

## Announcements

### 1. **Third hour exam – Wednesday, March 23, 2005 – covering Chapters 33-34.**

- 4 problems – show your work and reasoning for possible partial credit.
- May bring one 8½” x 11” sheet of paper to the exam (to be turned in with your exam papers).
- Should also bring clear head

### 2. Example exam available on [website](#).

### 3. Today's lecture –

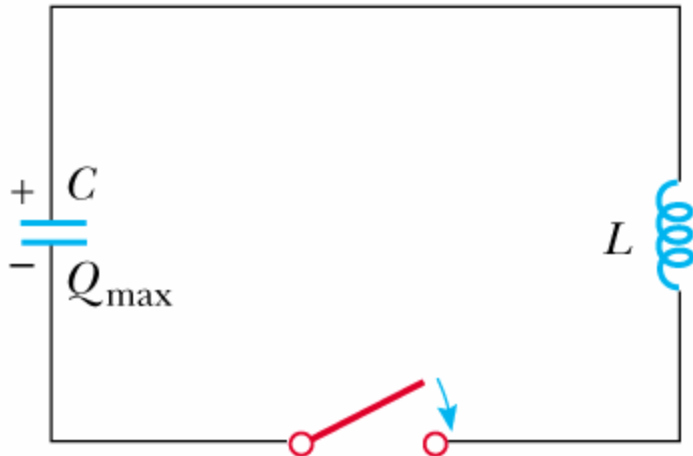
Advice for studying

Systematic review

## List of topics

1. Circuits with L, C, and R components and emf source (either constant or oscillating in time)
  - a. Current-voltage relationships
  - b. Transient and steady-state effects
  - c. Frequency dependent effects
  - d. Energy stored or dissipated
2. Electromagnetic waves as a consequence of Maxwell's equations
  - a. Relationship between Maxwell's equations and the wave equation
  - b. Vector and functional relationships of  $\mathbf{E}$ ,  $\mathbf{B}$ , and  $\mathbf{k}$  and  $c$ .
  - c. Poynting vector, intensity, and power.
  - d. Refractive index
  - e. Reflection and refraction of EM waves at surfaces

# LC circuit



**Resonant frequency:**

$$\omega_0 = \sqrt{\frac{1}{LC}} = 2\pi f_0 = \frac{2\pi}{T_0}$$

$$-\frac{q}{C} - L \frac{d^2 q}{dt^2} = 0$$

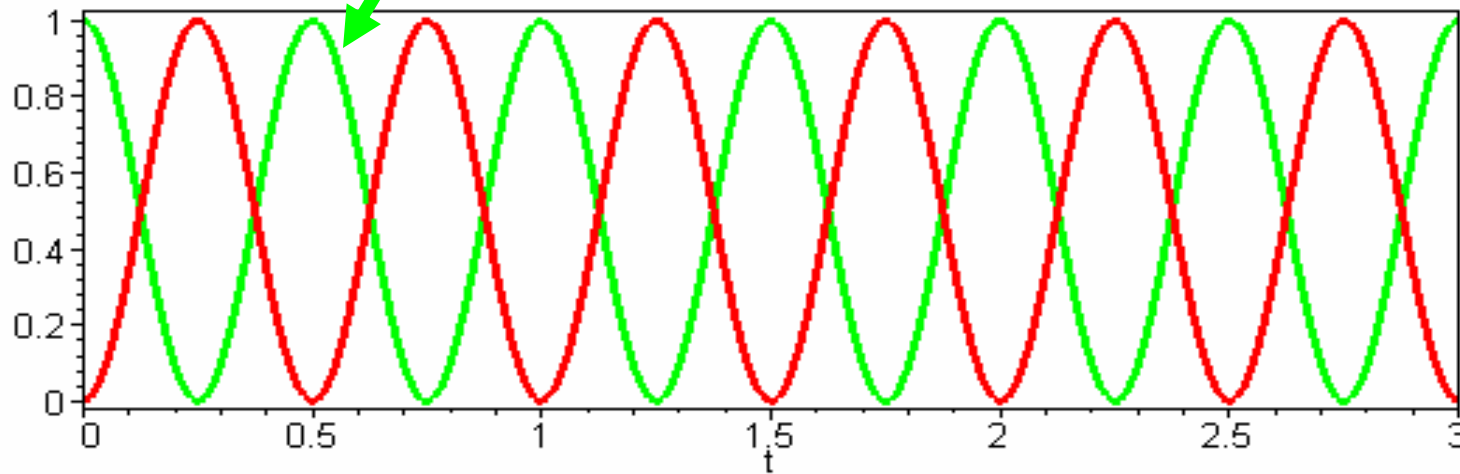
$$q(t) = Q_{\max} \cos(\omega_0 t + \varphi)$$

$$i(t) = \frac{dq}{dt} = -\underbrace{\omega_0 Q_{\max}}_{I_{\max}} \sin(\omega_0 t + \varphi)$$

# Energy in LC circuits

Energy in capacitor :  $U_E = \frac{q^2}{2C} = \frac{Q_{\max}^2 \cos^2(\omega_0 t + \varphi)}{2C}$

Energy in inductor :  $U_B = \frac{Li^2}{2} = \frac{\omega_0^2 L Q_{\max}^2 \sin^2(\omega_0 t + \varphi)}{2}$

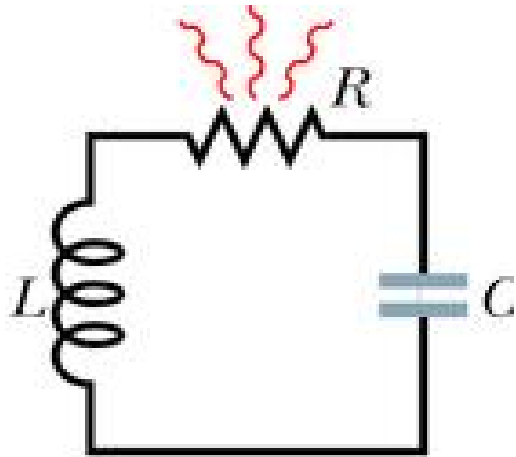


## Total energy in LC circuit:

$$\begin{aligned}U_E + U_B &= \frac{q^2}{2C} + \frac{Li^2}{2} \\&= \frac{Q_{\max}^2}{2C} (\cos^2(\omega_0 t + \varphi) + \sin^2(\omega_0 t + \varphi)) \\&= \frac{Q_{\max}^2}{2C} = \frac{LI_{\max}^2}{2}\end{aligned}$$

$$\text{Note : } \omega_0^2 L = \frac{1}{LC} L = \frac{1}{C}$$

## Effects of resistance in the circuit



Solution :

$$-\frac{q}{C} - L \frac{di}{dt} - Ri = 0$$

$$-\frac{1}{C} q - L \frac{d^2 q}{dt^2} - R \frac{dq}{dt} = 0$$

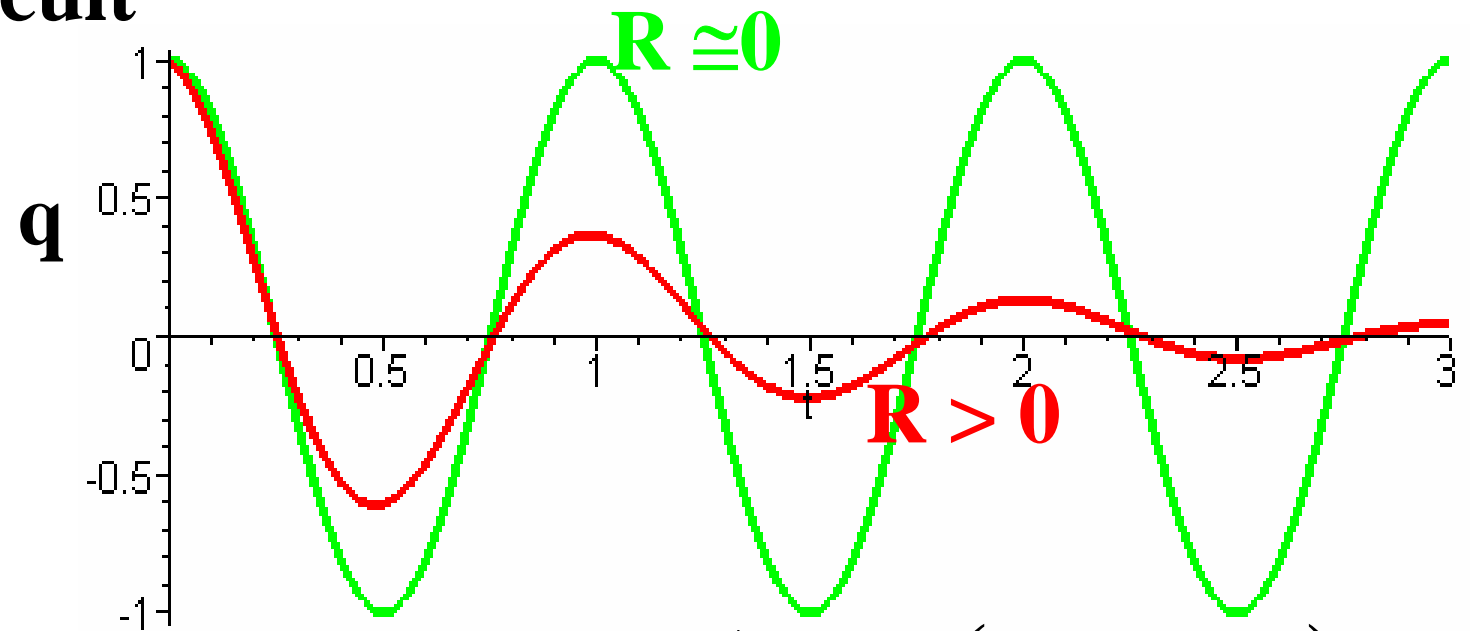
$$q(t) = Q_{\max} e^{-t/\tau} \cos(\omega_0 t + \varphi)$$

