

Announcements

1. Comments about Exam 4 and final exam format

2. Extra credit sessions:

Monday evening at 8 PM – 4 students so far

Tuesday evening at 8 PM – 3 students so far

Wednesday evening at 8 PM – 0 students so far

3. Today's topic: (Chapters 39 & 40 of your text)

Continued discussion of quantum mechanics

Lasers

4. Wednesday – general review

Quantum physics –

➤ Electromagnetic waves sometimes behave like particles

➔ one “photon” has a quantum of energy $E=hf$

momentum $p=h/\lambda=hf/c$

➤ Particles sometimes behave like waves

➔ “wavelength” of particle related to momentum:

$$\lambda=h/p$$

➔ quantum particles can “tunnel” to places classically “forbidden”

➔ Stationary quantum states have quantized energies

Wave equations

Electromagnetic waves:

$$\frac{\partial^2 \mathbf{E}}{\partial t^2} = c^2 \frac{\partial^2 \mathbf{E}}{\partial x^2}$$

Matter waves: (Schrödinger equation)

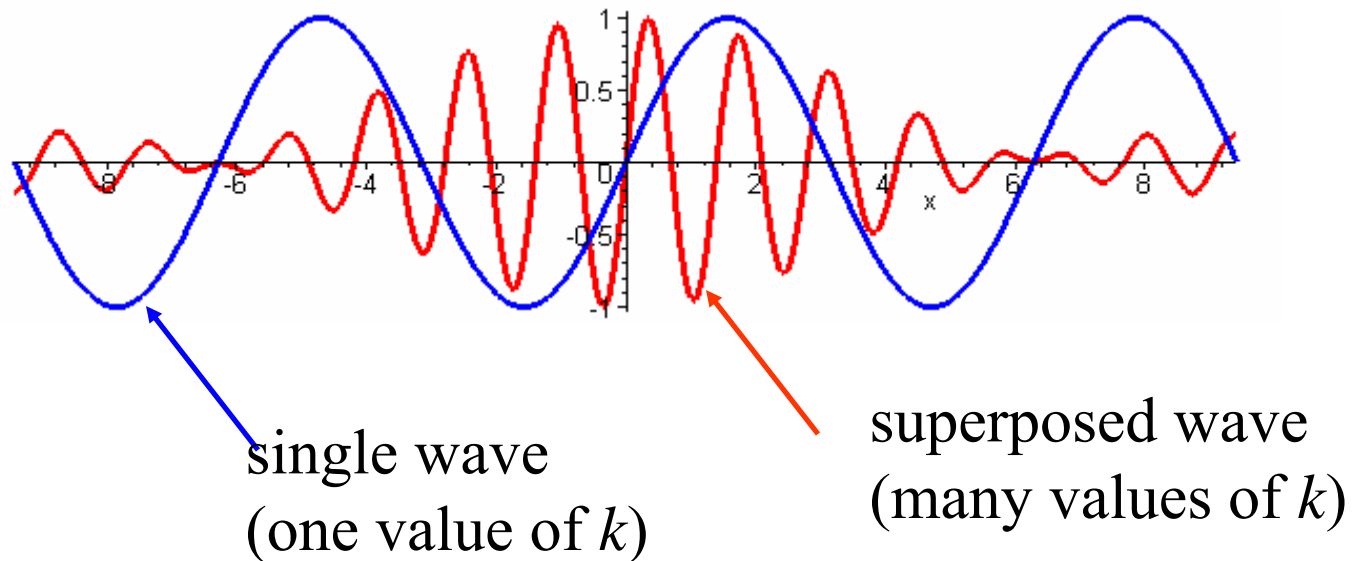
$$-i\hbar \frac{\partial}{\partial t} \Psi(x, t) = \left[-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + V(x) \right] \Psi(x, t)$$

Particle properties	Wave properties
Position as a function of time is known -- $\mathbf{r}(t)$	Phenomenon is spread out over many positions at an instant of time.
Particle is spatially confined when $E \leq U(r)$.	Finite probability that particle can be found in regions other than $E \leq U(r)$.
Particles are independent.	Interference effects.

Mathematical representation of particle and wave behaviors.

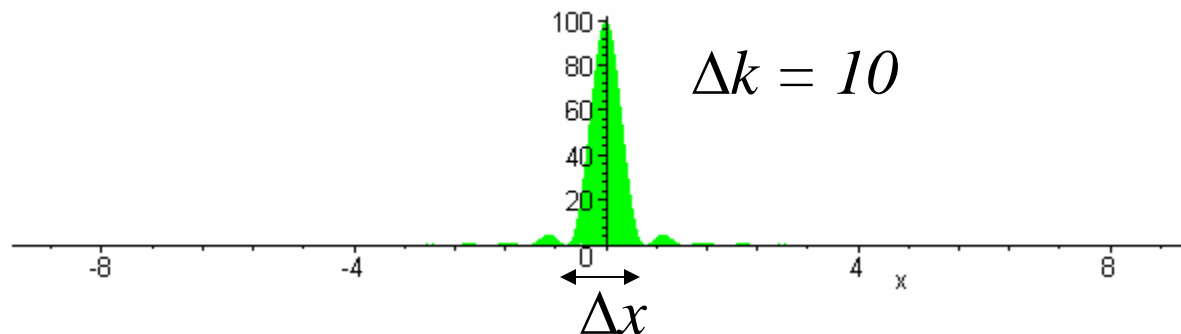
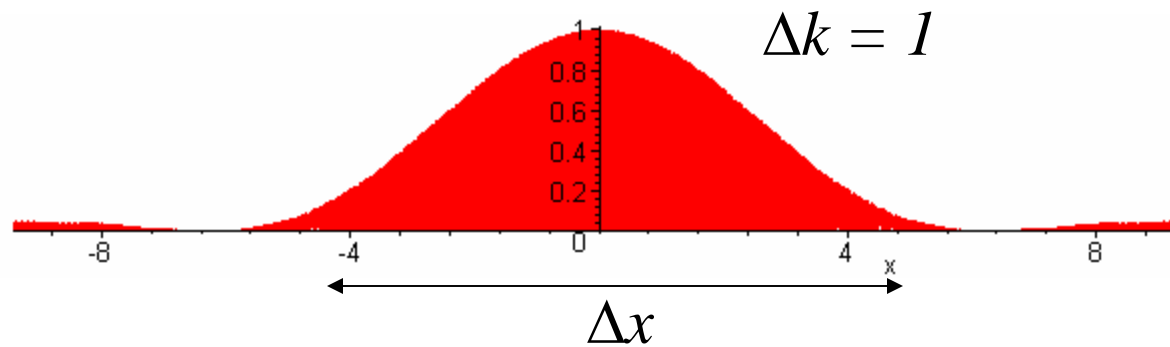
Consider a superposition of periodic waves at $t=0$:

$$E(x, t) = \sum_i E_{\max} \sin(k_i x)$$



$$[E(x,0)]^2 = \left(\sum_i E_{\max} \sin(k_i x) \right)^2$$

$$\Delta x \Delta k \approx 2\pi$$



Δx smaller \rightarrow more particle like

Δk smaller \rightarrow more wave like

$\Delta x \Delta k \approx 2\pi \quad \rightarrow$ Heisenberg's uncertainty principle

