

**PHY 711 Classical Mechanics and
Mathematical Methods
10-10:50 AM MWF Olin 103
Plan for Lecture 25:**

Motions of elastic membranes (Chap. 8)

- 1. Review of standing waves on a string**
- 2. Standing waves on a two dimensional membrane.**
- 3. Boundary value problems**

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**Colloquium: “Energy Landscapes for Protein
Folding and Misfolding” Dr. Peter Wolynes,
Wednesday, October 31, 2018, at 4:00 PM**

Dr. Peter G. Wolynes, Bullard-Welch Foundation Professor of Science;
Professor of Chemistry, MSNE, and Physics and Astronomy at Rice University,
George P. Williams, Jr. Lecture Hall, (Olin 101)
Wednesday, October 31, 2018, at 4:00 PM
(Colloquium sponsored jointly with WFU Dept. of Chemistry)

There will be a reception with refreshments at 3:30 PM in the lounge. All
interested persons are cordially invited to attend.

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18	Wed, 10/10/2018	Chap. 7	Wave equation		
	Fri, 10/12/2018	No class	Fall break		
19	Mon, 10/15/2018	Chap. 7	Wave equation		
20	Wed, 10/17/2018	Chap. 7	Fourier Transforms	#12	10/22/2018
21	Fri, 10/19/2018	Chap. 7	Laplace transforms; Contour integrals	#13	10/24/2018
22	Mon, 10/22/2018	Chap. 7	Contour integrals		
23	Wed, 10/24/2018	Chap. 5	Rigid body motion	#14	10/26/2018
24	Fri, 10/26/2018	Chap. 5	Rigid body motion	#15	10/31/2018
25	Mon, 10/29/2018	Chap. 8	Mechanics of elastic membranes	#16	11/02/2018
26	Wed, 10/31/2018				
27	Fri, 11/02/2018				
28	Mon, 11/05/2018				
29	Wed, 11/07/2018				
30	Fri, 11/09/2018				
31	Mon, 11/12/2018				
32	Wed, 11/14/2018				
33	Fri, 11/16/2018				
34	Mon, 11/19/2018				
	Wed, 11/21/2018	No class	Thanksgiving holiday		
	Fri, 11/23/2018	No class	Thanksgiving holiday		
35	Mon, 11/26/2018				

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Elastic media in two or more dimensions --

Review of wave equation in one-dimension – here $\mu(x,t)$ can describe either a longitudinal or transverse wave.

Traveling wave solutions --

$$\frac{\partial^2 \mu}{\partial t^2} - c^2 \frac{\partial^2 \mu}{\partial x^2} = 0$$

Note that for any function $f(q)$ or $g(q)$:

$$\mu(x,t) = f(x-ct) + g(x+ct)$$

satisfies the wave equation.

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Initial value problem : $\mu(x,0) = \phi(x)$ and $\frac{\partial \mu}{\partial t}(x,0) = \psi(x)$

then : $\mu(x,0) = \phi(x) = f(x) + g(x)$

$$\frac{\partial \mu}{\partial t}(x,0) = \psi(x) = -c \left(\frac{df(x)}{dx} - \frac{dg(x)}{dx} \right)$$

$$\Rightarrow f(x) - g(x) = -\frac{1}{c} \int \psi(x') dx'$$

For each x , find $f(x)$ and $g(x)$:

$$f(x) = \frac{1}{2} \left(\phi(x) - \frac{1}{c} \int \psi(x') dx' \right)$$

$$g(x) = \frac{1}{2} \left(\phi(x) + \frac{1}{c} \int \psi(x') dx' \right)$$

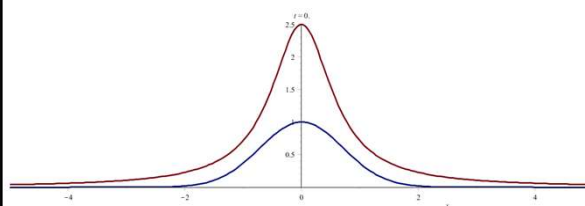
$$\Rightarrow \mu(x,t) = \frac{1}{2} (\phi(x-ct) + \phi(x+ct)) + \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(x') dx'$$

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Example with $\psi(x) = 0$ and $\phi(x) = \frac{1}{x^2 + 0.4}$



Example with $\psi(x) = 0$ and $\phi(x) = e^{-x^2}$

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Standing wave solutions of wave equation:

$$\frac{\partial^2 \mu}{\partial t^2} - c^2 \frac{\partial^2 \mu}{\partial x^2} = 0$$

with $\mu(0, t) = \mu(L, t) = 0$.

