



# **PHY 711 Classical Mechanics and Mathematical Methods**

**10-10:50 AM MWF in Olin 103**

## **Discussion notes for Lecture 3**

**Scattering theory – Coordinate frames**  
**Center of mass and laboratory frames**

# Your questions –



# PHY 711 Classical Mechanics and Mathematical Methods

MWF 10 AM-10:50 AM | OPL 103 | <http://www.wfu.edu/~natalie/f21phy711/>

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## Course schedule

(Preliminary schedule -- subject to frequent adjustment.)

	Date	F&W Reading	Topic	Assignment	Due
1	Mon, 8/23/2021	Chap. 1	Introduction	<a href="#">#1</a>	8/27/2021
2	Wed, 8/25/2021	Chap. 1	Scattering theory	<a href="#">#2</a>	8/30/2021
3	Fri, 8/27/2021	Chap. 1	Scattering theory		
4	Mon, 8/30/2021	Chap. 1	Scattering theory		
5	Wed, 9/01/2021	Chap. 1	Scattering theory		
6	Fri, 9/03/2021	Chap. 2	Non-inertial coordinate systems		

## Clarification on homework #2

The total energy of the system is a given constant and  $V_0$  happens to a particular energy value (that hopefully makes the algebra of the problem somewhat convenient.).

### PHY 711 – Assignment #2

**Assume particle M is at rest.**

08/25/2021

1. Consider a particle of mass  $m$  moving in the vicinity of another particle of mass  $M$  where  $m \ll M$ . The particles interact with a conservative central potential of the form

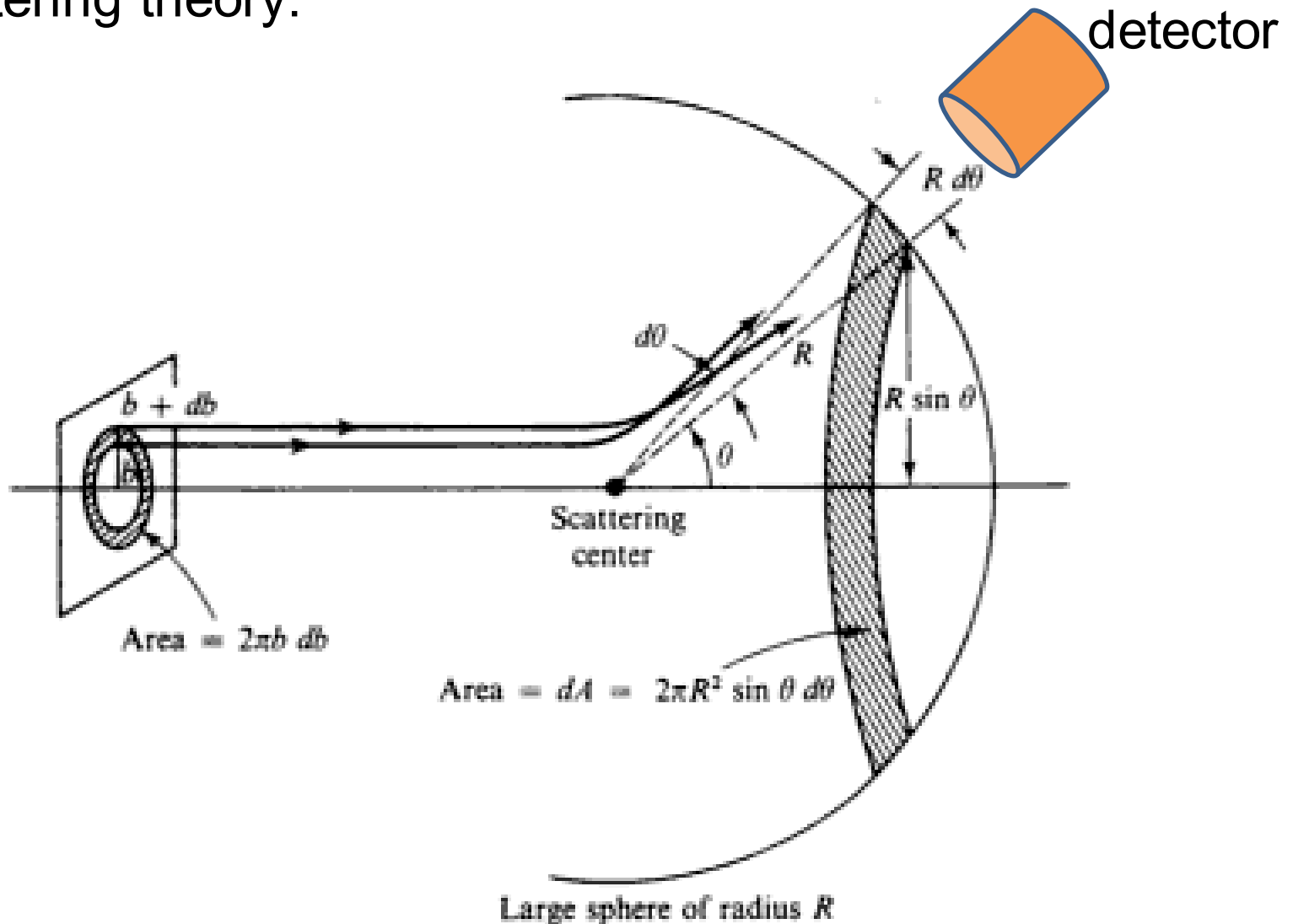
$$V(r) = V_0 \left( \left( \frac{r_0}{r} \right)^2 - \left( \frac{r_0}{r} \right) \right),$$

where  $r$  denotes the magnitude of the particle separation and  $V_0$  and  $r_0$  denote energy and length constants, respectively. The total energy of the system is  $V_0$ .

- (a) First consider the case where the impact parameter  $b = 0$ . Find the distance of closest approach of the particles.

- (b) Now consider the case where the impact parameter  $b = r_0$ . Find the distance of closest approach of the particles.

# Scattering theory:



**Figure 5.5** The scattering problem and relation of cross section to impact parameter.

Can you think of examples of such an experimental setup?

Other experimental designs –

At CERN <https://home.cern/science/experiments/totem> the study of highly energetic proton-proton scattering is designed in the center of mass frame of reference by accelerating two proton beams focused to collide head on in the Large Hadron Collider LHC facility.