De Broglie's particle moment – wavelength relation:

$$p = \frac{h}{\lambda} = \frac{h/2\pi}{\lambda/2\pi} = \hbar k$$

Heisenberg's hypotheses: $\Delta x \Delta p \geq \frac{\hbar}{2}$

$$\Delta t \Delta E \geq \frac{\hbar}{2}$$

$$h = 6.6 \times 10^{-34} \text{ Js} = 4.14 \times 10^{-15} \text{ eVs}$$

Electromagnetic waves	Matter waves
<p>Vector – E or B fields</p> <p>Second order t dependence</p> <p>Examples:</p> $E_y(x, t) = E_{\max} \sin(kx - \omega t)$ $B_z(x, t) = \frac{E_{\max}}{c} \sin(kx - \omega t)$	<p>Scalar – probability amplitude</p> <p>First order t dependence</p> <p>Examples:</p> $\Psi(x, t) = \Psi_0 \sin(kx) e^{-iEt / \hbar}$ <p>“bound” states</p>

What is the meaning of the matter wave function $\Psi(x,t)$?

➤ $\Psi(x,t)$ is not directly measurable

➤ $|\Psi(x,t)|^2$ is measurable – represents the density of particles at position x at time t .

➤ For a single particle system, $\int_{-\infty}^{\infty} |\Psi(x,t)|^2 dx = 1$

➤ For many systems of interest, the wave function can be written in the form $\Psi(x,t) = \psi(x)e^{-iEt/\hbar}$

$$|\Psi(x,t)|^2 = |\psi(x)|^2$$

Wave-like properties of particles

Louis de Broglie suggested that a wavelength could be associated with a particle's momentum

$$p = \frac{h}{\lambda}$$

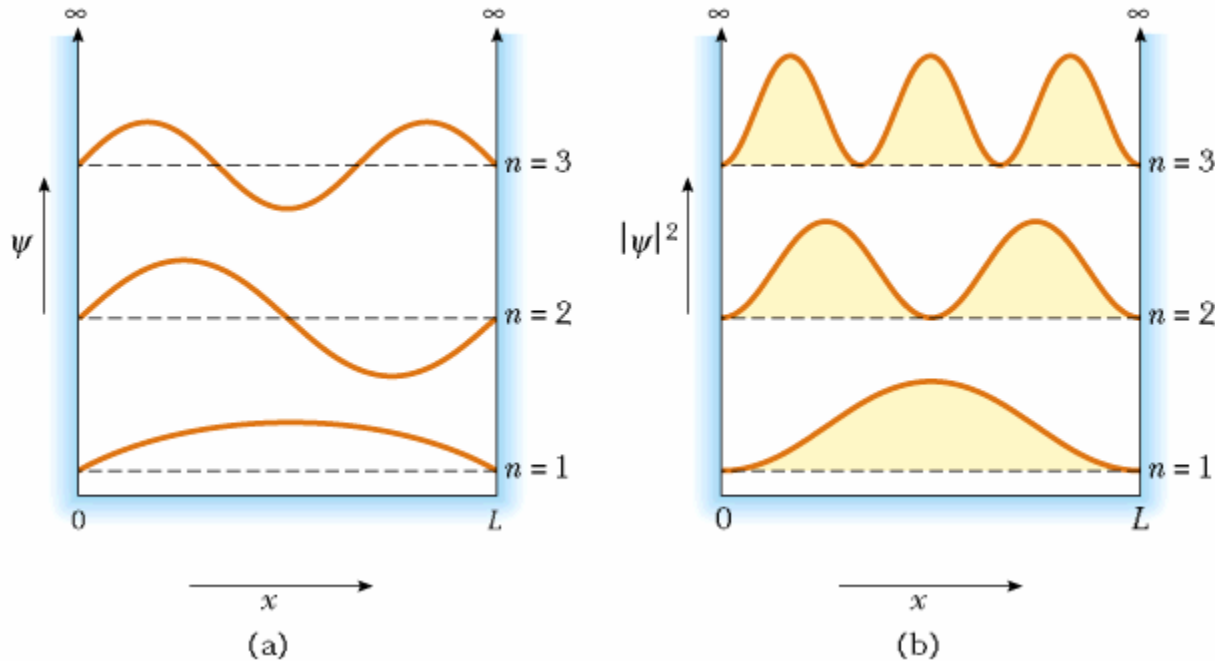
“Wave” equation for particles – Schrödinger equation

$$\left[-\frac{h^2}{(2\pi)^2 m} \frac{\partial^2}{\partial x^2} + V(x) \right] \Psi(x, t) = -i \frac{h}{2\pi} \frac{\partial}{\partial t} \Psi(x, t)$$

Stationary - state wavefunctions : $\Psi(\mathbf{r}, t) = \psi(\mathbf{r}) e^{-iEt/\hbar}$

$$\left[-\frac{h^2}{(2\pi)^2 m} \frac{\partial^2}{\partial x^2} + V(x) \right] \Psi(x, t) = E \Psi(x, t)$$

Electrons in an infinite box:

