PHY 711 Classical Mechanics and Mathematical Methods
10-10:50 AM MWF in Olin 103

Notes on Lecture 33:

Chapter 10 in F & W: Surface waves

1. **Water waves in a channel**
2. **Wave-like solutions; wave speed**

28	Mon, 10/31/2022	Chap. 9	Mechanics of 3 dimensional fluids	#21	11/02/2022
29	Wed, 11/02/2022	Chap. 9	Mechanics of 3 dimensional fluids	#22	11/04/2022
30	Fri, 11/04/2022	Chap. 9	Linearized hydrodynamics equations	#23	11/07/2022
31	Mon, 11/07/2022	Chap. 9	Linear sound waves	#24	11/09/2022
32	Wed, 11/09/2022	Chap. 9	Scattering of sound and non-linear effects	#25	11/11/2022
33	Fri, 11/11/2022	Chap. 10	Surface waves in fluids	#26	11/16/2022
34	Mon, 11/14/2022	Chap. 10	Surface waves in fluids; soliton solutions		
35	Wed, 11/16/2022	Chap. 11	Heat conduction		
36	Fri, 11/18/2022	Chap. 12	Viscous effects on hydrodynamics		
37	Mon, 11/21/2022	Chap 1-12	Review		
	Wed, 11/23/2022		Thanksgiving Holiday		
	Fri, 11/25/2022		Thanksgiving Holiday		
	Mon, 11/28/2022		Presentations I		
	Wed, 11/30/2022		Presentations II		
	Fri, 12/02/2022		Presentations III		

PHY 711 -- Assignment #26

Nov. 11, 2022

Start reading Chapter 10 in **Fetter & Walecka**.

1. Work Problem 10.3 at the end of Chapter 10 in **Fetter and Walecka**.

   Note that some of the ideas are discussed in today's lecture.

Reference: Chapter 10 of Fetter and Walecka

Physics of incompressible fluids and their surfaces

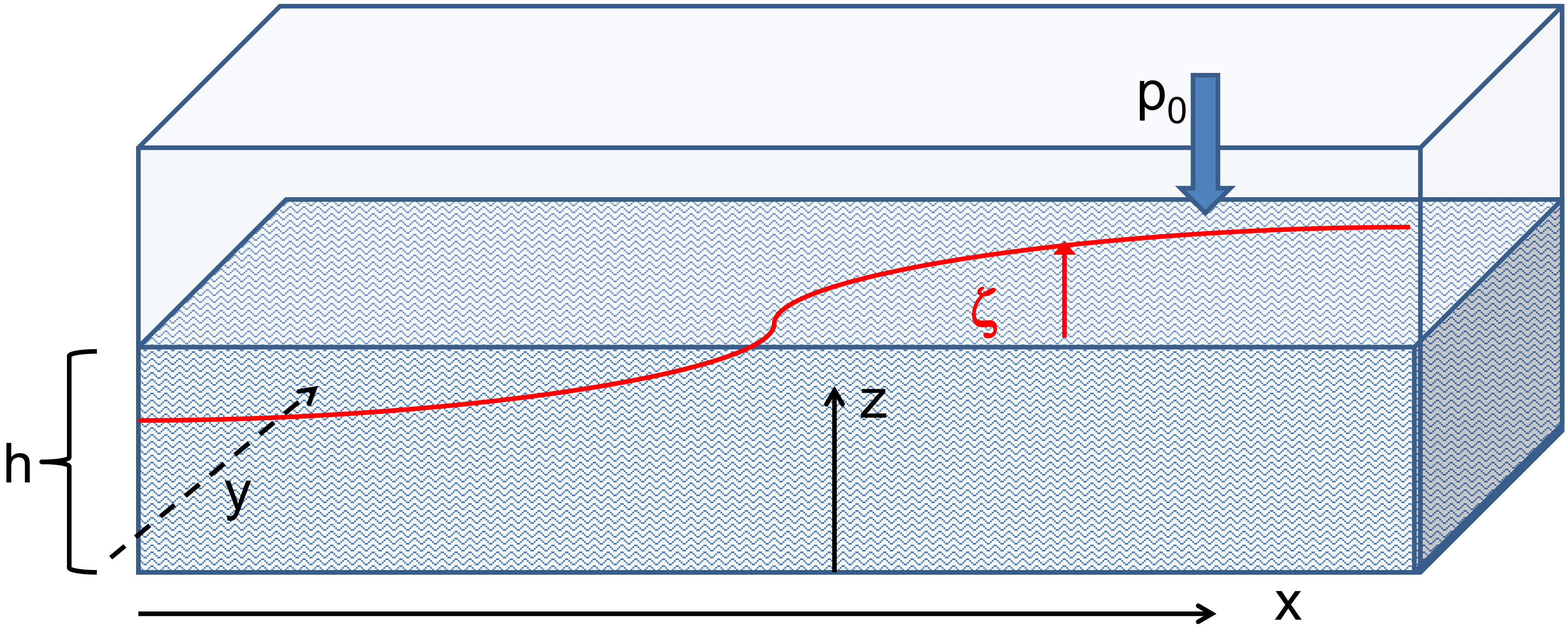
Consider a container of water with average height h and surface h+ζ(x,y,t); (h ←→ z_0 on some of the slides)

Atmospheric pressure is in equilibrium with the surface of water

Pressure at a height z above the bottom where the surface is at a height $h+\zeta$:

$$p(z) = \begin{cases} p_0 + \rho g\left(h+\zeta - z\right) & \text{For } z \le h+\zeta \\ p_0 & \text{For } z > h+\zeta \end{cases}$$

Here ρ represents density of water



Why do we not consider ρ_{air} in this analysis?

a. Because it is a reasonable approximation
b. Because it simplifies the analysis
c. Both of the above

Related question from Lee:
On slides 4 & 5, the air density does not appear in the equations. Isn't this because all considerations of the air column above the surface are accounted for in p0? If the situation were modified, let's say for water within a pressurized vessel where the density of air becomes more substantial, would the only change be to increase p0?