Figure from CERN website

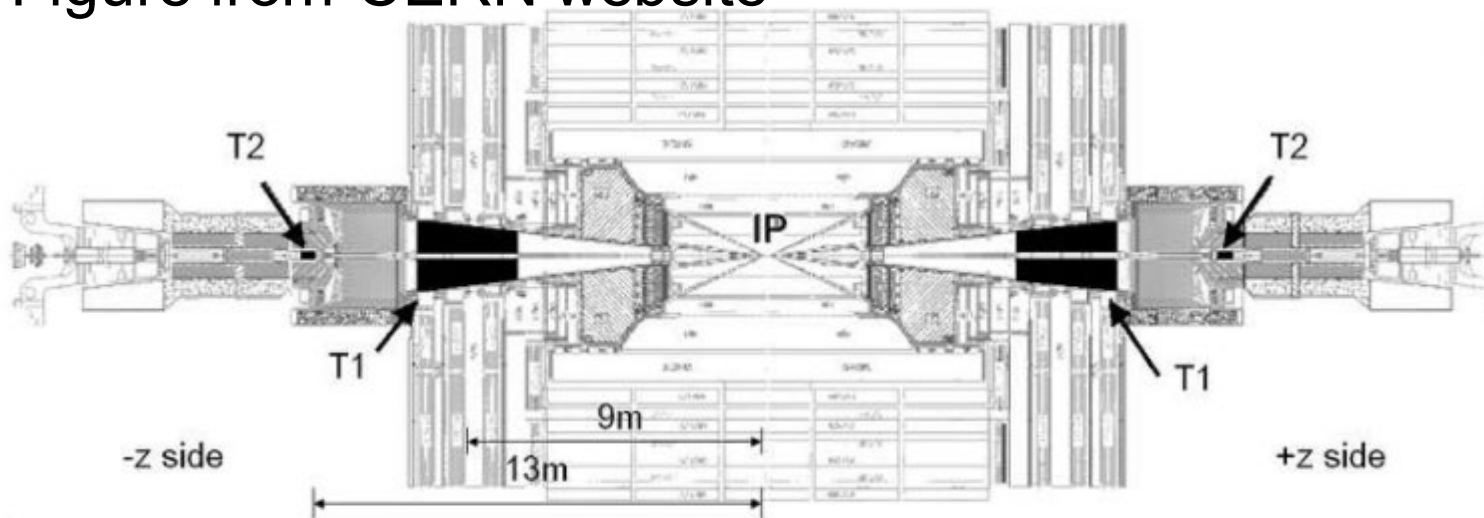
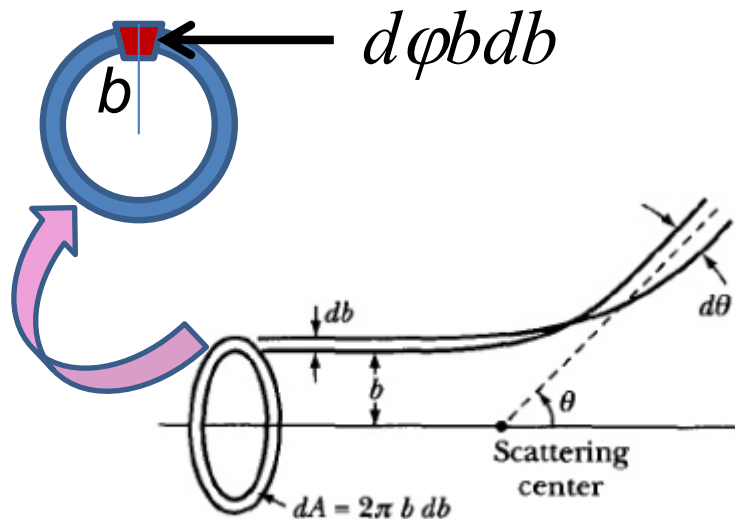


Figure 1.17: View of the inelastic forward trackers T1 and T2 inside the CMS detector.

## Differential cross section

$$\left( \frac{d\sigma}{d\Omega} \right) = \frac{\text{Number of detected particles at } \theta \text{ per target particle}}{\text{Number of incident particles per unit area}}$$

= Area of incident beam that is scattered into detector  
at angle  $\theta$

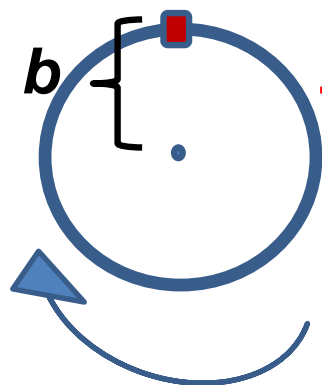


$$\left( \frac{d\sigma}{d\Omega} \right) = \frac{d\phi b db}{d\phi \sin\theta d\theta} = \frac{b}{\sin\theta} \left| \frac{db}{d\theta} \right|$$

Figure from Marion & Thorton, Classical Dynamics

More details --

We imagine that the beam of particles has a cylindrical geometry and that the physics is totally uniform in the azimuthal direction. The cross section of the beam is a circle.



← This piece of the beam scatters into the detector at angle  $\theta$

This logic leads to the notion that  $b$  is a function of  $\theta$  and we will try to find  $b(\theta)$  for various cases.

$\varphi \equiv$  azimuthal angle



**Note:** The notion of cross section is common to many areas of physics including classical mechanics, quantum mechanics, optics, etc. Only in the **classical mechanics** can we calculate it from a knowledge of the particle trajectory as it relates to the scattering geometry.