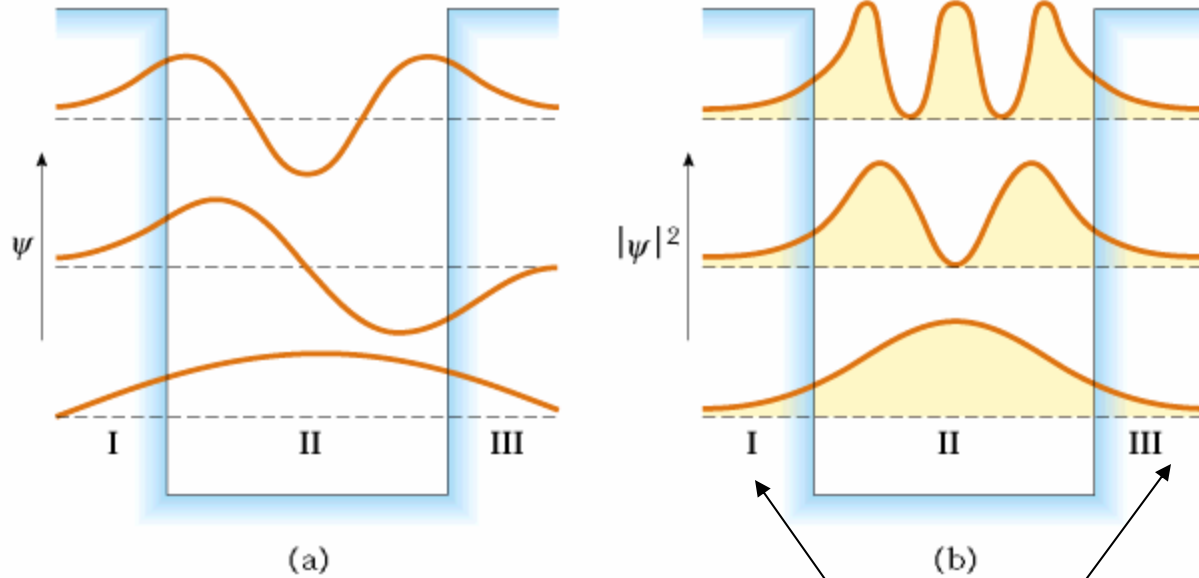


$$E \psi(x) = \left[-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \right] \psi(x) \quad \text{for } 0 \leq x \leq L$$

$$\psi(x) = \psi_0 \sin\left(\frac{n\pi x}{L}\right) \quad n = 1, 2, 3, \dots$$

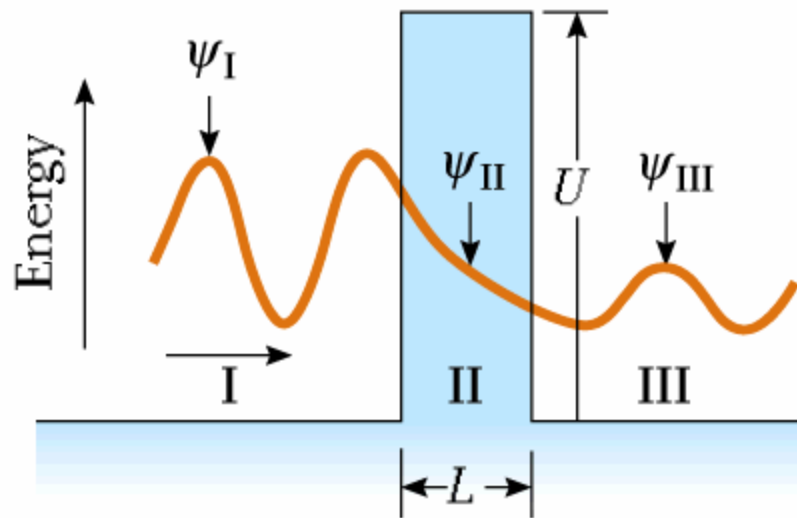
$$E_n = \frac{\hbar^2 \pi^2 n^2}{2m}$$

Electrons in a finite box:



