

Lecture 26

Chemical potentials and phase equilibria (Chapter 7 in STP)

- Phase diagrams
- Van der Waals equation of state

	<i>4/24/2012</i>	<i>AP² – no class, take-home exam</i>		
	3/2/2012	APS – no class, take-home exam		
18	3/05/2012	Exam due – Ising model	5.5	HW 17 03/07/2012
19	3/07/2012	Ising model	5.6-7	HW 18 03/09/2012
20	3/09/2012	Phase transformation	5.8-10	
	3/12/2012	Spring Break		
	3/14/2012	Spring Break		
	3/16/2012	Spring Break		
21	3/19/2012	Many particle systems	6.1-2	HW 19 03/23/2012
22	3/21/2012	Fermi and Bose particles	6.3-4	
23	3/23/2012	Bose and Fermi particles	6.5-11	HW 20 03/28/2012
24	3/26/2012	Bose and Fermi particles	6.5-11	
25	3/28/2012	Phase transformations	7.1-7.3	HW 21 03/30/2012
26	3/30/2012	Phase transformations, continued	7.1-7.3	
	4/02/2012			
	4/04/2012			
	4/06/2012	Good Friday, Holiday		

Reminder – second exam in April
-- student presentations 4/30, 5/2 (need to pick topics)

Review:

Clausius - Clapeyron Equation -- approximate $P(T)$ for liquid - gas coexistence

$$\frac{dP}{dT} = \frac{L_{AB}}{T(V_A - V_B)}$$

Example: $A \equiv \text{vapor}$ $B \Rightarrow \text{water}$

$$L_{AB} = 2.257 \times 10^6 \text{ J/kg} = 40.7 \times 10^3 \text{ J/mole} \quad T \geq 373.15 \text{ K}$$

$$V_A \approx \frac{R_M T}{P} \text{ (per mole)} \quad R_M = kN_{\text{Avo}}$$

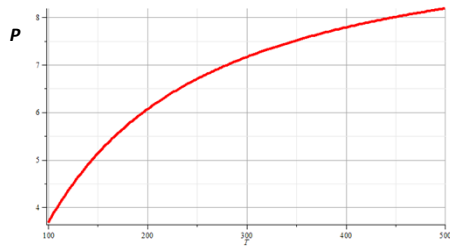
$$V_B \approx 0$$

$$\frac{dP}{dT} = \frac{L_{AB}}{R_M} \frac{P}{T^2} \Rightarrow \frac{dP}{P} = \frac{L_{AB}}{R_M} \frac{dT}{T^2}$$

$$\ln(P) = \text{constant} - \frac{L_{AB}}{R_M T} \Rightarrow P(T) = P_0 e^{-L_{AB}/R_M T}$$

Approximate liquid-gas vaporization curve from Clausius-Clapeyron equation:

$$P(T) = P_0 e^{-L_{\text{vap}}/R_u T}$$



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The van der Waals equation of state

-- More realistic than the ideal gas law; contains some of the correct attributes for liquid-gas phase transitions.

Ideal gas equation of state: $PV = NkT$

van der Waals equation of state: $\left(P + a \frac{N^2}{V^2}\right)(V - bN) = NkT$

here a, b are material - dependent parameters

Dimensionless variables:

$$\tilde{P} \equiv 27 \left(\frac{b^2}{a}\right) P \quad \tilde{T} \equiv \frac{27}{8} \left(\frac{b}{a}\right) kT \quad \tilde{\rho} \equiv 3b \frac{N}{V}$$

van der Waals equation of state: $\tilde{P} = \frac{8\tilde{\rho}\tilde{T}}{3-\tilde{\rho}} - 3\tilde{\rho}^2$

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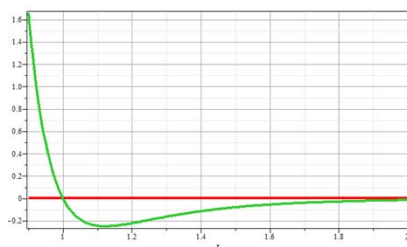
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Some motivation for van der Waals equation of state

Johannes Diderik van der Waals -- Ph. D. Thesis written in 1873 (The Netherlands) -- received Nobel Prize in 1910