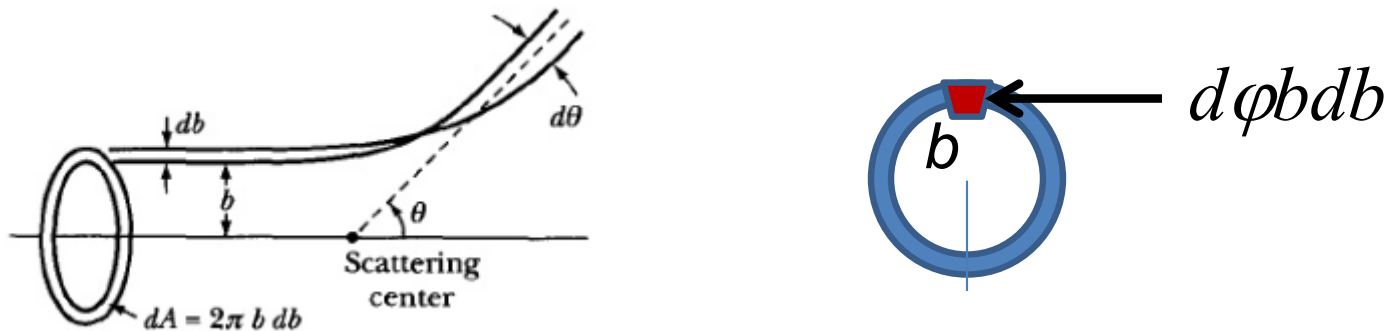
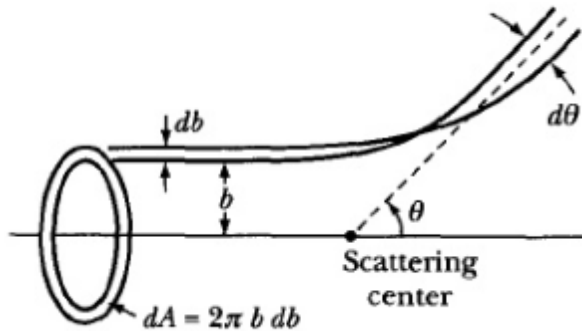


Figure from Marion & Thorton, Classical Dynamics

$$\left( \frac{d\sigma}{d\Omega} \right) = \frac{d\phi b db}{d\phi \sin\theta d\theta} = \frac{b}{\sin\theta} \left| \frac{db}{d\theta} \right|$$

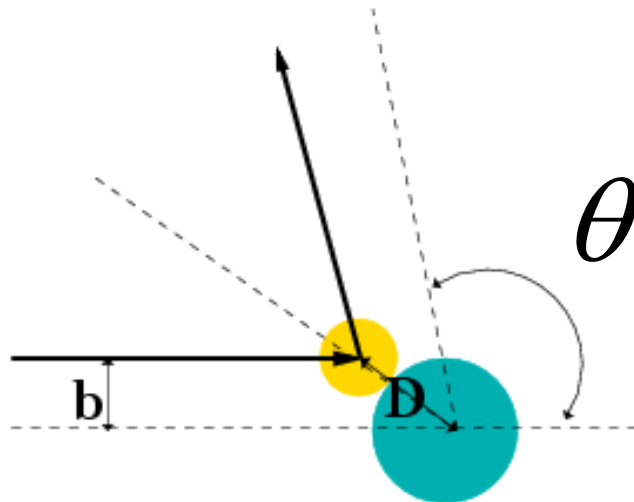
Note: We are assuming that the process is isotropic in  $\phi$

## Simple example – collision of hard spheres



$$\left( \frac{d\sigma}{d\Omega} \right) = \frac{b}{\sin\theta} \left| \frac{db}{d\theta} \right|$$

Microscopic view:

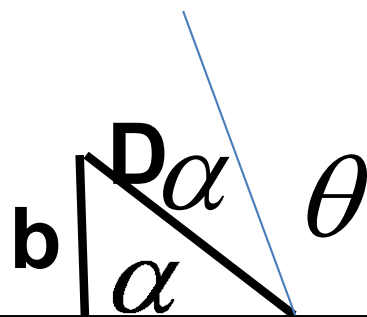
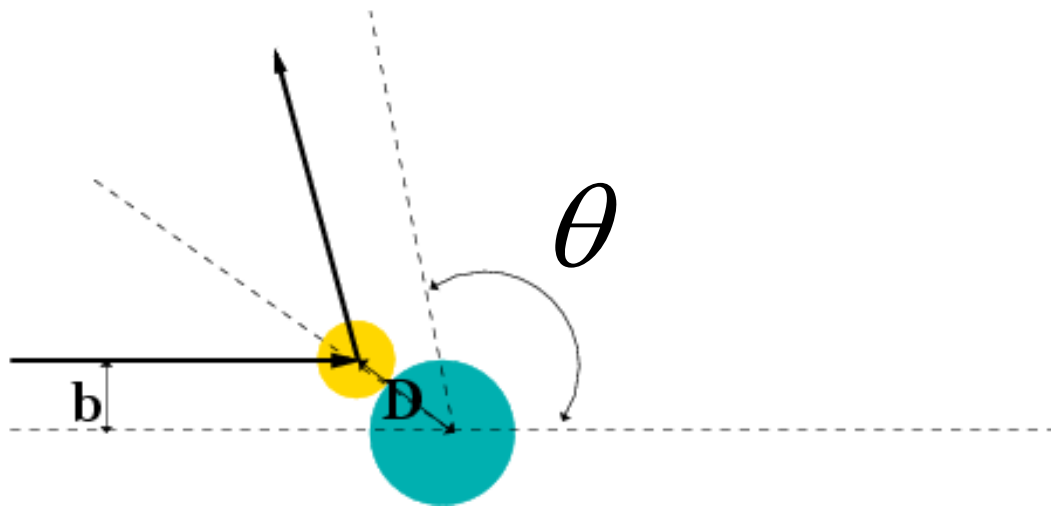


$$b(\theta) = ?$$

$$b(\theta) = D \sin\left(\frac{\pi}{2} - \frac{\theta}{2}\right)$$

$$\left( \frac{d\sigma}{d\Omega} \right) = \frac{D^2}{4}$$

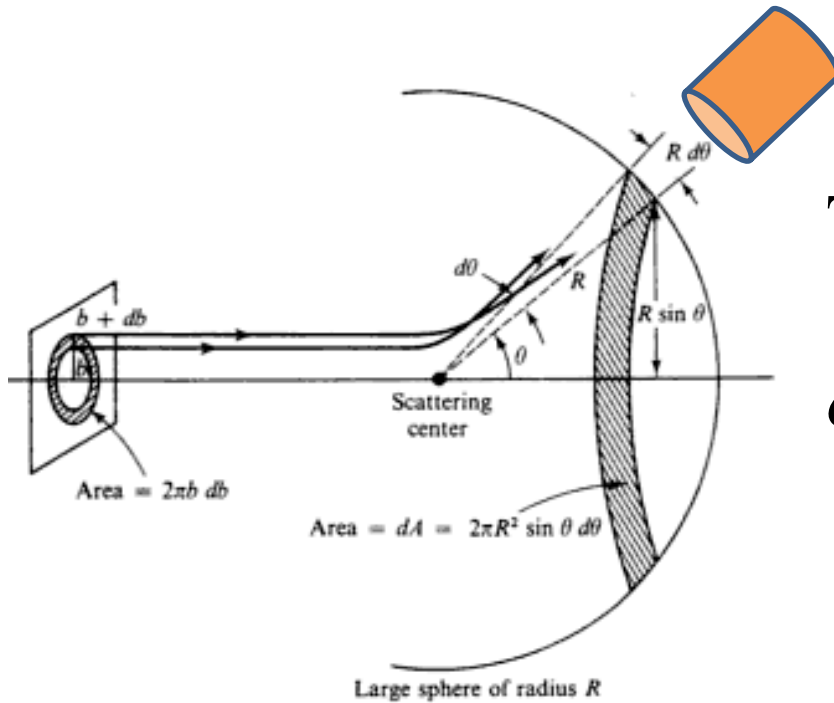
Some more details of form of  $b(\theta)$



