

PHY 341/641  
Thermodynamics and Statistical Physics

## Lecture 33

Analysis of classical gases and liquids (STP Chapt. 8)

- Virial coefficients
  - Radial distribution functions

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24	3/26/2012	Bose and Fermi particles	6.5-6.11
25	3/28/2012	Phase transformations	7.1-7.3 <a href="#">HW 21</a>
26	3/30/2012	Van der Waals equation	7.4
27	4/02/2012	Equilibrium constants	7.4-7.5 <a href="#">HW 22</a>
28	4/04/2012	Equilibrium constants	7.5
	4/06/2012	Good Friday Holiday	
29	4/09/2012	Review – begin take-home exam	5-7
4/11/2012		No class – work on exam	5-7
30	4/13/2012	Simulation of chemical potential	7.2 Exam continued
31	4/16/2012	Classical treatment of dense systems	8.1-8.2 Exam due
32	4/18/2012	Review exam; Vinal expansion	8.3-8.4
33	4/20/2012	Radial distribution function	8.5
34	4/23/2012	More topics on classical fluids	8.6-8.9
35	4/25/2012	Review	
36	4/27/2012	Review	
	4/30/2012	Student presentations I	
5/02/2012		Student presentations II	
5/09/2012		9 AM Final exam	

-- student presentations 4/30, 5/2 (need to pick topics)

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Part of SPS zone 5 conference  
April 20-21, 2012

## Time, Einstein, and the Coolest Stuff in the Universe

A free public lecture by Nobel Laureate

Dr. William Phillips

National Institute of Standards and Technology

8:00 PM Friday April 20

Broadbent Recital Hall

[Brendle Recital Hall](#)  
Wake Forest University

[www.wfu.edu/physics/sps/spszone52012cont/welcome.html](http://www.wfu.edu/physics/sps/spszone52012cont/welcome.html)

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Estimate of partition function for interacting gas or fluid;  
virial expansion:

Free energy for interacting particles :

$$F(T, V, N) = -kT \ln Z_N = -kT \ln Z_{IG} - kT \ln Z_C$$

$$= F_{IG}(T, V, N) + F_C(T, V, N)$$

$$P = -\left(\frac{\partial F}{\partial V}\right)_{T,N} = \frac{NkT}{V} - \left(\frac{\partial F_C}{\partial V}\right)_{T,N}$$

$$\frac{PV}{NkT} = 1 + \frac{V}{N} \left( \frac{\partial \ln Z_C}{\partial V} \right)_{T,N}$$

$$Z_C = \frac{1}{V^N} \int d^3 r_1 d^3 r_2 \cdots d^3 r_N e^{-\beta U}$$

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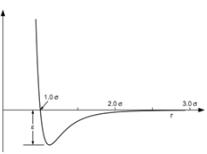
Model simulations of  $Z_C$

Isotropic pairwise interaction potentials :

$$U(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_N) = \sum_{i < j} u(|\mathbf{r}_i - \mathbf{r}_j|)$$

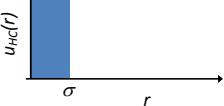
Lennard - Jones potential :

$$u_{LJ}(r) = 4\epsilon \left[ \left( \frac{\sigma}{r} \right)^{12} - \left( \frac{\sigma}{r} \right)^6 \right]$$



Hard - core potential :

$$u_{HC}(r) = \begin{cases} \infty & \text{for } r < \sigma \\ 0 & \text{for } r > \sigma \end{cases}$$



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For pair potential :

$$Z_C = \frac{1}{V^N} \int d^3 r_1 d^3 r_2 \cdots d^3 r_N e^{-\beta U}$$

$$= \frac{1}{V^N} \int d^3 r_1 d^3 r_2 \cdots d^3 r_N e^{-\beta u(r_{12}) - \beta u(r_{23}) - \beta u(r_{13}) - \cdots - \beta u(r_{(N-1)N})}$$