Assume: $\mu(x, t) = \Re(e^{-i\omega t} \rho(x))$

where $\frac{d^2 \rho(x)}{dx^2} + k^2 \rho(x) = 0$ $k = \frac{\omega}{c}$

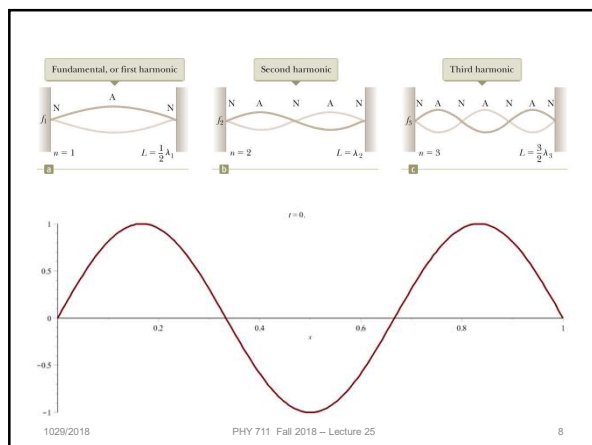
$$\rho_v(x) = A \sin\left(\frac{v\pi x}{L}\right)$$

$$k_v = \frac{v\pi}{L} \quad \omega_v = ck_v$$

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Wave motion on a two-dimensional surface – elastic membrane (transverse wave; linear regime).

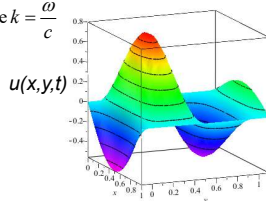
Two-dimensional wave equation:

$$\frac{\partial^2 u}{\partial t^2} - c^2 \nabla^2 u = 0 \quad \text{where } c^2 = \frac{\tau}{\sigma}$$

Standing wave solutions:

$$u(x, y, t) = \Re(e^{-i\omega t} \rho(x, y))$$

$$(\nabla^2 + k^2) \rho(x, y) = 0 \quad \text{where } k = \frac{\omega}{c}$$



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Lagrangian density: $\mathcal{L}\left(u, \frac{\partial u}{\partial x}, \frac{\partial u}{\partial y}, \frac{\partial u}{\partial t}; x, y, t\right)$

$$L = \int \mathcal{L}\left(u, \frac{\partial u}{\partial x}, \frac{\partial u}{\partial y}, \frac{\partial u}{\partial t}; x, y, t\right) dx dy$$

Hamilton's principle:

$$\delta \int_{t_1}^{t_2} L dt = 0$$

$$\frac{\partial \mathcal{L}}{\partial u} - \frac{\partial}{\partial t} \frac{\partial \mathcal{L}}{\partial (\partial u / \partial t)} - \frac{\partial}{\partial x} \frac{\partial \mathcal{L}}{\partial (\partial u / \partial x)} - \frac{\partial}{\partial y} \frac{\partial \mathcal{L}}{\partial (\partial u / \partial y)} = 0$$

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Lagrangian density for elastic membrane with constant σ and τ :

$$\mathcal{L}\left(u, \frac{\partial u}{\partial x}, \frac{\partial u}{\partial y}, \frac{\partial u}{\partial t}; x, y, t\right) = \frac{1}{2} \sigma \left(\frac{\partial u}{\partial t}\right)^2 - \frac{1}{2} \tau (\nabla u)^2$$

$$\frac{\partial \mathcal{L}}{\partial u} - \frac{\partial}{\partial t} \frac{\partial \mathcal{L}}{\partial (\partial u / \partial t)} - \frac{\partial}{\partial x} \frac{\partial \mathcal{L}}{\partial (\partial u / \partial x)} - \frac{\partial}{\partial y} \frac{\partial \mathcal{L}}{\partial (\partial u / \partial y)} = 0$$

$$\frac{\partial^2 u}{\partial t^2} - c^2 \nabla^2 u = 0 \quad \text{where } c^2 = \frac{\tau}{\sigma}$$

Two - dimensional wave equation:

$$\frac{\partial^2 u}{\partial t^2} - c^2 \nabla^2 u = 0 \quad \text{where } c^2 = \frac{\tau}{\sigma}$$