$$b = D \sin \alpha = D \sin \left( \frac{\pi}{2} - \frac{\theta}{2} \right)$$

$$2\alpha + \theta = \pi$$

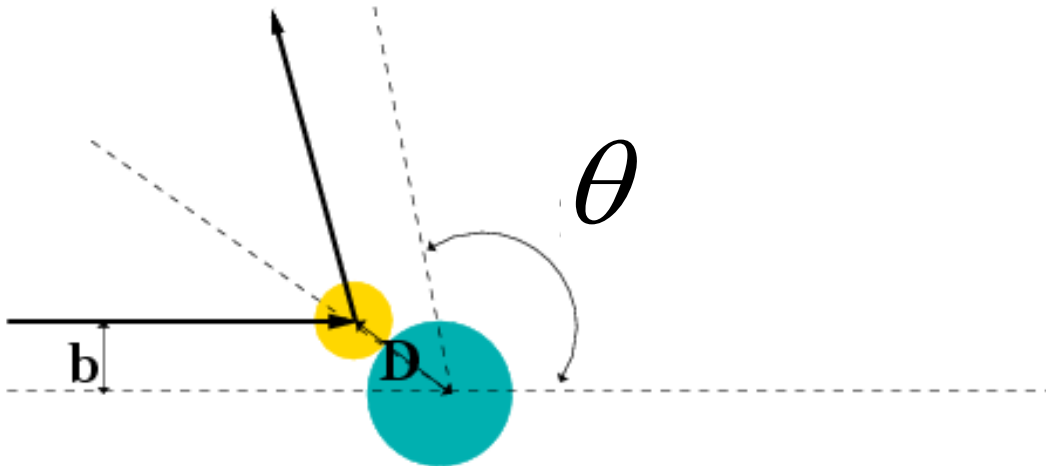
# Simple example – collision of hard spheres -- continued



Total scattering cross section:

$$\sigma = \int \left( \frac{d\sigma}{d\Omega} \right) d\Omega$$

Hard sphere:



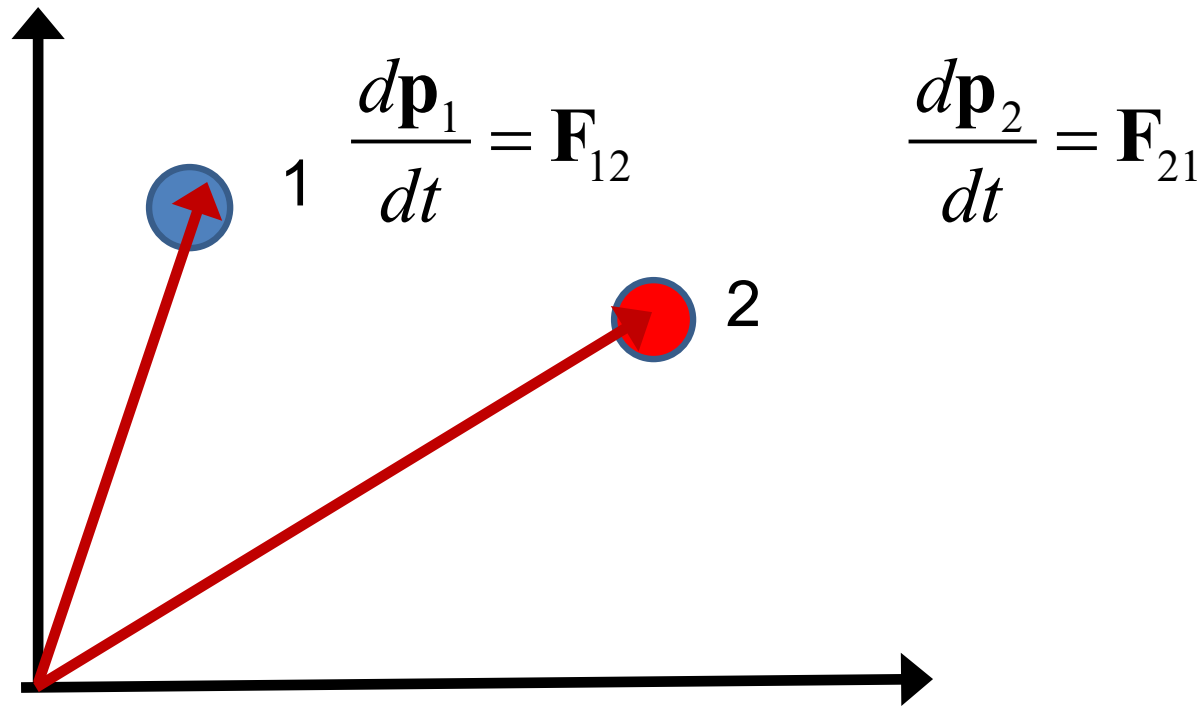
$$\left( \frac{d\sigma}{d\Omega} \right) = \frac{D^2}{4}$$

$$\sigma = \pi D^2$$

Now consider the more general case of particle interactions and the corresponding scattering analysis.

Scattering theory can help us analyze the interaction potential  $V(r)$ . First we need to simplify the number of variables.

