## Notes for Lecture \#4

## Examples of solutions of the one-dimensional Poisson equation

Consider the following one dimensional charge distribution:

$$
\rho(x)= \begin{cases}0 & \text { for } x<-a  \tag{1}\\ -\rho_{0} & \text { for }-a<x<0 \\ +\rho_{0} & \text { for } 0<x<a \\ 0 & \text { for } x>a\end{cases}
$$

We want to find the electrostatic potential such that

$$
\begin{equation*}
\frac{d^{2} \Phi(x)}{d x^{2}}=-\frac{\rho(x)}{\varepsilon_{0}}, \tag{2}
\end{equation*}
$$

with the boundary condition $\Phi(-\infty)=0$.
The solution to the differential equation is given by:

$$
\Phi(x)=\left\{\begin{array}{ll}
0 & \text { for } x<-a  \tag{3}\\
\frac{\rho_{0}}{2 \varepsilon_{0}}(x+a)^{2} & \text { for }-a<x<0 \\
\frac{-\rho_{0}}{2 \varepsilon_{0}}(x-a)^{2}+\frac{\rho_{0} a^{2}}{\varepsilon_{0}} & \text { for } 0<x<a \\
\frac{\rho_{0}}{\varepsilon_{0}} a^{2} & \text { for } x>a
\end{array} .\right.
$$

The electrostatic field is given by:

$$
E(x)=\left\{\begin{array}{ll}
0 & \text { for } x<-a  \tag{4}\\
-\frac{\rho_{0}}{\varepsilon_{0}}(x+a) & \text { for }-a<x<0 \\
\frac{\rho_{0}}{\varepsilon_{0}}(x-a) & \text { for } 0<x<a \\
0 & \text { for } x>a
\end{array} .\right.
$$

This particular example is one that is used to model semiconductor junctions where the charge density is controlled by introducing charged impurities near the junction. A plot of the results is given below.


# Electric charge density Electric potential <br> Electric field 

The solution of the above differential equation can be determined by piecewise solution within each of the four regions or by use of the Green's function $G\left(x, x^{\prime}\right)=4 \pi x_{<}$, where,

$$
\begin{equation*}
\Phi(x)=\frac{1}{4 \pi \varepsilon_{0}} \int_{-\infty}^{\infty} G\left(x, x^{\prime}\right) \rho\left(x^{\prime}\right) d x^{\prime} \tag{5}
\end{equation*}
$$

In the expression for $G\left(x, x^{\prime}\right), x_{<}$should be taken as the smaller of $x$ and $x^{\prime}$. It can be shown that Eq. 5 gives the identical result for $\Phi(x)$ as given in Eq. 3 .

## Notes on the one-dimensional Green's functions

The Green's function for the Poisson equation can be defined as a solution to the equation:

$$
\begin{equation*}
\nabla^{2} G\left(x, x^{\prime}\right)=-4 \pi \delta\left(x-x^{\prime}\right) \tag{6}
\end{equation*}
$$

Here the factor of $4 \pi$ is not really necessary, but ensures consistency with your text's treatment of the 3-dimensional case. The meaning of this expression is that $x^{\prime}$ is held fixed while taking the derivative with respect to $x$. It is easily shown that with this definition of the Green's function (6), Eq. (5) finds the electrostatic potential $\Phi(x)$ for an arbitrary charge density $\rho(x)$. In order to find the Green's function which satisfies Eq. (6), we notice that we can use two independent solutions to the homogeneous equation

$$
\begin{equation*}
\nabla^{2} \phi_{i}(x)=0 \tag{7}
\end{equation*}
$$

where $i=1$ or 2 , to form

$$
\begin{equation*}
G\left(x, x^{\prime}\right)=\frac{4 \pi}{W} \phi_{1}\left(x_{<}\right) \phi_{2}\left(x_{>}\right) . \tag{8}
\end{equation*}
$$

This notation means that $x_{<}$should be taken as the smaller of $x$ and $x^{\prime}$ and $x_{>}$should be taken as the larger. In this expression $W$ is the "Wronskian":

$$
\begin{equation*}
W \equiv \frac{d \phi_{1}(x)}{d x} \phi_{2}(x)-\phi_{1}(x) \frac{d \phi_{2}(x)}{d x} . \tag{9}
\end{equation*}
$$

We can check that this "recipe" works by noting that for $x \neq x^{\prime}$, Eq. (8) satisfies the defining equation (6) by virtue of the fact that it is equal to a product of solutions to the homogeneous equation 7. The defining equation is singular at $x=x^{\prime}$, but integrating Eq. (6) over $x$ in the neighborhood of $x^{\prime}\left(x^{\prime}-\epsilon<x<x^{\prime}+\epsilon\right)$, gives the result:

$$
\begin{equation*}
\left.\left.\frac{d G\left(x, x^{\prime}\right)}{d x}\right\rfloor_{x=x^{\prime}+\epsilon}-\frac{d G\left(x, x^{\prime}\right)}{d x}\right\rfloor_{x=x^{\prime}-\epsilon}=-4 \pi . \tag{10}
\end{equation*}
$$

In our present case, we can choose $\phi_{1}(x)=x$ and $\phi_{2}(x)=1$, so that $W=1$, and the Green's function is as given above. For this piecewise continuous form of the Green's function, the integration 5 can be evaluated:

$$
\begin{equation*}
\Phi(x)=\frac{1}{4 \pi \varepsilon_{0}}\left\{\int_{-\infty}^{x} G\left(x, x^{\prime}\right) \rho\left(x^{\prime}\right) d x^{\prime}+\int_{x}^{\infty} G\left(x, x^{\prime}\right) \rho\left(x^{\prime}\right) d x^{\prime}\right\} \tag{11}
\end{equation*}
$$

which becomes

$$
\begin{equation*}
\Phi(x)=\frac{1}{\varepsilon_{0}}\left\{\int_{-\infty}^{x} x^{\prime} \rho\left(x^{\prime}\right) d x^{\prime}+x \int_{x}^{\infty} \rho\left(x^{\prime}\right) d x^{\prime}\right\} \tag{12}
\end{equation*}
$$

Evaluating this expression, we find that we obtain the same result as given in Eq. (3).
In general, the Green's function $G\left(x, x^{\prime}\right)$ solution (5) depends upon the boundary conditions of the problem as well as on the charge density $\rho(x)$. In this example, the solution is valid for all neutral charge densities, that is $\int_{-\infty}^{\infty} \rho(x) d x=0$.