Standing wave solutions:

$$u(x, y, t) = \Re\left\{e^{-i\omega t} \rho(x, y)\right\}$$

$$(\nabla^2 + k^2) \rho(x, y) = 0 \quad \text{where } k = \frac{\omega}{c}$$

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Consider a rectangular boundary:

b



a

Clamped boundary conditions:

$$\rho(0, y) = \rho(a, y) = \rho(x, 0) = \rho(x, b) = 0 \quad (\nabla^2 + k^2) \rho(x, y) = 0$$

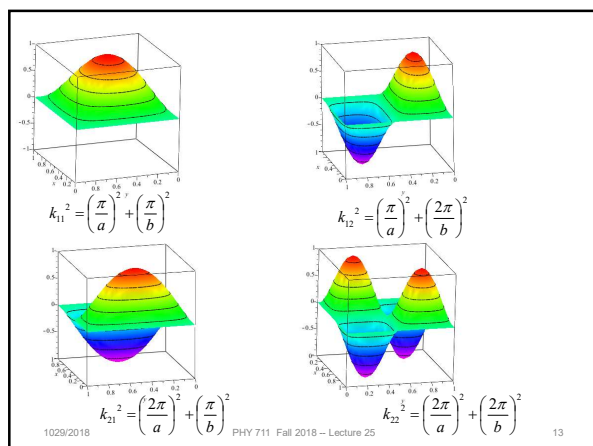
$$\Rightarrow \rho_{mn}(x, y) = A \sin\left(\frac{m\pi x}{a}\right) \sin\left(\frac{n\pi y}{b}\right) \quad \text{where } k = \frac{\omega}{c}$$

$$k_{mn}^2 = \left(\frac{m\pi}{a}\right)^2 + \left(\frac{n\pi}{b}\right)^2 \quad \omega_{mn} = ck_{mn}$$

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More general boundary conditions:

$\tau \nabla u|_b = \kappa u|_b$ represents bounded side constrained with spring

$\tau \nabla u|_b = 0$ represents "free" side

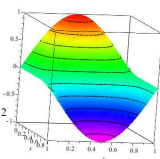
Mixed boundary conditions:

$$\rho(x,0) = \rho(x,b) = \frac{\partial \rho(0,y)}{\partial x} = \frac{\partial \rho(a,y)}{\partial x} = 0$$

$$\Rightarrow \rho_{mn}(x,y) = A \cos\left(\frac{m\pi x}{a}\right) \sin\left(\frac{n\pi y}{b}\right)$$

$$k_{mn}^2 = \left(\frac{m\pi}{a}\right)^2 + \left(\frac{n\pi}{b}\right)^2 \quad \omega_{mn} = ck_{mn}$$

$$k_{11}^2 = \left(\frac{\pi}{a}\right)^2 + \left(\frac{\pi}{b}\right)^2$$



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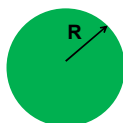
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Consider a circular boundary:

Clamped boundary conditions for $\rho(r,\varphi)$:

$$\rho(R,\varphi) = 0$$



$$(\nabla^2 + k^2)\rho(r,\varphi) = 0 \quad \text{where } k = \frac{\omega}{c}$$

In cylindrical coordinate system

$$\nabla^2 = \frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \varphi^2}$$

Assume: $\rho(r,\varphi) = f(r)\Phi(\varphi)$

Let: $\Phi(\varphi) = e^{im\varphi}$

Note: $\Phi(\varphi) = \Phi(\varphi + 2\pi)$

$$\Rightarrow m = \text{integer}$$

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Consider circular boundary -- continued

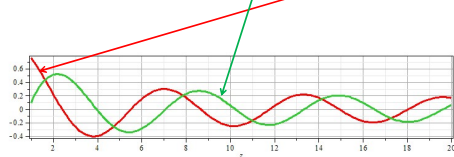
Differential equation for radial function:

$$\left(\frac{d^2}{dr^2} + \frac{1}{r} \frac{d}{dr} - \frac{m^2}{r^2} + k^2 \right) f(r) = 0$$

⇒ Bessel equation of integer order with transcendental solutions

Cylindrical Bessel function $J_m(z)$

Cylindrical Neumann function $N_m(z)$ also called $Y_m(z)$



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Some properties of Bessel functions

Ascending series: $J_m(z) = \left(\frac{z}{2} \right)^m \sum_{j=0}^{\infty} \frac{(-1)^j}{j!(j+m)!} \left(\frac{z}{2} \right)^{2j}$

Recursion relations: $J_{m-1}(z) + J_{m+1}(z) = \frac{2m}{z} J_m(z)$

$$J_{m-1}(z) - J_{m+1}(z) = 2 \frac{dJ_m(z)}{dz}$$

Asymptotic form: $J_m(z) \xrightarrow{z \gg 1} \sqrt{\frac{2}{\pi z}} \cos\left(z - \frac{m\pi}{2} - \frac{\pi}{4}\right)$