Relationship of scattering cross-section to particle interactions --  
Classical mechanics of a conservative 2-particle system.



$$\mathbf{F}_{12} = -\nabla_1 V(\mathbf{r}_1 - \mathbf{r}_2) \quad \Rightarrow \quad E = \frac{1}{2}m_1 v_1^2 + \frac{1}{2}m_2 v_2^2 + V(\mathbf{r}_1 - \mathbf{r}_2)$$



Relationship between center of mass and laboratory frames of reference. At and time  $t$ , the following relationships apply --

Definition of center of mass  $\mathbf{R}_{CM}$

$$m_1 \mathbf{r}_1 + m_2 \mathbf{r}_2 = (m_1 + m_2) \mathbf{R}_{CM}$$

$$m_1 \dot{\mathbf{r}}_1 + m_2 \dot{\mathbf{r}}_2 = (m_1 + m_2) \dot{\mathbf{R}}_{CM} = (m_1 + m_2) \mathbf{V}_{CM}$$

Note that  $\dot{\mathbf{r}} \equiv \frac{d\mathbf{r}}{dt}$

$$E = \frac{1}{2} m_1 v_1^2 + \frac{1}{2} m_2 v_2^2 + V(\mathbf{r}_1 - \mathbf{r}_2)$$

$$= \frac{1}{2} (m_1 + m_2) V_{CM}^2 + \frac{1}{2} \mu |\mathbf{v}_1 - \mathbf{v}_2|^2 + V(\mathbf{r}_1 - \mathbf{r}_2)$$

where:  $\mu \equiv \frac{m_1 m_2}{m_1 + m_2}$



Classical mechanics of a conservative 2-particle system -- continued

$$E = \frac{1}{2}(m_1 + m_2)V_{CM}^2 + \frac{1}{2}\mu|\mathbf{v}_1 - \mathbf{v}_2|^2 + V(\mathbf{r}_1 - \mathbf{r}_2)$$

For central potentials:  $V(\mathbf{r}_1 - \mathbf{r}_2) = V(|\mathbf{r}_1 - \mathbf{r}_2|) \equiv V(r_{12})$

Relative angular momentum is also conserved:

$$\mathbf{L}_{12} \equiv \mathbf{r}_{12} \times \mu \mathbf{v}_{12}$$

$$E = \frac{1}{2}(m_1 + m_2)V_{CM}^2 + \frac{1}{2}\mu v_{12}^2 + \frac{L_{12}^2}{2\mu r_{12}^2} + V(r_{12})$$

Simpler notation:

$$E = \frac{1}{2}(m_1 + m_2)V_{CM}^2 + \frac{1}{2}\mu\dot{r}^2 + \frac{\ell^2}{2\mu r^2} + V(r)$$



Simpler notation:

$$E = \frac{1}{2}(m_1 + m_2)V_{CM}^2 + \frac{1}{2}\mu\dot{r}^2 + \frac{\ell^2}{2\mu r^2} + V(r)$$

The diagram illustrates the equation for total energy  $E$  in a two-body system. The equation is:  $E = \frac{1}{2}(m_1 + m_2)V_{CM}^2 + \frac{1}{2}\mu\dot{r}^2 + \frac{\ell^2}{2\mu r^2} + V(r)$ . Arrows indicate the time dependence of each term: 

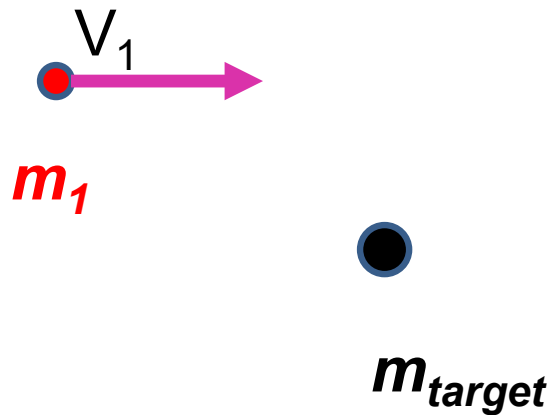
- Pink arrows point to  $(m_1 + m_2)$ ,  $V_{CM}^2$ , and  $\ell^2$ , which are labeled as "constants".
- Green arrows point to  $\dot{r}^2$ ,  $r^2$ , and  $V(r)$ , which are labeled as "vary in time".

For scattering analysis only need to know trajectory **before** and **after** the collision.

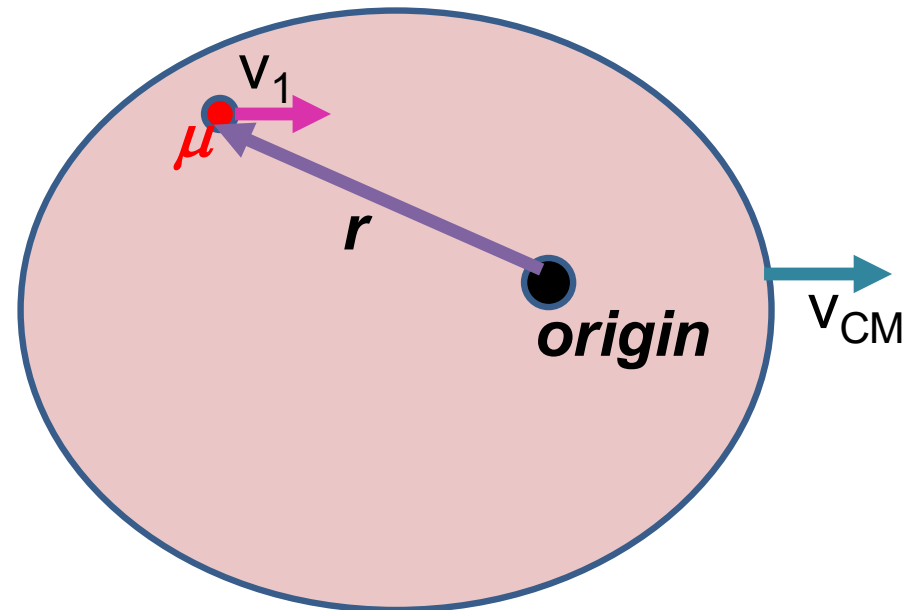
vary in time

Note: The following analysis will be carried out in the center of mass frame of reference.

In laboratory frame:



In center-of-mass frame:



Also note: We are assuming that the interaction between particle and target  $V(r)$  conserves energy and angular momentum.

$$\mu = \frac{m_1 m_{\text{target}}}{m_1 + m_{\text{target}}}$$

$$\ell = |\mathbf{r} \times \mu \mathbf{v}_1|$$

Typically, the laboratory frame is where the data is taken, but the center of mass frame is where the analysis is most straightforward.