finite probability of electron
existing outside of classical
region

Tunneling through a barrier

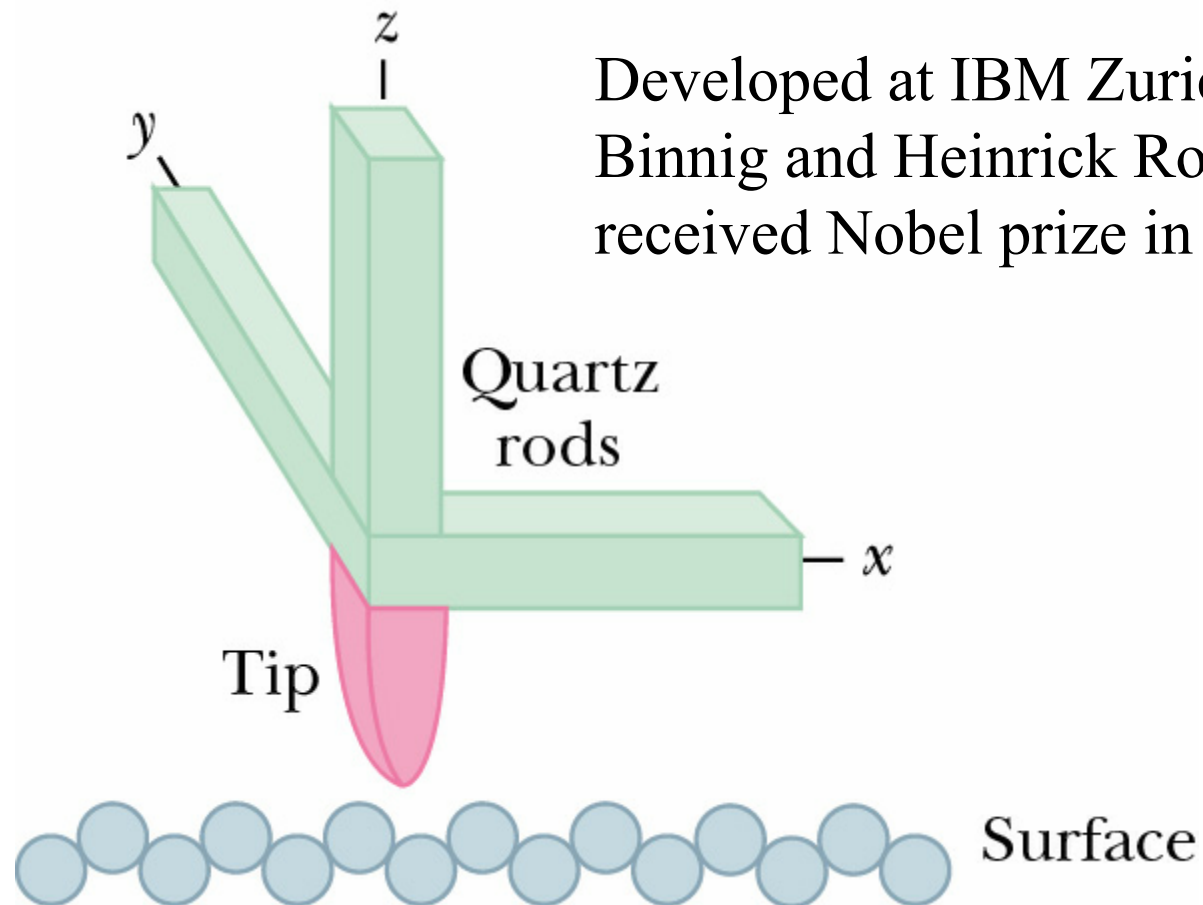


surface region

tip

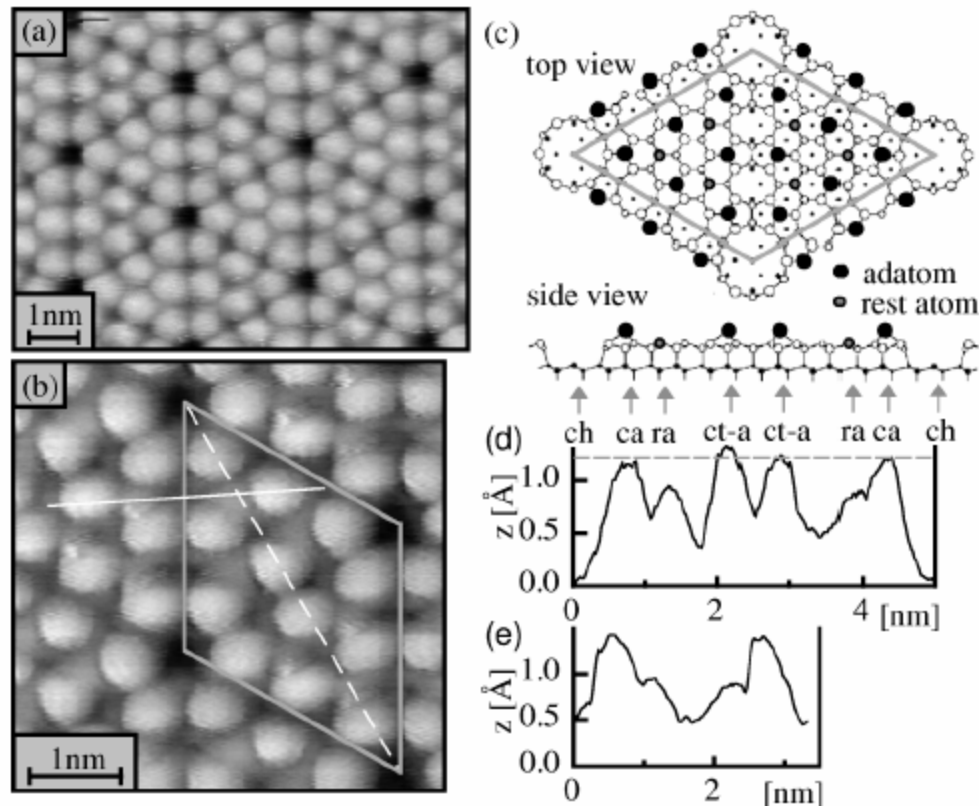
vacuum

How a scanning tunneling microscope works:



Visualization of $|\psi(x)|^2$

A surface if a nearly perfect Si crystal



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Issue 12, pp. 2642-2645

The physics of atoms –

Features are described by solutions to the matter wave equation – Schrödinger equation:

$$-i\hbar \frac{\partial}{\partial t} \Psi(\mathbf{r}, t) = \left[-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial \mathbf{r}^2} + V(\mathbf{r}) \right] \Psi(\mathbf{r}, t)$$

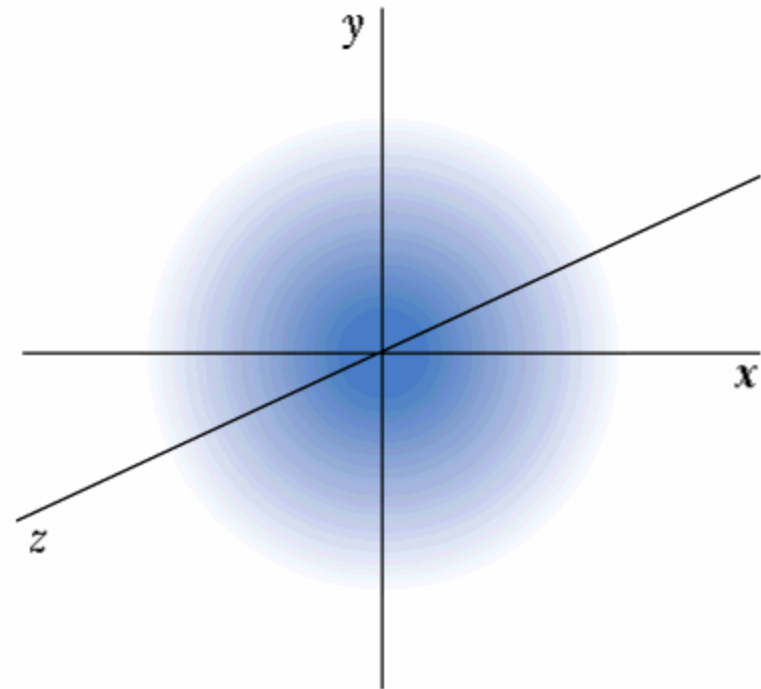
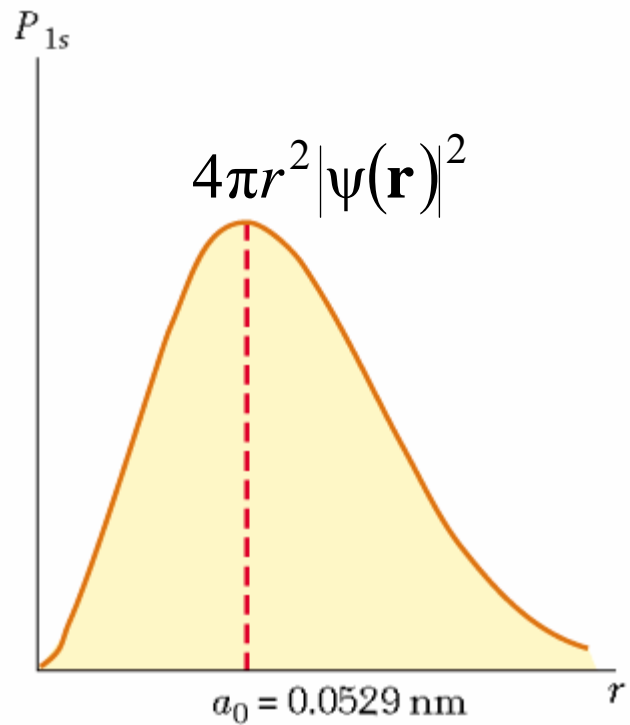
“reduced” mass of electron and proton \nearrow $-\frac{Ze^2}{4\pi\epsilon_0 r}$