General potential between two particles:



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van der Waals equation of state : $P = \frac{NkT}{(V - bN)} - a \frac{N^2}{V^2}$

Substance	a (10^{-30}eV m^3)	b (10^{-30}m^3)
water	9.55	50.7
SO ₂	11.9	94.7
CO ₂	6.3	71.3
O ₂	2.38	52.9
Ar	2.36	53.8
N ₂	2.36	64.3
H ₂	0.428	44.3
⁴ He	0.0597	39.4

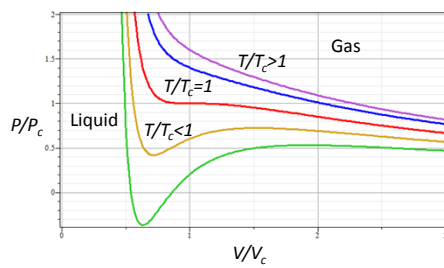
From: Baierlein, *Thermal Physics*

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van der Waals equation of state : $P = \frac{NkT}{(V - bN)} - a \frac{N^2}{V^2}$

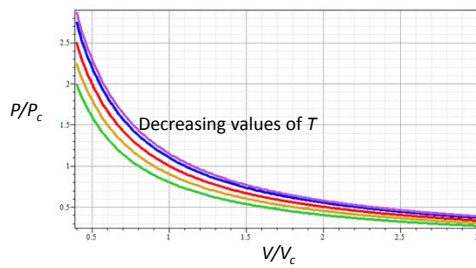


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For comparison -- ideal gas equation of state : $P = \frac{NkT}{V}$



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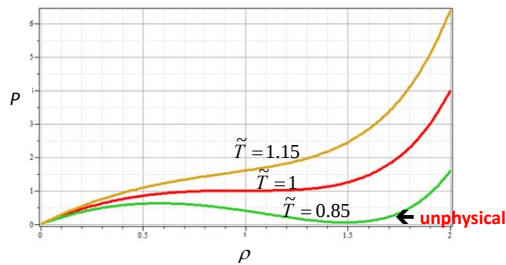
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Dimensionless variables:

$$\tilde{P} \equiv 27 \left(\frac{b^2}{a} \right) P \quad \tilde{T} \equiv \frac{27}{8} \left(\frac{b}{a} \right) kT \quad \tilde{\rho} \equiv 3b \frac{N}{V}$$

van der Waals equation of state: $\tilde{P} = \frac{8\tilde{\rho}\tilde{T}}{3-\tilde{\rho}} - 3\tilde{\rho}^2$



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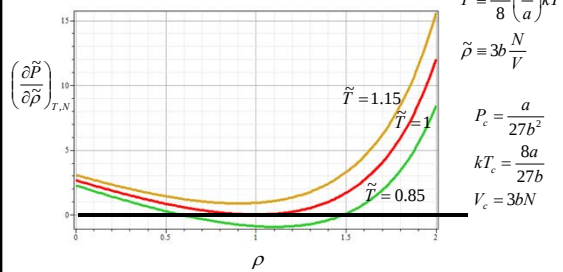
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Examination of compressibility in van der Waals material:

$$\kappa \equiv -\frac{1}{V} \left(\frac{\partial V}{\partial P} \right)_{T,N} = \frac{1}{\rho} \left(\frac{\partial \rho}{\partial P} \right)_{T,N} \geq 0$$

for: $\tilde{P} = \frac{8\tilde{\rho}\tilde{T}}{3-\tilde{\rho}} - 3\tilde{\rho}^2$



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$$\tilde{P} \equiv 27 \left(\frac{b^2}{a} \right) P$$

$$\tilde{T} \equiv \frac{27}{8} \left(\frac{b}{a} \right) kT$$

$$\tilde{\rho} \equiv 3b \frac{N}{V}$$

$$P_c = \frac{a}{27b^2}$$

$$kT_c = \frac{8a}{27b}$$

$$V_c = 3bN$$

Behavior of the thermodynamic potentials for the van der Waals equation of state

In terms of unscaled variables: $\left(P + a \frac{N^2}{V^2} \right) (V - bN) = NkT$

