



**PHY 711 Classical Mechanics and
Mathematical Methods
10-10:50 AM MWF in Olin 103**

Notes for Lecture 3 based on Chap. 1 of F&W

**Scattering theory – Coordinate frames
Center of mass and laboratory frames**

Your questions --

From Evan -- On slide 17, it is mentioned that we only need to consider what is happening before and after the collision. Why is it that we can ignore what is happening during the collision (the intermediate)? Is it because we have already accounted for this when considering what is happening after the collision?

From Zezong -- I am still a bit confused about why $V_{cm}/V_1 = m_1/m_2$ for elastic scattering when deriving the expression of $\tan \theta$.

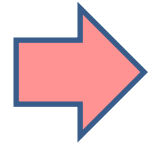
From Sam -- What exactly is the physical explanation of the term μ that occurs in the center of mass FOR? is it to account for movement that doesn't change the center of mass? I remember using it but can't remember why it exists.



Course schedule

(Preliminary schedule -- subject to frequent adjustment.)

	Date	F&W Reading	Topic	Assignment	Due
1	Mon, 8/22/2022		Introduction	#1	8/26/2022
2	Wed, 8/24/2022	Chap. 1	Scattering theory		
3	Fri, 8/26/2022	Chap. 1	Scattering theory	#2	8/29/2022
4	Mon, 8/29/2022	Chap. 1	Scattering theory		



PHY 711 – Assignment #2

08/24/2022

Consider a particle of mass m moving in the vicinity of another particle of mass M , initially at rest, where $m \ll M$. The particles interact with a conservative central potential of the form

$$V(r) = V_0 \left(\left(\frac{r_0}{r} \right)^2 - \left(\frac{r_0}{r} \right) \right),$$

where r denotes the magnitude of the particle separation and V_0 and r_0 denote energy and length constants, respectively. The total energy of the system E is constant and $E = V_0$.

- (a) First consider the case where the impact parameter $b = 0$. Find the distance of closest approach of the particles.
- (b) Now consider the case where the impact parameter $b = r_0$. Find the distance of closest approach of the particles.

Scattering theory:

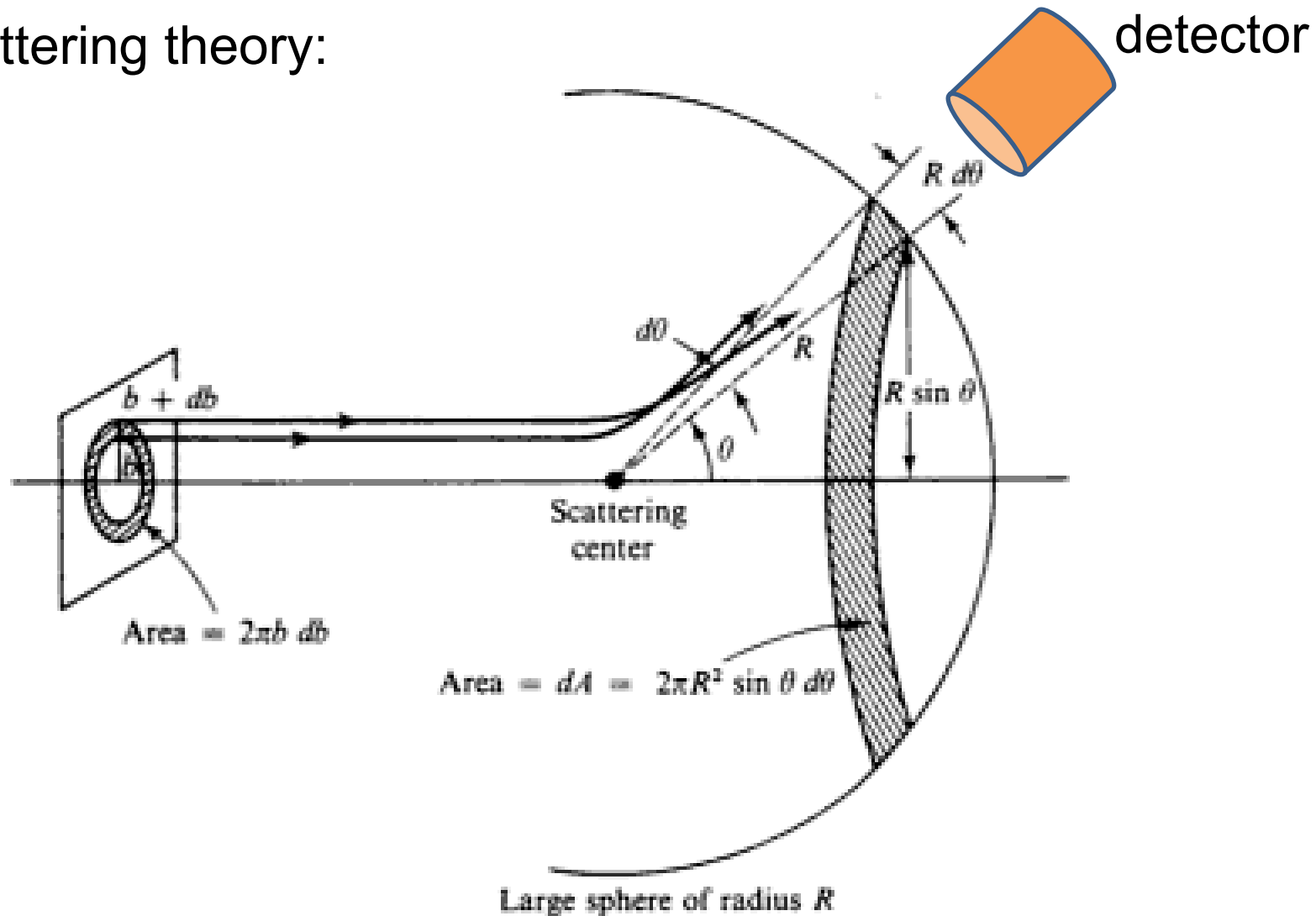


Figure 5.5 The scattering problem and relation of cross section to impact parameter.

Can you think of examples of such an experimental setup?

Other experimental designs –

At CERN <https://home.cern/science/experiments/totem> the study of highly energetic proton-proton scattering is designed in the center of mass frame of reference by accelerating two proton beams focused to collide head on in the Large Hadron Collider LHC facility.