Euler's equation inside a incompressible fluid:

$$\frac{d\mathbf{v}}{dt} = f_{applied} - \frac{\nabla p}{\rho} = -g\hat{\mathbf{z}} - \frac{\nabla p}{\rho}$$

Assume that $v_z << v_x, v_y$ $\quad \Rightarrow -g - \frac{1}{\rho}\frac{\partial p}{\partial z} \approx 0$

$$\Rightarrow p(x,y,z,t) = p_0 + \rho g\left(\zeta(x,y,t) + h - z\right)$$

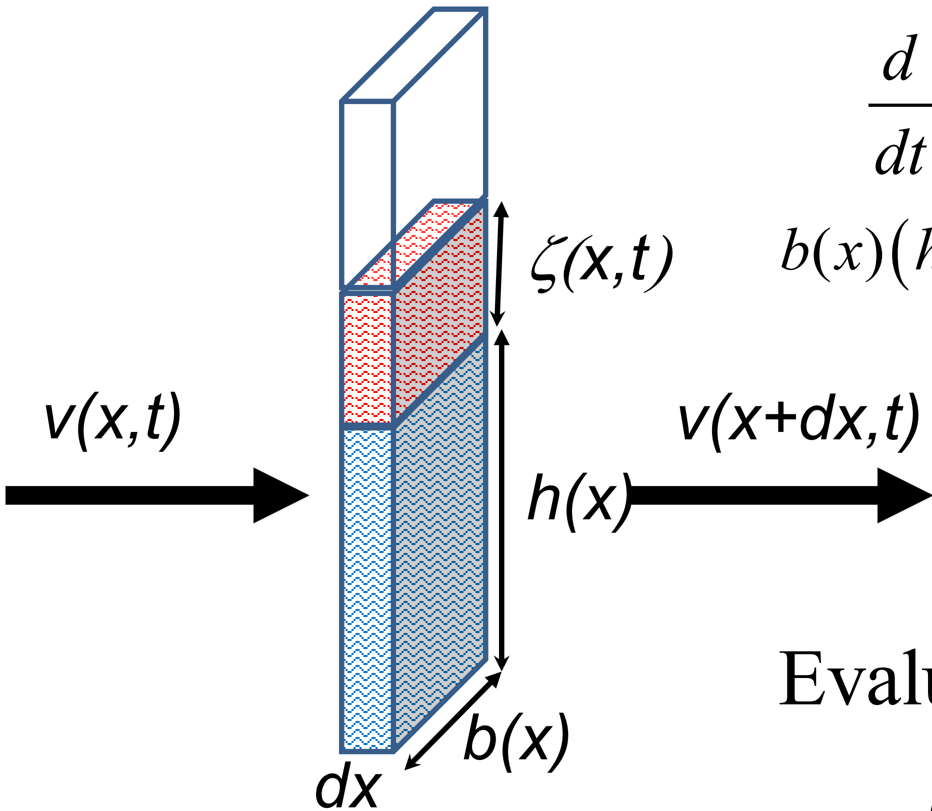
within the water

Horizontal fluid motions (keeping leading terms):

$$\frac{dv_x}{dt} \approx \frac{\partial v_x}{\partial t} = -\frac{1}{\rho}\frac{\partial p}{\partial x} = -g\frac{\partial \zeta}{\partial x}$$

$$\frac{dv_y}{dt} \approx \frac{\partial v_y}{\partial t} = -\frac{1}{\rho}\frac{\partial p}{\partial y} = -g\frac{\partial \zeta}{\partial y}$$

Consider a surface wave moving in the *x*-direction in a channel of width *b(x)* and height *h(x)* +*ζ(x,t)* :



Continuity condition in integral form:

$$\frac{d}{dt}\int_V \rho dV + \int_A \rho \mathbf{v}\cdot d\mathbf{A} = 0$$

$$b(x)\big(h(x)+\zeta(x,t)\big)\hat{\mathbf{x}}$$

$$b(x)\big(h(x)+\zeta(x,t)\big)dx$$

Evaluating continuity condition:

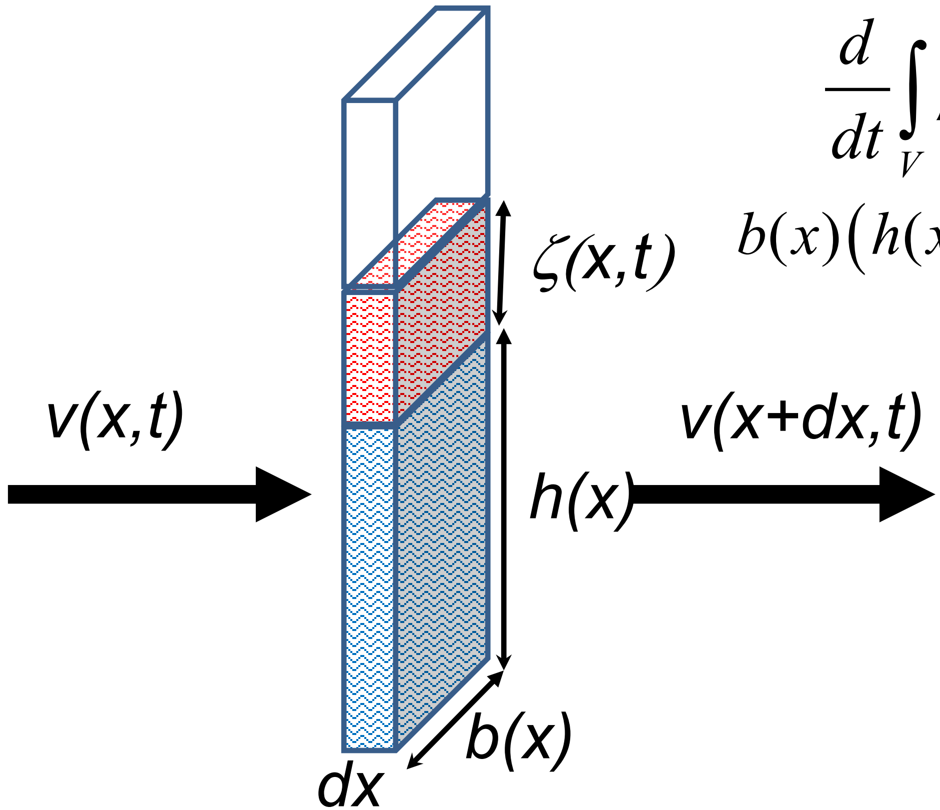
$$b(x)\frac{\partial \zeta}{\partial t} = -\frac{\partial}{\partial x}\big(h(x)b(x)v(x,t)\big)$$

