

PHY 741 Quantum Mechanics
12-12:50 AM MWF Olin 103

Plan for Lecture 25:

Chap. 17 in Shankar:
Time-Independent Perturbation Theory

- 1. Case of non-degenerate spectrum**
- 2. Treatment of degenerate spectrum**

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17 Mon, 10/09/2017	Chap. 13	Quantum mechanics of a hydrogen atom	#10	10/16/2017
18 Wed, 10/11/2017	Chap. 13	Quantum mechanics of multi-electron atoms		
Fri, 10/13/2017		Fall break -- No class		
19 Mon, 10/16/2017		Discuss exam questions and topics for presentations/Topic		10/18/2017
20 Wed, 10/18/2017	Chap. 14	Intrinsic spin	#11	10/20/2017
21 Fri, 10/20/2017	Chap. 15	Addition of Angular Momentum	#12	10/23/2017
22 Mon, 10/23/2017	Chap. 15	Multi-electron atoms	#13	10/25/2017
23 Wed, 10/25/2017	Chap. 15	Multi-electron atoms	#14	10/30/2017
24 Fri, 10/27/2017		Effects of nuclear motion		
25 Mon, 10/30/2017	Chap. 17	Time-independent perturbation theory	#15	11/3/2017
26 Wed, 11/01/2017				
27 Fri, 11/03/2017				
28 Mon, 11/06/2017				
29 Wed, 11/08/2017				
30 Fri, 11/10/2017				
31 Mon, 11/13/2017				
32 Wed, 11/15/2017				
33 Fri, 11/17/2017				
34 Mon, 11/20/2017				
Wed, 11/22/2017		Thanksgiving Holiday -- No class		

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Methods for finding approximate solutions to the time-independent Schrödinger equation

Bound states	Continuum states
Variational methods	Scattering theory
Perturbation theory	

$$H|n\rangle = E_n|n\rangle$$

$$H = H^0 + \epsilon H^1$$

In general, we approach the problem using the complete basis set of H^0 :

$$H^0|n^0\rangle = E_n^0|n^0\rangle$$

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$$H|n\rangle = E_n|n\rangle$$

$$H = H^0 + \epsilon H^1$$

$$\text{Assume: } |n\rangle = |n^0\rangle + \epsilon|n^1\rangle + \epsilon^2|n^2\rangle + \dots$$

$$E_n = E_n^0 + \epsilon E_n^1 + \epsilon^2 E_n^2 + \dots$$

Note: This method depends on the non-degenerate

case: $E_n \neq E_m$ for $n \neq m$

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First order corrections

$$\begin{aligned} (H^0 + \epsilon H^1)(|n^0\rangle + \epsilon|n^1\rangle + \epsilon^2|n^2\rangle + \dots) \\ = (E_n^0 + \epsilon E_n^1 + \epsilon^2 E_n^2 + \dots)(|n^0\rangle + \epsilon|n^1\rangle + \epsilon^2|n^2\rangle + \dots) \end{aligned}$$

Collecting terms according to their order:

$$\epsilon^0: H^0|n^0\rangle = E_n^0|n^0\rangle$$

$$\epsilon^1: H^0|n^1\rangle + H^1|n^0\rangle = E_n^0|n^1\rangle + E_n^1|n^0\rangle$$

$$\langle n^0|H^0|n^1\rangle + \langle n^0|H^1|n^0\rangle = \langle n^0|E_n^0|n^1\rangle + \langle n^0|E_n^1|n^0\rangle$$

$$\Rightarrow E_n^1 = \langle n^0|H^1|n^0\rangle$$

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First order corrections -- continued

Since the eigenstates $|n^0\rangle$ for a complete set of states:

$$\text{Assume } |n^1\rangle = \sum_{m \neq n} C_m |m^0\rangle$$

$$C_m = \frac{\langle m^0|H^1|n^0\rangle}{E_n^0 - E_m^0}$$

$$|n^1\rangle = \sum_{m \neq n} \left(\frac{\langle m^0|H^1|n^0\rangle}{E_n^0 - E_m^0} \right) |m^0\rangle$$