Figure from CERN website

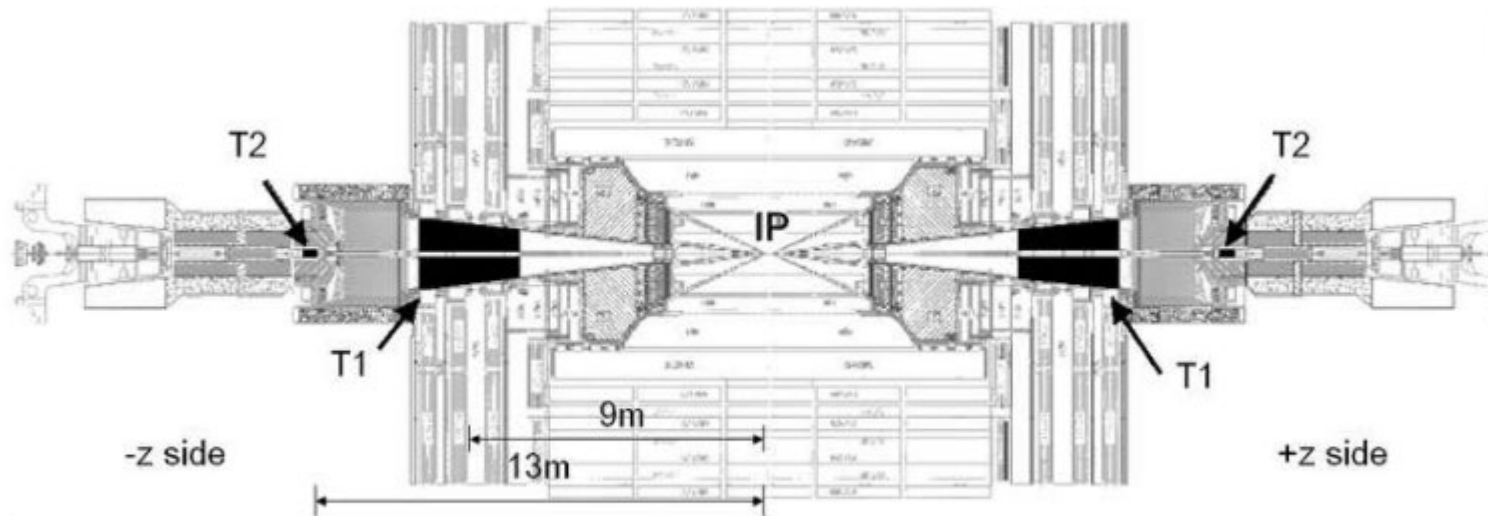


Figure 1.17: View of the inelastic forward trackers T1 and T2 inside the CMS detector.

What might be the advantage/disadvantage of this design?

What are the benefits/disadvantages of expressing the scattering cross section in the laboratory frame of reference vs center of mass frame of reference? (When or why to use a particular frame of reference)

Advantages of Lab frame

1. Natural experimental design.
2. Some targets are more naturally at rest.
3. ??

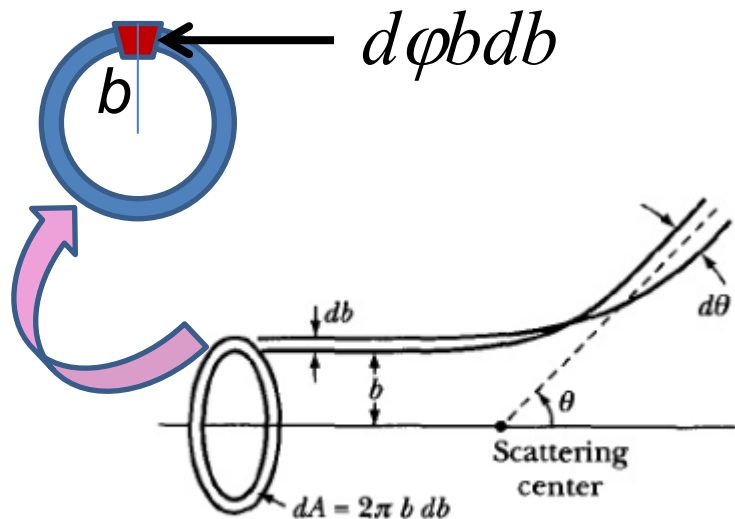
Advantages of CM frame

1. Analysis is done in CM frame.
2. Experiment is more energy efficient.
3. ??

Differential cross section

$$\left(\frac{d\sigma}{d\Omega}\right) = \frac{\text{Number of detected particles at } \theta \text{ per target particle}}{\text{Number of incident particles per unit area}}$$

= Area of incident beam that is scattered into detector
at angle θ



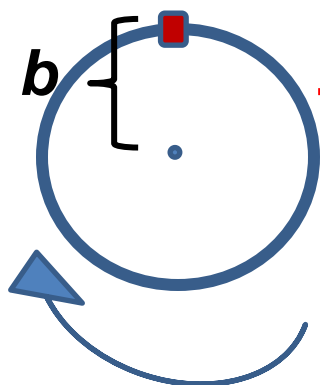
$$\left(\frac{d\sigma}{d\Omega}\right) = \frac{d\phi b db}{d\phi \sin\theta d\theta} = \frac{b}{\sin\theta} \left|\frac{db}{d\theta}\right|$$

Figure from Marion & Thorton, Classical Dynamics

More details --

We imagine that the beam of particles has a cylindrical geometry and that the physics is totally uniform in the azimuthal direction. The cross section of the beam is a circle.

View of beam
cross section:



← This piece of the beam scatters into the detector at angle θ

This logic leads to the notion that b is a function of θ and we will try to find $b(\theta)$ for various cases.

$\varphi \equiv$ azimuthal angle

Note: The notion of cross section is common to many areas of physics including classical mechanics, quantum mechanics, optics, etc. Only in the **classical mechanics** can we calculate it from a knowledge of the particle trajectory as it relates to the scattering geometry.

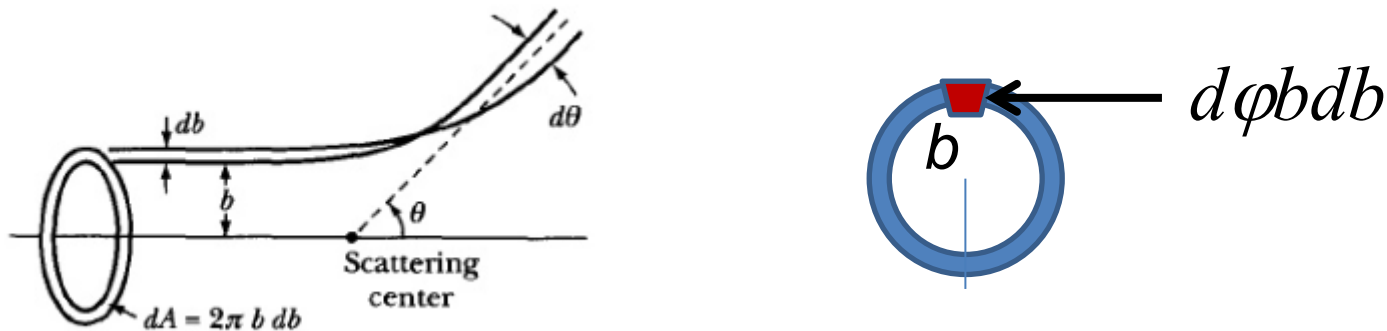
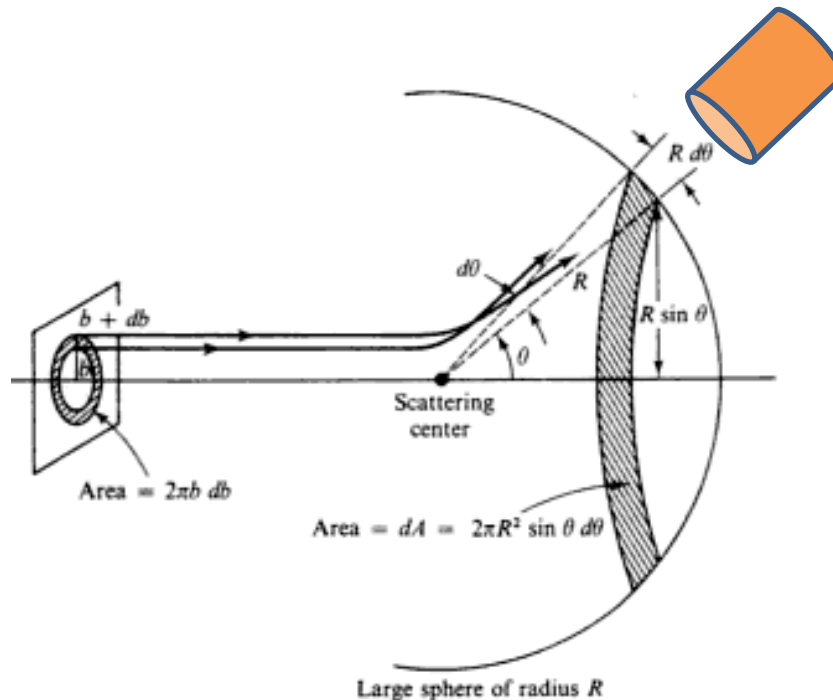


Figure from Marion & Thorton, Classical Dynamics

$$\left(\frac{d\sigma}{d\Omega} \right) = \frac{d\phi b db}{d\phi \sin\theta d\theta} = \frac{b}{\sin\theta} \left| \frac{db}{d\theta} \right|$$

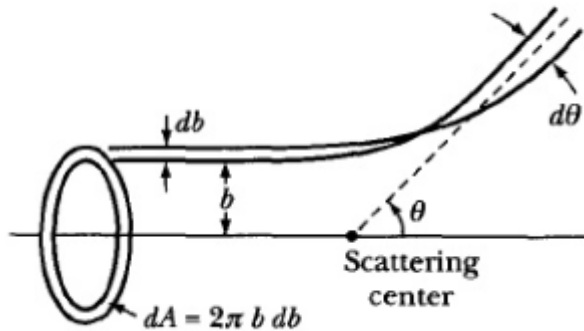
Note: We are assuming that the process is isotropic in ϕ

Elaboration on how we know that $b db d\phi$ is the relevant piece of beam ending up in our detector?



