PHY 711 Classical Mechanics and Mathematical Methods
10-10:50 AM MWF Online or (occasionally) in Olin 103

Discussion notes for Lecture 3

Scattering theory;

Center of mass and laboratory coordinate frames

Schedule for weekly one-on-one meetings

Nick – 11 AM Monday (ED/ST)
Tim – 9 AM Tuesday
Bamidele – 7 PM Tuesday
Zhi – 9 PM Tuesday
Jeanette – 11 AM Friday
Derek – 12 PM Friday

Your questions –

From Jeanette

1. Why do we start with scattering? I'm just curious, I don't think scattering was specifically addressed in my undergrad course.
2. Slide #6 Last sentence - What is the function b(theta)? Is it simply referring to solving that equation for b?
3. Slide 19 - Is the whole yellow boxed equation for elastic scattering or only the far-right term?

From Nick

1. What is \varphi representing? Is it the angle of a sector of the beam?
2. Why does b(\theta)=Dsin(\pi/2 - \theta/2)? I'm not sure I follow that.
3. Can we spend some time going over the geometry of the center of mass frame?

From Tim

1. For Assignment #2 are we assuming the larger mass, M is at rest or does it have an initial velocity?

From Derek

1. Assignment #2 says that the total energy of the system in the problem is V0 but shouldn't it be V0 plus the kinetic energy of the moving particle?

From Zhi–

1. About the lecture 3. I have a question: What is laboratory reference frame' strong points compared to center of mass reference frame when studying scattering?

PHY 711 Classical Mechanics and Mathematical Methods

MWF 10 AM-10:50 AM	**OPL 103**	**http://www.wfu.edu/~natalie/f20phy711/**

Instructor: Natalie Holzwarth	**Office:**300 OPL	**e-mail:**natalie@wfu.edu

Course schedule

(Preliminary schedule -- subject to frequent adjustment.)

	Date	F&W Reading	Topic	Assignment	Due
1	Wed, 8/26/2020	Chap. 1	Introduction	#1	8/31/2020
2	Fri, 8/28/2020	Chap. 1	Scattering theory	#2	9/02/2020
3	Mon, 8/31/2020	Chap. 1	Scattering theory	#3	9/04/2020
4	Wed, 9/02/2020	Chap. 1	Scattering theory		
5	Fri, 9/04/2020	Chap. 2	Non-inertial coordinate systems		
6	Mon, 9/07/2020	Chap. 3	Calculus of Variation		

PHY 711 -- Assignment #3

Aug. 31, 2020

Read Chapter 1 in **Fetter & Walecka**.

1. In Lecture 3, we derived equations relating the laboratory scattering angle to the scattering angle in the center of mass reference frame. We also worked out the relationship between the differential scattering cross sections in the laboratory and center of mass frames. After you have convinced yourselves of the validity of those derivations, evaluate both the lab and center of mass scattering angles and the corresponding cross section factors for the following mass ratios
   a. $m_1/m_2=1$
   b. $m_1/m_2=1/10$
   c. $m_1/m_2=10/1$

From your questions -- About homework #2

1. For Assignment #2 are we assuming the larger mass, M is at rest or does it have an initial velocity?
2. Assignment #2 says that the total energy of the system in the problem is V0 but shouldn't it be V0 plus the kinetic energy of the moving particle?

The total energy of the system is a given constant and V_0 happens to a particular energy value (that hopefully makes the algebra of the problem somewhat convenient.).

1. Consider a particle of mass m moving in the vicinity of another particle of mass M where $m << M$. The particles interact with a conservative central potential of the form

Assume particle M is at rest.

$$V(r) = V_0\left(\left(\frac{r_0}{r}\right)^2 - \left(\frac{r_0}{r}\right)\right),$$

where r denotes the magnitude of the particle separation and V_0 and r_0 denote energy and length constants, respectively. The total energy of the system is V_0.

(a) First consider the case where the impact parameter $b = 0$. Find the distance of closest approach of the particles.

(b) Now consider the case where the impact parameter $b = r_0$. Find the distance of closest approach of the particles.

Your question -- Why do we start with scattering? I'm just curious, I don't think scattering was specifically addressed in my undergrad course.

Comment – Short answer – Fetter and Walecka chose scattering as the first chapter, so I thought that we should play along. Longer answer – scattering ideas come up a lot in physics in different contexts. In fact, the analysis for classical mechanics is especially rewarding (at least in principle). This is because the particle trajectory can be directly related to the scattering behavior. More generally, cross section analysis is a standard way to compare “theory” and experiment.

Scattering theory:

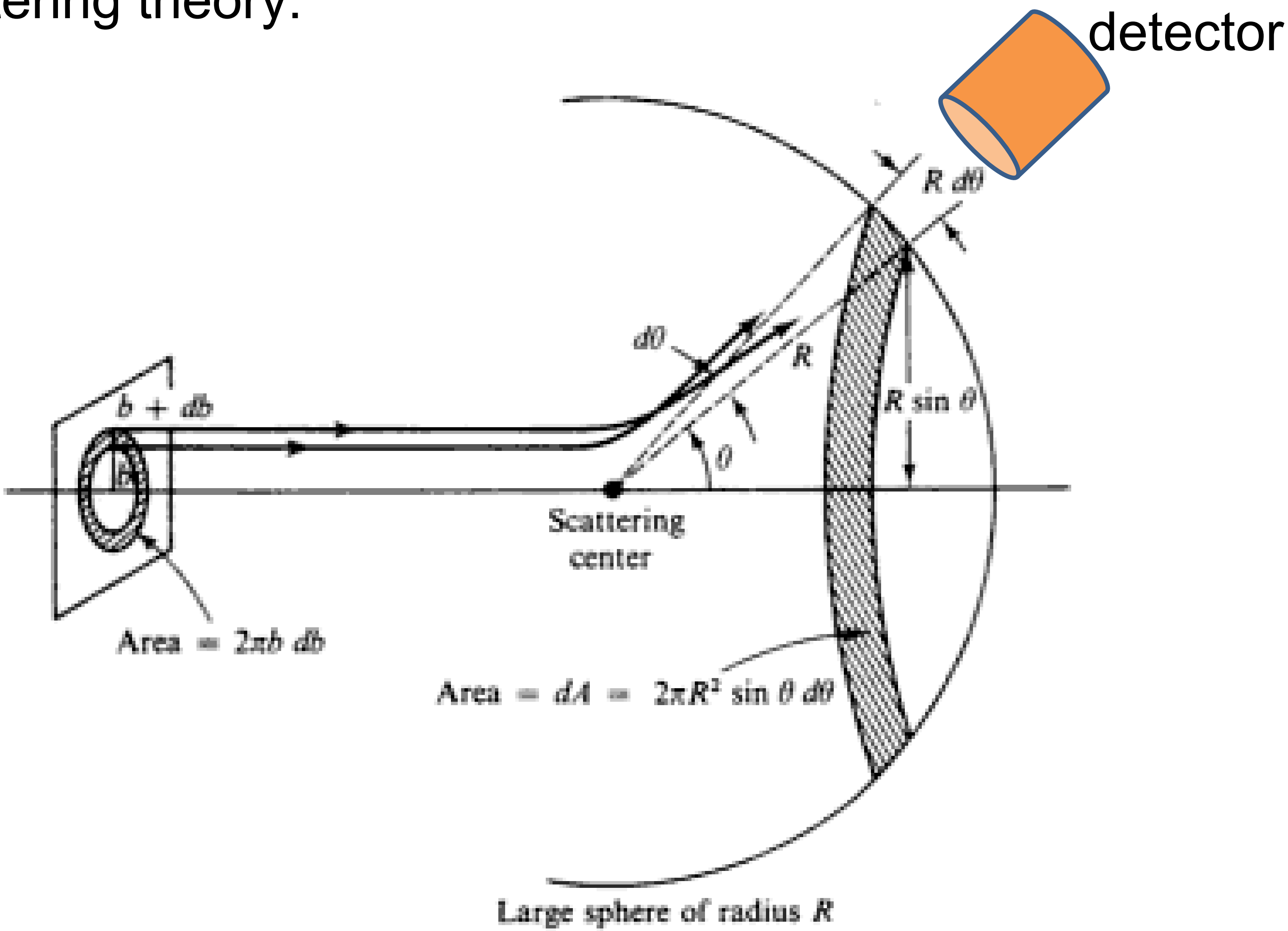


Figure 5.5 The scattering problem and relation of cross section to impact parameter.

Can you think of examples of such an experimental setup?

Other experimental designs –
At CERN https://home.cern/science/experiments/totem
the study of highly energetic proton-proton scattering is designed in the center of mass frame of reference by accelerating two proton beams focused to collide head on in the Large Hadron Collider LHC facility.

Figure from CERN website

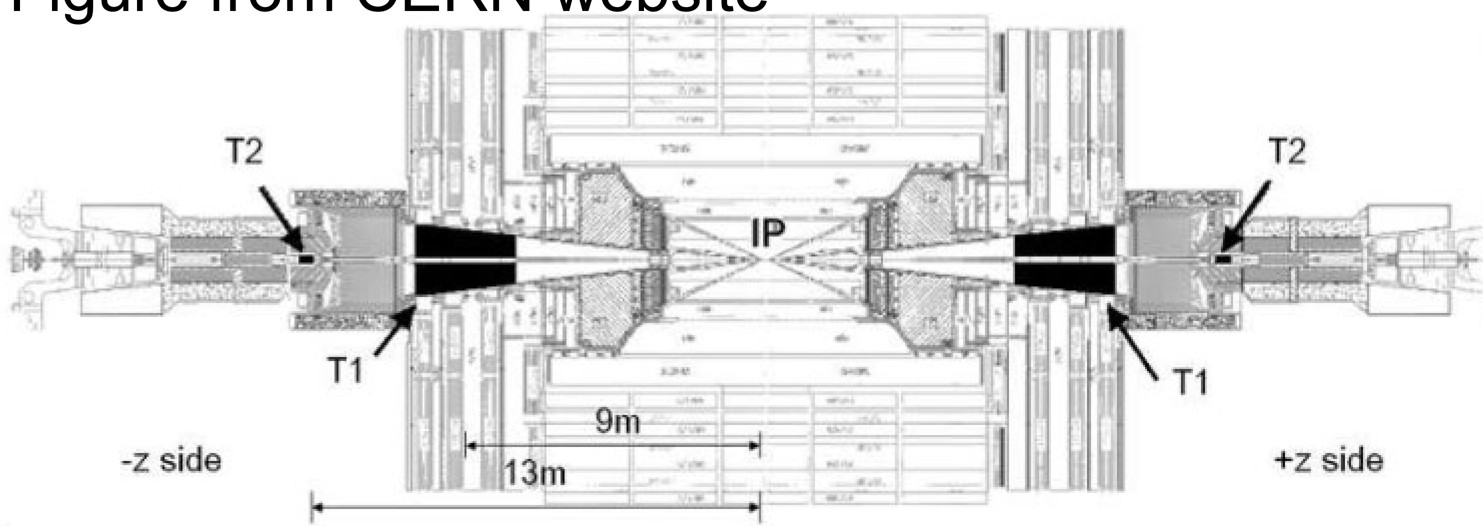


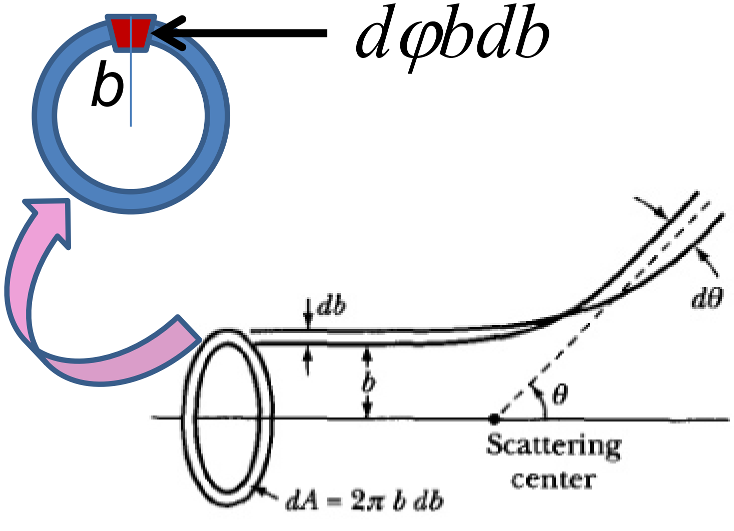
Figure 1.17: View of the inelastic forward trackers T1 and T2 inside the CMS detector.

Differential cross section

$$\left(\frac{d\sigma}{d\Omega}\right)=\frac{\text{Number of detected particles at }\theta\text{ per target particle}}{\text{Number of incident particles per unit area}}$$

$$=\text{Area of incident beam that is scattered into detector at angle }\theta$$

$$d\varphi b db$$

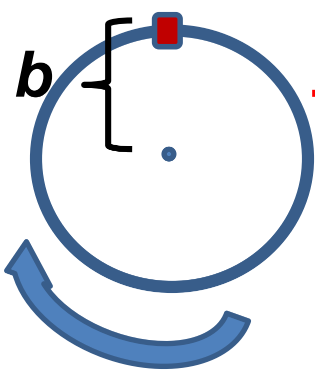


$$\left(\frac{d\sigma}{d\Omega}\right)=\frac{d\varphi\, b\, db}{d\varphi \sin\theta\, d\theta}=\frac{b}{\sin\theta}\left|\frac{db}{d\theta}\right|$$

Figure from Marion & Thorton, Classical Dynamics

Your questions -- Slide #6 Last sentence - What is the function b(theta)? Is it simply referring to solving that equation for b? What is \varphi representing? Is it the angle of a sector of the beam?

Comment -- We imagine that the beam of particles has a cylindrical geometry and that the physics is totally uniform in the azimuthal direction. The cross section of the beam is a circle.