Zeros of Bessel functions $J_m(z_{mn}) = 0$

$m=0$: $z_{0n} = 2.406, 5.520, 8.654, \dots$

$m=1$: $z_{1n} = 3.832, 7.016, 10.173, \dots$

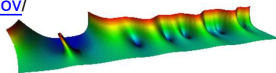
$m=2$: $z_{2n} = 5.136, 8.417, 11.620, \dots$

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<http://dlmf.nist.gov/>



NIST Digital Library of Mathematical Functions

Project News

2014-08-29 DLMF Update: Version 1.0.2

2014-04-25 DLMF Update: Version 1.0.1b: errata & improved MathML

2014-03-21 DLMF Update: Version 1.0.1: New Features: Improved Math & 3D Graphics

2013-08-16 Billie C. Carlson, DLMF Author, dies at age 89

[More news](#)

Foreword

Preface

Mathematical Introduction

1 Algebraic and Analytic Methods

2 Asymptotic Approximations

3 Numerical Methods

4 Elementary Functions

5 Gamma Function

6 Exponential, Logarithmic, Sine, and Cosine Integrals

7 Error Functions, Dawson's and Fresnel Integrals

8 Incomplete Gamma and Related Functions

9 Airy and Related Functions

10 Bessel Functions

11 Struve and Related Functions

12 Parabolic Cylinder Functions

19 Elliptic Integrals

20 Theta Functions

21 Multidimensional Theta Functions

22 Jacobian Elliptic Functions

23 Weierstrass Elliptic and Modular Functions

24 Bernoulli and Euler Polynomials

25 Zeta and Related Functions

26 Combinatorial Analysis

27 Functions of Number Theory

28 Mathieu Functions and Hill's Equation

29 Lamé Functions

30 Spheroidal Wave Functions

31 Heun Functions

32 Painlevé Transcendents

33 Coulomb Functions

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Series expansions of Bessel and Neumann functions

$$J_\nu(z) = \left(\frac{1}{2}z\right)^\nu \sum_{k=0}^{\infty} \frac{(-1)^k}{k! \Gamma(\nu + k + 1)} \left(\frac{1}{4}z^2\right)^k.$$

$$Y_n(z) = -\frac{\left(\frac{1}{2}z\right)^{-n}}{\pi} \sum_{k=0}^{n-1} \frac{(n-k-1)!}{k!} \left(\frac{1}{4}z^2\right)^k + \frac{2}{\pi} \ln\left(\frac{1}{2}z\right) J_n(z) \\ - \frac{\left(\frac{1}{2}z\right)^n}{\pi} \sum_{k=0}^{\infty} \left(\psi(k+1) + \psi(n+k+1)\right) \frac{\left(-\frac{1}{4}z^2\right)^k}{k!(n+k)!}.$$

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Some properties of Bessel functions -- continued

Note : It is possible to prove the following

identity for the functions $J_m\left(\frac{z_{mn}}{R}r\right)$:

$$\int_0^R J_m\left(\frac{z_{mn}}{R}r\right) J_m\left(\frac{z_{mn'}}{R}r\right) r dr = \frac{R^2}{2} (J_{m+1}(z_{mn}))^2 \delta_{nn'}.$$

Returning to differential equation for radial function :

$$\left(\frac{d^2}{dr^2} + \frac{1}{r} \frac{d}{dr} - \frac{m^2}{r^2} + k^2\right) f(r) = 0$$

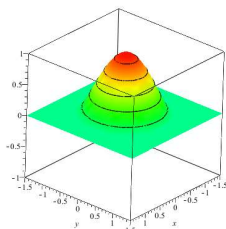
$$\Rightarrow f_{mn}(r) = A J_m\left(\frac{z_{mn}}{R}r\right); \quad k_{mn} = \frac{z_{mn}}{R}$$