**constants**

$$\tau = \frac{2L}{R}$$

$$\omega_0 = \sqrt{\omega_0^2 - \frac{1}{\tau^2}}$$

# LCR circuit

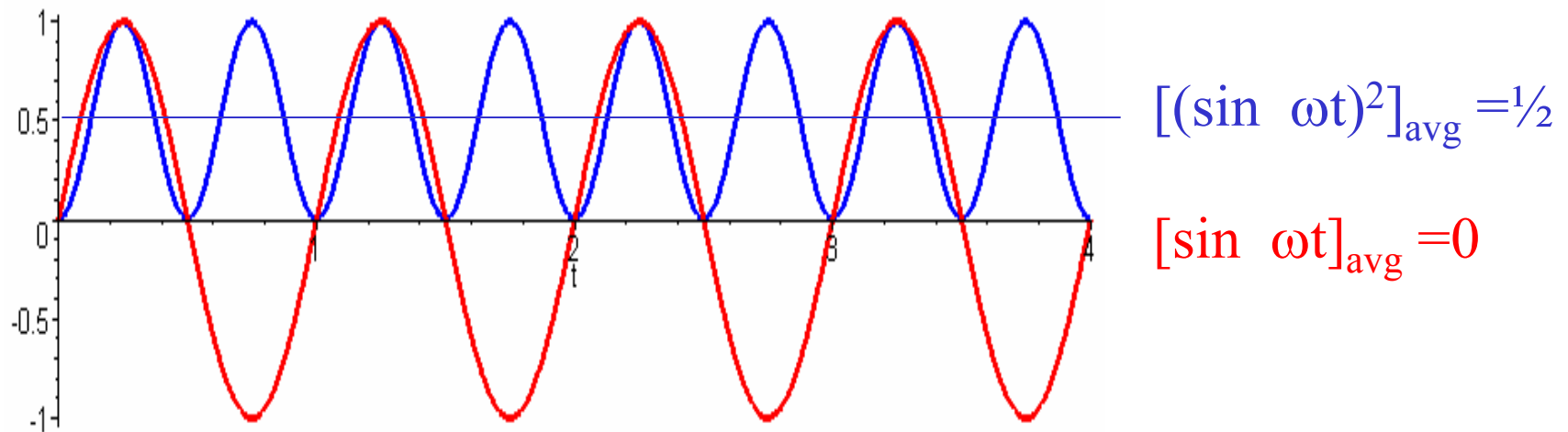


$$q(t) = Q_{\max} e^{-t/\tau} \cos(\varpi_0 t + \varphi)$$

$$\tau = \frac{2L}{R} \quad \varpi_0 = \sqrt{\omega_0^2 - \frac{1}{\tau^2}}$$

## Properties of AC circuits

$$\mathcal{E} = \mathcal{E}_{\max} \sin \omega t \quad \text{or} \quad \mathcal{E}_{\max} \cos \omega t$$



$$\langle |\mathcal{E}|^2 \rangle_{\text{avg}} = \frac{1}{2} |\mathcal{E}_{\max}|^2$$

$$\mathcal{E}_{\text{rms}} = \sqrt{\langle |\mathcal{E}|^2 \rangle_{\text{avg}}}$$

$$\mathcal{E}_{\text{rms}} = \sqrt{\frac{1}{2}} \mathcal{E}_{\max}$$

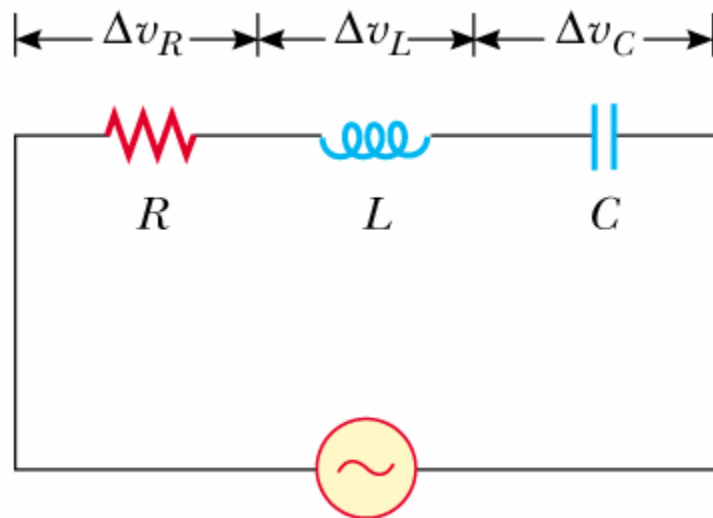
similarly,

$$I_{\text{rms}} = \sqrt{\frac{1}{2}} I_{\max}$$



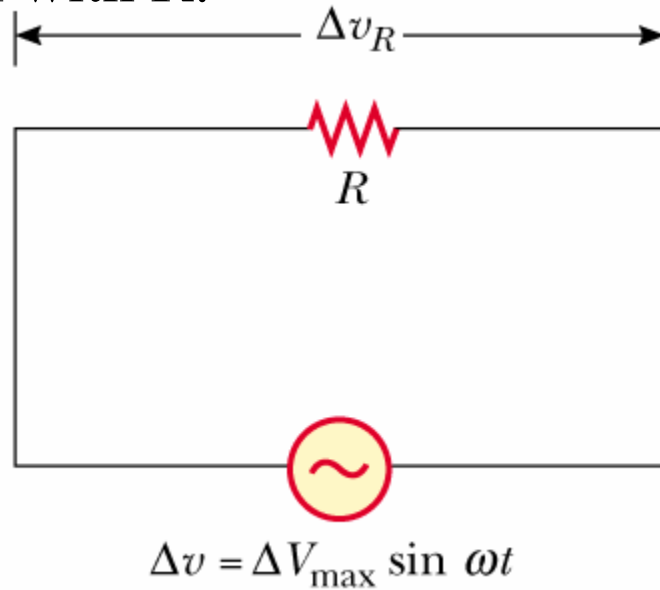
DC circuits	AC circuits
<p><math>\mathcal{E} = \text{constant in time}</math></p> <p> <math>I = \begin{cases} \text{constant in time} \\ \text{constant} + I_0 e^{-t/\tau} \\ \text{damped oscillations} \end{cases}</math> </p> <p>Kirchhoff's rules apply</p>	<p><math>\mathcal{E} = \mathcal{E}_{\text{max}} \sin \omega t \text{ or } \mathcal{E}_{\text{max}} \cos \omega t</math></p> <p><math>I = \text{transients} + I_0 \sin (\omega t - \phi)</math></p> <p>Kirchhoff's rules apply</p>

## Analysis of LCR circuits with AC emf:



Consider each component separately with AC emf --

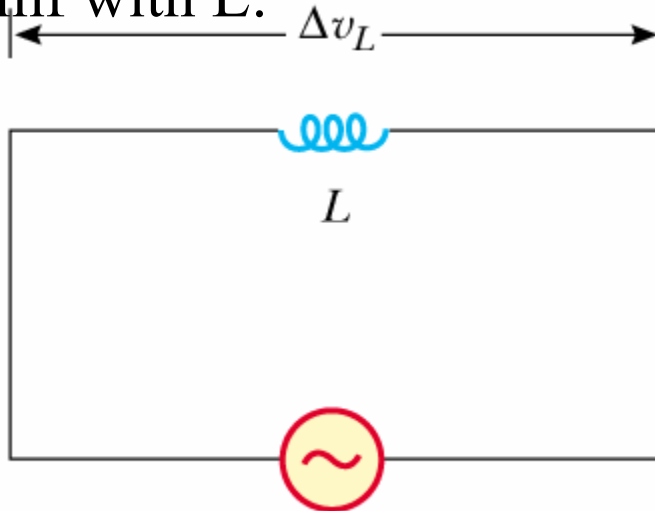
AC emf with R:



$$-IR + \Delta V_{\max} \sin \omega t = 0 \quad \text{Note: } \mathcal{E}_{\max} \leftrightarrow \Delta V_{\max}$$

$$I = \frac{\Delta V_{\max}}{R} \sin \omega t$$

AC emf with L:

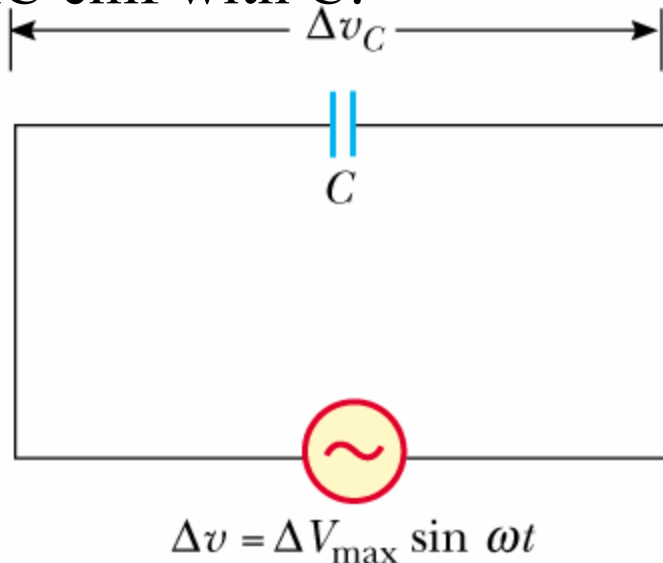


$$\Delta v = \Delta V_{\max} \sin \omega t$$
$$-L \frac{dI}{dt} + \Delta V_{\max} \sin \omega t = 0$$

$$\frac{dI}{dt} = \frac{\Delta V_{\max}}{L} \sin \omega t$$

$$I = -\frac{\Delta V_{\max}}{\omega L} \cos \omega t$$

AC emf with C:

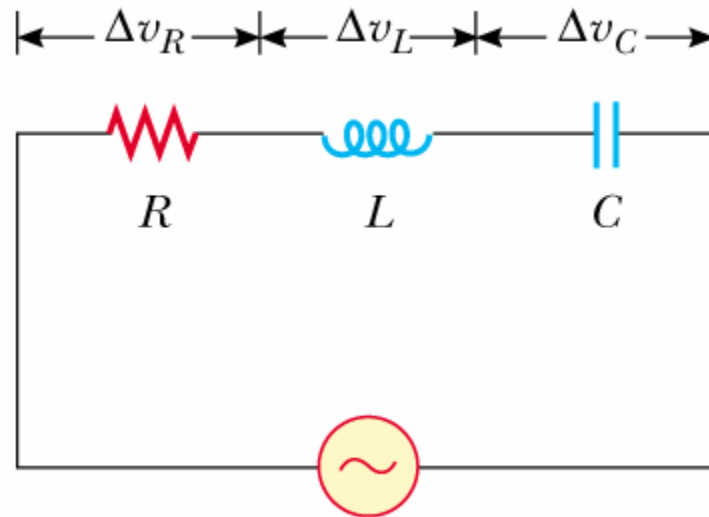


$$-\frac{q}{C} + \Delta V_{\max} \sin \omega t = 0$$

$$q = C \Delta V_{\max} \sin \omega t$$

$$I = \omega C \Delta V_{\max} \cos \omega t$$

AC emf with series LCR circuit:



Differential eq:  $-RI - L \frac{dI}{dt} - \frac{q}{C} + \mathcal{E}_{\max} \sin \omega t = 0$

Steady-state term

Transient term

Solution for  $I(t)$ :  $I = I_{\max} \sin(\omega t - \varphi) + K e^{-t/\tau} \cos(\omega_0 t + \alpha)$

$$\tau = \frac{2L}{R}$$

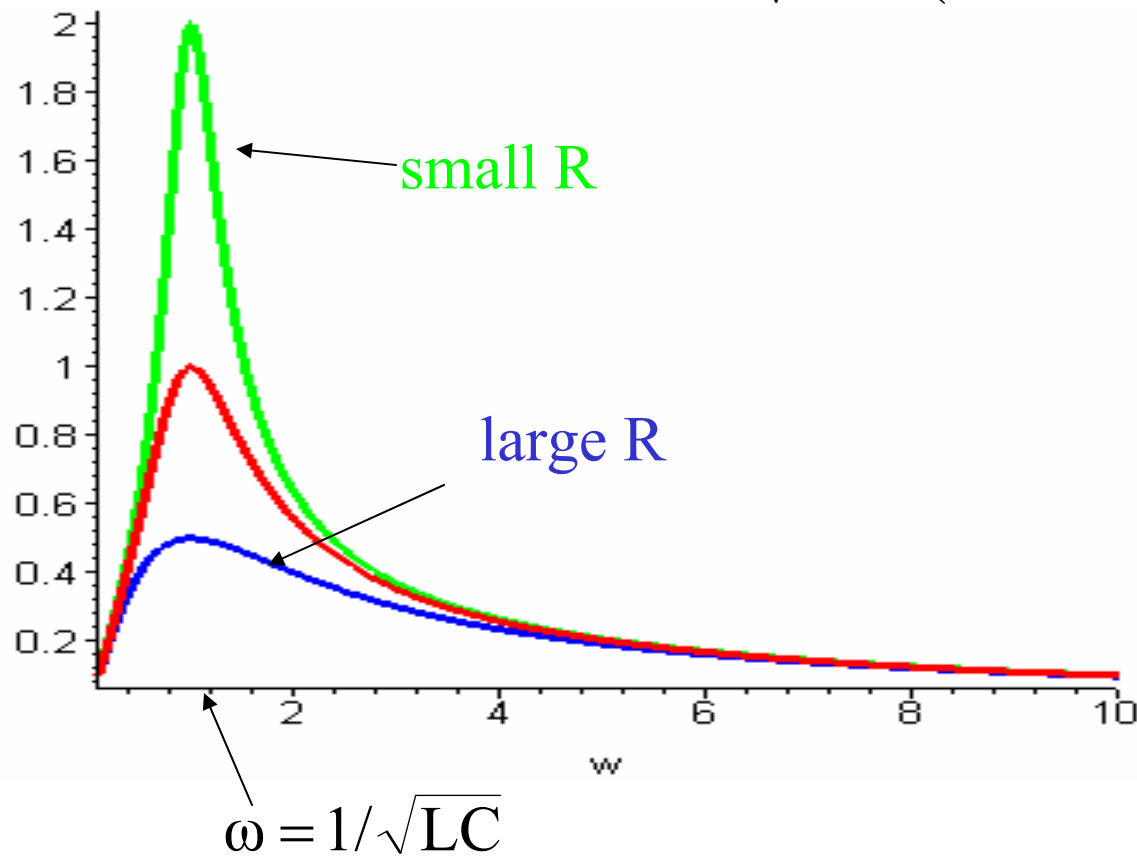
where:  $I_{\max} = \frac{\mathcal{E}_{\max}}{Z}$   $Z \equiv \sqrt{R^2 + (\omega L - 1/\omega C)^2}$