Comment: The interaction potential will determine the detailed shape of the particle trajectory which we can express as $r(\theta)$, which in principle can be related to the impact parameter as a function of scattering angle $b(\theta)$.

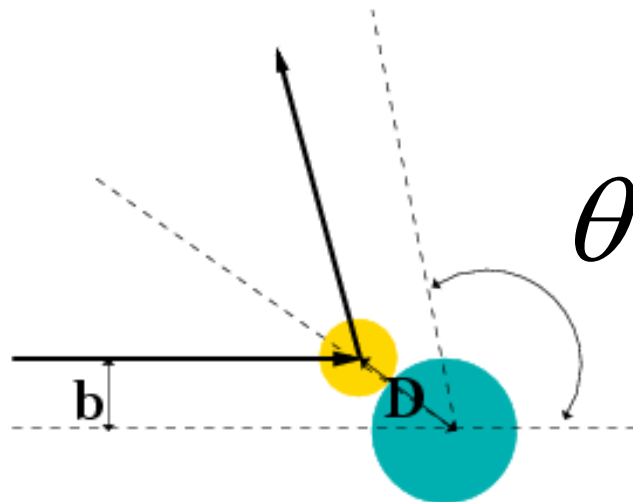
Simple example – collision of hard spheres



$$\left(\frac{d\sigma}{d\Omega} \right) = \frac{b}{\sin\theta} \left| \frac{db}{d\theta} \right|$$

Microscopic view:

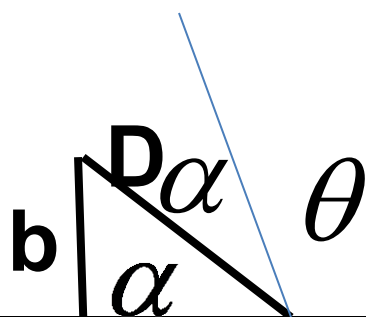
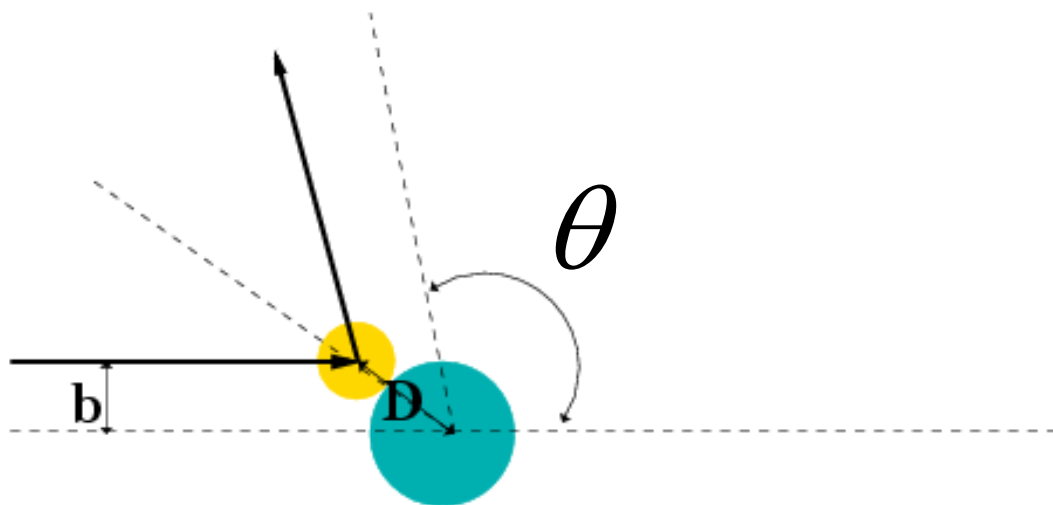
$$b(\theta) = ?$$



$$b(\theta) = D \sin\left(\frac{\pi}{2} - \frac{\theta}{2}\right)$$

$$\left(\frac{d\sigma}{d\Omega} \right) = \frac{D^2}{4}$$

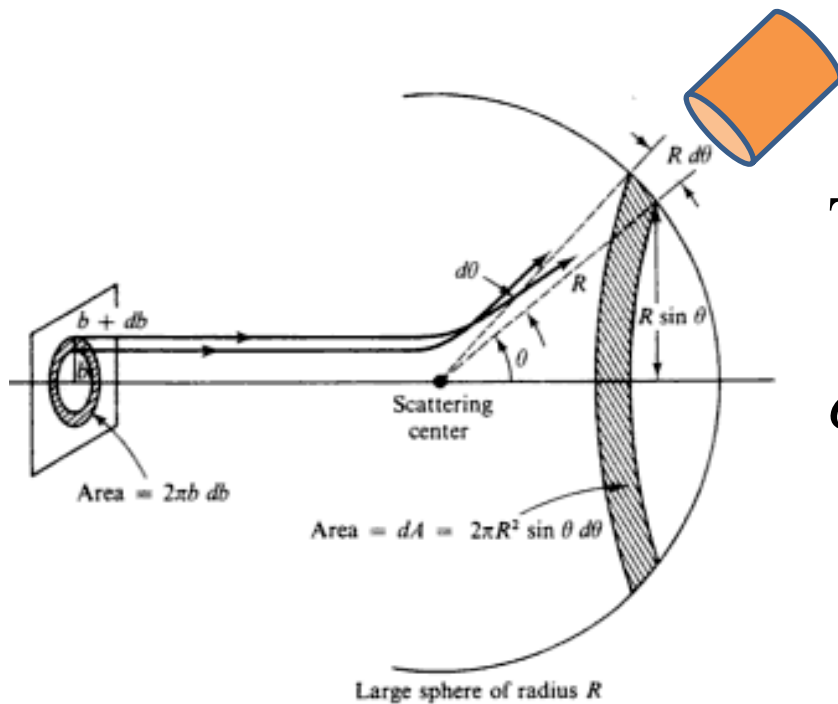
Some more details of form of $b(\theta)$



$$b = D \sin \alpha = D \sin \left(\frac{\pi}{2} - \frac{\theta}{2} \right)$$

$$2\alpha + \theta = \pi$$

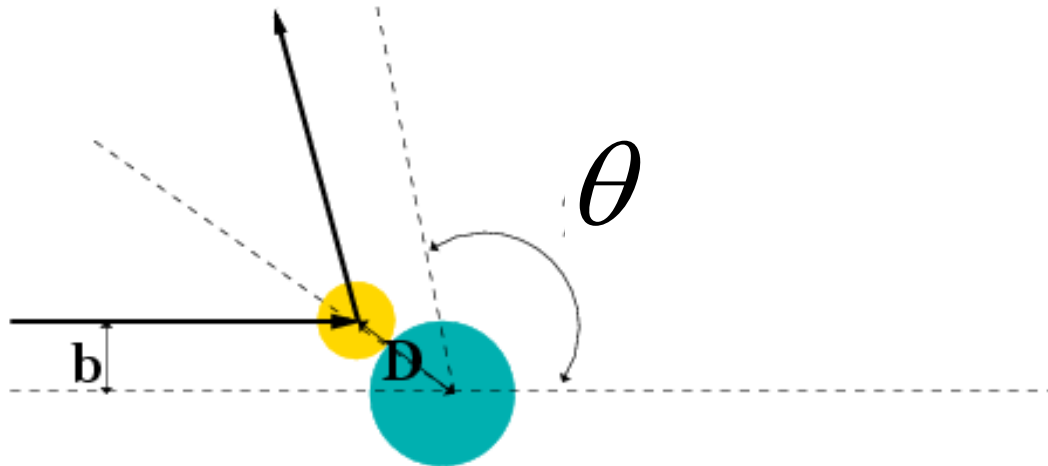
Simple example – collision of hard spheres -- continued