Helmholz free energy F : $P = - \left(\frac{\partial F}{\partial V} \right)_{T,N}$

$$F = - \int^V P dV + NkTw(T)$$

$$= - \int^V \left(\frac{NkT}{V-bN} - a \frac{N^2}{V^2} \right) dV + NkTw(T)$$

$$= -NkT \ln(V-bN) - a \frac{N^2}{V} + NkTw(T)$$

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Behavior of the thermodynamic potentials for the van der Waals equation of state

$$F = -NkT \ln(V - bN) - a \frac{N^2}{V} + NkTw(T)$$

Gibb's free energy :

$$G = F + PV = -NkT \ln(V - bN) - a \frac{N^2}{V} + PV + NkTw(T)$$

In scaled variables with

$$\tilde{g} \equiv \frac{G}{N} \frac{8}{3kT_c} = \frac{G}{N} \frac{9b}{a} \quad \tilde{P} \equiv 27 \left(\frac{b^2}{a} \right) P \quad \tilde{T} \equiv \frac{27}{8} \left(\frac{b}{a} \right) kT \quad \tilde{\rho} \equiv 3b \frac{N}{V}$$

$$G = F + PV = -NkT \ln(V - bN) - a \frac{N^2}{V} + PV + NkTw(T)$$

$$\tilde{g} = -\frac{8}{3} \tilde{T} \ln \left(\left(\frac{3}{\tilde{\rho}} - 1 \right) bN \right) - 3\tilde{\rho} + \frac{\tilde{P}}{\tilde{\rho}} + w'(\tilde{T})$$

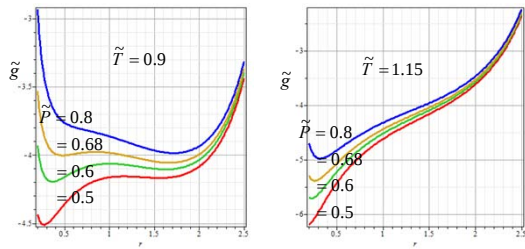
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Behavior of Gibbs chemical potential :

$$\begin{aligned} \tilde{g}(\tilde{T}, \tilde{P}, \tilde{\rho}) &= -\frac{8}{3} \tilde{T} \ln \left(\left(\frac{3}{\tilde{\rho}} - 1 \right) bN \right) - 3\tilde{\rho} + \frac{\tilde{P}}{\tilde{\rho}} + w'(\tilde{T}) \\ &= -\frac{8}{3} \tilde{T} \ln \left(\frac{3}{\tilde{\rho}} - 1 \right) - 3\tilde{\rho} + \frac{\tilde{P}}{\tilde{\rho}} + w''(\tilde{T}, N) \end{aligned}$$

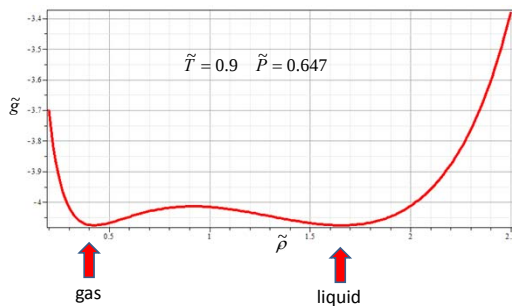


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Gibbs chemical potential at coexistence point



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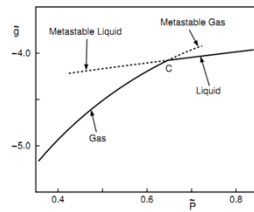


Figure 7.7: Plot of the dimensionless Gibbs free energy per particle $\tilde{g}$ as a function of $\tilde{P}$ at $T = 0.9T_c$ and fixed density. The system is metastable along the dashed curves.

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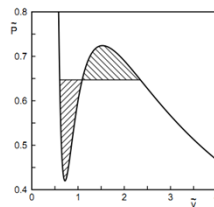


Figure 7.8: Maxwell equal area construction. The pressure $\tilde{P}$ where two phase coexistence begins for $\tilde{T} = 0.9$ is determined so that the areas above and below the horizontal line are equal. In this case $\tilde{P} \approx 0.647$.

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