Previous equations --

$$E = \frac{1}{2}(m_1 + m_2)V_{CM}^2 + \frac{1}{2}\mu\dot{r}^2 + \frac{\ell^2}{2\mu r^2} + V(r)$$



constant



relative coordinate system;  
visualize as “in” CM frame

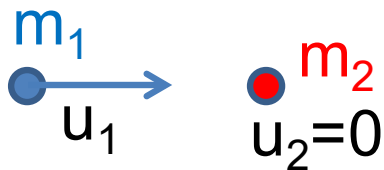


It is often convenient to analyze the scattering cross section in the center of mass reference frame.

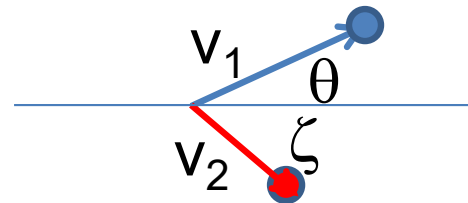
Relationship between normal laboratory reference and center of mass:

Laboratory reference frame:

Before

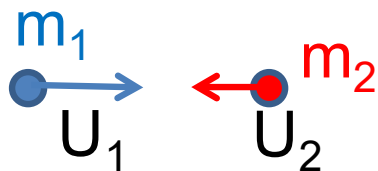


After

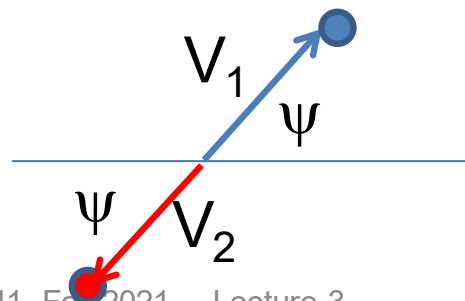


Center of mass reference frame:

Before



After





## Relationship between center of mass and laboratory frames of reference -- continued

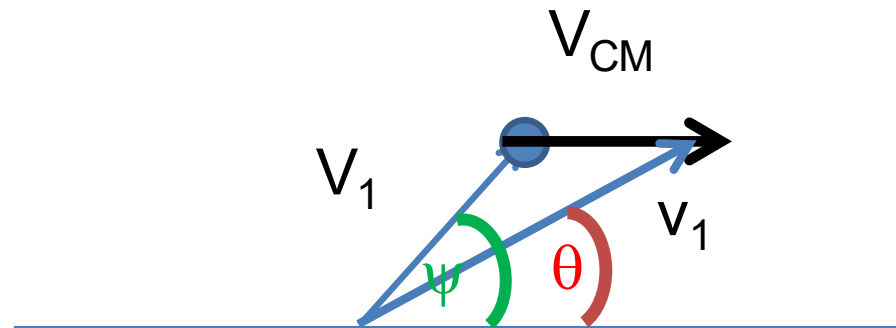
Since  $m_2$  is initially at rest :

$$\mathbf{V}_{CM} = \frac{m_1}{m_1 + m_2} \mathbf{u}_1 \quad \mathbf{u}_1 = \mathbf{U}_1 + \mathbf{V}_{CM} \quad \Rightarrow \quad \mathbf{U}_1 = \frac{m_2}{m_1 + m_2} \mathbf{u}_1 = \frac{m_2}{m_1} \mathbf{V}_{CM}$$
$$\mathbf{u}_2 = \mathbf{U}_2 + \mathbf{V}_{CM} \quad \Rightarrow \quad \mathbf{U}_2 = -\frac{m_1}{m_1 + m_2} \mathbf{u}_1 = -\mathbf{V}_{CM}$$

$$\mathbf{v}_1 = \mathbf{V}_1 + \mathbf{V}_{CM}$$

$$\mathbf{v}_2 = \mathbf{V}_2 + \mathbf{V}_{CM}$$

# Relationship between center of mass and laboratory frames of reference for the scattering particle 1



$$\mathbf{v}_1 = \mathbf{V}_1 + \mathbf{V}_{CM}$$

$$v_1 \sin \theta = V_1 \sin \psi$$

$$v_1 \cos \theta = V_1 \cos \psi + V_{CM}$$

$$\tan \theta = \frac{\sin \psi}{\cos \psi + V_{CM} / V_1} = \frac{\sin \psi}{\cos \psi + m_1 / m_2}$$

For elastic scattering

## Digression – elastic scattering

$$\begin{aligned} \frac{1}{2} m_1 U_1^2 + \frac{1}{2} m_2 U_2^2 + \frac{1}{2} (m_1 + m_2) V_{CM}^2 \\ = \frac{1}{2} m_1 V_1^2 + \frac{1}{2} m_2 V_2^2 + \frac{1}{2} (m_1 + m_2) V_{CM}^2 \end{aligned}$$

Also note:

$$m_1 \mathbf{U}_1 + m_2 \mathbf{U}_2 = 0 \qquad m_1 \mathbf{V}_1 + m_2 \mathbf{V}_2 = 0$$

$$\mathbf{U}_1 = \frac{m_2}{m_1} \mathbf{V}_{CM} \qquad \mathbf{U}_2 = -\mathbf{V}_{CM}$$

$$\Rightarrow |\mathbf{U}_1| = |\mathbf{V}_1| \quad \text{and} \quad |\mathbf{U}_2| = |\mathbf{V}_2| = |\mathbf{V}_{CM}|$$

$$\text{Also note that: } m_1 |\mathbf{U}_1| = m_2 |\mathbf{U}_2|$$

$$\text{So that: } V_{CM}/V_1 = V_{CM}/U_1 = m_1/m_2$$

## Summary of results --

$$\mathbf{v}_1 = \mathbf{V}_1 + \mathbf{V}_{CM}$$

$$v_1 \sin \theta = V_1 \sin \psi$$

$$v_1 \cos \theta = V_1 \cos \psi + V_{CM}$$

$$\tan \theta = \frac{\sin \psi}{\cos \psi + V_{CM} / V_1} = \frac{\sin \psi}{\cos \psi + m_1 / m_2}$$