Some details

Continuity condition in integral form:

$$\frac{d}{dt}\int_V \rho dV + \int_A \rho \mathbf{v}\cdot d\mathbf{A} = 0$$

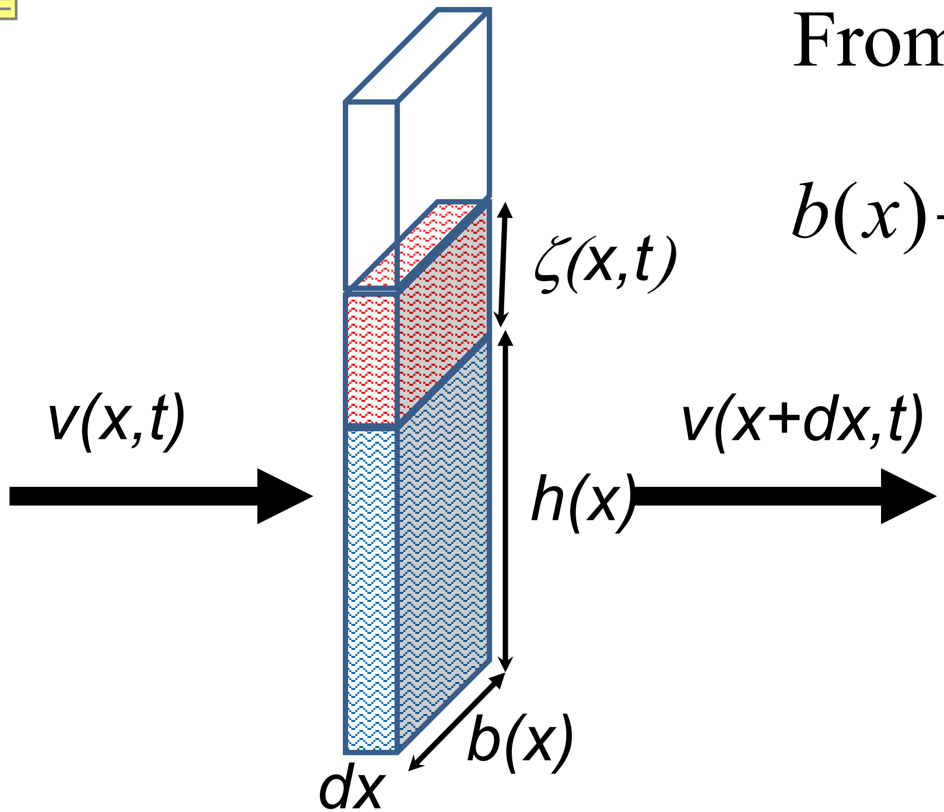
$$b(x)\left(h(x)+\zeta(x,t)\right)dx \qquad\qquad b(x)\left(h(x)+\zeta(x,t)\right)\hat{\mathbf{x}}$$



Here, we are assuming that ρ is constant

$$\frac{d}{dt}\int_V \rho dV + \int_A \rho \mathbf{v}\cdot d\mathbf{A} = \rho\int b(x)\frac{\partial \zeta}{\partial t}dx + \rho\int \frac{\partial}{\partial x}\left(b(x)(h(x)+\zeta(x,t))v(x,t)\right)dx = 0$$

$$\Rightarrow b(x)\frac{\partial \zeta}{\partial t} = -\frac{\partial}{\partial x}\left(h(x)b(x)v(x,t)\right)$$



From continuity condition:

$$b(x)\frac{\partial \zeta}{\partial t} = -\frac{\partial}{\partial x}\left(h(x)b(x)v(x,t)\right)$$

Example (Problem 10.3):

$$b(x) = b_0 \qquad h(x) = \kappa x$$

(A special case sometimes found at a beach.)

$$b_0\frac{\partial \zeta}{\partial t} = -\frac{\partial}{\partial x}\left((\kappa x)b_0 v(x,t)\right)$$

$$\frac{\partial \zeta}{\partial t} = -\kappa\left(v + x\frac{\partial v}{\partial x}\right)$$

From Newton-Euler equation:

$$\frac{dv}{dt} \approx \frac{\partial v}{\partial t} = -g\frac{\partial \zeta}{\partial x}$$

Example continued

$$\frac{\partial \zeta}{\partial t} = -\kappa\left(v + x\frac{\partial v}{\partial x}\right) \quad \Rightarrow \quad \frac{\partial^2 \zeta}{\partial t^2} = -\kappa\left(\frac{\partial v}{\partial t} + x\frac{\partial^2 v}{\partial x \partial t}\right)$$

$$\frac{\partial v}{\partial t} = -g\frac{\partial \zeta}{\partial x} \quad \Rightarrow \quad \frac{\partial^2 \zeta}{\partial t^2} = \kappa g\left(\frac{\partial \zeta}{\partial x} + x\frac{\partial^2 \zeta}{\partial x^2}\right)$$

It can be shown that a solution can take the form:

$$\zeta(x,t) = CJ_0\left(\frac{2\omega}{\sqrt{\kappa g}}\sqrt{x}\right)\cos(\omega t)$$

Note that $J_0(u)$ satisfies the equation: $\left(\frac{d^2}{du^2} + \frac{1}{u}\frac{d}{du} + 1\right)J_0(u) = 0$

Therefore, for $u = \frac{2\omega}{\sqrt{\kappa g}}\sqrt{x}$

$$\left(x\frac{d^2}{dx^2} + \frac{d}{dx}\right)J_0(u) = \frac{\omega^2}{\kappa g}\left(\frac{d^2}{du^2} + \frac{1}{u}\frac{d}{du}\right)J_0(u) = -\frac{\omega^2}{\kappa g}J_0(u)$$