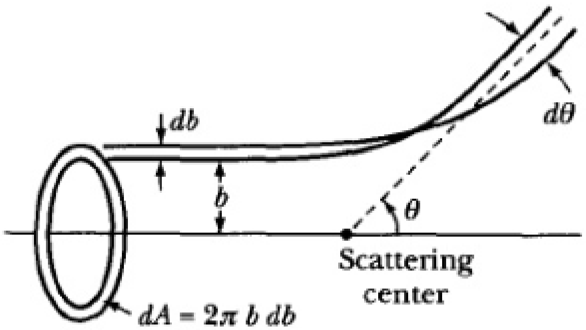


←This piece of the beam scatters into the detector at angle θ

This logic leads to the notion that b is a function of theta and we will try to find $b(\theta)$ for various cases.

$$\varphi \equiv \text{azimuthal angle}$$

Note: The notion of cross section is common to many areas of physics including classical mechanics, quantum mechanics, optics, etc. Only in the **classical mechanics** can we calculate it from a knowledge of the particle trajectory as it relates to the scattering geometry.



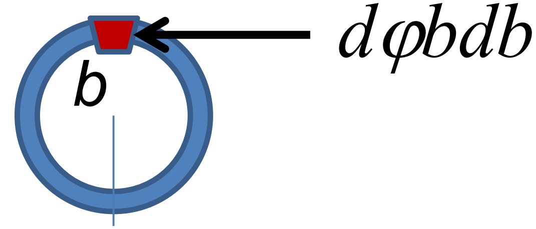
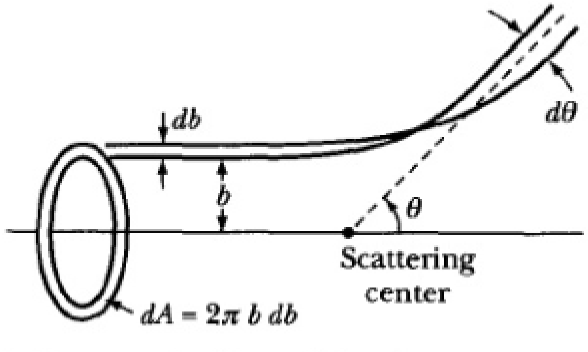


Figure from Marion & Thorton, Classical Dynamics

$$\left(\frac{d\sigma}{d\Omega}\right)=\frac{d\varphi\, b\, db}{d\varphi \sin\theta\, d\theta}=\frac{b}{\sin\theta}\left|\frac{db}{d\theta}\right|$$

Note: We are assuming that the process is isotropic in ϕ

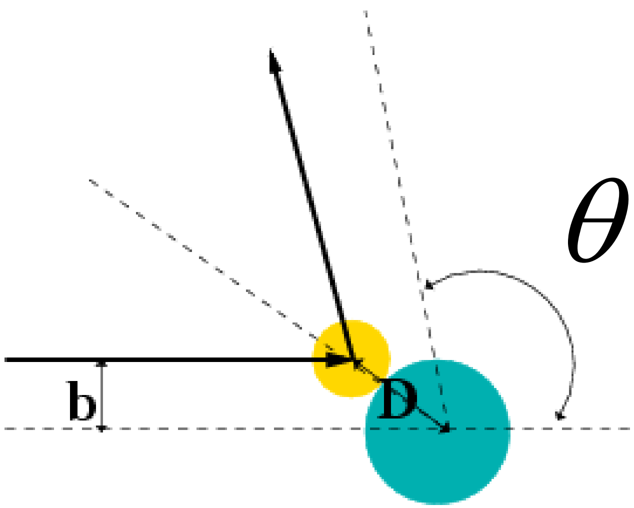
Simple example – collision of hard spheres



$$\left(\frac{d\sigma}{d\Omega}\right) = \frac{b}{\sin\theta}\left|\frac{db}{d\theta}\right|$$

Microscopic view:

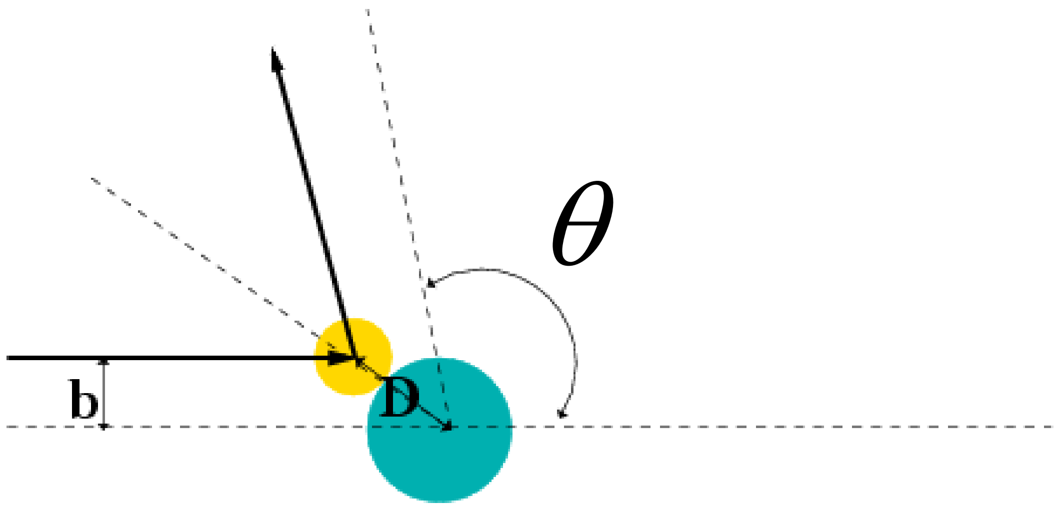
$$b(\theta) = ?$$

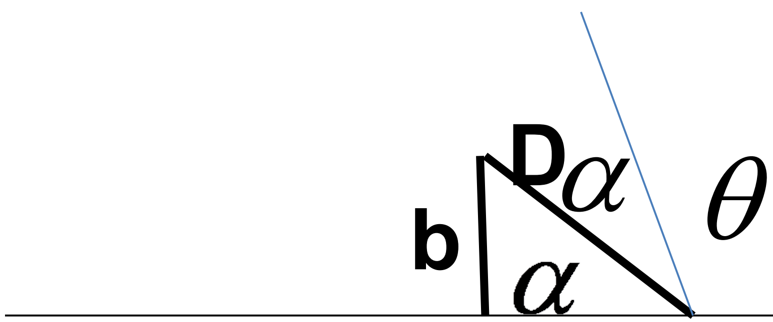


$$b(\theta) = D\sin\left(\frac{\pi}{2} - \frac{\theta}{2}\right)$$

$$\left(\frac{d\sigma}{d\Omega}\right) = \frac{D^2}{4}$$

Your question -- Why does b(\theta)=Dsin(\pi/2 - \theta/2)? I'm not sure I follow that.