$$\omega_0 = \sqrt{\omega_0^2 - 1/\tau^2}$$

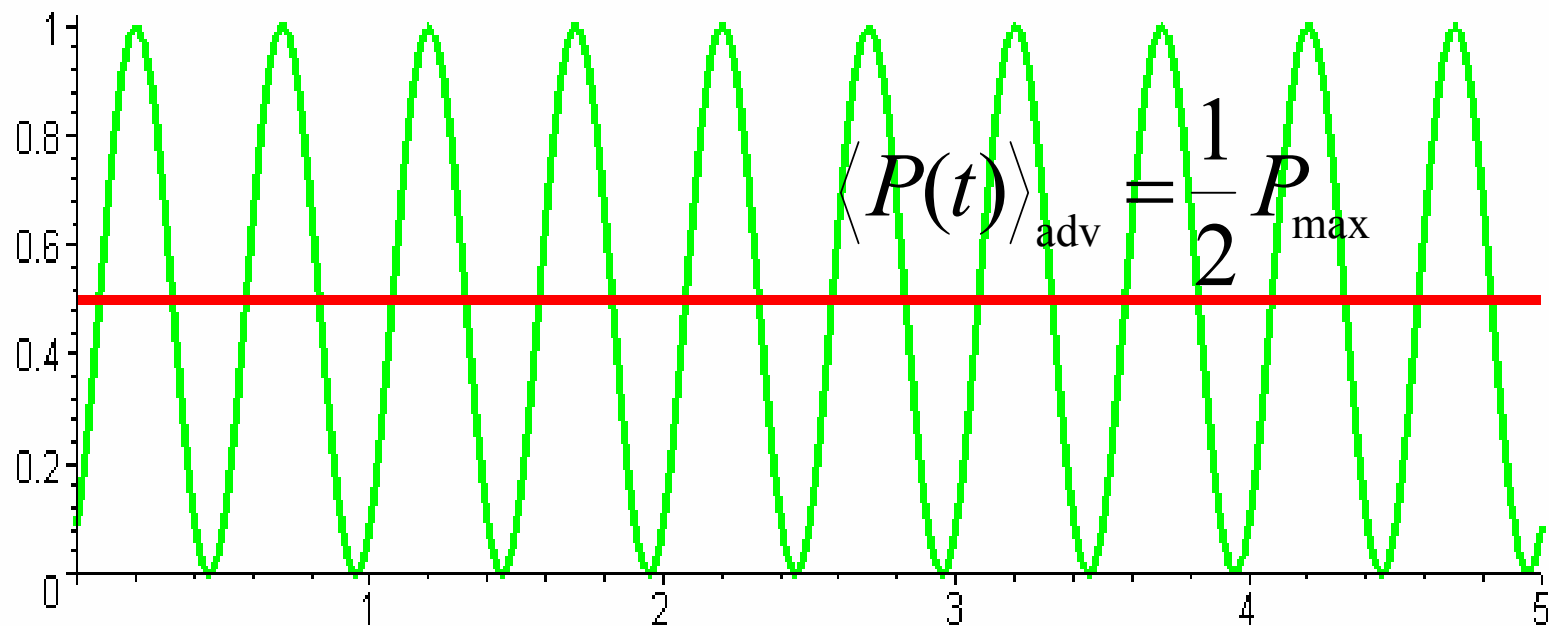
$$\tan \varphi = \frac{\omega L - 1/\omega C}{R}$$

Behavior of  $I_{\max}$  in steady-state term --

$$I_{\max} = \frac{\Delta V_{\max}}{\sqrt{R^2 + (\omega L - 1/\omega C)^2}}$$



Power in AC circuit:  $P = I^2 R = (I_{\max} \sin(\omega t - \phi))^2 R$

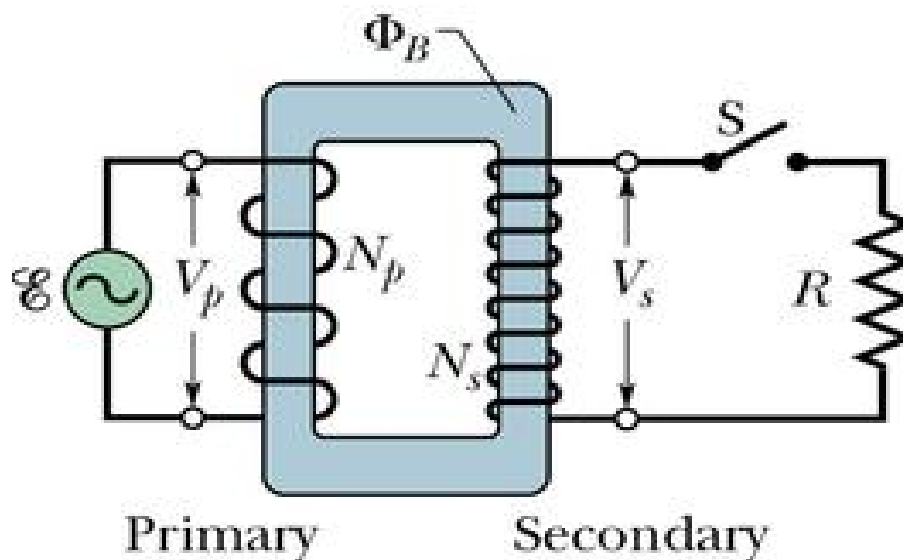


$$\langle P(t) \rangle_{\text{adv}} = \frac{1}{2} P_{\max} = \frac{1}{2} I_{\max}^2 R = \frac{1}{2} \frac{\mathcal{E}_{\max}^2}{Z^2} R$$

$$I_{\text{rms}} = \sqrt{\frac{1}{2} I_{\max}^2} = \sqrt{\frac{1}{2}} I_{\max}$$



## AC transformer



Faraday's law for single wire loop :  $\mathcal{E} = -\frac{d\Phi_B}{dt}$

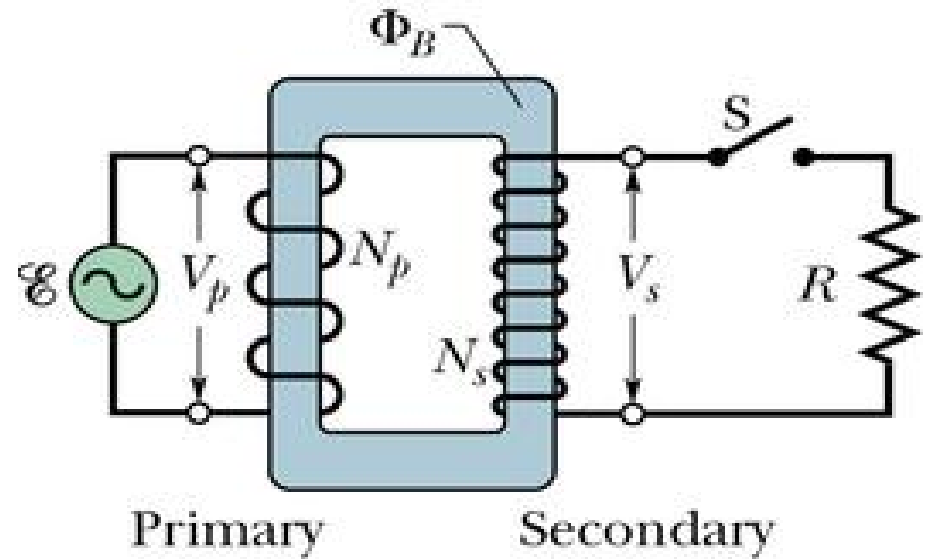
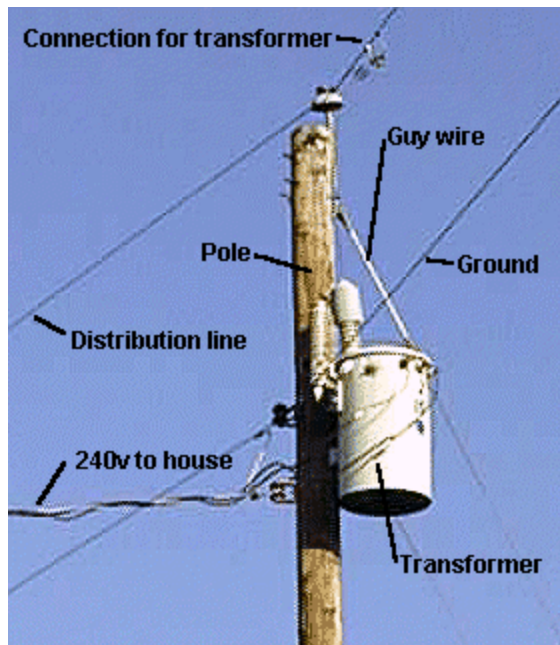
For transformer :

$$\mathcal{E}_p = -N_p \frac{d\Phi_B}{dt} \qquad \mathcal{E}_s = -N_s \frac{d\Phi_B}{dt}$$

$$-\frac{d\Phi_B}{dt} = \frac{\mathcal{E}_p}{N_p} = \frac{\mathcal{E}_s}{N_s} \qquad \Rightarrow \mathcal{E}_s = \frac{N_s}{N_p} \mathcal{E}_p$$

Practical use of transformer:

$$\mathcal{E}_S = \frac{N_S}{N_P} \mathcal{E}_P$$



Example:

$$\mathcal{E}_S = \frac{N_S}{N_P} \mathcal{E}_P = \frac{100}{6000} 7200V = 120V$$