Therefore, for $u = \frac{2\omega}{\sqrt{\kappa g}}\sqrt{x} \quad \Rightarrow \frac{1}{\sqrt{x}} = \frac{2\omega}{\sqrt{\kappa g}}\frac{1}{u}$

$$\left(x\frac{d^2}{dx^2}+\frac{d}{dx}\right)J_0(u) = \frac{\omega^2}{\kappa g}\left(\frac{d^2}{du^2}+\frac{1}{u}\frac{d}{du}\right)J_0(u) = -\frac{\omega^2}{\kappa g}J_0(u)$$

Detail: $\frac{dJ_0(u)}{dx} = \frac{dJ_0(u)}{du}\frac{\omega}{\sqrt{\kappa g}}\frac{1}{\sqrt{x}}$

$$\frac{d^2J_0(u)}{dx^2} = \frac{d^2J_0(u)}{du^2}\left(\frac{\omega}{\sqrt{\kappa g}}\frac{1}{\sqrt{x}}\right)^2 - \frac{dJ_0(u)}{du}\frac{\omega}{2\sqrt{\kappa g}}\frac{1}{x\sqrt{x}}$$

Therefore: $\left(x\frac{d^2}{dx^2}+\frac{d}{dx}\right)J_0(u) = \left(\frac{\omega^2}{\kappa g}\frac{d^2J_0(u)}{du^2}+\frac{dJ_0(u)}{du}\frac{\omega}{2\sqrt{\kappa g}}\frac{1}{\sqrt{x}}\right)$

$$= \frac{\omega^2}{\kappa g}\left(\frac{d^2J_0(u)}{du^2}+\frac{dJ_0(u)}{du}\frac{1}{u}\right)$$

Example continued

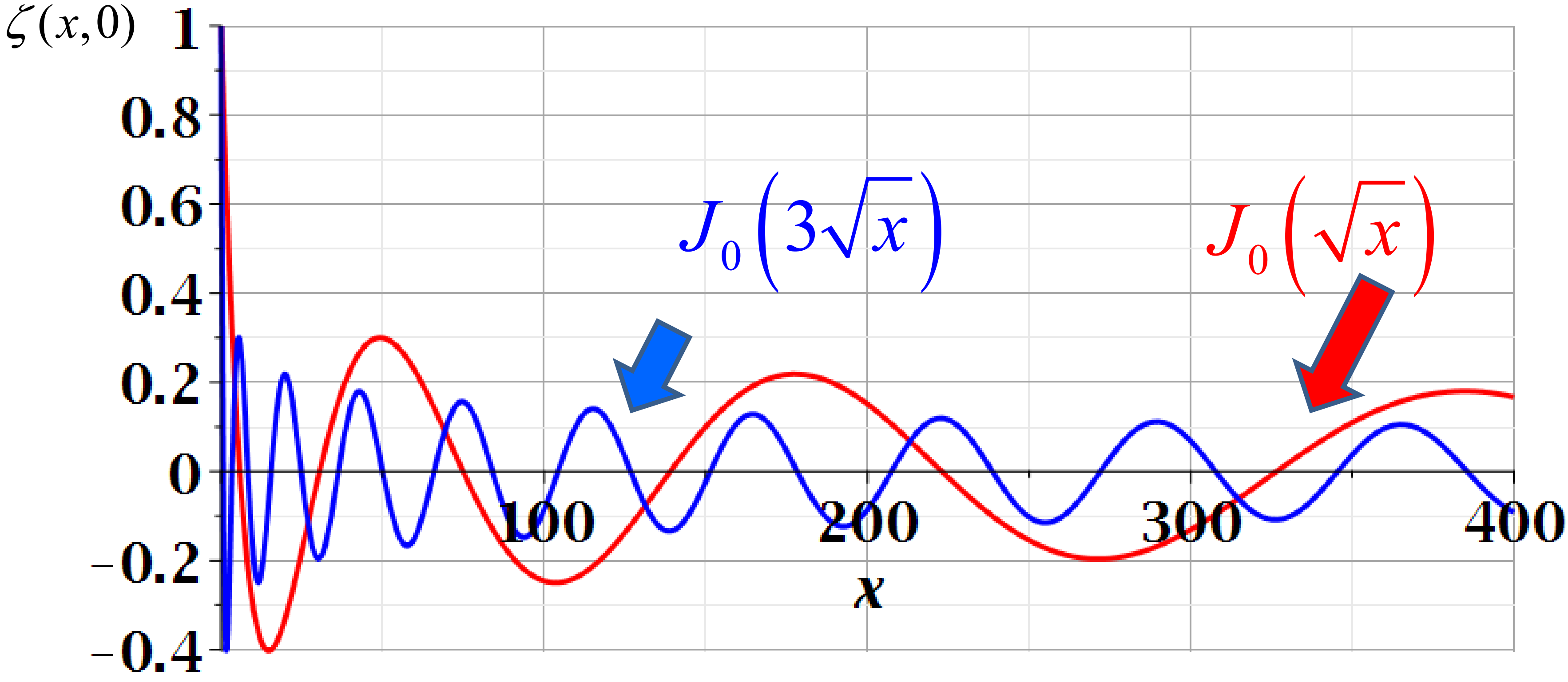
$$\frac{\partial^2 \zeta}{\partial t^2} = \kappa g \left(\frac{\partial \zeta}{\partial x} + x \frac{\partial^2 \zeta}{\partial x^2} \right)$$

$$\Rightarrow \zeta(x,t) = C J_0 \left(\frac{2\omega\sqrt{x}}{\sqrt{\kappa g}} \right) \cos(\omega t)$$

Check:

$$-\omega^2 C J_0 \left(\frac{2\omega\sqrt{x}}{\sqrt{\kappa g}} \right) \cos(\omega t) = \kappa g \left(\frac{\partial}{\partial x} + x \frac{\partial^2}{\partial x^2} \right) C J_0 \left(\frac{2\omega\sqrt{x}}{\sqrt{\kappa g}} \right) \cos(\omega t)$$