Total scattering cross section:

$$\sigma = \int \left(\frac{d\sigma}{d\Omega} \right) d\Omega$$

Hard sphere:



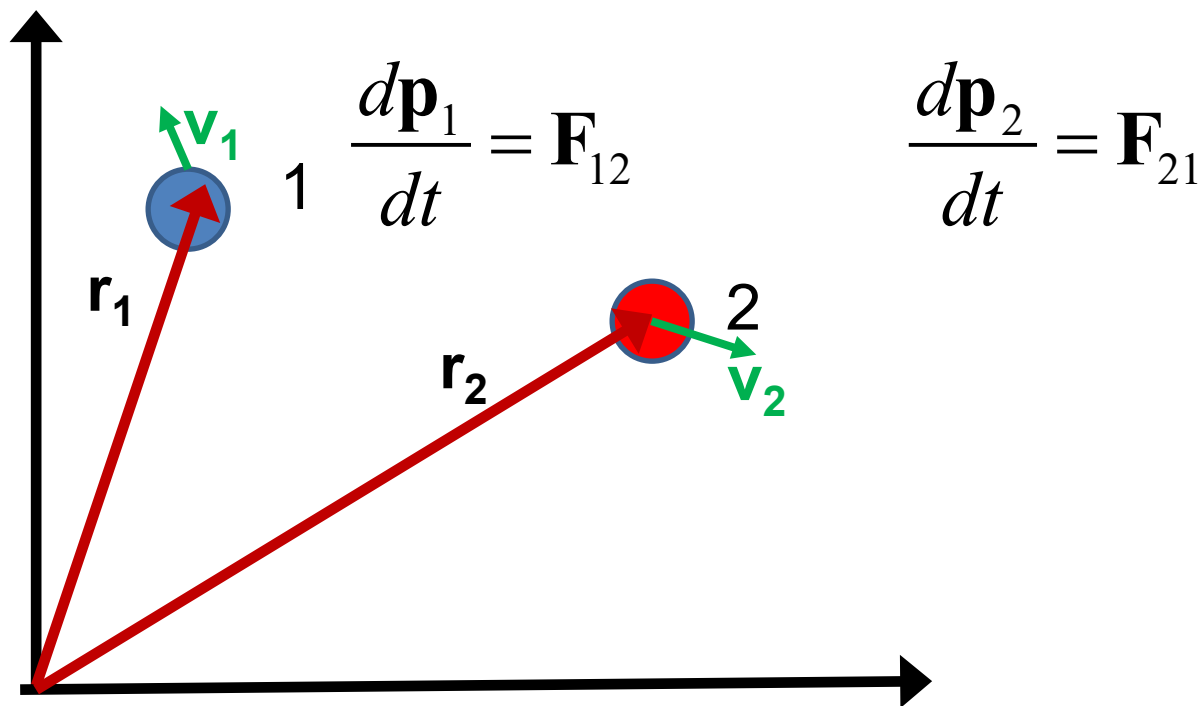
$$\left(\frac{d\sigma}{d\Omega} \right) = \frac{D^2}{4}$$

$$\sigma = \pi D^2$$

Now consider the more general case of particle interactions and the corresponding scattering analysis.

Scattering theory can help us analyze the interaction potential $V(r)$. First, we need to simplify the number of variables.

Relationship of scattering cross-section to particle interactions --
Classical mechanics of a conservative 2-particle system.



$$\mathbf{F}_{12} = -\nabla_1 V(\mathbf{r}_1 - \mathbf{r}_2) \quad \Rightarrow \quad E = \frac{1}{2} m_1 v_1^2 + \frac{1}{2} m_2 v_2^2 + V(\mathbf{r}_1 - \mathbf{r}_2)$$

Relationship between center of mass and laboratory frames of reference. At a time t , the following relationships apply --

Definition of center of mass \mathbf{R}_{CM}

$$m_1 \mathbf{r}_1 + m_2 \mathbf{r}_2 = (m_1 + m_2) \mathbf{R}_{CM}$$

$$m_1 \dot{\mathbf{r}}_1 + m_2 \dot{\mathbf{r}}_2 = (m_1 + m_2) \dot{\mathbf{R}}_{CM} = (m_1 + m_2) \mathbf{V}_{CM}$$

Note that $\dot{\mathbf{r}} \equiv \frac{d\mathbf{r}}{dt}$

$$E = \frac{1}{2} m_1 v_1^2 + \frac{1}{2} m_2 v_2^2 + V(\mathbf{r}_1 - \mathbf{r}_2)$$

$$= \frac{1}{2} (m_1 + m_2) V_{CM}^2 + \frac{1}{2} \mu |\mathbf{v}_1 - \mathbf{v}_2|^2 + V(\mathbf{r}_1 - \mathbf{r}_2)$$

where: $\mu \equiv \frac{m_1 m_2}{m_1 + m_2}$

Why do this? We need to make the mathematics tractable...

Classical mechanics of a conservative 2-particle system -- continued

$$E = \frac{1}{2}(m_1 + m_2)V_{CM}^2 + \frac{1}{2}\mu|\mathbf{v}_1 - \mathbf{v}_2|^2 + V(\mathbf{r}_1 - \mathbf{r}_2)$$

For central potentials: $V(\mathbf{r}_1 - \mathbf{r}_2) = V(|\mathbf{r}_1 - \mathbf{r}_2|) \equiv V(r_{12})$

Relative angular momentum is also conserved:

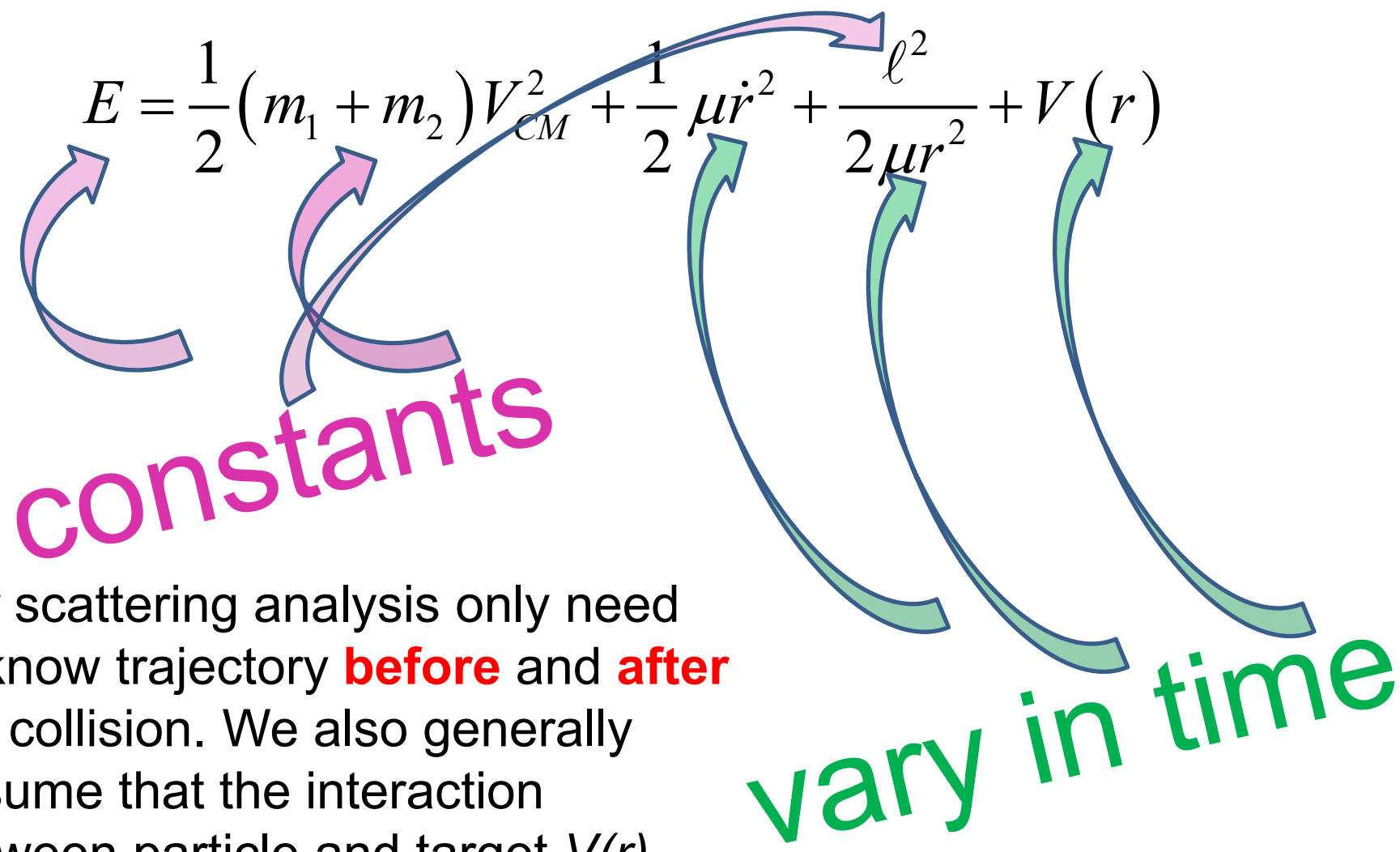
$$\mathbf{L}_{12} \equiv \mathbf{r}_{12} \times \mu\mathbf{v}_{12}$$

$$E = \frac{1}{2}(m_1 + m_2)V_{CM}^2 + \frac{1}{2}\mu v_{12}^2 + \frac{L_{12}^2}{2\mu r_{12}^2} + V(r_{12})$$

Simpler notation:

$$E = \frac{1}{2}(m_1 + m_2)V_{CM}^2 + \frac{1}{2}\mu\dot{r}^2 + \frac{\ell^2}{2\mu r^2} + V(r)$$

Simpler notation:

$$E = \frac{1}{2}(m_1 + m_2)V_{CM}^2 + \frac{1}{2}\mu\dot{r}^2 + \frac{\ell^2}{2\mu r^2} + V(r)$$
The diagram features several curved arrows. Three pink arrows originate from the first three terms of the equation: the first from the mass sum, the second from the reduced mass term, and the third from the angular momentum term. These arrows point towards the word 'constants'. Three green arrows originate from the last two terms: the first from the radial velocity term, the second from the angular momentum term, and the third from the potential energy term. These arrows point towards the phrase 'vary in time'.

constants

For scattering analysis only need to know trajectory **before** and **after** the collision. We also generally assume that the interaction between particle and target $V(r)$ conserves energy and angular momentum.

vary in time

Comment: The impact parameter b is a useful concept in the general case.

$$E_{total} = \underbrace{\frac{1}{2}(m_1 + m_2)V_{CM}^2}_{E_{CM}} + \underbrace{\frac{1}{2}\mu\dot{r}^2 + \frac{\ell^2}{2\mu r^2} + V(r)}_{E_{rel}}$$

$$E_{rel} = \frac{1}{2}\mu\dot{r}^2 + \frac{E_{rel}b^2}{r^2} + V(r)$$

In what situations do particles undergo inelastic scattering, rather than elastic scattering?

Comment – elastic scattering means $E_{\text{initial}} = E_{\text{final}}$

Typically, elastic scattering occurs when two fundamental particles interact (as long as the final kinetic energy of both particles is taken into account).

Elastically bouncing ball

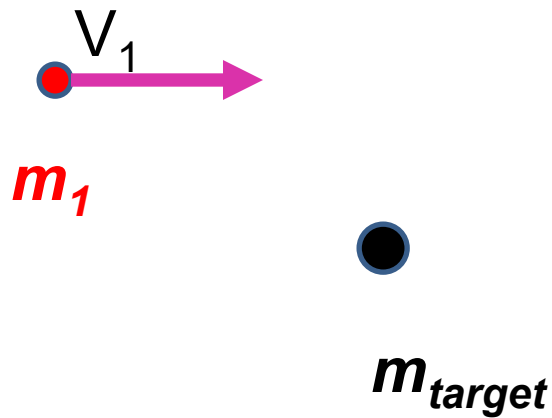


Inelastically collision

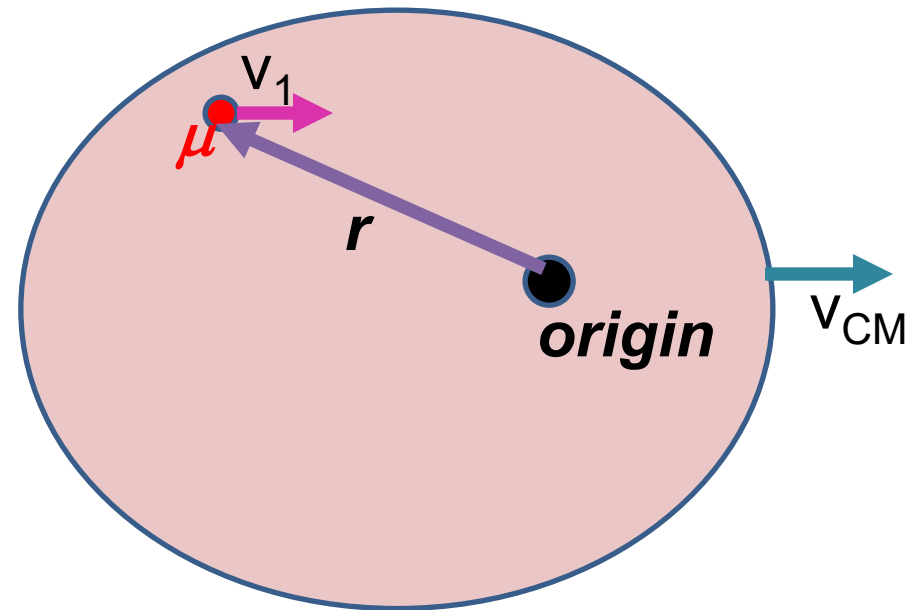


Note: The following analysis will be carried out in the center of mass frame of reference.

In laboratory frame:



In center-of-mass frame:



$$\mu = \frac{m_1 m_{\text{target}}}{m_1 + m_{\text{target}}}$$

$$\ell = |\mathbf{r} \times \mu \mathbf{v}_1|$$

Typically, the laboratory frame is where the data is taken, but the center of mass frame is where the analysis is most straightforward.