<http://science.howstuffworks.com/power9.htm>

The mathematical form of the “wave equation”:

$$\frac{\partial^2 y}{\partial t^2} = v^2 \frac{\partial^2 y}{\partial x^2}$$

velocity of wave

Solution form for a periodic solution:

$$y(x, t) = A \sin\left(\frac{2\pi}{\lambda}(x - vt) + \varphi\right) \equiv A \sin\left(\frac{2\pi x}{\lambda} - 2\pi ft + \varphi\right) = A \sin(kx - \omega t + \varphi)$$

amplitude

phase

$$v = \lambda f = \frac{\omega}{k}$$

## Comparison of mechanical and electromagnetic waves

Mechanical	Electromagnetic
<p>Satisfy wave equation</p> <p><math>v</math> depends upon propagation material.</p> <p>Can be transverse or longitudinal</p> <p>Can only propagate within materials (solids, liquids, gases, strings, etc.)</p> <p>Doppler effect :</p> $f' = f \frac{1 \pm u_O / v}{1 \mp u_S / v}$	<p>Satisfy wave equation</p> <p><math>v</math> depends upon propagation material (or vacuum).</p> <p>Can only be transverse</p> <p>Can propagate within a vacuum and within some materials.</p> <p>Doppler effect :</p> $f' = f \sqrt{\frac{1 + u / v}{1 - u / v}}$

# Maxwell's equations:

$$\oiint \mathbf{E} \cdot d\mathbf{A} = Q / \epsilon_0$$

**Coulomb-Gauss law**

$$\oiint \mathbf{B} \cdot d\mathbf{A} = 0$$

**Gauss's for magnetic field**

$$\oint \mathbf{B} \cdot d\mathbf{s} = \mu_0 I + \mu_0 \epsilon_0 \frac{d\Phi_E}{dt}$$

**Ampere-Maxwell law**

$$\oint \mathbf{E} \cdot d\mathbf{s} = -\frac{d\Phi_B}{dt}$$

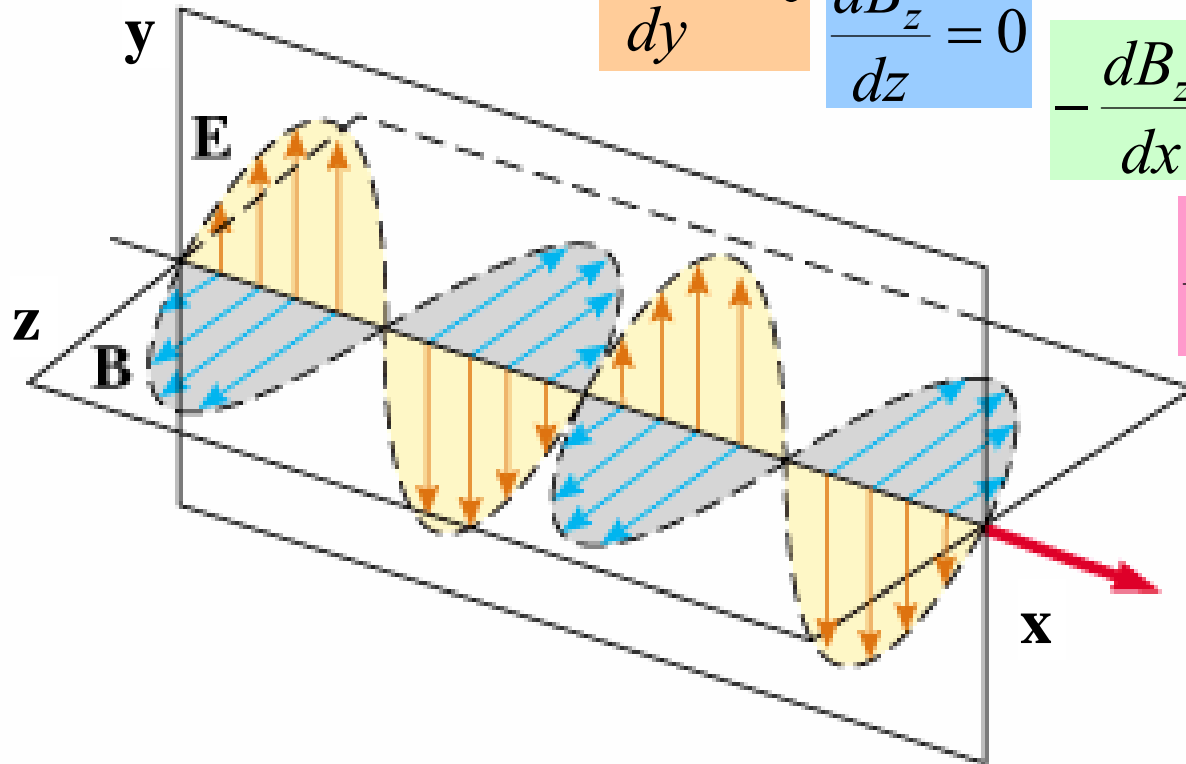
**Faraday's law**

Maxwell's equations far away from sources ( $Q=0$ ,  $I=0$ ):

$$\oiint \mathbf{E} \cdot d\mathbf{A} = 0 \quad \oiint \mathbf{B} \cdot d\mathbf{A} = 0 \quad \oint \mathbf{B} \cdot d\mathbf{s} = \mu_0 \epsilon_0 \frac{d(\oiint \mathbf{E} \cdot d\mathbf{A})}{dt} \quad \oint \mathbf{E} \cdot d\mathbf{s} = -\frac{d(\oiint \mathbf{B} \cdot d\mathbf{A})}{dt}$$

Plane-polarized solution:

$$\frac{dE_y}{dy} = 0 \quad \frac{dB_z}{dz} = 0 \quad -\frac{dB_z}{dx} = \mu_0 \epsilon_0 \frac{dE_y}{dt} \quad \frac{dE_y}{dx} = -\frac{dB_z}{dt}$$



Results:

$$-\frac{dB_z}{dx} = \mu_0 \epsilon_0 \frac{dE_y}{dt}$$

$$\frac{dE_y}{dx} = -\frac{dB_z}{dt}$$

Wave equations for  $E_y$  and  $B_z$ :

$$\frac{\partial^2 E_y}{\partial t^2} = \frac{1}{\mu_0 \epsilon_0} \frac{\partial^2 E_y}{\partial x^2}$$

$$\frac{\partial^2 B_z}{\partial t^2} = \frac{1}{\mu_0 \epsilon_0} \frac{\partial^2 B_z}{\partial x^2}$$

$$\Rightarrow v \equiv c = \sqrt{\frac{1}{\mu_0 \epsilon_0}} = 2.99792458 \times 10^8 \text{ m/s}$$

Periodic solution:

$$E_y(x, t) = E_{\max} \sin\left(\frac{2\pi}{\lambda}(x - ct)\right) = E_{\max} \sin(kx - \omega t)$$
$$B_z(x, t) = \frac{E_{\max}}{c} \sin\left(\frac{2\pi}{\lambda}(x - ct)\right) = \frac{E_{\max}}{c} \sin(kx - \omega t)$$

## Summary of significant properties of electromagnetic waves:

➡ “Self-sustaining” electric and magnetic fields which can propagate in vacuum at a velocity of  $c = 2.99792458 \times 10^8$  m/s (or within matter at a velocity of  $v = c/n$ ).

➡ **E** and **B** fields are perpendicular to each other and perpendicular to propagation direction (transverse waves).

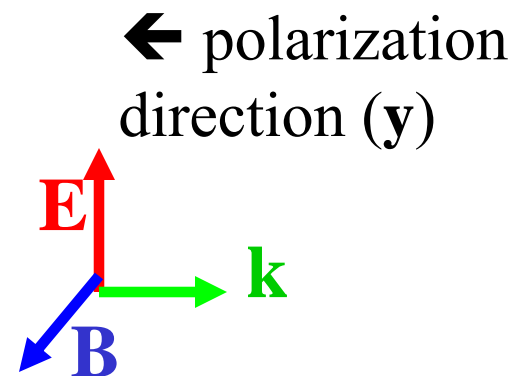
➡  $|\mathbf{E}(x, t)| = \frac{|\mathbf{B}(x, t)|}{v}$

➡ Periodic waves have the form:

$$E_y(x, t) = E_{\max} \sin(kx - \omega t)$$

$$B_z(x, t) = \frac{E_{\max}}{v} \sin(kx - \omega t)$$

$$\frac{\omega}{k} = v$$





## Radiation power

Poynting vector:  $\mathbf{S} = \frac{1}{\mu_0} \mathbf{E} \times \mathbf{B}$  units: Watts/area  
pointing in propagation direction