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Second order corrections

$$e^2: \quad H^0 |n^2\rangle + H^1 |n^1\rangle = E_n^0 |n^2\rangle + E_n^1 |n^1\rangle + E_n^2 |n^0\rangle$$

$$\langle n^0 | H^0 |n^2\rangle + \langle n^0 | H^1 |n^1\rangle = \langle n^0 | E_n^0 |n^2\rangle + \langle n^0 | E_n^1 |n^1\rangle + \langle n^0 | E_n^2 |n^0\rangle$$

$$\Rightarrow E_n^2 = \langle n^0 | H^1 |n^1\rangle = \sum_{m \neq n} \frac{\langle n^0 | H^1 |m^0\rangle \langle m^0 | H^1 |n^0\rangle}{E_n^0 - E_m^0}$$

$$|n^2\rangle = \sum_{m \neq n} |m^0\rangle \sum_{l \neq n} \frac{\langle m^0 | H^1 |l^0\rangle \langle l^0 | H^1 |n^0\rangle}{(E_n^0 - E_m^0)(E_n^0 - E_l^0)}$$

$$- \sum_{m \neq n} |m^0\rangle \frac{\langle m^0 | H^1 |n^0\rangle \langle n^0 | H^1 |n^0\rangle}{(E_n^0 - E_m^0)^2} - \frac{1}{2} |n^0\rangle \sum_{m \neq n} \frac{|\langle m^0 | H^1 |n^0\rangle|^2}{(E_n^0 - E_m^0)^2}$$

Due to normalization

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Example – Harmonic oscillator

$$H^0 = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + \frac{1}{2} m \omega^2 x^2$$

$$H^1 = \frac{1}{2} m \Omega^2 x^2$$

$$\text{Exact eigenstates: } E_n^{\text{exact}} = \hbar \sqrt{\omega^2 + \Omega^2} \left(n + \frac{1}{2} \right)$$

$$= \hbar \omega \left(1 + \frac{1}{2} \frac{\Omega^2}{\omega^2} - \frac{1}{8} \frac{\Omega^4}{\omega^4} \dots \right) \left(n + \frac{1}{2} \right)$$

Perturbation theory results:

$$E_n^1 = \langle n^0 | \frac{1}{2} m \Omega^2 x^2 | n^0 \rangle = \hbar \omega \frac{\Omega^2}{2 \omega^2} \left(n + \frac{1}{2} \right)$$

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Example: He atom

$$H(\mathbf{r}_1, \mathbf{r}_2) = -\frac{\hbar^2}{2m} \nabla_1^2 - \frac{\hbar^2}{2m} \nabla_2^2 - \frac{2e^2}{r_1} - \frac{2e^2}{r_2} + \frac{e^2}{|\mathbf{r}_1 - \mathbf{r}_2|}$$

$$H^0(\mathbf{r}_1, \mathbf{r}_2) = -\frac{\hbar^2}{2m} \nabla_1^2 - \frac{\hbar^2}{2m} \nabla_2^2 - \frac{2e^2}{r_1} - \frac{2e^2}{r_2}$$

$$H^1(\mathbf{r}_1, \mathbf{r}_2) = \frac{e^2}{|\mathbf{r}_1 - \mathbf{r}_2|}$$

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What happens when the spectrum of H^0 has degeneracy?

$$\begin{aligned} & (H^0 + \epsilon H^1) \left(|n^0\rangle + \epsilon |n^1\rangle + \epsilon^2 |n^2\rangle + \dots \right) \\ &= (E_n^0 + \epsilon E_n^1 + \epsilon^2 E_n^2 + \dots) \left(|n^0\rangle + \epsilon |n^1\rangle + \epsilon^2 |n^2\rangle + \dots \right) \end{aligned}$$

Collecting terms according to their order:

$$\epsilon^0: H^0 |n^0\rangle = E_n^0 |n^0\rangle$$

$$\epsilon^1: H^0 |n^1\rangle + H^1 |n^0\rangle = E_n^0 |n^1\rangle + E_n^1 |n^0\rangle$$

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Since the eigenstates $|n^0\rangle$ for a complete set of states:

Assume $|n^1\rangle = \sum_{m \neq n} C_m |m^0\rangle$

$$C_m = \frac{\langle m^0 | H_1 | n^0 \rangle}{E_n^0 - E_m^0}$$

$$|n^1\rangle = \sum_{m \neq n} \left(\frac{\langle m^0 | H_1 | n^0 \rangle}{E_n^0 - E_m^0} \right) |m^0\rangle$$

- problem when $E_n^0 = E_m^0$
- solution – consider the degenerate states separately

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