Define :

$$f_{ij} \equiv e^{-\beta u(r_{ij})} - 1$$

$$Z_C = \frac{1}{V^N} \int d^3 r_1 d^3 r_2 \cdots d^3 r_N (1 + f_{12})(1 + f_{23})(1 + f_{13}) \cdots (1 + f_{(N-1)N})$$

$$= 1 + \frac{N(N-1)}{2V^2} \int d^3 r_1 d^3 r_2 f_{12} + \frac{N(N-1)(N-2)}{6V^3} \int d^3 r_1 d^3 r_2 d^3 r_3 f_{12} f_{13} f_{23} \cdots$$

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"Virial" expansion of equation of state :

$$\frac{PV}{NkT} = 1 + \frac{V}{N} \left( \frac{\partial \ln Z_C}{\partial V} \right)_{T,N}$$

$$\approx 1 + \rho B_2(T) + \rho^2 B_3(T) + \dots$$

where  $\rho \equiv \frac{N}{V}$

Evaluation of  $B_2(T)$ :

$$\frac{N(N-1)}{2V^2} \int d^3 r_1 d^3 r_2 f_{12} = \frac{N(N-1)}{2V} \int d^3 r_{12} (e^{-\beta u(r_{12})} - 1)$$

$$\left( \frac{\partial}{\partial V} \left( \frac{N(N-1)}{2V} \int d^3 r_{12} (e^{-\beta u(r_{12})} - 1) \right) \right)_{T,N} \approx -\frac{\rho^2}{2} \int d^3 r_{12} (e^{-\beta u(r_{12})} - 1)$$

$$\Rightarrow B_2(T) = -\frac{1}{2} \int d^3 r_{12} (e^{-\beta u(r_{12})} - 1)$$

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Evaluation of  $B_2(T)$ :

$$B_2(T) = -\frac{1}{2} \int d^3 r_{12} (e^{-\beta u(r_{12})} - 1)$$

For hard sphere :  $B_2(T) = \frac{2\pi\sigma^3}{3}$

For square well :  $B_2(T) = \frac{2\pi\sigma^3}{3} \left( 1 - (e^{\rho\epsilon} - 1)(R^3 - 1) \right)$

Lennard - Jones potential :

$$u_{LJ}(r) = 4\epsilon \left[ \left( \frac{\sigma}{r} \right)^{12} - \left( \frac{\sigma}{r} \right)^6 \right]$$

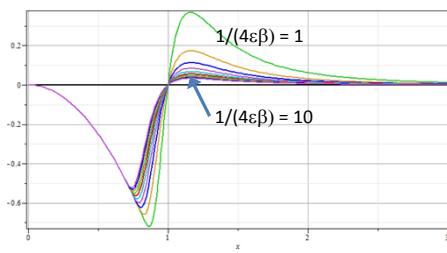
$$B_2(T) = -2\pi\sigma^3 \int_0^\infty x^2 dx \left( e^{-4\epsilon\beta \left( \frac{1}{x^{12}} - \frac{1}{x^6} \right)} - 1 \right)$$

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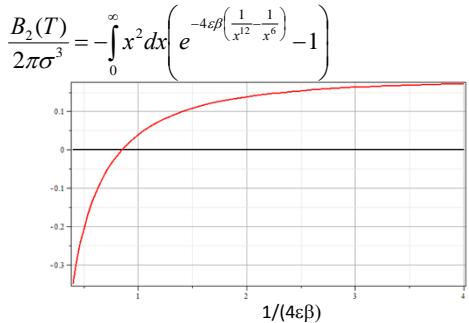
$$x^2 \left( e^{-4\epsilon\beta \left( \frac{1}{x^{12}} - \frac{1}{x^6} \right)} - 1 \right)$$



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Note: Correlates well with experimental measurements on N<sub>2</sub>, Ne, Ar, and He (less well)

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### Radial distribution function $g(r)$

Previously we have defined the correlation contribution to the partition function:

$$Z_C = \frac{1}{V^N} \int d^3r_1 d^3r_2 \cdots d^3r_N e^{-\beta U}$$

This can be rewritten :

$$Z_C = \int d^3r_1 d^3r_2 \left( \frac{1}{V^N} \int d^3r_3 \cdots d^3r_N e^{-\beta U} \right)$$