Stationary - state wavefunctions : $\Psi(\mathbf{r}, t) = \psi(\mathbf{r}) e^{-iEt/\hbar}$

$$\text{Solutions : } E_n = -\frac{Z^2 e^2}{8\pi\epsilon_0 a_0} \frac{1}{n^2} = -13.6 \frac{Z^2}{n^2} \text{ eV}$$

$$a_0 = \frac{4\pi\epsilon_0 \hbar^2}{me^2} = 0.0529 \text{ nm}$$

Form of probability density for ground state ($n = 1$)



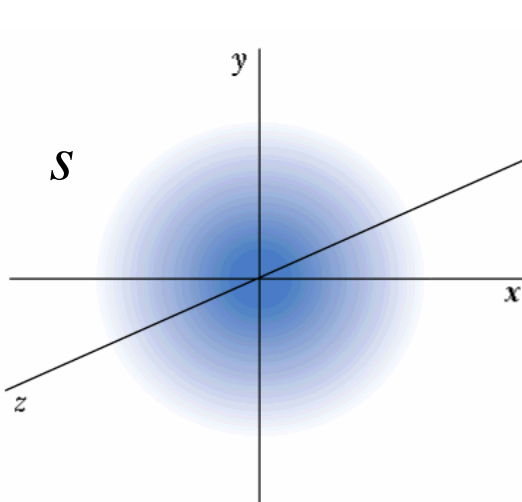
Angular degrees of freedom

-- since the force between the electron and nucleus depends only on distance and not on angle, angular momentum $\mathbf{L} \equiv \mathbf{r} \times \mathbf{p}$ is conserved. Quantum numbers associated with angular

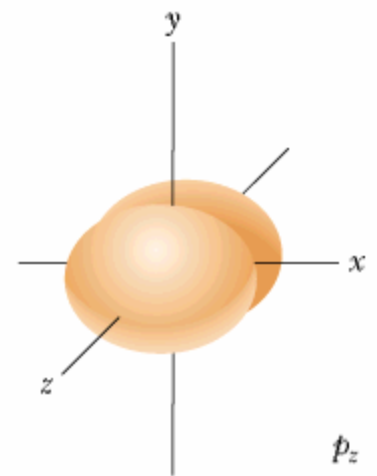
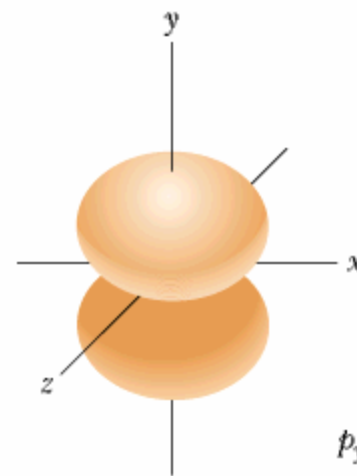
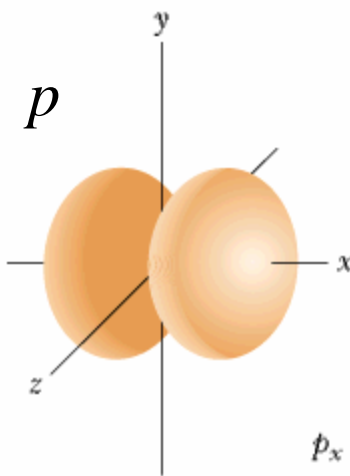
momentum: $\mathbf{L}^2 = \hbar^2 \ell(\ell + 1)$ $\ell = 0, 1, 2, \dots, (n - 1)$

$L_z = \hbar m$ $-\ell \leq m \leq \ell$ total of $2\ell + 1$ states

Notation: $\ell = 0 \Rightarrow s$, $1 \Rightarrow p$, $2 \Rightarrow d$



4/25/2005

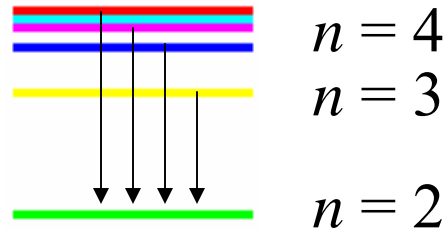


PHY 114 -- Lecture 32

18

Summary of results for H-atom:

$$E_n = -13.6 \frac{Z^2}{n^2} \text{ eV}$$



Balmer series
spectra

degeneracy associated
with each n : $2n^2$



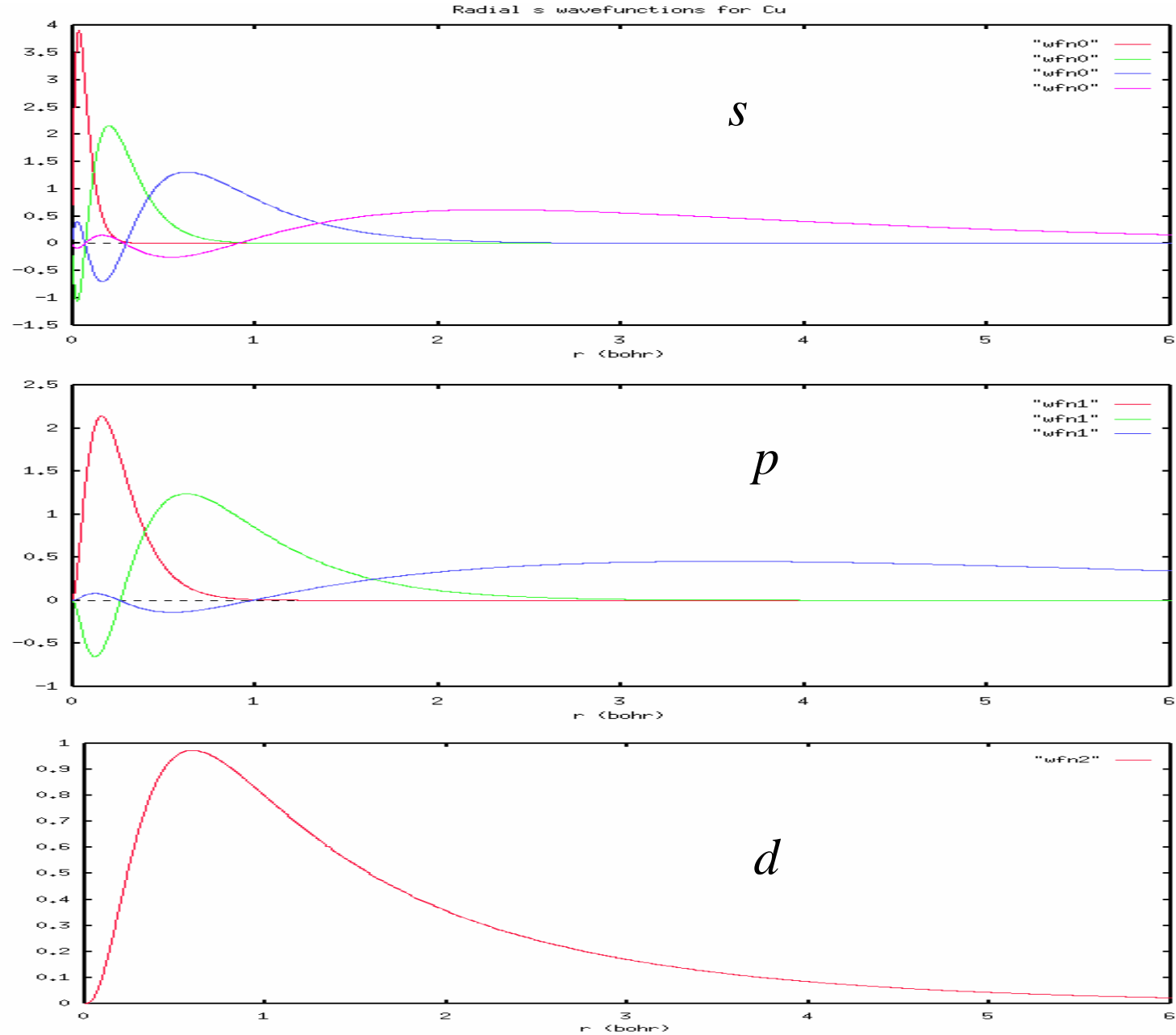
Atomic states of atoms throughout periodic table:

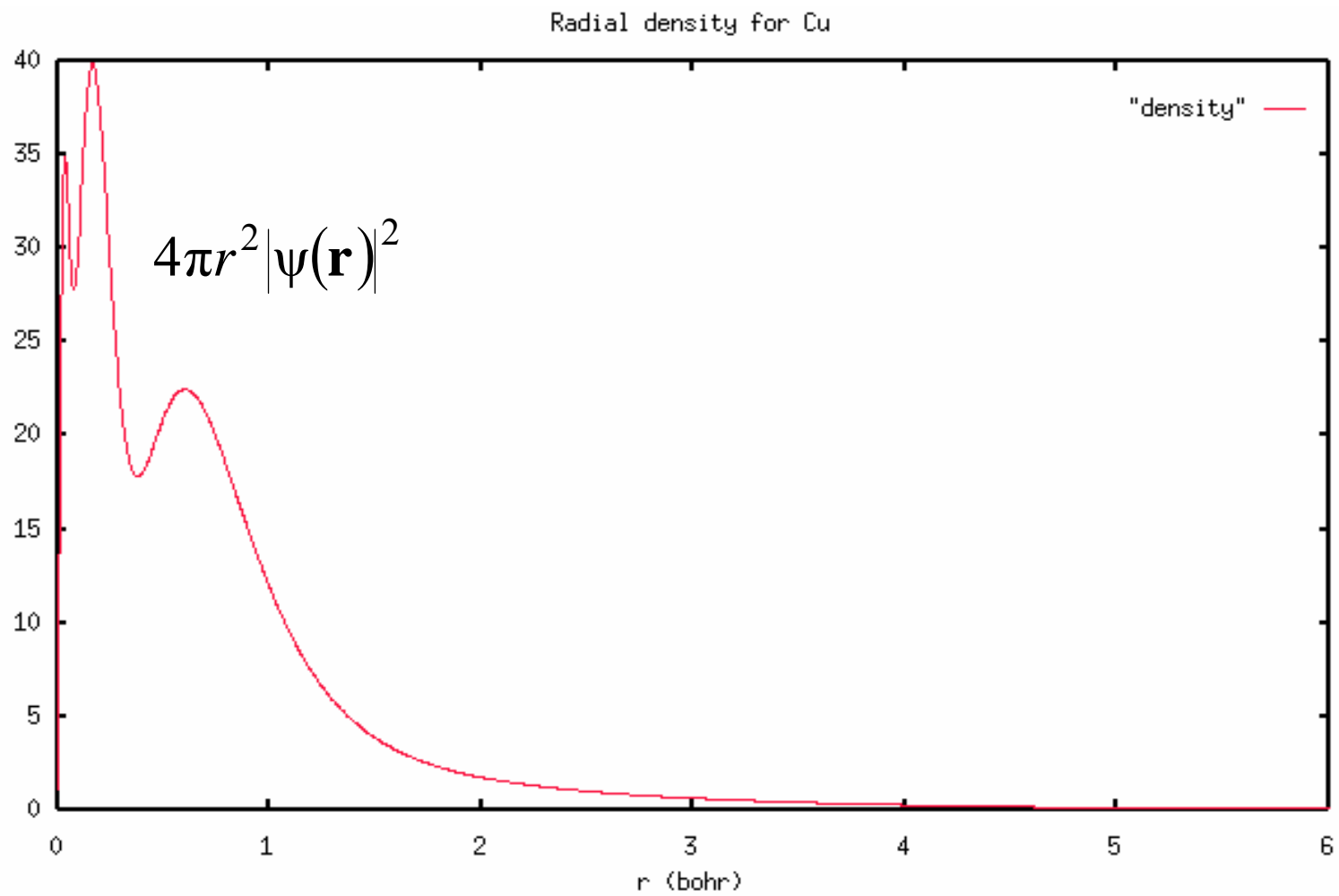
1 H																	2 He
3 Li	4 Be											5 B	6 C	7 N	8 O	9 F	10 Ne
11 Na	12 Mg											13 Al	14 Si	15 P	16 S	17 Cl	18 Ar
19 K	20 Ca	21 Sc	22 Ti	23 V	24 Cr	25 Mn	26 Fe	27 Co	28 Ni	29 Cu	30 Zn	31 Ga	32 Ge	33 As	34 Se	35 Br	36 Kr
37 Rb	38 Sr	39 Y	40 Zr	41 Nb	42 Mo	43 Tc	44 Ru	45 Rh	46 Pd	47 Ag	48 Cd	49 In	50 Sn	51 Sb	52 Te	53 I	54 Xe
55 Cs	56 Ba	57 La	72 Hf	73 Ta	74 W	75 Re	76 Os	77 Ir	78 Pt	79 Au	80 Hg	81 Tl	82 Pb	83 Bi	84 Po	85 At	86 Rn
87 Fr	88 Ra	89 Ac															