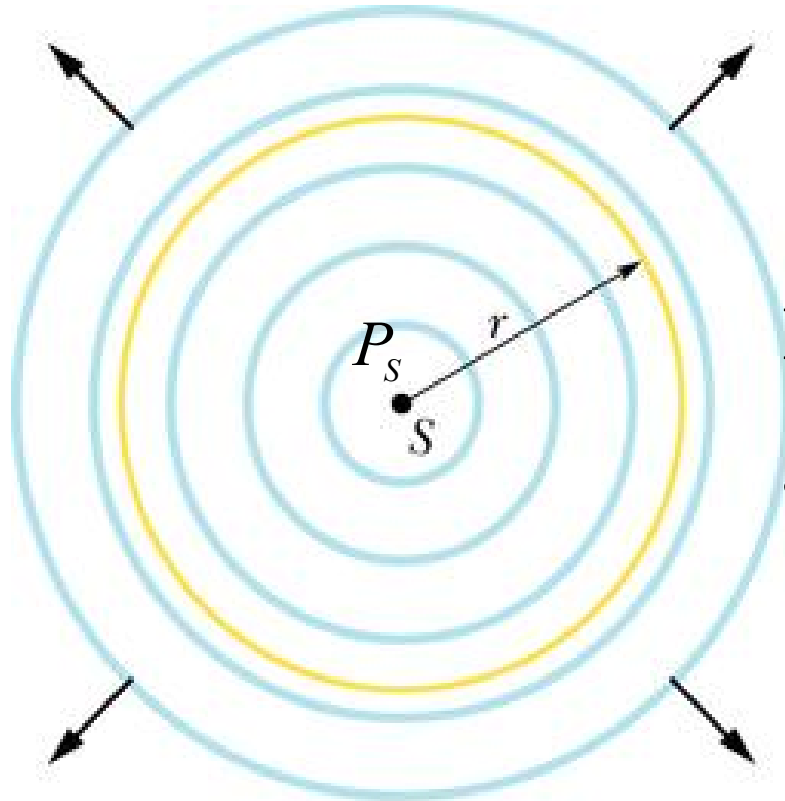
Intensity :  $I \equiv \langle |\mathbf{S}| \rangle_{avg}$