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$$Z_C = \int d^3r_1 d^3r_2 \left( \frac{1}{V^N} \int d^3r_3 \cdots d^3r_N e^{-\beta U} \right)$$

The quantity in () is related to the “pair distribution function”  $g(\mathbf{r}_1, \mathbf{r}_2)$  – given a particle at  $\mathbf{r}_1$  what is the probability of finding a second particle at  $\mathbf{r}_2$ .

$$Z_C = \int d^3r_1 d^3r_2 \left( \frac{1}{V^N} \int d^3r_3 \cdots d^3r_N e^{-\beta U} \right)$$

$$\rho^2 g(\mathbf{r}_1, \mathbf{r}_2) = \frac{N(N-1)}{Z_C V^N} \int d^3r_3 \cdots d^3r_N e^{-\beta U}$$

$$\text{where } \rho \equiv \frac{N}{V}$$

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Radial distribution function – continued --

$$\rho^2 g(\mathbf{r}_1, \mathbf{r}_2) = \frac{N(N-1)}{Z_c V^N} \int d^3 r_3 \cdots d^3 r_N e^{-\beta U}$$

If  $g(\mathbf{r}_1, \mathbf{r}_2) = g(|\mathbf{r}_1 - \mathbf{r}_2|)$  then

$$\rho \int 4\pi r^2 dr g(r) = N - 1 \approx N$$

Measurement of  $g(r)$  by neutron scattering:

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PHYSICAL REVIEW A VOLUME 7, NUMBER 6 JUNE 1973

**Structure Factor and Radial Distribution Function for Liquid Argon at 85 K<sup>†</sup>**

J. L. Yarnell, M. J. Katz, and R. G. Wenzel  
Los Alamos Scientific Laboratory, University of California, Los Alamos, New Mexico 87544

S. H. Koenig  
IBM Thomas J. Watson Research Center, Yorktown Heights, New York 10598

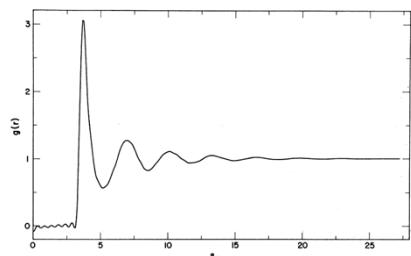


FIG. 4. Radial distribution function  $g(r)$  for  $^{36}\text{Ar}$  at 85 K. This is the Fourier transform of the smoothed and extended  $S(Q)$  shown as a solid line in Fig. 3.

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Experiment measures neutron scattering signal as a function of the change in neutron wave number  $\mathbf{Q}$ :

$$S(Q) = 1 + \frac{4\pi\rho}{Q} \int_0^\infty (g(r) - 1)r \sin(Qr) dr$$

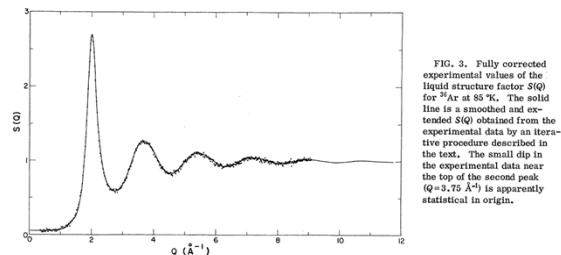


FIG. 3. Fully corrected experimental values of the liquid structure factor  $S(Q)$  for  $^{36}\text{Ar}$  at 85 K. The solid line is a smoothed and extended version of the experimental data by an iterative procedure described in the text. The small dip in the experimental data near the end of the second peak ( $Q=3.75 \text{ \AA}^{-1}$ ) is apparently statistical in origin.

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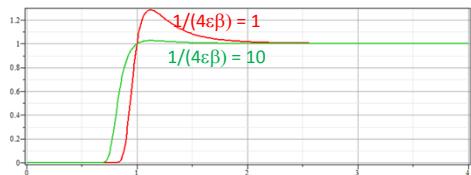
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Estimate of  $g(r)$  from theory:

For systems described by pair potentials in the dilute limit:

$$g(r) \approx e^{-\beta u(r)}$$



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