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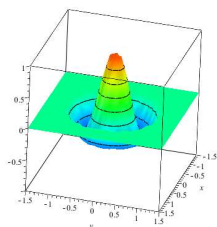
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$$\rho_{01}(r, \varphi) = f_{01}(r) = A J_0\left(\frac{z_{01}}{R}r\right)$$



$$k_{01} = \frac{2.406}{R}$$

$$\rho_{02}(r, \varphi) = f_{02}(r) = A J_0\left(\frac{z_{02}}{R}r\right)$$

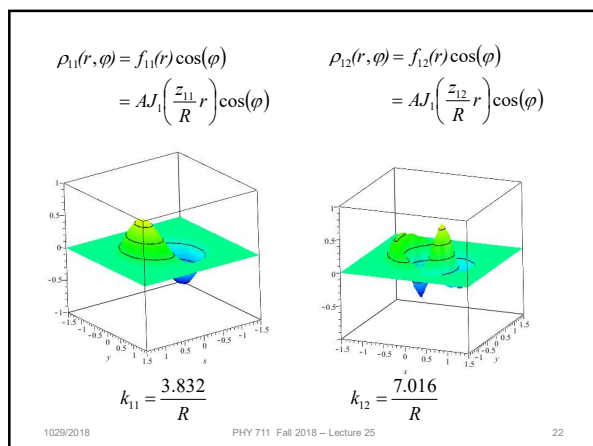


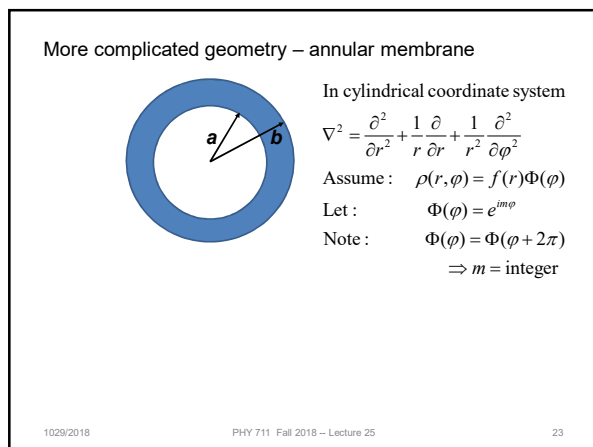
$$k_{02} = \frac{5.520}{R}$$

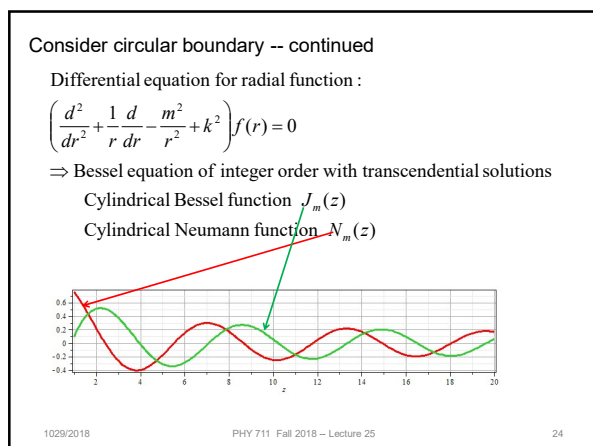
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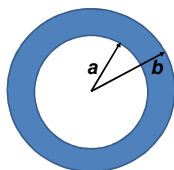
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Normal modes of an annular membrane -- continued



Differential equation for radial function:

$$\left(\frac{d^2}{dr^2} + \frac{1}{r} \frac{d}{dr} - \frac{m^2}{r^2} + k^2 \right) f(r) = 0$$

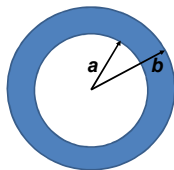
General form of radial function: $f(r) = AJ_m(kr) + BN_m(kr)$

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Normal modes of an annular membrane -- continued



Boundary conditions:

$$f(a) = 0 \quad f(b) = 0$$

$$AJ_m(ka) + BN_m(ka) = 0$$

$$AJ_m(kb) + BN_m(kb) = 0$$

 $\Rightarrow$ 2 equations and 2 unknowns -- k and $\frac{B}{A}$

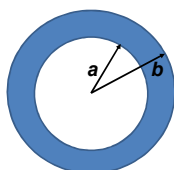
$$\frac{B}{A} = \frac{-J_m(ka)}{N_m(ka)} = \frac{-J_m(kb)}{N_m(kb)} \quad (\text{transcendental equation for } k)$$

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Normal modes of an annular membrane -- continued



Boundary conditions:

$$f(a) = 0 \quad f(b) = 0$$

$$\frac{B}{A} = \frac{-J_m(ka)}{N_m(ka)} = \frac{-J_m(kb)}{N_m(kb)} \quad \text{-- in terms of solution } k_{mn} :$$

$$f(r) = A \left(J_m(k_{mn}r) - \frac{J_m(k_{mn}a)}{N_m(k_{mn}a)} N_m(k_{mn}r) \right)$$

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