General case



Special case of  
elastic scattering

For elastic scattering

$$V_{CM} / V_1 = m_1 / m_2$$



## Relationship between center of mass and laboratory frames of reference – continued (elastic scattering)

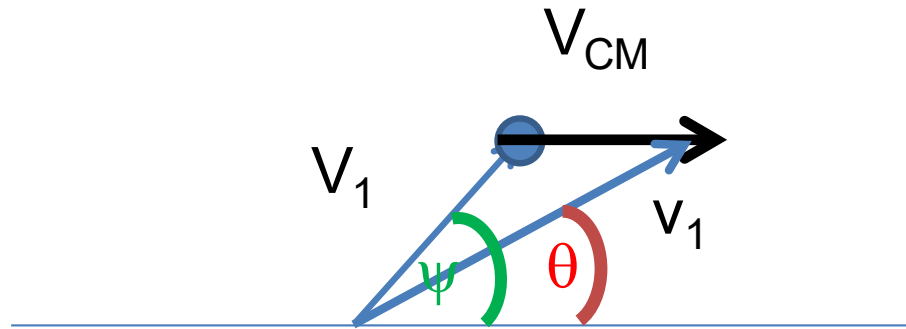
$$\mathbf{v}_1 = \mathbf{V}_1 + \mathbf{V}_{CM}$$

$$v_1 \sin \theta = V_1 \sin \psi$$

$$v_1 \cos \theta = V_1 \cos \psi + V_{CM}$$

$$\tan \theta = \frac{\sin \psi}{\cos \psi + V_{CM} / V_1} = \frac{\sin \psi}{\cos \psi + m_1 / m_2}$$

$$\text{Also: } \cos \theta = \frac{\cos \psi + m_1 / m_2}{\sqrt{1 + 2m_1 / m_2 \cos \psi + (m_1 / m_2)^2}}$$

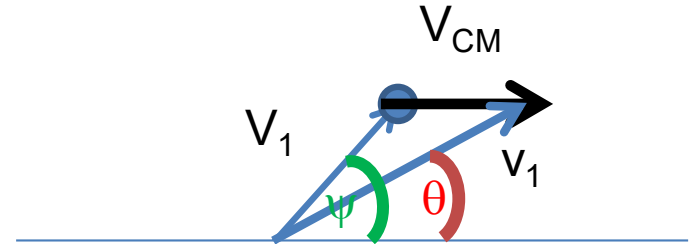


More details -- from the diagram and equations --

$$\mathbf{v}_1 = \mathbf{V}_1 + \mathbf{V}_{CM}$$

$$v_1 \sin \theta = V_1 \sin \psi$$

$$v_1 \cos \theta = V_1 \cos \psi + V_{CM}$$



Take the dot product of the first equation with itself

$$v_1^2 = V_1^2 + 2V_1V_{CM} \cos \psi + V_{CM}^2$$

$$\text{or } \frac{v_1}{V_1} = \sqrt{1 + 2 \frac{V_{CM}}{V_1} \cos \psi + \frac{V_{CM}^2}{V_1^2}} = \sqrt{1 + 2 \frac{m_1}{m_2} \cos \psi + \left( \frac{m_1}{m_2} \right)^2}$$

$$\Rightarrow \cos \theta = \frac{\cos \psi + m_1 / m_2}{\sqrt{1 + 2m_1 / m_2 \cos \psi + (m_1 / m_2)^2}}$$



## Differential cross sections in different reference frames

$$\left( \frac{d\sigma_{LAB}(\theta)}{d\Omega_{LAB}} \right) = \left( \frac{d\sigma_{CM}(\psi)}{d\Omega_{CM}} \right) \frac{d\Omega_{CM}}{d\Omega_{LAB}}$$

$$\frac{d\Omega_{CM}}{d\Omega_{LAB}} = \left| \frac{\sin \psi}{\sin \theta} \frac{d\psi}{d\theta} \right| = \left| \frac{d \cos \psi}{d \cos \theta} \right|$$

Using:

$$\cos \theta = \frac{\cos \psi + m_1 / m_2}{\sqrt{1 + 2(m_1 / m_2) \cos \psi + (m_1 / m_2)^2}}$$

$$\left| \frac{d \cos \theta}{d \cos \psi} \right| = \frac{(m_1 / m_2) \cos \psi + 1}{\left( 1 + 2(m_1 / m_2) \cos \psi + (m_1 / m_2)^2 \right)^{3/2}}$$



Differential cross sections in different reference frames – continued:

$$\left( \frac{d\sigma_{LAB}(\theta)}{d\Omega_{LAB}} \right) = \left( \frac{d\sigma_{CM}(\psi)}{d\Omega_{CM}} \right) \left| \frac{d\cos\psi}{d\cos\theta} \right|$$

$$\left( \frac{d\sigma_{LAB}(\theta)}{d\Omega_{LAB}} \right) = \left( \frac{d\sigma_{CM}(\psi)}{d\Omega_{CM}} \right) \frac{\left( 1 + 2m_1 / m_2 \cos\psi + (m_1 / m_2)^2 \right)^{3/2}}{(m_1 / m_2) \cos\psi + 1}$$

where:  $\tan\theta = \frac{\sin\psi}{\cos\psi + m_1 / m_2}$

$$\left( \frac{d\sigma_{LAB}(\theta)}{d\Omega_{LAB}} \right) = \left( \frac{d\sigma_{CM}(\psi)}{d\Omega_{CM}} \right) \frac{\left( 1 + 2m_1 / m_2 \cos \psi + (m_1 / m_2)^2 \right)^{3/2}}{(m_1 / m_2) \cos \psi + 1}$$

where:  $\tan \theta = \frac{\sin \psi}{\cos \psi + m_1 / m_2}$

Example: suppose  $m_1 = m_2$

In this case:  $\tan \theta = \frac{\sin \psi}{\cos \psi + 1} \Rightarrow \theta = \frac{\psi}{2}$

note that  $0 \leq \theta \leq \frac{\pi}{2}$

$$\left( \frac{d\sigma_{LAB}(\theta)}{d\Omega_{LAB}} \right) = \left( \frac{d\sigma_{CM}(2\theta)}{d\Omega_{CM}} \right) \cdot 4 \cos \theta$$