Previous equations --

$$E = \frac{1}{2}(m_1 + m_2)V_{CM}^2 + \frac{1}{2}\mu\dot{r}^2 + \frac{\ell^2}{2\mu r^2} + V(r)$$



constant



relative coordinate system;
visualize as “in” CM frame

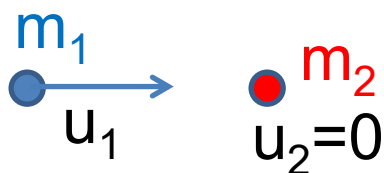


It is often convenient to analyze the scattering cross section in the center of mass reference frame.

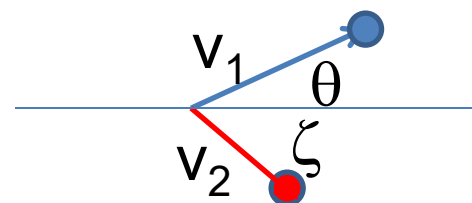
Relationship between normal laboratory reference and center of mass:

Laboratory reference frame:

Before

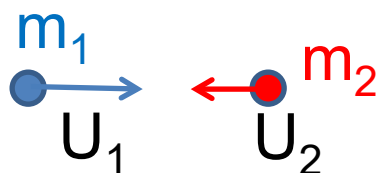


After

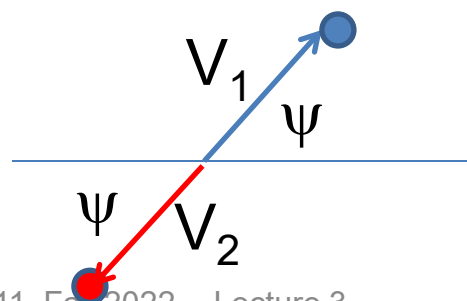


Center of mass reference frame:

Before



After





Relationship between center of mass and laboratory frames of reference -- continued

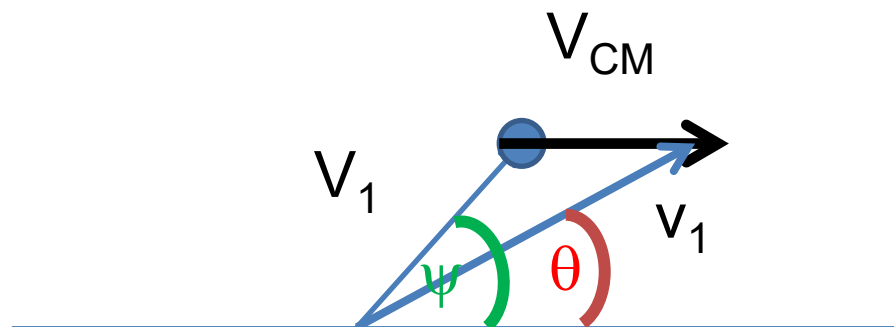
Since m_2 is initially at rest in lab frame:

$$\mathbf{V}_{CM} = \frac{m_1}{m_1 + m_2} \mathbf{u}_1 \quad \mathbf{u}_1 = \mathbf{U}_1 + \mathbf{V}_{CM} \quad \Rightarrow \quad \mathbf{U}_1 = \frac{m_2}{m_1 + m_2} \mathbf{u}_1 = \frac{m_2}{m_1} \mathbf{V}_{CM}$$
$$\mathbf{u}_2 = \mathbf{U}_2 + \mathbf{V}_{CM} \quad \Rightarrow \quad \mathbf{U}_2 = -\frac{m_1}{m_1 + m_2} \mathbf{u}_1 = -\mathbf{V}_{CM}$$

$$\mathbf{v}_1 = \mathbf{V}_1 + \mathbf{V}_{CM}$$

$$\mathbf{v}_2 = \mathbf{V}_2 + \mathbf{V}_{CM}$$

Relationship between center of mass and laboratory frames of reference for the scattering particle 1



$$\mathbf{v}_1 = \mathbf{V}_1 + \mathbf{V}_{CM}$$

$$v_1 \sin \theta = V_1 \sin \psi$$

$$v_1 \cos \theta = V_1 \cos \psi + V_{CM}$$

$$\tan \theta = \frac{\sin \psi}{\cos \psi + V_{CM} / V_1} = \frac{\sin \psi}{\cos \psi + m_1 / m_2}$$

For elastic scattering

Digression – elastic scattering

$$\begin{aligned} \frac{1}{2} m_1 U_1^2 + \frac{1}{2} m_2 U_2^2 + \frac{1}{2} (m_1 + m_2) V_{CM}^2 \\ = \frac{1}{2} m_1 V_1^2 + \frac{1}{2} m_2 V_2^2 + \frac{1}{2} (m_1 + m_2) V_{CM}^2 \end{aligned}$$

Also note:

$$m_1 \mathbf{U}_1 + m_2 \mathbf{U}_2 = 0 \qquad m_1 \mathbf{V}_1 + m_2 \mathbf{V}_2 = 0$$

$$\mathbf{U}_1 = \frac{m_2}{m_1} \mathbf{V}_{CM} \qquad \mathbf{U}_2 = -\mathbf{V}_{CM}$$

$$\Rightarrow |\mathbf{U}_1| = |\mathbf{V}_1| \quad \text{and} \quad |\mathbf{U}_2| = |\mathbf{V}_2| = |\mathbf{V}_{CM}|$$

$$\text{Also note that: } m_1 |\mathbf{U}_1| = m_2 |\mathbf{U}_2|$$

$$\text{So that: } V_{CM}/V_1 = V_{CM}/U_1 = m_1/m_2$$

Summary of results --

$$\mathbf{v}_1 = \mathbf{V}_1 + \mathbf{V}_{CM}$$

$$v_1 \sin \theta = V_1 \sin \psi$$

$$v_1 \cos \theta = V_1 \cos \psi + V_{CM}$$

$$\tan \theta = \frac{\sin \psi}{\cos \psi + V_{CM} / V_1} = \frac{\sin \psi}{\cos \psi + m_1 / m_2}$$



General case



Special case of
elastic scattering

For elastic scattering

$$V_{CM} / V_1 = m_1 / m_2$$

Relationship between center of mass and laboratory frames of reference – continued (elastic scattering)

$$\mathbf{v}_1 = \mathbf{V}_1 + \mathbf{V}_{CM}$$

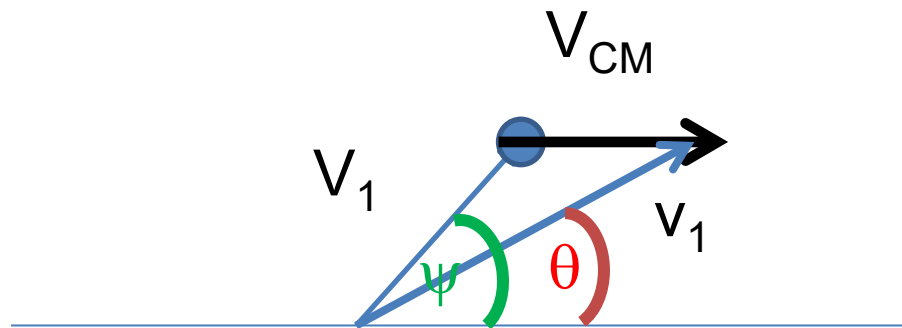
$$v_1 \sin \theta = V_1 \sin \psi$$

$$v_1 \cos \theta = V_1 \cos \psi + V_{CM}$$

$$\tan \theta = \frac{\sin \psi}{\cos \psi + V_{CM} / V_1} = \frac{\sin \psi}{\cos \psi + m_1 / m_2}$$

Also:

$$\cos \theta = \frac{\cos \psi + m_1 / m_2}{\sqrt{1 + 2m_1 / m_2 \cos \psi + (m_1 / m_2)^2}}$$