Intensity for plane wave :  $I = \frac{E_{\max}^2}{2c\mu_0}$

due to averaging



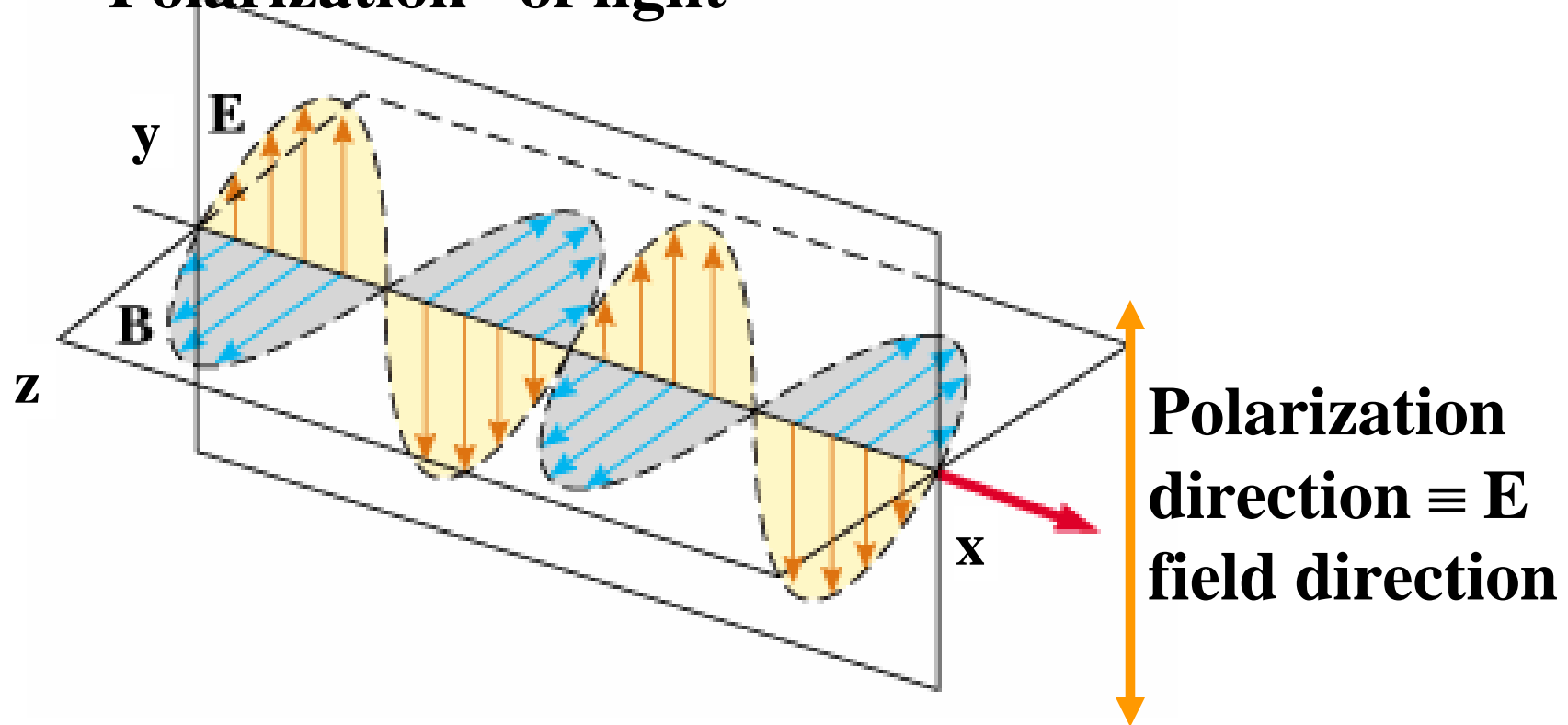
# Radiation power from point source



**Intensity from point source  
at a distance  $r$ :**

$$I = \frac{P_s}{4\pi r^2}$$

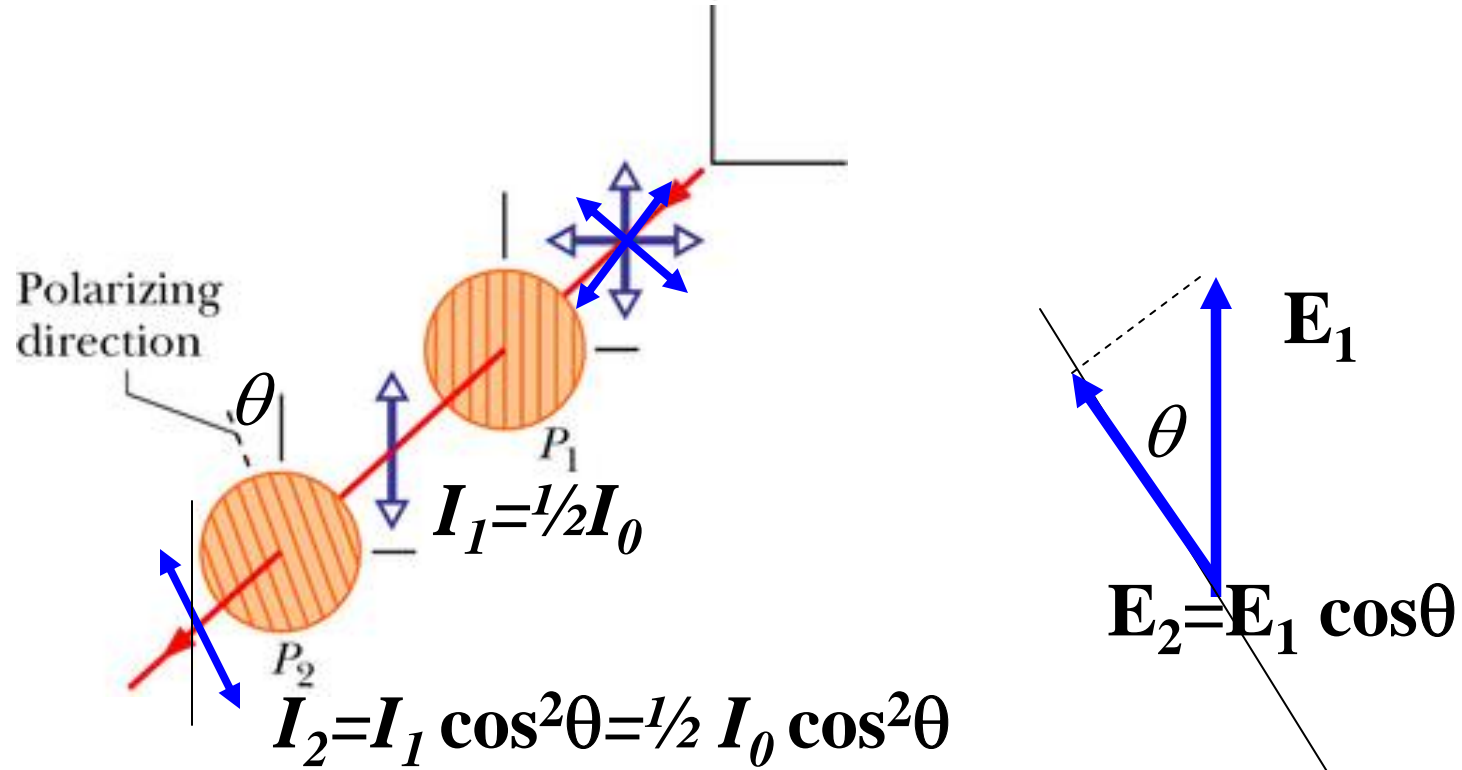
## “Polarization” of light



**Some radiation sources produce polarized light –**

**For example: lasers, antennas**

# Filtration of polarized light as a function of angle



# Electromagnetic waves in materials

$$-\frac{dB_z}{dx} = \mu\epsilon \frac{dE_y}{dt}$$

$$\frac{dE_y}{dx} = -\frac{dB_z}{dt}$$

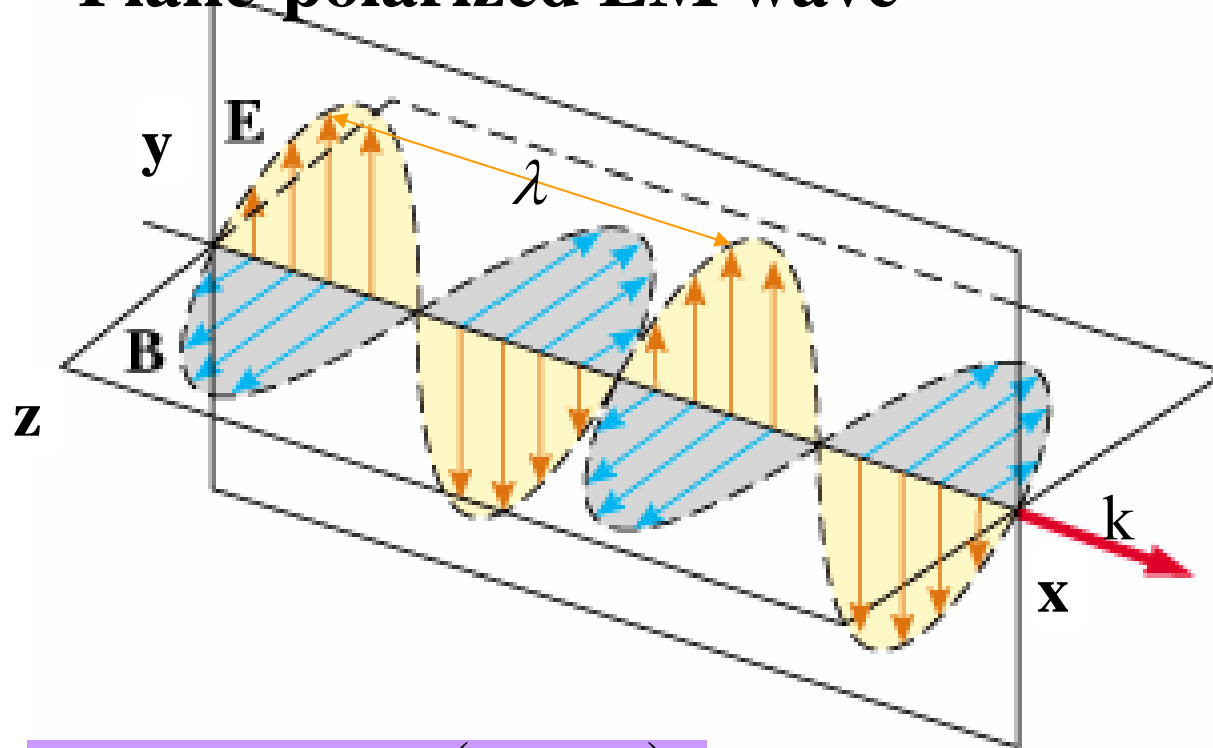
$$\frac{\partial^2 E_y}{\partial t^2} = \frac{1}{\mu\epsilon} \frac{\partial^2 E_y}{\partial x^2}$$

$$\frac{\partial^2 B_z}{\partial t^2} = \frac{1}{\mu\epsilon} \frac{\partial^2 B_z}{\partial x^2}$$

$$\Rightarrow v = \sqrt{\frac{1}{\mu\epsilon}} = \sqrt{\frac{\mu_0\epsilon_0}{\mu\epsilon}} \sqrt{\frac{1}{\mu_0\epsilon_0}} \equiv \frac{c}{n}$$

refractive index

# Plane-polarized EM wave



$$\mathbf{E}(x, t) = E_{\max} \sin(kx - \omega t) \hat{\mathbf{y}}$$

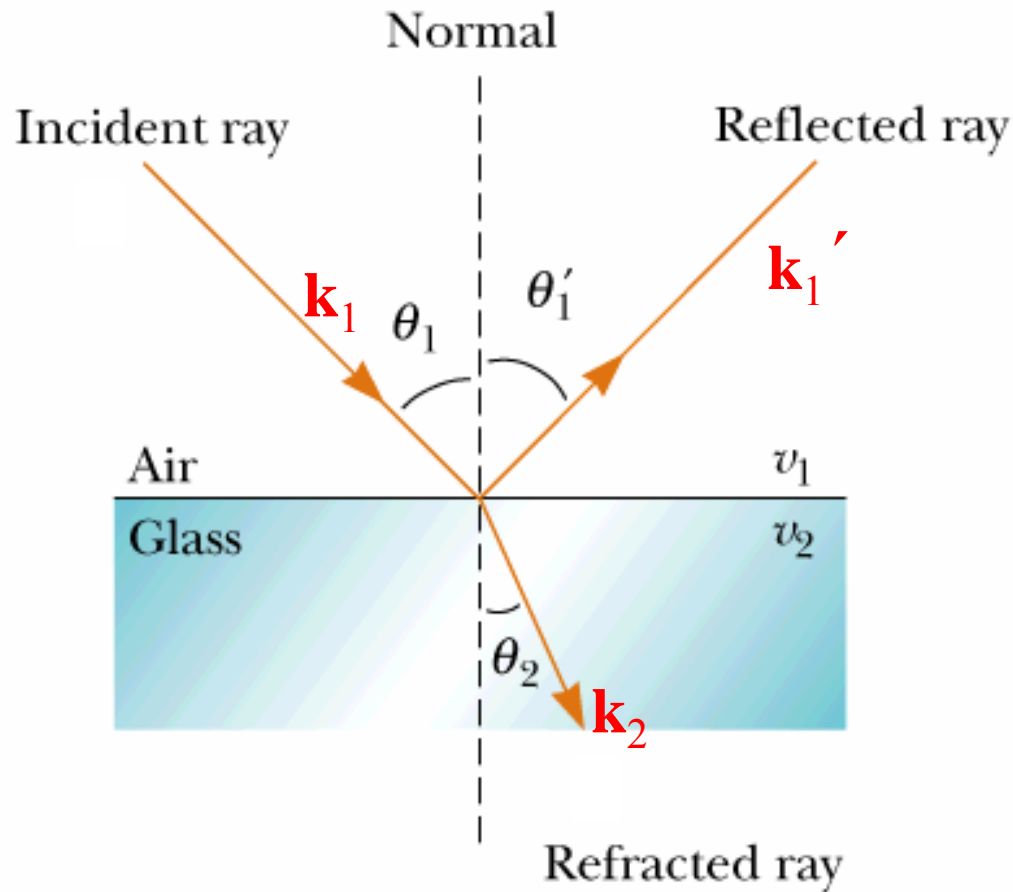
$$\mathbf{B}(x, t) = \frac{E_{\max}}{c} \sin(kx - \omega t) \hat{\mathbf{z}}$$

$$\frac{\omega}{k} = \frac{c}{n} = \frac{\text{speed of light in vacuum}}{\text{refractive index of medium}}$$

$$\Rightarrow k = \frac{n\omega}{c} \quad \text{or} \quad \lambda = \frac{c}{nf}$$