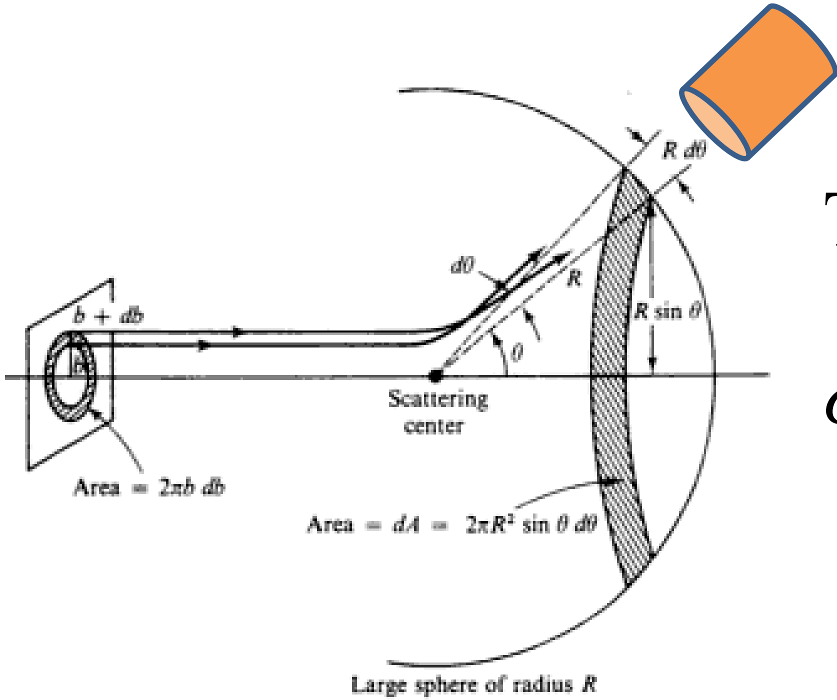




$$b = D\sin\alpha = D\sin\left(\frac{\pi}{2} - \frac{\theta}{2}\right)$$

$$2\alpha + \theta = \pi$$

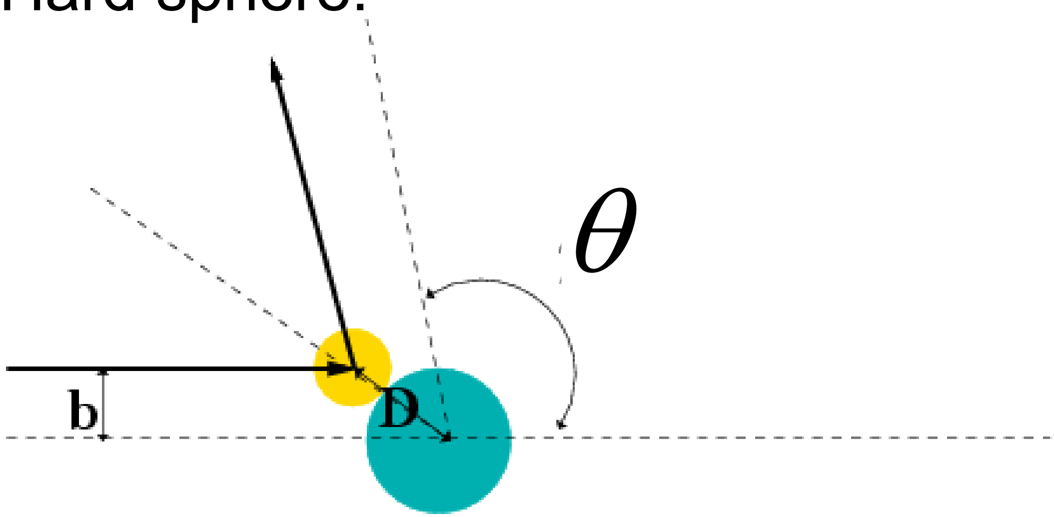
Simple example – collision of hard spheres -- continued



Total scattering cross section:

$$\sigma = \int \left(\frac{d\sigma}{d\Omega} \right) d\Omega$$

Hard sphere:

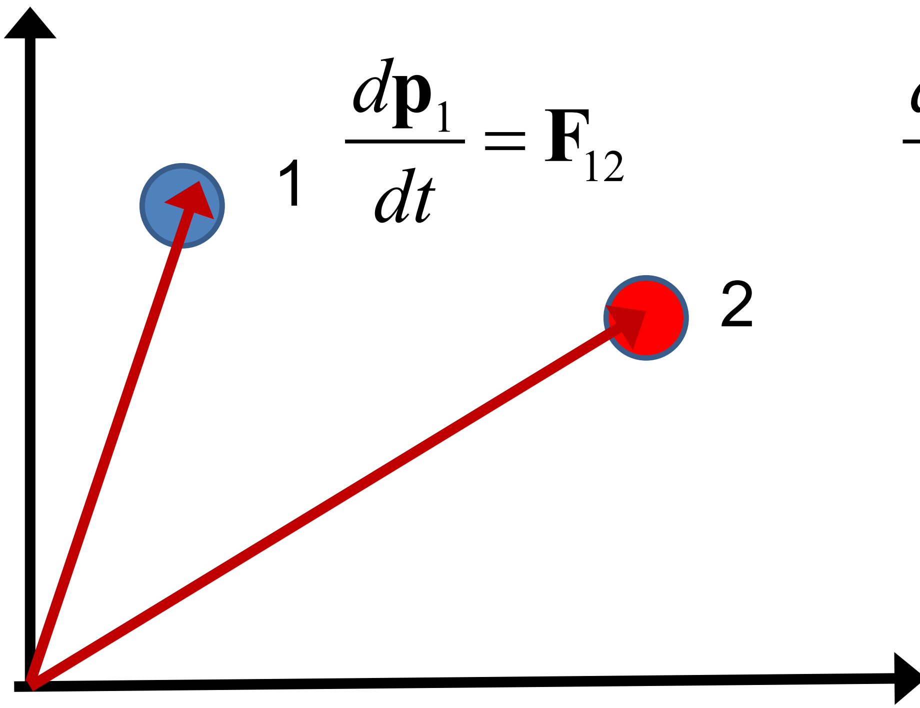


$$\left(\frac{d\sigma}{d\Omega} \right) = \frac{D^2}{4}$$

$$\sigma = \pi D^2$$

Now consider the more general case of particle interactions and the corresponding scattering analysis.

Relationship of scattering cross-section to particle interactions -- Classical mechanics of a conservative 2-particle system.



$$\frac{d\mathbf{p}_1}{dt} = \mathbf{F}_{12} \qquad \frac{d\mathbf{p}_2}{dt} = \mathbf{F}_{21}$$

$$\mathbf{F}_{12} = -\nabla_1 V\left(\mathbf{r}_1 - \mathbf{r}_2\right) \quad \Rightarrow E = \frac{1}{2}m_1 v_1^2 + \frac{1}{2}m_2 v_2^2 + V\left(\mathbf{r}_1 - \mathbf{r}_2\right)$$

Typical two-particle interactions –

Central potential: $$V(\mathbf{r}_1 - \mathbf{r}_2) = V(|\mathbf{r}_1 - \mathbf{r}_2|) \equiv V(r)$$

Hard sphere: $$V(r) = \begin{cases} \infty & r \le a \\ 0 & r > a \end{cases}$$

Coulomb or gravitational: $$V(r) = \frac{K}{r}$$

Lennard-Jones: $$V(r) = \frac{A}{r^{12}} - \frac{B}{r^6}$$

Scattering theory can help us analyze the interaction potential *V(r).* First we need to simply the number of variables.

