}_2 = 0$$

$$\mathbf{U}_1 = \frac{m_2}{m_1} \mathbf{V}_{CM} \quad \mathbf{U}_2 = -\mathbf{V}_{CM}$$

$$\Rightarrow |\mathbf{U}_1| = |\mathbf{V}_1| \quad \text{and} \quad |\mathbf{U}_2| = |\mathbf{V}_2| = |\mathbf{V}_{CM}|$$

$$\text{Also note that: } m_1 |\mathbf{U}_1| = m_2 |\mathbf{U}_2|$$

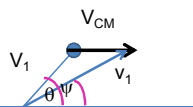
$$\text{So that: } V_{CM} / V_1 = V_{CM} / U_1 = m_1 / m_2$$

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Relationship between center of mass and laboratory frames of reference – continued (elastic scattering)



$$\mathbf{v}_1 = \mathbf{V}_1 + \mathbf{V}_{CM}$$

$$v_1 \sin \psi = V_1 \sin \theta$$

$$v_1 \cos \psi = V_1 \cos \theta + V_{CM}$$

$$\tan \psi = \frac{\sin \theta}{\cos \theta + V_{CM} / V_1} = \frac{\sin \theta}{\cos \theta + m_1 / m_2}$$

$$\text{Also: } \cos \psi = \frac{\cos \theta + m_1 / m_2}{\sqrt{1 + 2m_1 / m_2 \cos \theta + (m_1 / m_2)^2}}$$

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Differential cross sections in different reference frames

$$\left(\frac{d\sigma_{LAB}(\psi)}{d\Omega_{LAB}} \right) = \left(\frac{d\sigma_{CM}(\theta)}{d\Omega_{CM}} \right) \frac{d\Omega_{CM}}{d\Omega_{LAB}}$$

$$\frac{d\Omega_{CM}}{d\Omega_{LAB}} = \left| \frac{\sin \theta}{\sin \psi} \frac{d\theta}{d\psi} \right| = \left| \frac{d \cos \theta}{d \cos \psi} \right|$$

Using :

$$\cos \psi = \frac{\cos \theta + m_1 / m_2}{\sqrt{1 + 2(m_1 / m_2) \cos \theta + (m_1 / m_2)^2}}$$

$$\left| \frac{d \cos \psi}{d \cos \theta} \right| = \frac{(m_1 / m_2) \cos \theta + 1}{(1 + 2(m_1 / m_2) \cos \theta + (m_1 / m_2)^2)^{3/2}}$$

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Differential cross sections in different reference frames – continued:

$$\left(\frac{d\sigma_{LAB}(\psi)}{d\Omega_{LAB}} \right) = \left(\frac{d\sigma_{CM}(\theta)}{d\Omega_{CM}} \right) \left| \frac{d \cos \theta}{d \cos \psi} \right|$$

$$\left(\frac{d\sigma_{LAB}(\psi)}{d\Omega_{LAB}} \right) = \left(\frac{d\sigma_{CM}(\theta)}{d\Omega_{CM}} \right) \frac{(1 + 2m_1 / m_2 \cos \theta + (m_1 / m_2)^2)^{3/2}}{(m_1 / m_2) \cos \theta + 1}$$

For elastic scattering

$$\text{where : } \tan \psi = \frac{\sin \theta}{\cos \theta + m_1 / m_2}$$

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$$\left(\frac{d\sigma_{LAB}(\psi)}{d\Omega_{LAB}} \right) = \left(\frac{d\sigma_{CM}(\theta)}{d\Omega_{CM}} \right) \frac{(1 + 2m_1 / m_2 \cos \theta + (m_1 / m_2)^2)^{3/2}}{(m_1 / m_2) \cos \theta + 1}$$

$$\text{where : } \tan \psi = \frac{\sin \theta}{\cos \theta + m_1 / m_2}$$

Example: suppose $m_1 = m_2$

$$\text{In this case : } \tan \psi = \frac{\sin \theta}{\cos \theta + 1} \Rightarrow \psi = \frac{\theta}{2}$$

$$\text{note that } 0 \leq \psi \leq \frac{\pi}{2}$$

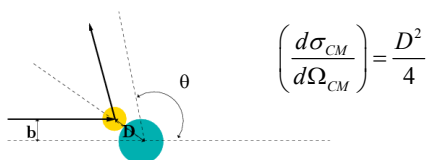
$$\left(\frac{d\sigma_{LAB}(\psi)}{d\Omega_{LAB}} \right) = \left(\frac{d\sigma_{CM}(2\psi)}{d\Omega_{CM}} \right) \cdot 4 \cos \psi$$

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Example of hard spheres

Cross section in lab frame when $m_1 = m_2$

$$\tan \psi = \frac{\sin \theta}{\cos \theta + 1} \Rightarrow \psi = \frac{\theta}{2} \Rightarrow \text{note that } 0 \leq \psi \leq \frac{\pi}{2}$$

$$\left(\frac{d\sigma_{LAB}(\psi)}{d\Omega_{LAB}} \right) = \left(\frac{d\sigma_{CM}(2\psi)}{d\Omega_{CM}} \right) \cdot 4 \cos \psi = D^2 \cos \psi$$

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Hard sphere example – continued

$$m_1 = m_2$$

Center of mass frame

Lab frame

$$\left(\frac{d\sigma_{CM}(\theta)}{d\Omega_{CM}} \right) = \frac{D^2}{4} \quad \left(\frac{d\sigma_{LAB}(\psi)}{d\Omega_{LAB}} \right) = D^2 \cos \psi \quad \psi = \frac{\theta}{2}$$

$$\int \frac{d\sigma_{CM}(\theta)}{d\Omega_{CM}} d\Omega_{CM} = \int \frac{d\sigma_{LAB}(\psi)}{d\Omega_{LAB}} d\Omega_{LAB} =$$

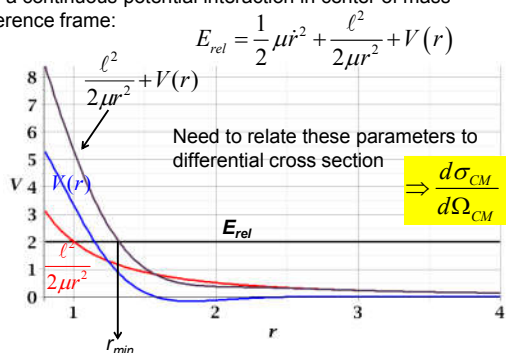
$$\frac{D^2}{4} 4\pi = \pi D^2 \quad 2\pi D^2 \int_0^{\pi/2} \cos \psi \sin \psi d\psi = \pi D^2$$

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For a continuous potential interaction in center of mass reference frame:



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