PHY 711 Classical Mechanics and Mathematical Methods 10-10:50 AM MWF in Olin 103

Discussion on Lecture 20 – Review of Chap. 1-7 (except 5 in F&W)

15						
	4	Mon, 8/30/2021	Chap. 1	Scattering theory	<u>#3</u>	9/01/2021
	5	Wed, 9/01/2021	Chap. 1	Summary of scattering theory	<u>#4</u>	9/03/2021
	6	Fri, 9/03/2021	Chap. 2	Non-inertial coordinate systems	<u>#5</u>	9/06/2021
	7	Mon, 9/06/2021	Chap. 3	Calculus of Variation	<u>#6</u>	9/10/2021
	8	Wed, 9/08/2021	Chap. 3	Calculus of Variation		
	9	Fri, 9/10/2021	Chap. 3 & 6	Lagrangian Mechanics	<u>#7</u>	9/13/2021
	10	Mon, 9/13/2021	Chap. 3 & 6	Lagrangian Mechanics	<u>#8</u>	9/17/2021
	11	Wed, 9/15/2021	Chap. 3 & 6	Constants of the motion		
	12	Fri, 9/17/2021	Chap. 3 & 6	Hamiltonian equations of motion	<u>#9</u>	9/20/2021
	13	Mon, 9/20/2021	Chap. 3 & 6	Liouville theorm	<u>#10</u>	9/22/2021
	14	Wed, 9/22/2021	Chap. 3 & 6	Canonical transformations		
	15	Fri, 9/24/2021	Chap. 4	Small oscillations about equilibrium	<u>#11</u>	9/27/2021
	16	Mon, 9/27/2021	Chap. 4	Normal modes of vibration	<u>#12</u>	9/29/2021
	17	Wed, 9/29/2021	Chap. 4	Normal modes of more complicated systems	<u>#13</u>	10/04/2021
	18	Fri, 10/01/2021	Chap. 7	Motion of strings	<u>#14</u>	10/06/2021
	19	Mon, 10/04/2021	Chap. 7	Sturm-Liouville equations		
	20	Wed, 10/06/2021	Chap.1-7	Review		
		Fri, 10/08/2021	No class	Fall break		
		Mon, 10/11/2021	No class	Take home exam		
		Wed, 10/13/2021	No class	Take home exam		
	21	Fri, 10/15/2021	Chap. 7	Sturm-Liouville equations exam due		
1.5						

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Comments on exam

- The purpose of the exam is to help with your understanding of the material
- In accordance with the honor code, the solutions you hand in must be your own work. That is, if you have any questions, please consult with me, but no one else.
- You will get credit for the reasoning and derivations as well as for the right answer, including Mathematica, Maple, etc work sheets.
- This is an open "book" exam which means that you can consult your textbook and lecture notes as long as you cite them. (Of course, if you find a source that works out the same problem, hopefully you will refrain from looking at that...).
- It is often helpful approach problems in more than one way – recalling that undergraduate physics is still true.

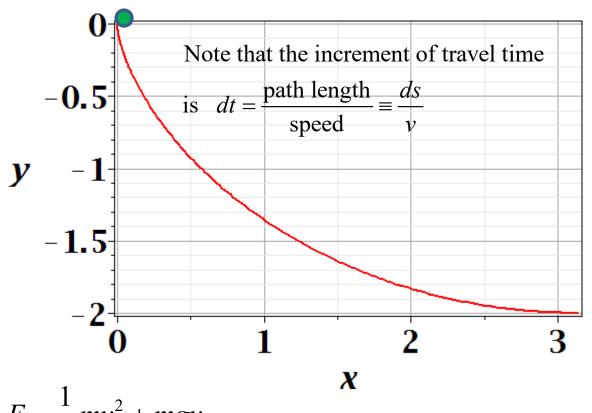
Your suggestions – From Can – Pertubation expansions and how to use them