Relationship between center of mass and laboratory frames of reference. At and time *t*, the following relationships apply --

Definition of center of mass $\mathbf{R}_{CM}$

$$m_1\mathbf{r}_1 + m_2\mathbf{r}_2 = \left(m_1 + m_2\right)\mathbf{R}_{CM}$$

$$m_1\dot{\mathbf{r}}_1 + m_2\dot{\mathbf{r}}_2 = \left(m_1 + m_2\right)\dot{\mathbf{R}}_{CM} = \left(m_1 + m_2\right)\mathbf{V}_{CM}$$

Note that $\dot{\mathbf{r}} \equiv \frac{d\mathbf{r}}{dt}$

$$E = \frac{1}{2}m_1 v_1^2 + \frac{1}{2}m_2 v_2^2 + V\left(\mathbf{r}_1 - \mathbf{r}_2\right)$$

$$= \frac{1}{2}\left(m_1 + m_2\right)V_{CM}^2 + \frac{1}{2}\mu\left|\mathbf{v}_1 - \mathbf{v}_2\right|^2 + V\left(\mathbf{r}_1 - \mathbf{r}_2\right)$$

where: $\mu \equiv \dfrac{m_1 m_2}{m_1 + m_2}$

Classical mechanics of a conservative 2-particle system -- continued

$$E = \frac{1}{2}(m_1 + m_2)V_{CM}^2 + \frac{1}{2}\mu|\mathbf{v}_1 - \mathbf{v}_2|^2 + V(\mathbf{r}_1 - \mathbf{r}_2)$$

For central potentials: $V(\mathbf{r}_1 - \mathbf{r}_2) = V(|\mathbf{r}_1 - \mathbf{r}_2|) \equiv V(r_{12})$

Relative angular momentum is also conserved:

$$\mathbf{L}_{12} \equiv \mathbf{r}_{12} \times \mu \mathbf{v}_{12}$$

$$E = \frac{1}{2}(m_1 + m_2)V_{CM}^2 + \frac{1}{2}\mu v_{12}^2 + \frac{L_{12}^2}{2\mu r_{12}^2} + V(r_{12})$$

Simpler notation:

$$E = \frac{1}{2}(m_1 + m_2)V_{CM}^2 + \frac{1}{2}\mu \dot{r}^2 + \frac{\ell^2}{2\mu r^2} + V(r)$$

Simpler notation:

$$E = \frac{1}{2}(m_1 + m_2)V_{CM}^2 + \frac{1}{2}\mu\dot{r}^2 + \frac{\ell^2}{2\mu r^2} + V(r)$$

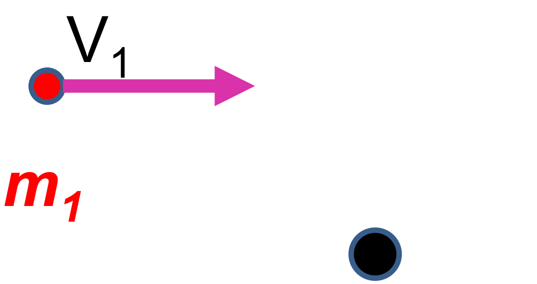
constants

vary in time

For scattering analysis only need to know trajectory **before** and **after** the collision.

Note: The following analysis will be carried out in the center of mass frame of reference.

In laboratory frame:

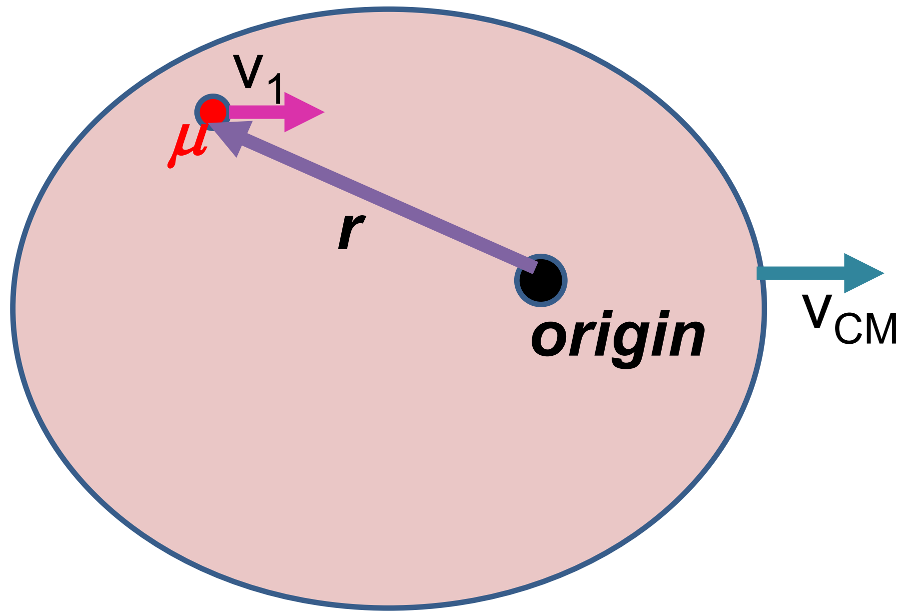


$$\mu = \frac{m_1 m_{\text{target}}}{m_1 + m_{\text{target}}}$$

$$\ell = \left| \mathbf{r} \times \mu \mathbf{v_1} \right|$$

***m**target*

In center-of-mass frame:



Also note: We are assuming that the interaction between particle and target *V(r)* conserves energy and angular momentum.

Your question – What is laboratory reference frame' strong points compared to center of mass reference frame when studying scattering?

Comment -- Typically, the laboratory frame is where the data is taken, but the center of mass frame is where the analysis is most straightforward.

Previous equations --

$$E = \frac{1}{2}(m_1 + m_2)V_{CM}^2 + \frac{1}{2}\mu\dot{r}^2 + \frac{\ell^2}{2\mu r^2} + V(r)$$

constant

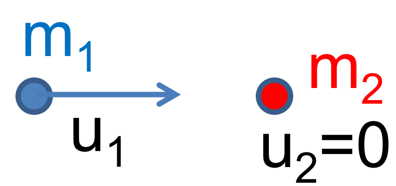
relative coordinate system; visualize as “in” CM frame

It is often convenient to analyze the scattering cross section in the center of mass reference frame.