Consider the behavior of a plane-polarized electromagnetic wave near the surface of two materials:

$$k_1 \sin \theta_1 = k'_1 \sin \theta'_1 = k_2 \sin \theta_2$$



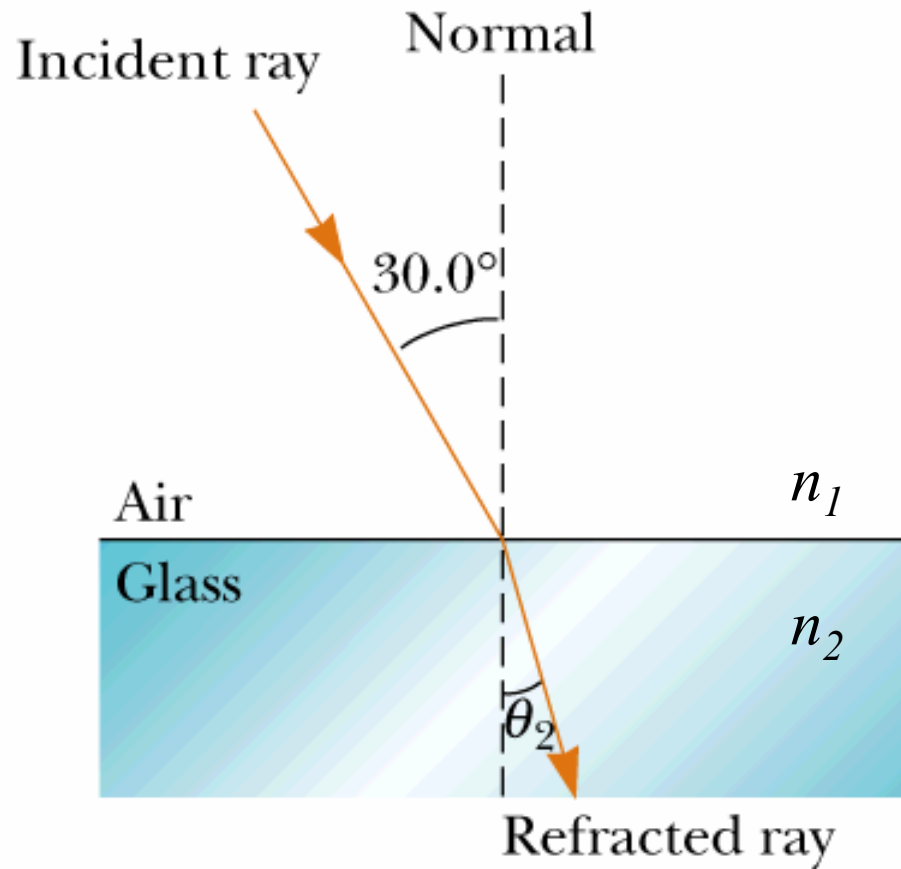
$$k_1 = \frac{n_1 \omega}{c} = k'_1$$

$$k_2 = \frac{n_2 \omega}{c}$$

$$\Rightarrow \theta_1 = \theta'_1$$

$$\Rightarrow n_1 \sin \theta_1 = n_2 \sin \theta_2$$

## Refraction



Snell's law:

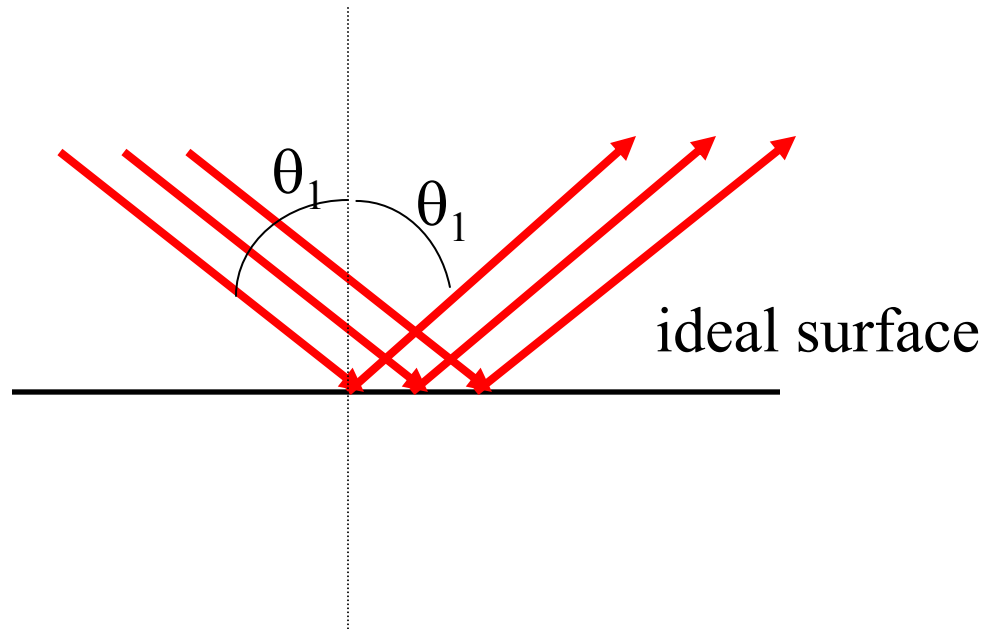
$$n_1 \sin \theta_1 = n_2 \sin \theta_2$$

$$\theta_2 = \sin^{-1} \left( \frac{n_1}{n_2} \sin \theta_1 \right)$$
$$= 17.5^\circ$$

for  $n_1 = 1, n_2 = 1.66$

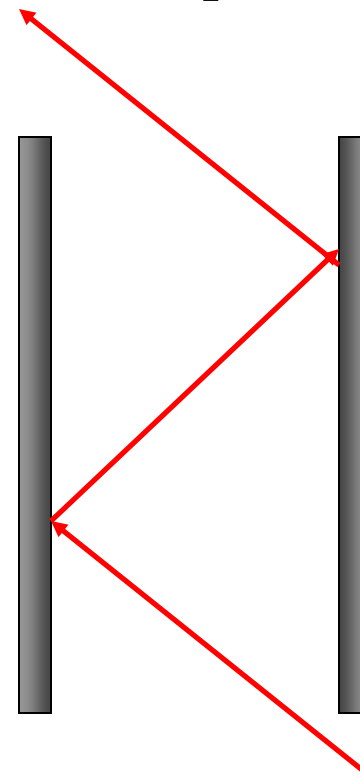


## Reflection

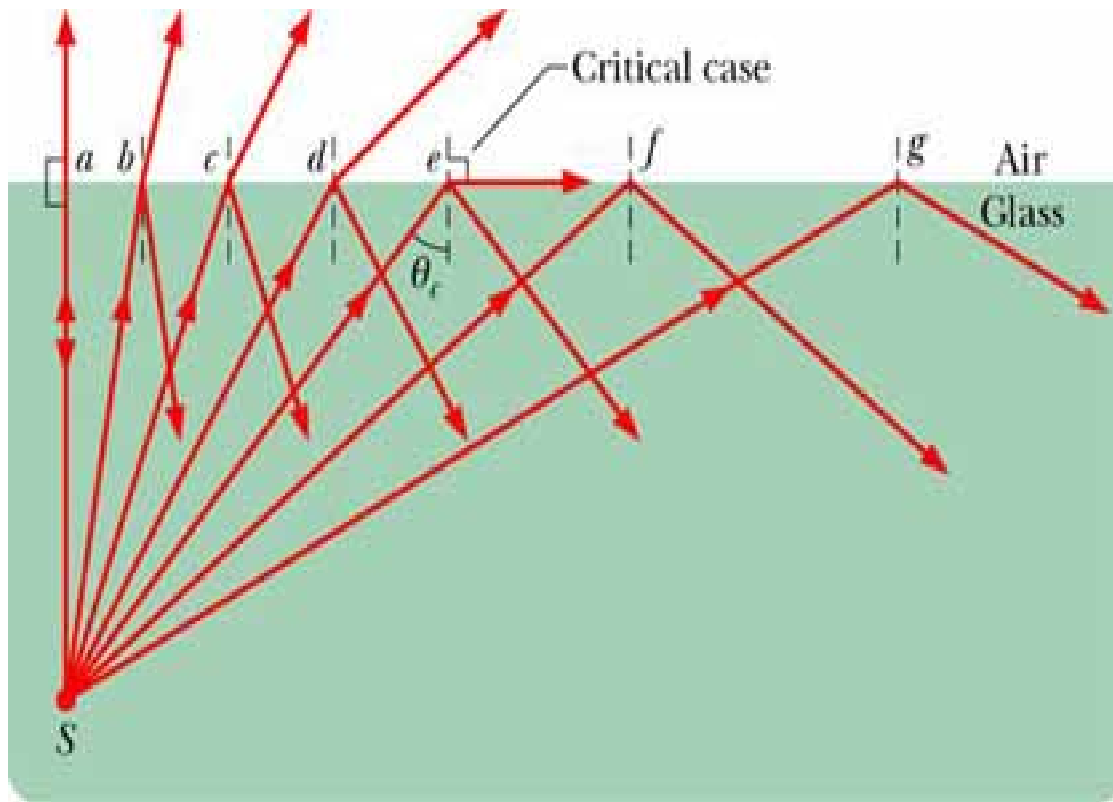


## Ray tracing

Example: 2 mirrors



## Total internal reflection



$$\sin \theta_1 = \frac{n_2}{n_1} \sin \theta_2 \leq 1$$

$$\theta_{2|\max} \equiv \theta_c = \sin^{-1} \left( \frac{n_1}{n_2} \right)$$