From Ramesh -- The Brachistochrone problem

From Natalie – Some homework problems

Brachistochrone problem: (solved by Newton in 1696)

http://mathworld.wolfram.com/BrachistochroneProblem.html

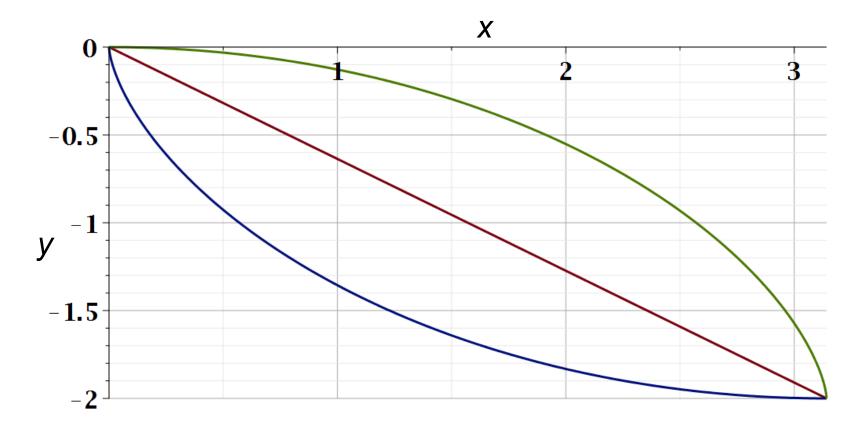


A particle of weight mg travels frictionlessly down a path of shape y(x). What is the shape of the path y(x) that minimizes the travel time from y(0)=0 to $y(\pi)=-2$?

$$E = \frac{1}{2}mv^2 + mgy$$

With the choice of initial conditions, E = 0

Vote for your favorite path



Which gives the shortest time?

- a. Green
- b. Red
- c. Blue

$$T = \int_{x_i y_i}^{x_f y_f} \frac{ds}{v} = \int_{x_i}^{x_f} \frac{\sqrt{1 + \left(\frac{dy}{dx}\right)^2}}{\sqrt{-2gy}} dx \quad \text{because} \quad \frac{1}{2}mv^2 = -mgy$$

$$f\left(\left\{y(x), \frac{dy}{dx}\right\}, x\right) = \sqrt{\frac{1 + \left(\frac{dy}{dx}\right)^2}{-y}}$$
 Note that for the original form of Euler-Lagrange equation:

$$\frac{d}{dx}\left(f - \frac{\partial f}{\partial(dy/dx)}\frac{dy}{dx}\right) = 0$$

$$\frac{d}{dx} \left[\frac{1}{\sqrt{-y \left(1 + \left(\frac{dy}{dx} \right)^2 \right)}} \right] = 0$$

$$\left(\frac{\partial f}{\partial y}\right)_{x,\frac{dy}{dx}} - \frac{d}{dx} \left[\left(\frac{\partial f}{\partial (dy/dx)}\right)_{x,y} \right] = 0,$$

differential equation is more complicated:

$$-\frac{1}{2}\sqrt{\frac{1+\left(\frac{dy}{dx}\right)^2}{-y^3}} - \frac{d}{dx}\left(\frac{\frac{dy}{dx}}{\sqrt{-y\left(1+\left(\frac{dy}{dx}\right)^2\right)}}\right) = 0$$

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$$f\left(\left\{y(x), \frac{dy}{dx}\right\}, x\right) = \sqrt{\frac{1 + \left(\frac{dy}{dx}\right)^2}{-y}}$$

$$\frac{d}{dx}\left(f - \frac{\partial f}{\partial (dy/dx)} \frac{dy}{dx}\right) = \left(\frac{\partial f}{\partial x}\right)$$

$$\Rightarrow \frac{d}{dx} \left(\frac{1}{\sqrt{-y \left(1 + \left(\frac{dy}{dx} \right)^2 \right)}} \right) = 0$$

$$\frac{dy}{dy^2} = 0 - y \left(1 + \left(\frac{dy}{dx} \right)^2 \right) = K \equiv 2a$$

Question – why this choice? Answer – because the answer will be more beautiful. (Be sure that was not my cleverness.)