## Summary --

Differential cross sections in different reference frames – continued:

$$\left( \frac{d\sigma_{LAB}(\theta)}{d\Omega_{LAB}} \right) = \left( \frac{d\sigma_{CM}(\psi)}{d\Omega_{CM}} \right) \left| \frac{d\cos\psi}{d\cos\theta} \right|$$

$$\left( \frac{d\sigma_{LAB}(\theta)}{d\Omega_{LAB}} \right) = \left( \frac{d\sigma_{CM}(\psi)}{d\Omega_{CM}} \right) \frac{\left( 1 + 2m_1 / m_2 \cos\psi + (m_1 / m_2)^2 \right)^{3/2}}{(m_1 / m_2) \cos\psi + 1}$$

where:  $\tan\theta = \frac{\sin\psi}{\cos\psi + m_1 / m_2}$

For elastic scattering

## Hard sphere example – continued

$$m_1 = m_2$$

Center of mass frame

Lab frame

$$\left( \frac{d\sigma_{CM}(\psi)}{d\Omega_{CM}} \right) = \frac{D^2}{4}$$

$$\left( \frac{d\sigma_{LAB}(\theta)}{d\Omega_{LAB}} \right) = D^2 \cos \theta \quad \theta = \frac{\psi}{2}$$

$$\int \frac{d\sigma_{CM}(\psi)}{d\Omega_{CM}} d\Omega_{CM} =$$

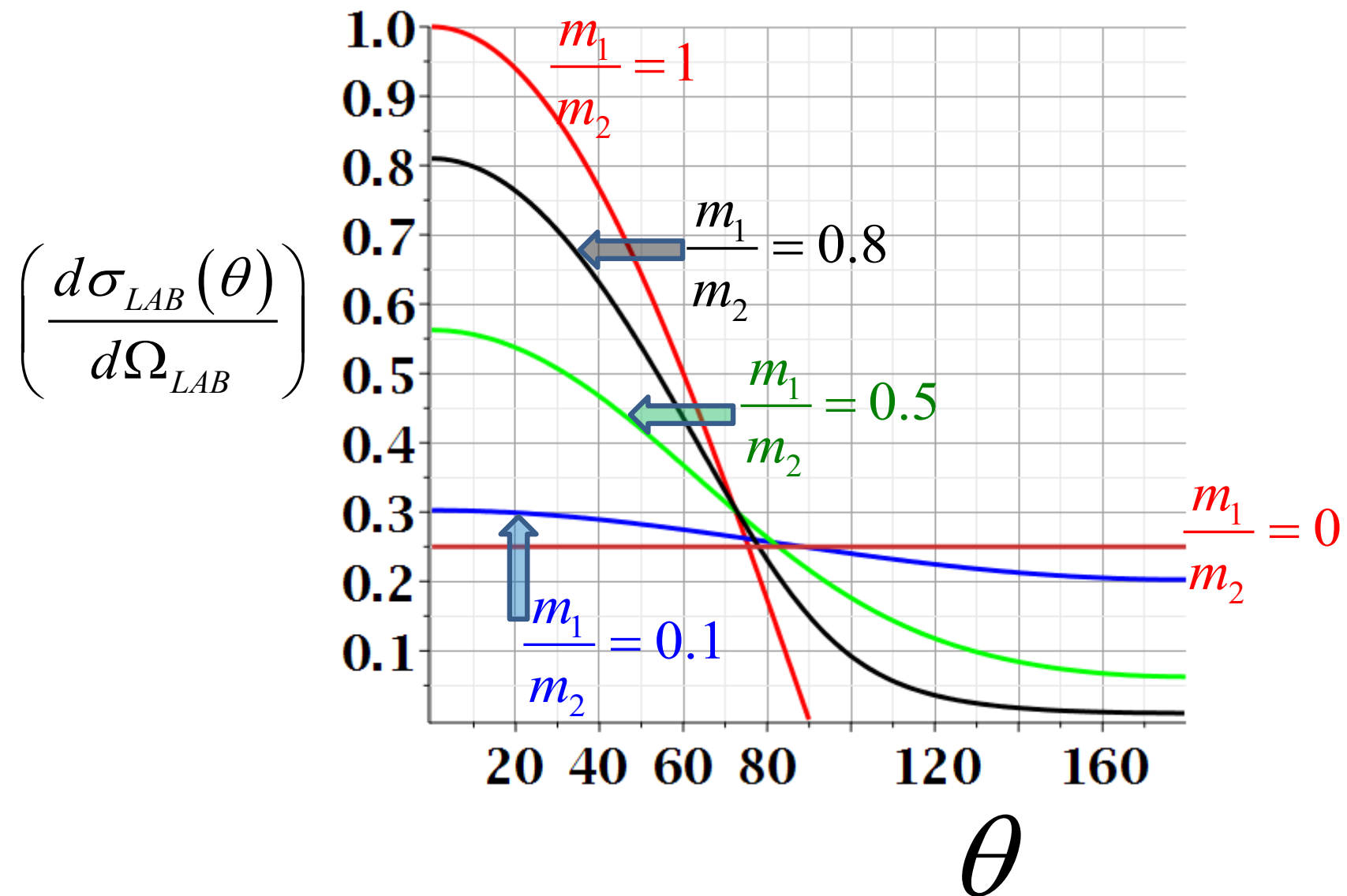
$$\frac{D^2}{4} 4\pi = \pi D^2$$

$$\int \frac{d\sigma_{lab}(\theta)}{d\Omega_{lab}} d\Omega_{lab} =$$

$$2\pi D^2 \int_0^{\pi/2} \cos \theta \sin \theta d\theta = \pi D^2$$



Scattering cross section for hard sphere in lab frame  
for various mass ratios:





For visualization, is convenient to make a "parametric" plot of

$$\left( \frac{d\sigma_{LAB}}{d\Omega}(\theta) \right) \text{ vs } \theta(\psi)$$

$$\left( \frac{d\sigma_{LAB}(\theta)}{d\Omega_{LAB}} \right) = \left( \frac{d\sigma_{CM}(\psi)}{d\Omega_{CM}} \right) \frac{\left( 1 + 2m_1 / m_2 \cos \psi + (m_1 / m_2)^2 \right)^{3/2}}{(m_1 / m_2) \cos \psi + 1}$$

where:  $\tan \theta = \frac{\sin \psi}{\cos \psi + m_1 / m_2}$

Maple syntax:

```
=
> plot( { [psi(theta, 0), sigma(theta, 0), theta = 0.001 ..3.14], [psi(theta, .1), sigma(theta, .1), theta
= 0.001 ..3.14], [psi(theta, .5), sigma(theta, .5), theta = 0.001 ..3.14], [psi(theta, .8),
sigma(theta, .8), theta = 0.001 ..3.14], [psi(theta, 1), sigma(theta, 1), theta = 0.001 ..3.14] },
thickness = 3, font = ['Times','bold', 24], gridlines = true, color = [red, blue, green, black,
orange])
```

For a continuous potential interaction in center of mass reference frame:

$$E_{rel} = \frac{1}{2} \mu \dot{r}^2 + \frac{\ell^2}{2\mu r^2} + V(r)$$