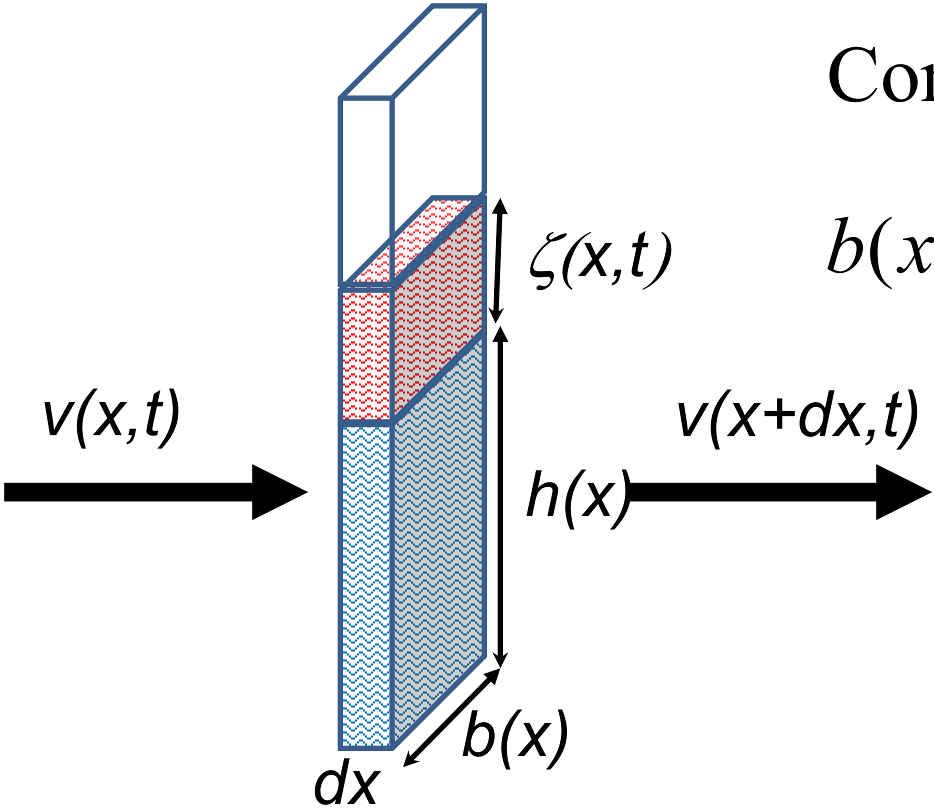
$$\zeta(x,t) = CJ_0\left(\frac{2\omega}{\sqrt{\kappa g}}\sqrt{x}\right)\cos(\omega t)$$



Imagine watching the waves at a beach – can you visualize the configuration for the surface wave pattern to approximation this situation?

a. Long flat beach
b. Beach in which average water level increases
c. Beach in which average water level decreases

A simplier example:



Continuity condition:

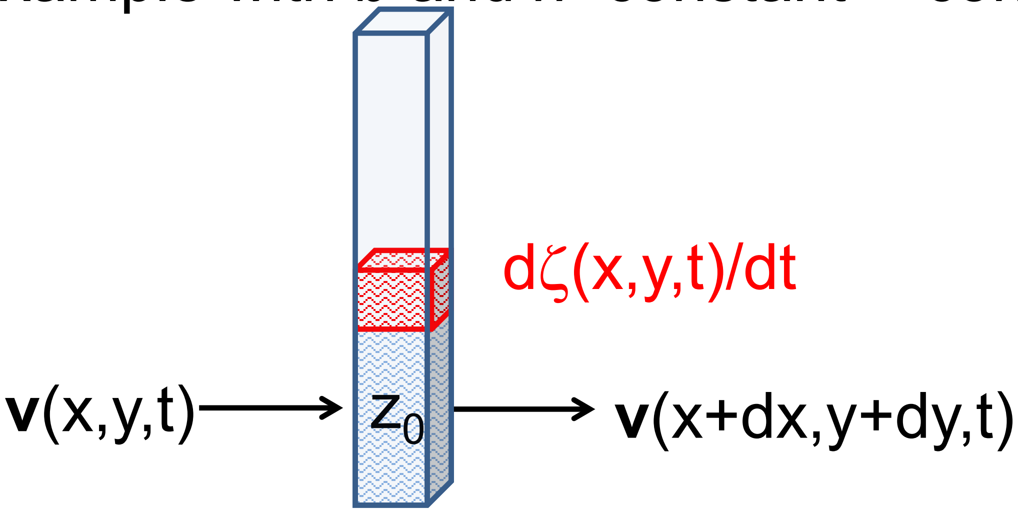
$$b(x)\frac{\partial \zeta}{\partial t} = -\frac{\partial}{\partial x}\Big(h(x)b(x)v(x,t)\Big)$$

Special case, where *b* and *h* are constant --

For constant b and h:

$$\frac{\partial \zeta}{\partial t} = -h\frac{\partial}{\partial x}\Big(v(x,t)\Big)$$

Example with *b* and *h* constant -- continued



Continuity condition for flow of incompressible fluid:

$$\frac{\partial \zeta}{\partial t} + h\nabla \cdot \mathbf{v} = 0$$

From horizontal flow relations: $$\frac{\partial \mathbf{v}}{\partial t} = -g\nabla \zeta$$

Equation for surface function: $$\frac{\partial^2 \zeta}{\partial t^2} - gh\nabla^2 \zeta = 0$$

For uniform channel:

Surface wave equation:

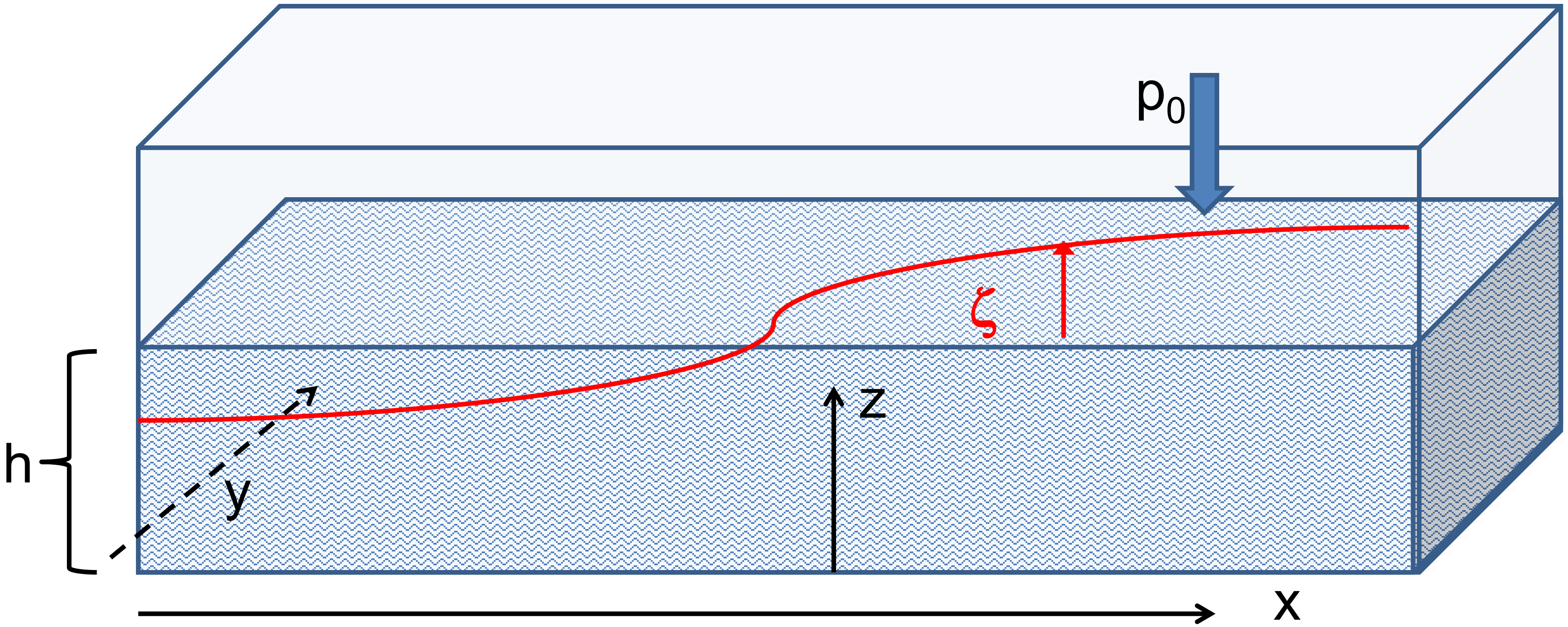
$$\frac{\partial^2 \zeta}{\partial t^2} - c^2 \nabla^2 \zeta = 0 \qquad c^2 = gh$$

More complete analysis finds:

$$c^2 = \frac{g}{k} \tanh(kh) \qquad \text{where } k = \frac{2\pi}{\lambda}$$

More details: -- recall setup --

Consider a container of water with average height h and surface h+$\zeta(x,y,t)$



Equations describing fluid itself (without boundaries)

Euler's equation for incompressible fluid:

$$\frac{d\mathbf{v}}{dt} = \frac{\partial \mathbf{v}}{\partial t} + \mathbf{v}\cdot\nabla\mathbf{v} = \frac{\partial \mathbf{v}}{\partial t} + \nabla\left(\tfrac{1}{2}v^2\right) + \mathbf{v}\times\left(\nabla\times\mathbf{v}\right) = -\nabla U - \frac{\nabla p}{\rho}$$

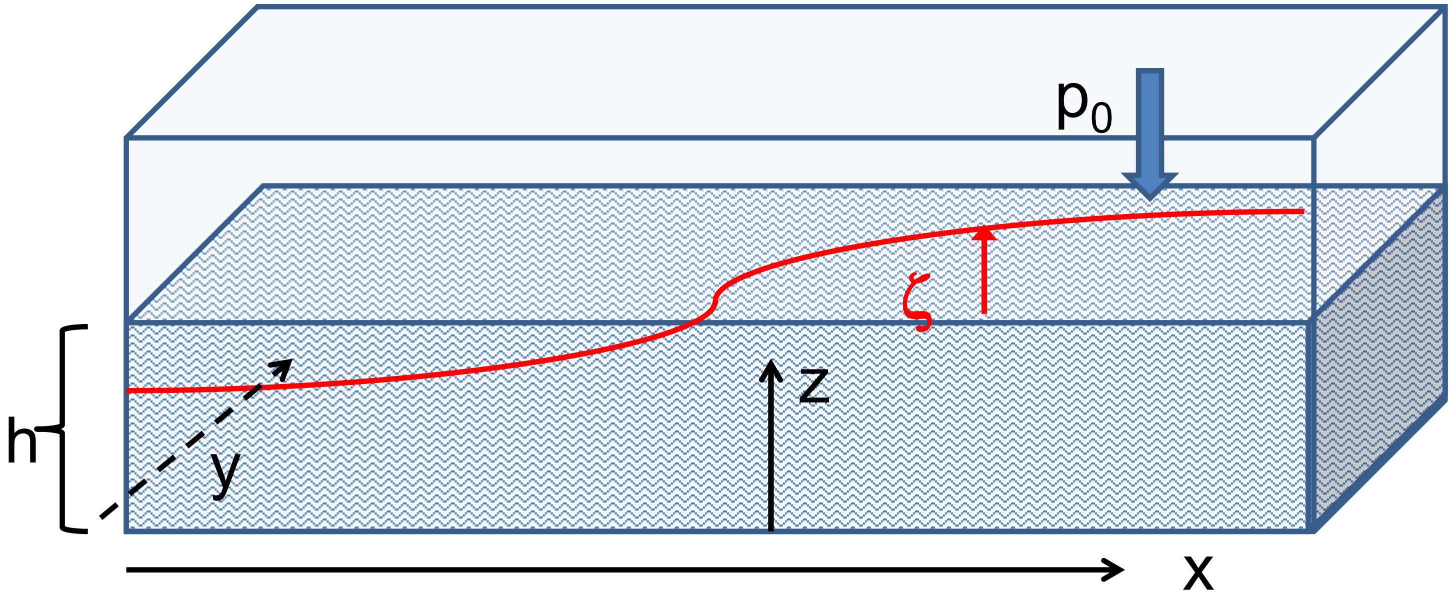
Assume that $\nabla\times\mathbf{v} = 0$ (irrotational flow) $\Rightarrow \mathbf{v} = -\nabla\Phi$

$$\Rightarrow \nabla\left(-\frac{\partial\Phi}{\partial t} + \tfrac{1}{2}v^2 + U + \frac{p}{\rho}\right) = 0$$

$$\Rightarrow -\frac{\partial\Phi}{\partial t} + \tfrac{1}{2}v^2 + U + \frac{p}{\rho} = \text{constant (within the fluid)}$$