$$-y\left(1+\left(\frac{dy}{dx}\right)^{2}\right) = K \equiv 2a$$

$$\frac{dy}{dx} = -\sqrt{\frac{2a}{-y}} - 1$$

$$-\frac{dy}{\sqrt{\frac{2a}{-1}}} = dx$$

Let
$$y = -2a\sin^2\frac{\theta}{2} = a(\cos\theta - 1)$$

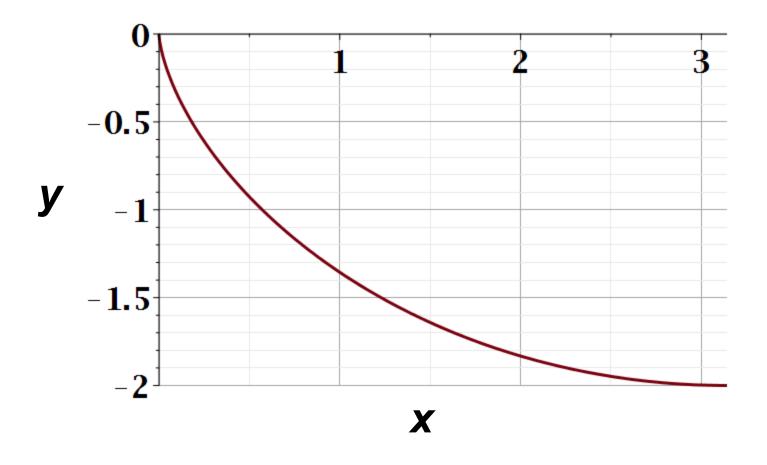
$$-\frac{dy}{\sqrt{\frac{2a}{-y} - 1}} = \frac{2a\sin\frac{\theta}{2}\cos\frac{\theta}{2}d\theta}{\sqrt{\frac{2a}{2a\sin^2\frac{\theta}{2}} - 1}} = dx$$

$$x = \int_0^\theta a(1 - \cos\theta')d\theta' = a(\theta - \sin\theta)$$

Parametric equations for Brachistochrone:

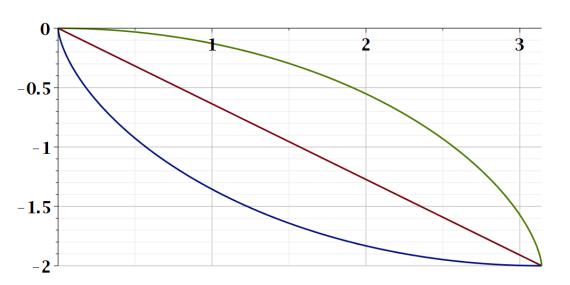
$$x = a(\theta - \sin \theta)$$
$$y = a(\cos \theta - 1)$$

Parametric plot -plot([theta-sin(theta), cos(theta)-1, theta = 0 .. Pi])



Checking the results

$$T = \int_{x_i y_i}^{x_f y_f} \frac{ds}{v} = \int_{x_i}^{x_f} \frac{\sqrt{1 + \left(\frac{dy}{dx}\right)^2}}{\sqrt{-2gy}} dx$$



(units of
$$\frac{1}{\sqrt{(2g)}}$$
)

Some comments on perturbation theory –

Perturbation theory methods are useful for solving differential equations when some of the contributions are well known but other "smaller" contributions are also present. The concept is used in many contexts although the details vary.

An example --

$$\frac{d^2y(t)}{dt^2} = -\omega^2 y(t) + \epsilon (y(t))^2$$
 where ϵ is small in some measure

Perturbation theory expansion:

Suppose:
$$y(t) = y_0(t) + \epsilon y_1(t) + \epsilon^2 y_2(t) ...$$

$$\frac{d^{2}\left(y_{0}(t) + \epsilon y_{1}(t) + \epsilon^{2} y_{2}(t)...\right)}{dt^{2}} = -\omega^{2}\left(y_{0}(t) + \epsilon y_{1}(t) + \epsilon^{2} y_{2}(t)...\right) + \epsilon\left(\left(y_{0}(t) + \epsilon y_{1}(t) + \epsilon^{2} y_{2}(t)...\right)\right)^{2}$$

Collecting terms by power of ϵ :

$$\epsilon^0: \frac{d^2y_0(t)}{dt^2} = -\omega^2y_0(t)$$

$$\rightarrow$$
 solve for $y_0(t)$

$$\epsilon^{1}: \frac{d^{2}y_{1}(t)}{dt^{2}} = -\omega^{2}y_{1}(t) + (y_{0}(t))^{2}$$
 solve for $y_{1}(t)$

$$\rightarrow$$
 solve for $y_1(t)$

PHY 711 – Assignment #8

September 13, 2021

The material for this exercise is covered in the lecture notes and in Chapters 3 and 6 of Fetter and Walecka.