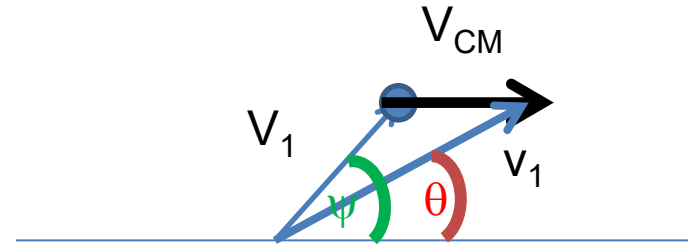


More details -- from the diagram and equations --

$$\mathbf{v}_1 = \mathbf{V}_1 + \mathbf{V}_{CM}$$

$$v_1 \sin \theta = V_1 \sin \psi$$

$$v_1 \cos \theta = V_1 \cos \psi + V_{CM}$$



Take the dot product of the first equation with itself

$$v_1^2 = V_1^2 + 2V_1V_{CM} \cos \psi + V_{CM}^2$$

$$\text{or } \frac{v_1}{V_1} = \sqrt{1 + 2 \frac{V_{CM}}{V_1} \cos \psi + \frac{V_{CM}^2}{V_1^2}} = \sqrt{1 + 2 \frac{m_1}{m_2} \cos \psi + \left(\frac{m_1}{m_2} \right)^2}$$

$$\Rightarrow \cos \theta = \frac{\cos \psi + m_1 / m_2}{\sqrt{1 + 2m_1 / m_2 \cos \psi + (m_1 / m_2)^2}}$$

Differential cross sections in different reference frames

$$\left(\frac{d\sigma_{LAB}(\theta)}{d\Omega_{LAB}} \right) = \left(\frac{d\sigma_{CM}(\psi)}{d\Omega_{CM}} \right) \frac{d\Omega_{CM}}{d\Omega_{LAB}}$$

$$\frac{d\Omega_{CM}}{d\Omega_{LAB}} = \left| \frac{\sin \psi}{\sin \theta} \frac{d\psi}{d\theta} \right| = \left| \frac{d \cos \psi}{d \cos \theta} \right|$$

Using:

$$\cos \theta = \frac{\cos \psi + m_1 / m_2}{\sqrt{1 + 2(m_1 / m_2) \cos \psi + (m_1 / m_2)^2}}$$

$$\left| \frac{d \cos \theta}{d \cos \psi} \right| = \frac{(m_1 / m_2) \cos \psi + 1}{\left(1 + 2(m_1 / m_2) \cos \psi + (m_1 / m_2)^2 \right)^{3/2}}$$

Differential cross sections in different reference frames – continued:

$$\left(\frac{d\sigma_{LAB}(\theta)}{d\Omega_{LAB}} \right) = \left(\frac{d\sigma_{CM}(\psi)}{d\Omega_{CM}} \right) \left| \frac{d\cos\psi}{d\cos\theta} \right|$$

$$\left(\frac{d\sigma_{LAB}(\theta)}{d\Omega_{LAB}} \right) = \left(\frac{d\sigma_{CM}(\psi)}{d\Omega_{CM}} \right) \frac{\left(1 + 2m_1/m_2 \cos\psi + (m_1/m_2)^2 \right)^{3/2}}{(m_1/m_2)\cos\psi + 1}$$

where: $\tan\theta = \frac{\sin\psi}{\cos\psi + m_1/m_2}$

$$\left(\frac{d\sigma_{LAB}(\theta)}{d\Omega_{LAB}} \right) = \left(\frac{d\sigma_{CM}(\psi)}{d\Omega_{CM}} \right) \frac{\left(1 + 2m_1/m_2 \cos\psi + (m_1/m_2)^2 \right)^{3/2}}{(m_1/m_2)\cos\psi + 1}$$

where: $\tan\theta = \frac{\sin\psi}{\cos\psi + m_1/m_2}$

Example: suppose $m_1 = m_2$

In this case: $\tan\theta = \frac{\sin\psi}{\cos\psi + 1} \Rightarrow \theta = \frac{\psi}{2}$

note that $0 \leq \theta \leq \frac{\pi}{2}$

$$\left(\frac{d\sigma_{LAB}(\theta)}{d\Omega_{LAB}} \right) = \left(\frac{d\sigma_{CM}(2\theta)}{d\Omega_{CM}} \right) \cdot 4 \cos\theta$$

Summary --

Differential cross sections in different reference frames – continued:

$$\left(\frac{d\sigma_{LAB}(\theta)}{d\Omega_{LAB}} \right) = \left(\frac{d\sigma_{CM}(\psi)}{d\Omega_{CM}} \right) \left| \frac{d\cos\psi}{d\cos\theta} \right|$$

$$\left(\frac{d\sigma_{LAB}(\theta)}{d\Omega_{LAB}} \right) = \left(\frac{d\sigma_{CM}(\psi)}{d\Omega_{CM}} \right) \frac{\left(1 + 2m_1/m_2 \cos\psi + (m_1/m_2)^2 \right)^{3/2}}{(m_1/m_2)\cos\psi + 1}$$

where: $\tan\theta = \frac{\sin\psi}{\cos\psi + m_1/m_2}$ For elastic scattering

Hard sphere example – continued

$$m_1 = m_2$$

Center of mass frame

Lab frame

$$\left(\frac{d\sigma_{CM}(\psi)}{d\Omega_{CM}} \right) = \frac{D^2}{4}$$

$$\left(\frac{d\sigma_{LAB}(\theta)}{d\Omega_{LAB}} \right) = D^2 \cos\theta \quad \theta = \frac{\psi}{2}$$

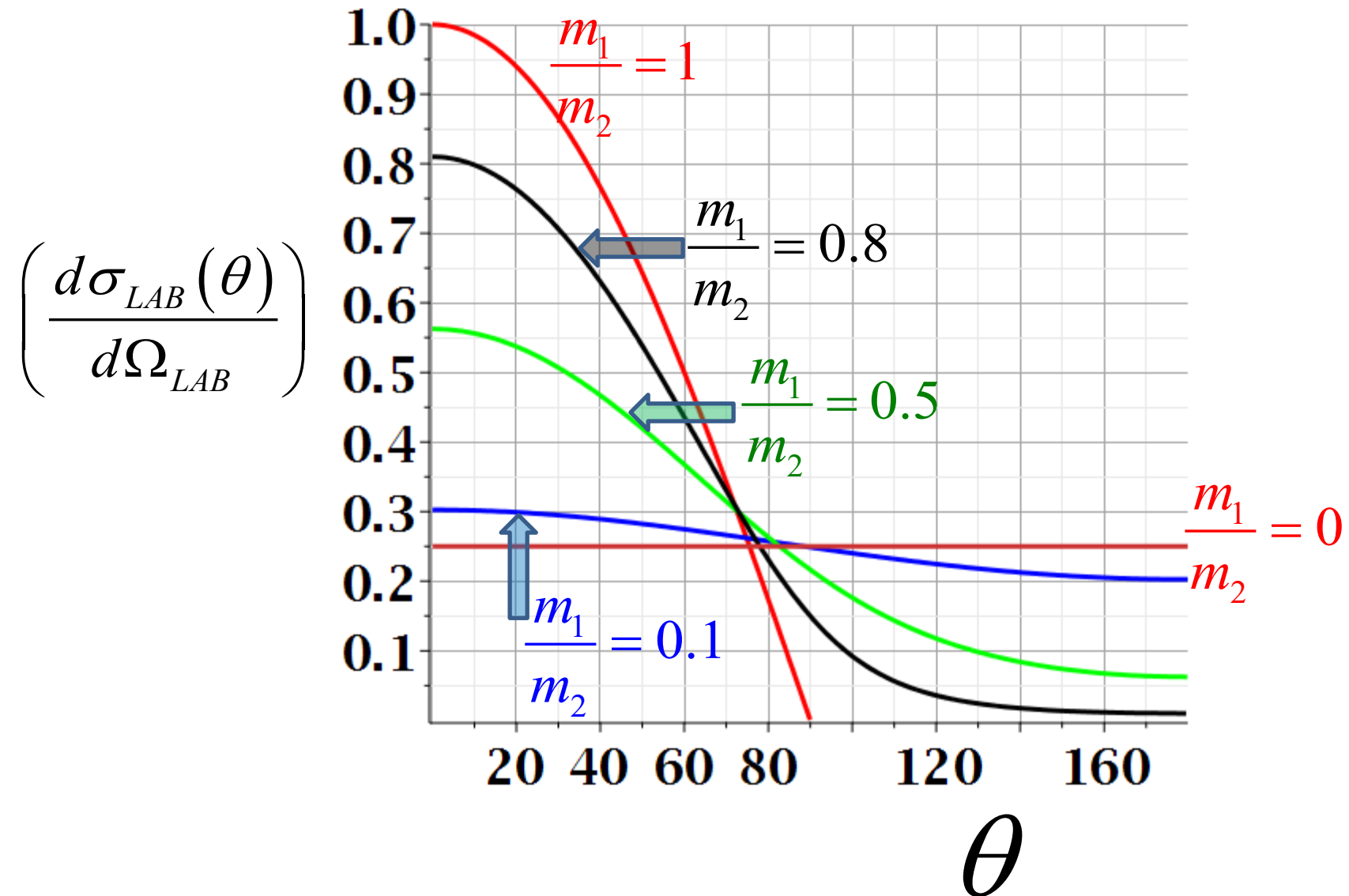
$$\int \frac{d\sigma_{CM}(\psi)}{d\Omega_{CM}} d\Omega_{CM} =$$