For the same system, the continuity condition becomes

$$\nabla\cdot\mathbf{v} = -\nabla^2\Phi = 0$$



Within fluid: $0 \leq z \leq h + \zeta$

$$-\frac{\partial \Phi}{\partial t} + \tfrac{1}{2}v^2 + g(z-h) = \text{constant}$$

(We have absorbed p_0 in "constant")

$$-\nabla^2 \Phi = 0$$

At surface: $z = h + \zeta$ with $\zeta = \zeta(x, y, t)$

$$\frac{d\zeta}{dt} = \frac{\partial \zeta}{\partial t} + v_x \frac{\partial \zeta}{\partial x} + v_y \frac{\partial \zeta}{\partial y}$$

where $v_{x,y} = v_{x,y}(x, y, h+\zeta, t)$

Full equations:

Within fluid: $0 \le z \le h+\zeta$

$$-\frac{\partial \Phi}{\partial t} + \tfrac{1}{2}v^2 + g(z-h) = \text{constant}$$

(We have absorbed p_0 in “constant”)

$$-\nabla^2 \Phi = 0$$

At surface: $z = h+\zeta$ with $\zeta = \zeta(x,y,t)$

$$\frac{d\zeta}{dt} = \frac{\partial \zeta}{\partial t} + v_x \frac{\partial \zeta}{\partial x} + v_y \frac{\partial \zeta}{\partial y} \quad \text{where } v_{x,y} = v_{x,y}(x,y,h+\zeta,t)$$

Linearized equations:

For $0 \le z \le h+\zeta$: $-\frac{\partial \Phi}{\partial t} + g(z-h) = 0 \qquad -\nabla^2 \Phi = 0$

At surface: $z = h+\zeta \qquad \frac{d\zeta}{dt} = \frac{\partial \zeta}{\partial t} = v_z(x,y,h+\zeta,t)$

$$-\frac{\partial \Phi(x,y,h+\zeta,t)}{\partial t} + g\zeta = 0$$

For simplicity, keep only linear terms and assume that horizontal variation is only along *x*:

For $\quad 0 \leq z \leq h + \zeta : \qquad \nabla^2 \Phi = \left(\frac{\partial^2}{\partial z^2} + \frac{\partial^2}{\partial x^2} \right) \Phi(x, z, t) = 0$

Consider and periodic waveform: $\quad \Phi(x, z, t) = Z(z) \cos\big(k(x - ct)\big)$

$$\Rightarrow \left(\frac{d^2}{dz^2} - k^2 \right) Z(z) = 0$$

Boundary condition at bottom of tank: $\quad v_z(x, 0, t) = 0$

$$\Rightarrow \frac{dZ}{dz}\big(0\big) = 0 \qquad\qquad Z(z) = A\cosh(kz)$$

For simplicity, keep only linear terms and assume that horizontal variation is only along *x* – continued:

$$\text{At surface:} \quad z = h + \zeta \qquad \frac{\partial \zeta}{\partial t} = v_z\left(x, h+\zeta, t\right) = -\frac{\partial \Phi\left(x, h+\zeta, t\right)}{\partial z}$$

$$-\frac{\partial \Phi\left(x, h+\zeta, t\right)}{\partial t} + g\zeta = 0$$

$$-\frac{\partial^2 \Phi\left(x, h+\zeta, t\right)}{\partial t^2} + g\frac{\partial \zeta}{\partial t} = -\frac{\partial^2 \Phi\left(x, h+\zeta, t\right)}{\partial t^2} - g\frac{\partial \Phi\left(x, h+\zeta, t\right)}{\partial z} = 0$$

$$\text{For} \quad \Phi(x, \left(h+\zeta\right), t) = A\cosh(k\left(h+\zeta\right))\cos\left(k(x-ct)\right)$$

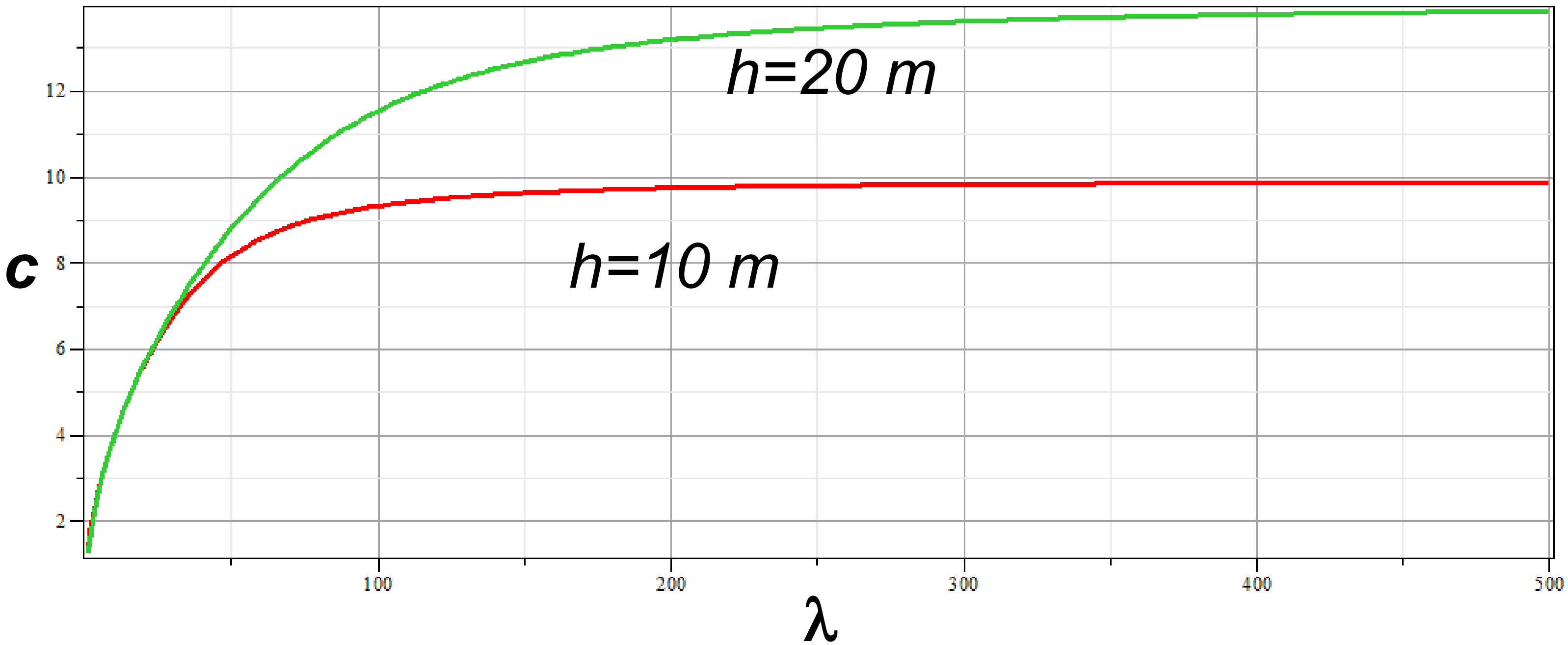
$$A\cosh(k\left(h+\zeta\right))\cos\left(k(x-ct)\right)\left(k^2c^2 - gk\frac{\sinh(k\left(h+\zeta\right))}{\cosh(k\left(h+\zeta\right))}\right) = 0$$