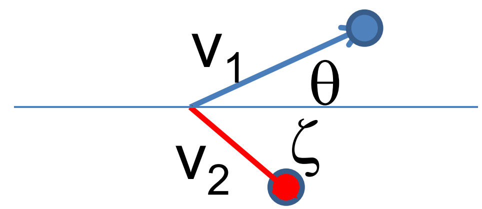
Relationship between normal laboratory reference and center of mass:

Laboratory reference frame:

Before

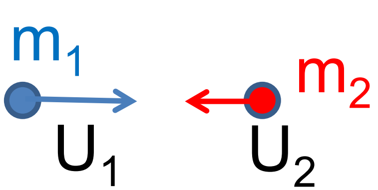


After

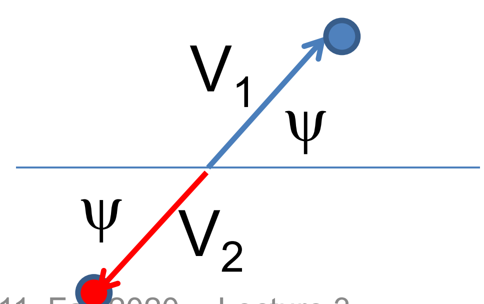


Center of mass reference frame:

Before



After



Relationship between center of mass and laboratory frames of reference -- continued

Since m_2 is initially at rest :

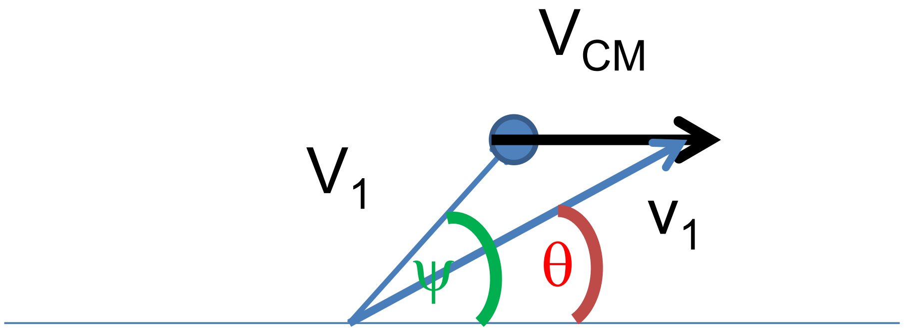
$$\mathbf{V}_{CM} = \frac{m_1}{m_1+m_2}\mathbf{u}_1 \qquad \mathbf{u}_1 = \mathbf{U}_1 + \mathbf{V}_{CM} \quad \Rightarrow \mathbf{U}_1 = \frac{m_2}{m_1+m_2}\mathbf{u}_1 = \frac{m_2}{m_1}\mathbf{V}_{CM}$$

$$\mathbf{u}_2 = \mathbf{U}_2 + \mathbf{V}_{CM} \quad \Rightarrow \mathbf{U}_2 = -\frac{m_1}{m_1+m_2}\mathbf{u}_1 = -\mathbf{V}_{CM}$$

$$\mathbf{v}_1 = \mathbf{V}_1 + \mathbf{V}_{CM}$$

$$\mathbf{v}_2 = \mathbf{V}_2 + \mathbf{V}_{CM}$$

Relationship between center of mass and laboratory frames of reference for the scattering particle 1



$$\mathbf{v}_1 = \mathbf{V}_1 + \mathbf{V}_{CM}$$

$$v_1 \sin\theta = V_1 \sin\psi$$

$$v_1 \cos\theta = V_1 \cos\psi + V_{CM}$$

$$\tan\theta = \frac{\sin\psi}{\cos\psi + V_{CM}/V_1} = \frac{\sin\psi}{\cos\psi + m_1/m_2}$$

For elastic scattering

Your question – Slide 19 - Is the whole yellow boxed equation for elastic scattering or only the far-right term?

$$\mathbf{v}_1 = \mathbf{V}_1 + \mathbf{V}_{CM}$$

$$v_1 \sin\theta = V_1 \sin\psi$$

$$v_1 \cos\theta = V_1 \cos\psi + V_{CM}$$

$$\tan\theta = \frac{\sin\psi}{\cos\psi + V_{CM} / V_1} = \frac{\sin\psi}{\cos\psi + m_1 / m_2}$$

General case

Special case of elastic scattering

For elastic scattering

$$V_{CM} / V_1 = m_1 / m_2$$

Digression – elastic scattering

$$\tfrac{1}{2}m_1U_1^2 + \tfrac{1}{2}m_2U_2^2 + \cancel{\tfrac{1}{2}(m_1 + m_2)V_{CM}^2}$$
$$= \tfrac{1}{2}m_1V_1^2 + \tfrac{1}{2}m_2V_2^2 + \cancel{\tfrac{1}{2}(m_1 + m_2)V_{CM}^2}$$

Also note:

$$m_1\mathbf{U}_1 + m_2\mathbf{U}_2 = 0 \qquad\qquad m_1\mathbf{V}_1 + m_2\mathbf{V}_2 = 0$$

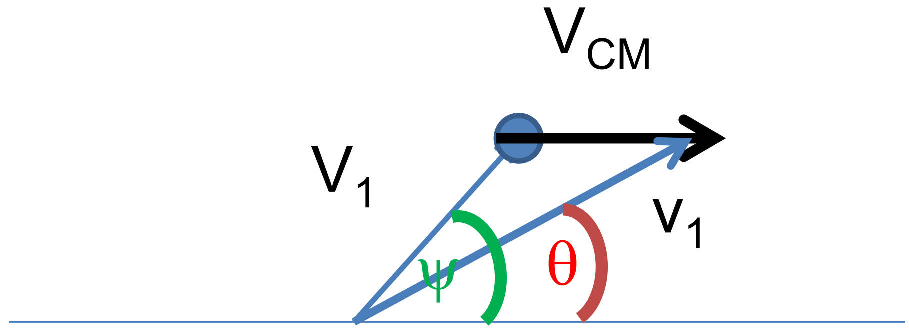
$$\mathbf{U}_1 = \frac{m_2}{m_1}\mathbf{V}_{CM} \qquad\qquad \mathbf{U}_2 = -\mathbf{V}_{CM}$$

$$\Rightarrow |\mathbf{U_1}| = |\mathbf{V_1}| \quad \text{and} \quad |\mathbf{U_2}| = |\mathbf{V_2}| = |\mathbf{V_{CM}}|$$

$$\text{Also note that :} \quad m_1|\mathbf{U_1}| = m_2|\mathbf{U_2}|$$

$$\text{So that :} \qquad V_{CM}/V_1 = V_{CM}/U_1 = m_1/m_2$$

Relationship between center of mass and laboratory frames of reference – continued (elastic scattering)



$$\mathbf{v}_1 = \mathbf{V}_1 + \mathbf{V}_{CM}$$

$$v_1 \sin\theta = V_1 \sin\psi$$

$$v_1 \cos\theta = V_1 \cos\psi + V_{CM}$$

$$\tan\theta = \frac{\sin\psi}{\cos\psi + V_{CM}/V_1} = \frac{\sin\psi}{\cos\psi + m_1/m_2}$$

Also: $$\cos\theta = \frac{\cos\psi + m_1/m_2}{\sqrt{1 + 2m_1/m_2\cos\psi + \left(m_1/m_2\right)^2}}$$

Differential cross sections in different reference frames

$$\left(\frac{d\sigma_{LAB}(\theta)}{d\Omega_{LAB}}\right)=\left(\frac{d\sigma_{CM}(\psi)}{d\Omega_{CM}}\right)\frac{d\Omega_{CM}}{d\Omega_{LAB}}$$

$$\frac{d\Omega_{CM}}{d\Omega_{LAB}}=\left|\frac{\sin\psi}{\sin\theta}\frac{d\psi}{d\theta}\right|=\left|\frac{d\cos\psi}{d\cos\theta}\right|$$