$$\int \frac{d\sigma_{lab}(\theta)}{d\Omega_{lab}} d\Omega_{lab} =$$

$$\frac{D^2}{4} 4\pi = \pi D^2$$

$$2\pi D^2 \int_0^{\pi/2} \cos\theta \sin\theta d\theta = \pi D^2$$

Scattering cross section for hard sphere in lab frame for various mass ratios:



For visualization, is convenient to make a "parametric" plot of

$$\left(\frac{d\sigma_{LAB}}{d\Omega}(\theta) \right) \text{ vs } \theta(\psi)$$

$$\left(\frac{d\sigma_{LAB}(\theta)}{d\Omega_{LAB}} \right) = \left(\frac{d\sigma_{CM}(\psi)}{d\Omega_{CM}} \right) \frac{\left(1 + 2m_1 / m_2 \cos\psi + (m_1 / m_2)^2 \right)^{3/2}}{(m_1 / m_2) \cos\psi + 1}$$

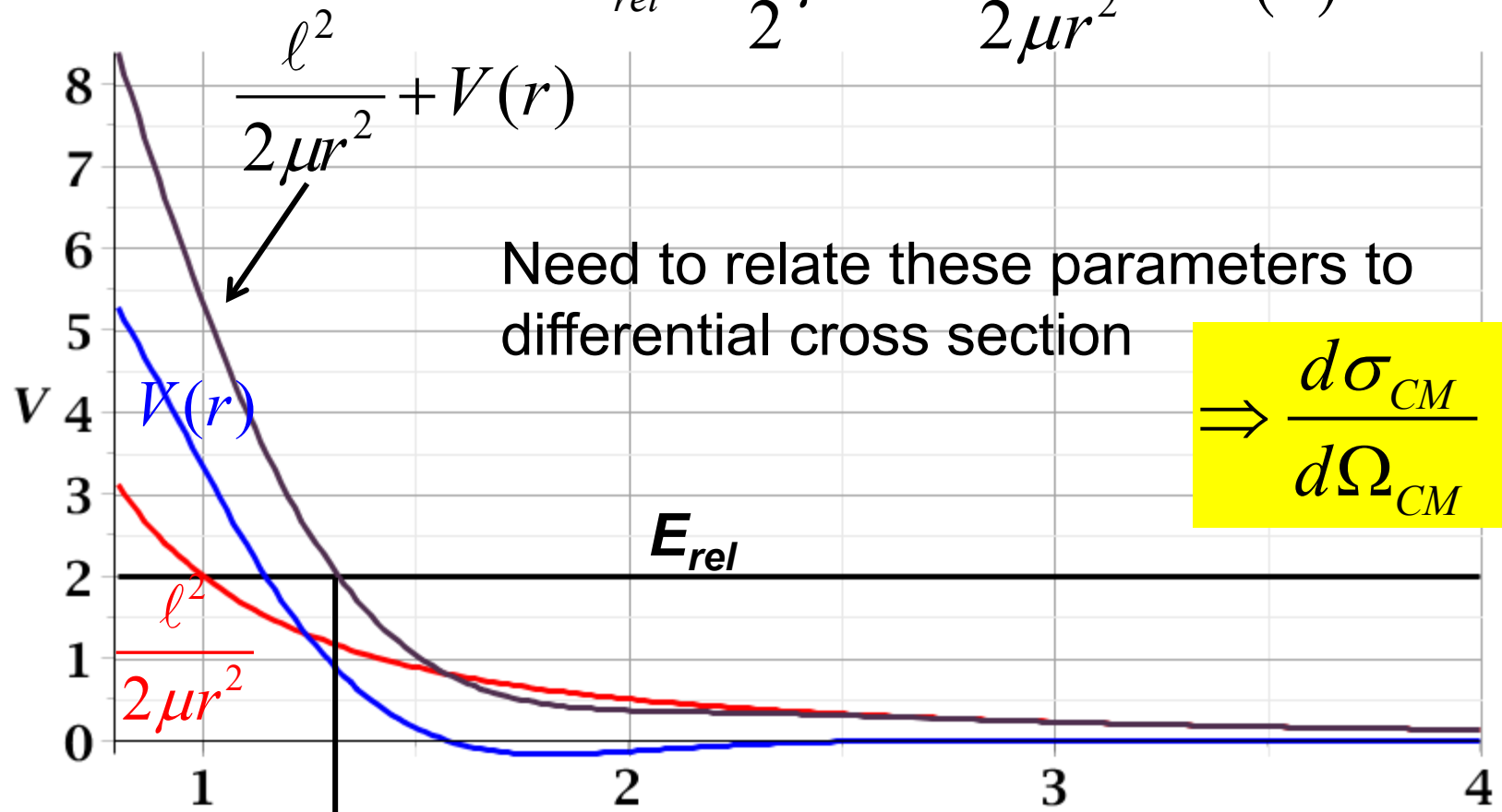
where: $\tan \theta = \frac{\sin \psi}{\cos \psi + m_1 / m_2}$

Maple syntax:

```
> plot( { [psi(theta, 0), sigma(theta, 0), theta = 0.001 ..3.14], [psi(theta, .1), sigma(theta, .1), theta = 0.001 ..3.14], [psi(theta, .5), sigma(theta, .5), theta = 0.001 ..3.14], [psi(theta, .8), sigma(theta, .8), theta = 0.001 ..3.14], [psi(theta, 1), sigma(theta, 1), theta = 0.001 ..3.14] }, thickness = 3, font = ['Times','bold', 24], gridlines = true, color = [red, blue, green, black, orange])
```

For a continuous potential interaction in center of mass reference frame:

$$E_{rel} = \frac{1}{2} \mu \dot{r}^2 + \frac{\ell^2}{2\mu r^2} + V(r)$$



$$\mu \equiv \frac{m_1 m_2}{m_1 + m_2} r$$

ℓ = angular momentum