$$\Rightarrow \mathrm{c}^2 = \frac{g}{k}\frac{\sinh(k\left(h+\zeta\right))}{\cosh(k\left(h+\zeta\right))}$$

For simplicity, keep only linear terms and assume that horizontal variation is only along *x* – continued:

$$c^2 = \frac{g}{k}\frac{\sinh(k(h+\zeta))}{\cosh(k(h+\zeta))} = \frac{g}{k}\tanh(k(h+\zeta))$$

Assuming $\zeta << h$: $$c^2 = \frac{g}{k}\tanh(kh) \qquad \lambda = \frac{2\pi}{k}$$



For simplicity, keep only linear terms and assume that horizontal variation is only along *x* – continued:

$$c^2 \approx \frac{g}{k}\tanh(kh) \qquad \text{For } \lambda >> h,\ c^2 \approx gh$$

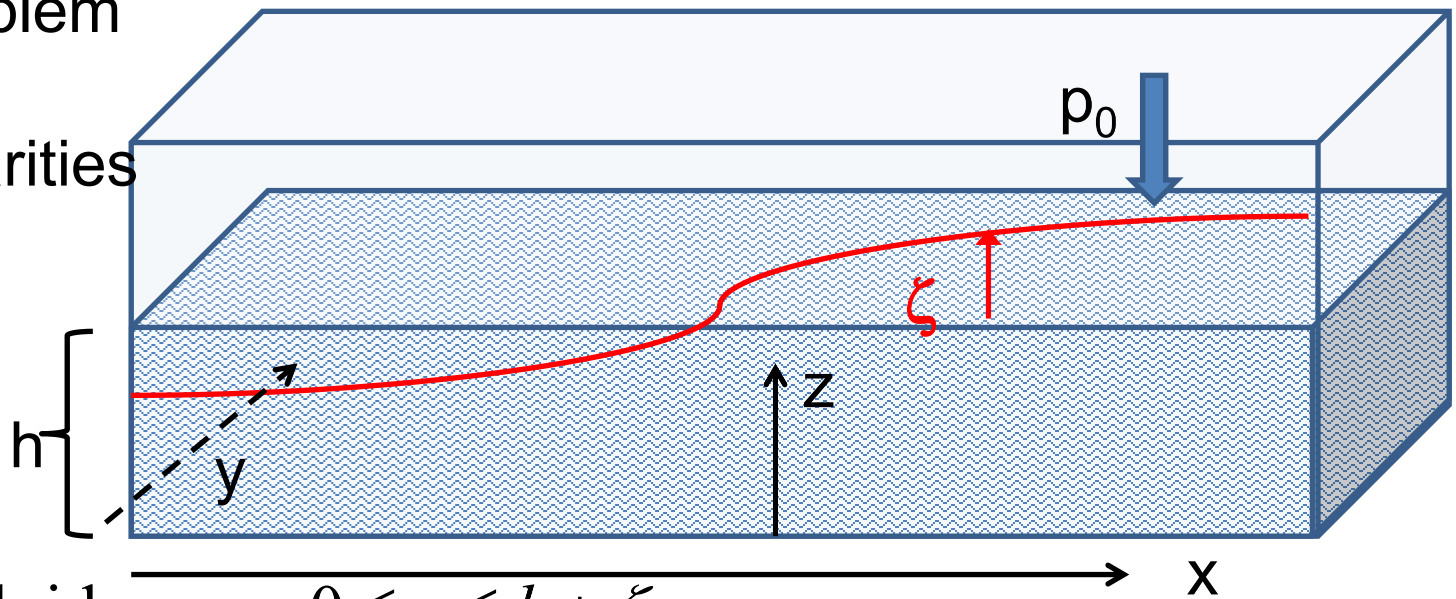
$$\Phi(x,z,t) = A\cosh(kz)\cos(k(x-ct))$$

$$\zeta(x,t) = \frac{1}{g}\frac{\partial\Phi(x,h+\zeta,t)}{\partial t} \approx \frac{kc}{g}A\cosh(kh)\sin(k(x-ct))$$

Note that for $\lambda >> h, \quad c^2 \approx gh$

(solutions are consistent with previous analysis)

General problem including non-linearities



Within fluid: $0 \le z \le h + \zeta$

$$-\frac{\partial \Phi}{\partial t} + \tfrac{1}{2} v^2 + g(z-h) = \text{constant}$$

(We have absorbed p_0 in our constant.)

$$-\nabla^2 \Phi = 0$$

At surface: $z = h + \zeta$ with $\zeta = \zeta(x, y, t)$

$$\frac{d\zeta}{dt} = \frac{\partial \zeta}{\partial t} + v_x \frac{\partial \zeta}{\partial x} + v_y \frac{\partial \zeta}{\partial y}$$

where $v_{x,y} = v_{x,y}(x, y, h+\zeta, t)$