$$E\psi(\mathbf{r}) = \left[-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial \mathbf{r}^2} + V(\mathbf{r}) \right] \psi(\mathbf{r})$$

effective potential for
an electron in atom

Example: Cu (Z=29) $1s^2 2s^2 2p^6 3s^2 3p^6 3d^{10} 4s^1$





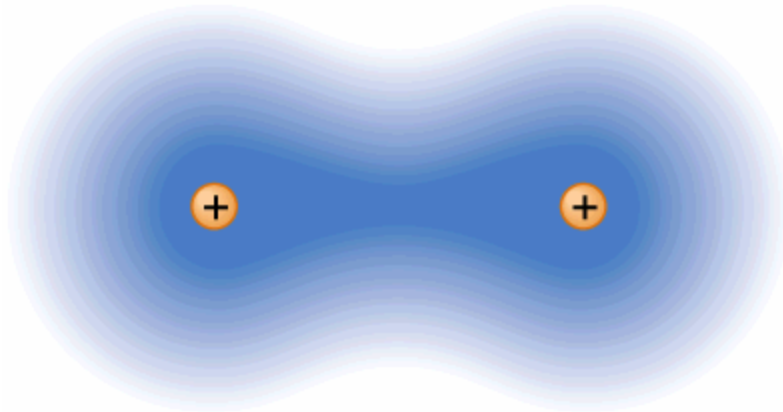
Physics of molecules and solids

$$E\psi(\mathbf{r}) = \left[-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial \mathbf{r}^2} + V(\mathbf{r}) \right] \psi(\mathbf{r})$$

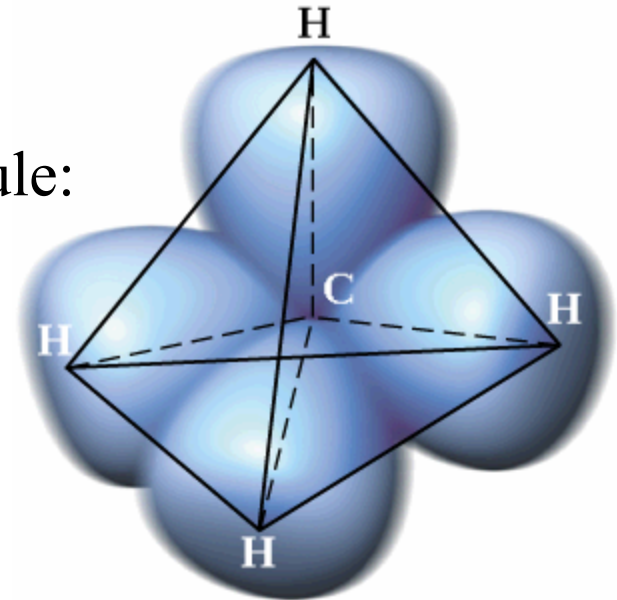
effective potential for
electron in molecule
or solid

Example: electron density associated with

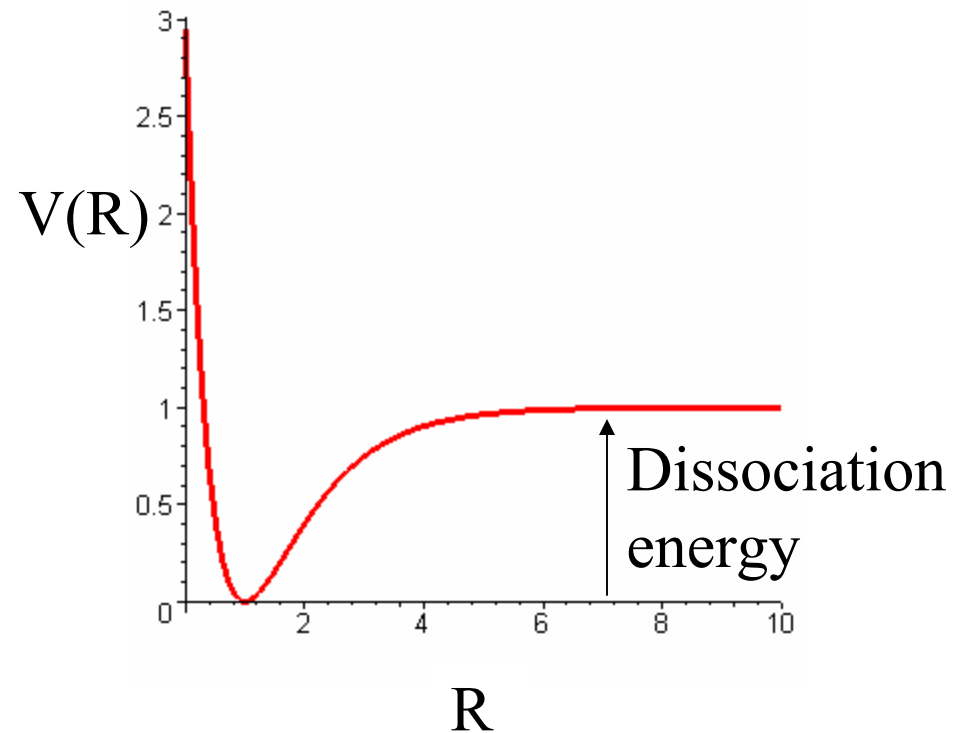
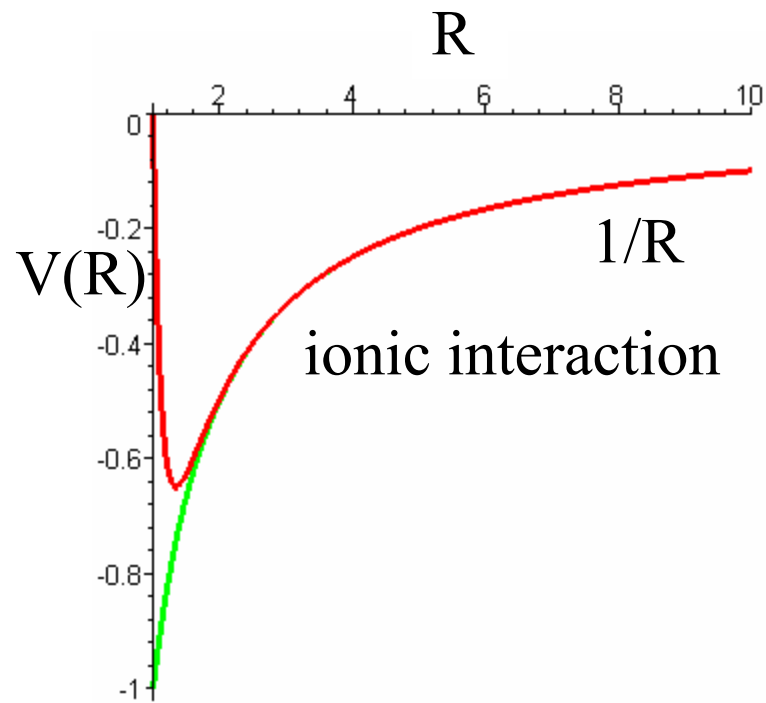
H₂ molecule:



CH₄
molecule:



Molecular binding of nuclei due to electron “glue”:

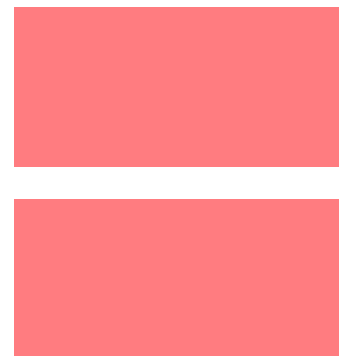
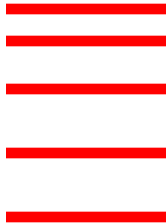


Physics of solids

$$E\psi(\mathbf{r}) = \left[-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial \mathbf{r}^2} + V(\mathbf{r}) \right] \psi(\mathbf{r})$$

effective potential for
electron in molecule
or solid

Energy spectrum of
atom or molecule :

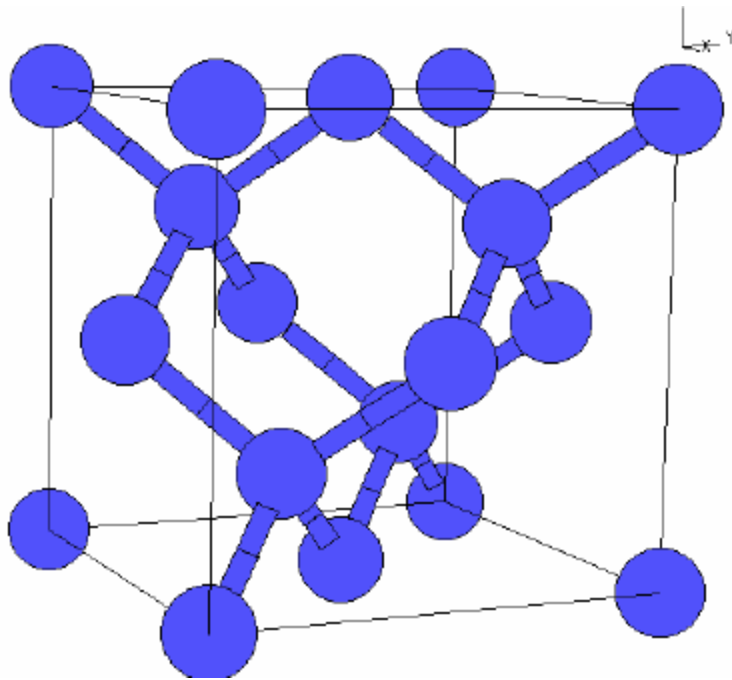


highest filled state
for metal

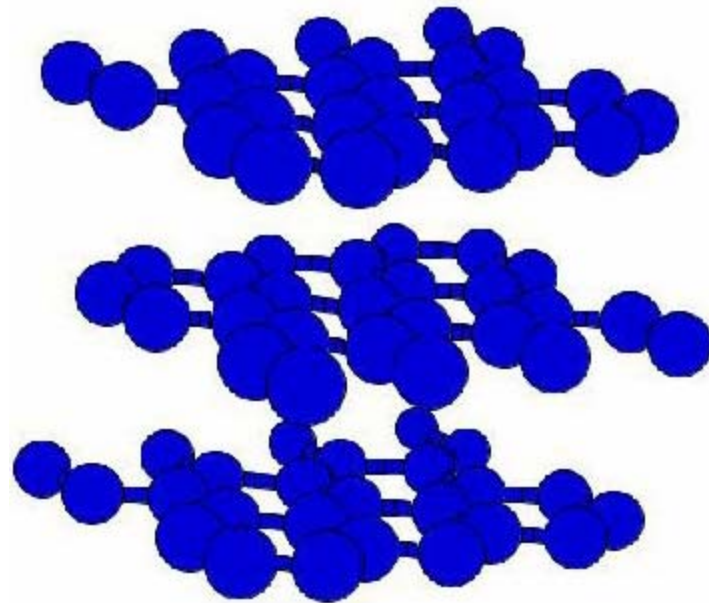
highest filled state
for insulator

Example: 2 materials made of pure carbon:

diamond (insulator)



graphite (semi-metal)



Laser technology:

- System with ground and excited state at desired λ ($E_{\text{ex}} - E_{\text{g}} = hc/\lambda$).
- Standing EM wave
- Mechanism for “population inversion”

The Helium–Neon Gas Laser

Figure 41-21 shows a type of **laser** commonly found in student laboratories. It was developed in 1961 by Ali Javan and his coworkers. The glass discharge tube is filled with a 20:80 mixture of helium and neon gases, neon being the medium in which laser action occurs.

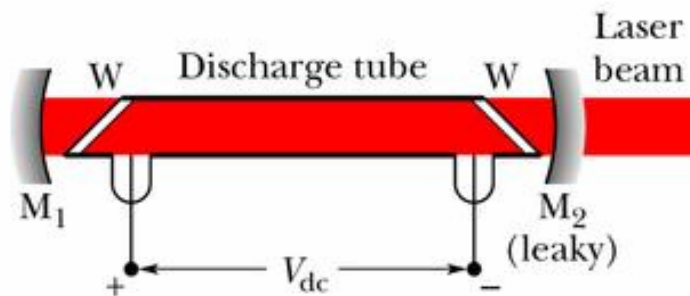


Fig. 41-21 The elements of a helium–neon gas laser. An applied **potential** V_{dc} sends electrons through a discharge tube containing a mixture of helium gas and neon gas. Electrons collide with helium atoms, which then collide with neon atoms, which emit **light** along the length of the tube. The light passes through transparent windows **W** and reflects back and forth through the tube from mirrors M_1 and M_2 to cause more neon **atom** emissions. Some of the light leaks through **mirror** M_2 to form the **laser** beam.

Figure 41-22 shows simplified energy-level diagrams for the two atoms. An electric current passed through the helium–neon gas mixture serves—through collisions between helium atoms and electrons of the current—to raise many helium atoms to state E_3 , which is metastable.