- 1. A particle of mass m and charge q is subjected to a vector potential A(r,t) = -(E₀ct + B₀x) \(\hat{z}\). (Note that we are using the cgs Gaussian units of your text book.) Here E₀ denotes a constant electric field amplitude and B₀ denotes a constant magnetic field amplitude. The initial particle position is r(0) = 0 and the initial particle velocity is \(\hat{r}(0) = 0\).
 - (a) Determine the Lagrangian $L(x, y, z, \dot{x}, \dot{y}, \dot{z}, t)$ which describes the particle's motion.
 - (b) Write the Euler-Lagrange equations for this system.
 - (c) Find and evaluate the constants of motion for this system.
 - (d) Find the particle trajectories x(t), y(t), z(t) by solving the equations and imposing the given initial conditions.

Steps for tackling a problem –

- 1. What are the basic concepts the apply to this problem.
- 2. Write down the the fundamental equations
- 3. Solve
- 4. Check.

In this case, we expect that we should use the Lagrangian formalism and thus we need to know how to represent electric and magnetic fields in the Lagrangian.

$$L(\mathbf{r}(t), \dot{\mathbf{r}}(t), t) = L_{mech}(\mathbf{r}(t), \dot{\mathbf{r}}(t), t) - q \left(\Phi(\mathbf{r}, t) - \frac{1}{c} \mathbf{A}(\mathbf{r}, t) \right)$$

In this example:
$$\Phi(\mathbf{r},t) = 0$$

$$\mathbf{A}(\mathbf{r},t) = -\hat{\mathbf{z}} \left(E_0 c t + B_0 x \right)$$

$$L(\mathbf{r}(t), \dot{\mathbf{r}}(t), t) = L_{mech}(\mathbf{r}(t), \dot{\mathbf{r}}(t), t) - q \left(\Phi(\mathbf{r}, t) - \frac{1}{c} \dot{\mathbf{r}} \cdot \mathbf{A}(\mathbf{r}, t) \right)$$

In this example: $\Phi(\mathbf{r},t) = 0$

$$\mathbf{A}(\mathbf{r},t) = -\hat{\mathbf{z}} \left(E_0 c t + B_0 x \right)$$

Note that this corresponds to an electric field:

$$\mathbf{E}(\mathbf{r},t) = -\frac{1}{c} \frac{\partial \mathbf{A}(\mathbf{r},t)}{\partial t} = E_0 \hat{\mathbf{z}}$$

and a magnetic field:

$$\mathbf{B}(\mathbf{r},t) = \nabla \times \mathbf{A}(\mathbf{r},t) = B_0 \hat{\mathbf{y}}$$

$$L(\mathbf{r}(t), \dot{\mathbf{r}}(t), t) = \frac{1}{2} m \left| \dot{\mathbf{r}}(t) \right|^2 - \frac{q\dot{z}}{c} \left(E_0 c t + B_0 x \right)$$

$$L(\mathbf{r}(t), \dot{\mathbf{r}}(t), t) = \frac{1}{2} m \left| \dot{\mathbf{r}}(t) \right|^2 - \frac{q\dot{z}}{c} \left(E_0 c t + B_0 x \right)$$

Digression on forming the Hamiltonian for this case:

$$\begin{split} p_{x} &= \frac{\partial L}{\partial \dot{x}} = m\dot{x} \\ p_{y} &= \frac{\partial L}{\partial \dot{y}} = m\dot{y} \\ p_{z} &= \frac{\partial L}{\partial \dot{z}} = m\dot{z} - \frac{q}{c} \left(E_{0}ct + B_{0}x \right) \\ H &= p_{x}\dot{x} + p_{y}\dot{y} + p_{z}\dot{z} - L = \frac{1}{2}m \left| \dot{\mathbf{r}}(t) \right|^{2} \end{split}$$

Is this correct?

$$H(\mathbf{r}(t), \mathbf{p}(t), t) = \frac{1}{2m} p_x^2 + \frac{1}{2m} p_y^2 + \frac{1}{2m} \left(p_z + qE_0 t + \frac{q}{c} B_0 x \right)^2$$