Using:

$$\cos\theta=\frac{\cos\psi+m_1/m_2}{\sqrt{1+2(m_1/m_2)\cos\psi+(m_1/m_2)^2}}$$

$$\left|\frac{d\cos\theta}{d\cos\psi}\right|=\frac{(m_1/m_2)\cos\psi+1}{\left(1+2(m_1/m_2)\cos\psi+(m_1/m_2)^2\right)^{3/2}}$$

Differential cross sections in different reference frames – continued:

$$\left(\frac{d\sigma_{LAB}(\theta)}{d\Omega_{LAB}}\right)=\left(\frac{d\sigma_{CM}(\psi)}{d\Omega_{CM}}\right)\left|\frac{d\cos\psi}{d\cos\theta}\right|$$

$$\left(\frac{d\sigma_{LAB}(\theta)}{d\Omega_{LAB}}\right)=\left(\frac{d\sigma_{CM}(\psi)}{d\Omega_{CM}}\right)\frac{\left(1+2m_1/m_2\cos\psi+\left(m_1/m_2\right)^2\right)^{3/2}}{\left(m_1/m_2\right)\cos\psi+1}$$

where: $\tan\theta=\dfrac{\sin\psi}{\cos\psi+m_1/m_2}$

$$\left(\frac{d\sigma_{LAB}(\theta)}{d\Omega_{LAB}}\right)=\left(\frac{d\sigma_{CM}(\psi)}{d\Omega_{CM}}\right)\frac{\left(1+2m_1/m_2\cos\psi+(m_1/m_2)^2\right)^{3/2}}{(m_1/m_2)\cos\psi+1}$$

where: $$\tan\theta=\frac{\sin\psi}{\cos\psi+m_1/m_2}$$

Example: suppose $m_1 = m_2$

In this case: $$\tan\theta=\frac{\sin\psi}{\cos\psi+1} \quad \Rightarrow \theta=\frac{\psi}{2}$$

note that $$0\le\theta\le\frac{\pi}{2}$$

$$\left(\frac{d\sigma_{LAB}(\theta)}{d\Omega_{LAB}}\right)=\left(\frac{d\sigma_{CM}(2\theta)}{d\Omega_{CM}}\right)\cdot 4\cos\theta$$

Summary --
Differential cross sections in different reference frames – continued:

$$\left(\frac{d\sigma_{LAB}(\theta)}{d\Omega_{LAB}}\right) = \left(\frac{d\sigma_{CM}(\psi)}{d\Omega_{CM}}\right)\left|\frac{d\cos\psi}{d\cos\theta}\right|$$

$$\left(\frac{d\sigma_{LAB}(\theta)}{d\Omega_{LAB}}\right) = \left(\frac{d\sigma_{CM}(\psi)}{d\Omega_{CM}}\right)\frac{\left(1+2m_1/m_2\cos\psi+\left(m_1/m_2\right)^2\right)^{3/2}}{\left(m_1/m_2\right)\cos\psi+1}$$

where: $$\tan\theta = \frac{\sin\psi}{\cos\psi + m_1/m_2}$$

For elastic scattering

Hard sphere example – continued

$m_1=m_2$

Center of mass frame

Lab frame

$$\left(\frac{d\sigma_{CM}(\psi)}{d\Omega_{CM}}\right)=\frac{D^2}{4}$$

$$\left(\frac{d\sigma_{LAB}(\theta)}{d\Omega_{LAB}}\right)=D^2\cos\theta \quad \theta=\frac{\psi}{2}$$

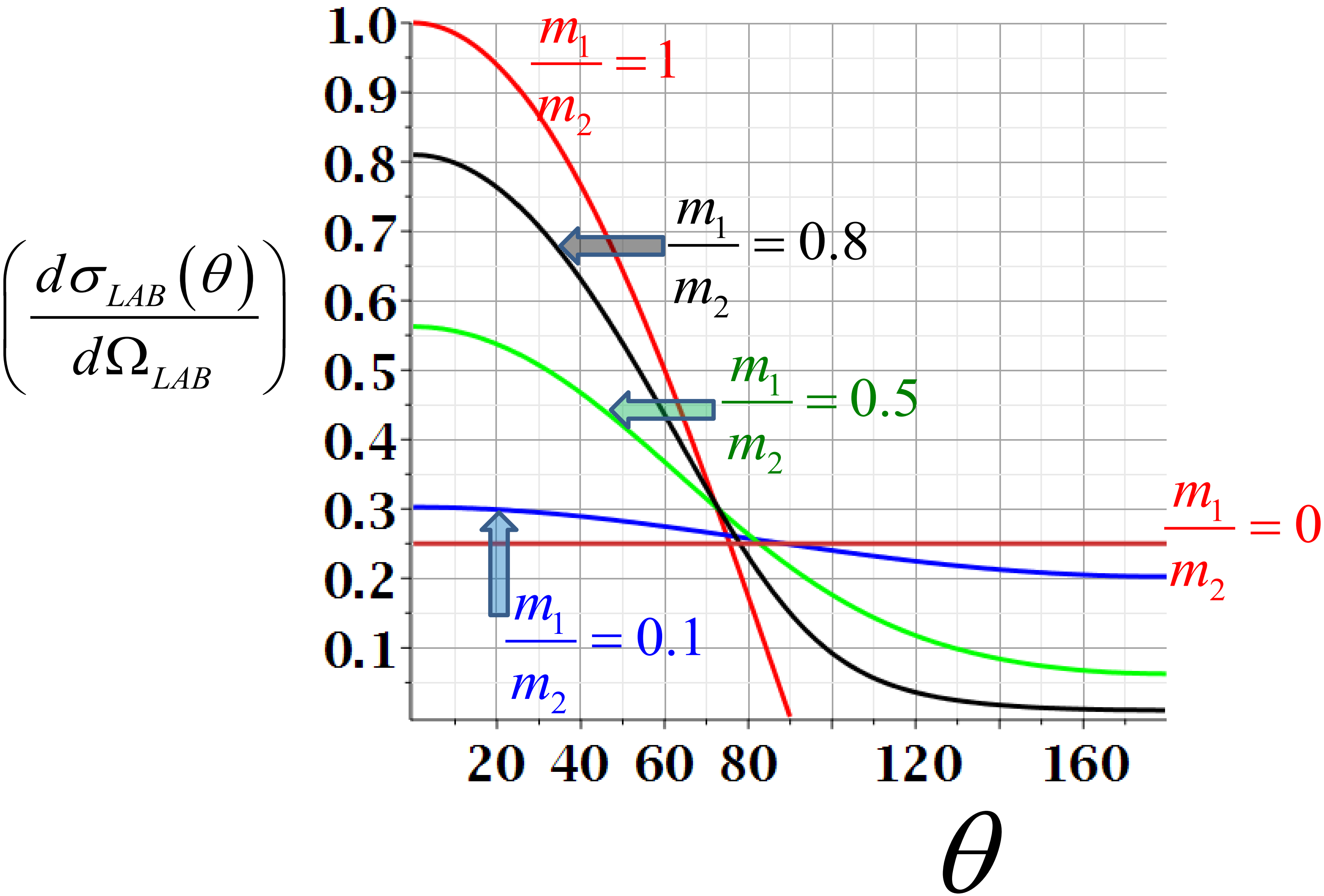
$$\int\frac{d\sigma_{CM}(\psi)}{d\Omega_{CM}}d\Omega_{CM}=$$

$$\frac{D^2}{4}4\pi=\pi D^2$$

$$\int\frac{d\sigma_{lab}(\theta)}{d\Omega_{lab}}d\Omega_{lab}=$$

$$2\pi D^2\int_0^{\pi/2}\cos\theta\,\sin\theta d\theta=\pi D^2$$

Scattering cross section for hard sphere in lab frame for various mass ratios:



For visualization, is convenient to make a "parametric" plot of

$$\left(\frac{d\sigma_{LAB}}{d\Omega}(\theta)\right) \text{ vs } \theta(\psi)$$

$$\left(\frac{d\sigma_{LAB}(\theta)}{d\Omega_{LAB}}\right) = \left(\frac{d\sigma_{CM}(\psi)}{d\Omega_{CM}}\right)\frac{\left(1+2m_1/m_2\cos\psi+(m_1/m_2)^2\right)^{3/2}}{(m_1/m_2)\cos\psi+1}$$

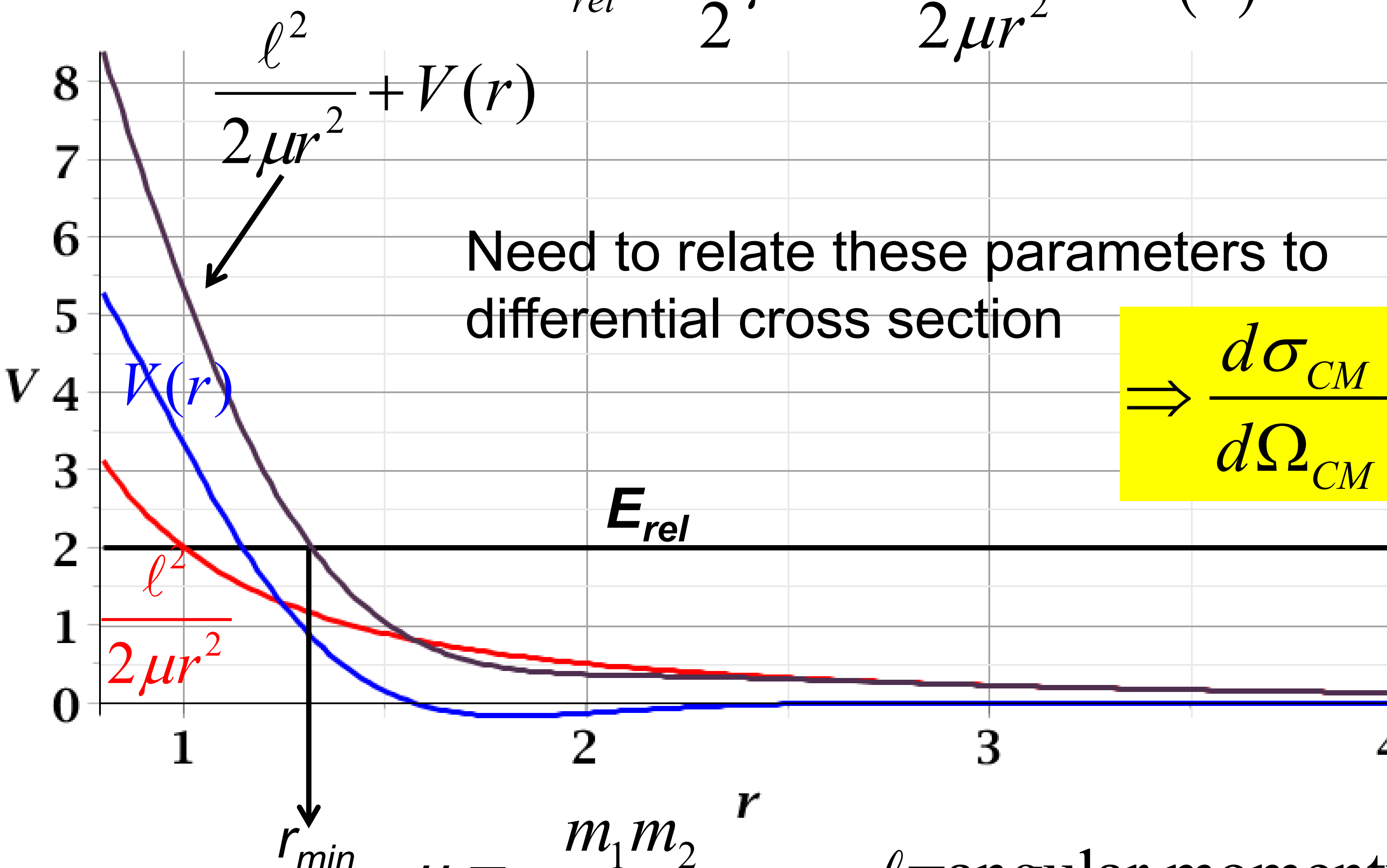
where: $\tan\theta = \dfrac{\sin\psi}{\cos\psi + m_1/m_2}$

Maple syntax:

```
> plot({[psi(theta, 0), sigma(theta, 0), theta = 0.001 ..3.14], [psi(theta, .1), sigma(theta, .1), theta
      = 0.001 ..3.14], [psi(theta, .5), sigma(theta, .5), theta = 0.001 ..3.14], [psi(theta, .8),
      sigma(theta, .8), theta = 0.001 ..3.14], [psi(theta, 1), sigma(theta, 1), theta = 0.001 ..3.14]},
      thickness = 3, font = ['Times','bold', 24], gridlines = true, color = [red, blue, green, black,
      orange])
```

For a continuous potential interaction in center of mass reference frame:

$$E_{rel} = \frac{1}{2}\mu\dot{r}^2 + \frac{\ell^2}{2\mu r^2} + V(r)$$



Need to relate these parameters to differential cross section

$$\Rightarrow \frac{d\sigma_{CM}}{d\Omega_{CM}}$$

$$\mu \equiv \frac{m_1 m_2}{m_1 + m_2}$$

ℓ=angular momentum

Focusing on the center of mass frame of reference:

Typical two-particle interactions –

Central potential: $$V(\mathbf{r}_1 - \mathbf{r}_2) = V(|\mathbf{r}_1 - \mathbf{r}_2|) \equiv V(r)$$

Hard sphere: $$V(r) = \begin{cases} \infty & r \leq D \\ 0 & r > D \end{cases}$$ X

Coulomb or gravitational: $$V(r) = \frac{K}{r}$$ ←

Lennard-Jones: $$V(r) = \frac{A}{r^{12}} - \frac{B}{r^{6}}$$