PHY 712 Electrodynamics
10-10:50 AM MWF Olin 107

Plan for Lecture 11:

Start reading Chapter 5

A. Magnetostatics

B. Vector potential

C. Example: current loop

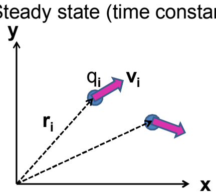
Course schedule for Spring 2014

(Preliminary schedule -- subject to frequent adjustment.) Please note that makeup lectures (indicated in red) are scheduled for Tuesdays or Thursdays at 11 AM - 12:15 PM in Olin 107.

	Lecture date	JDJ Reading	Topic	Assign.	Due date
1	Wed: 01/15/2014	Chap. 1	Introduction, units and Poisson equation	#1	01/31/2014
2	Thu: 01/16/2014	Chap. 1	Electrostatic energy calculations	#2	01/31/2014
3	Fri: 01/17/2014	Chap. 1	Poisson equation and Green's theorm	#3	01/31/2014
	Mon: 01/20/2014		MLK Holiday - no class		
4	Wed: 01/22/2014	Chap. 1	Green's functions for cartesian coordinates	#4	01/31/2014
5	Thu: 01/23/2014	Chap. 1	Brief introduction to numerical methods	#5	01/31/2014
6	Fri: 01/24/2014	Chap. 2	Method of images	#6	01/31/2014
	Mon: 01/27/2014		NAWH out of town - no class		
	Wed: 01/29/2014		NAWH out of town - no class		
7	Fri: 01/31/2014	Chap. 3	Cylindrical and spherical geometries	#7	02/05/2014
8	Mon: 02/03/2014	Chap. 4	Multipole analysis of charge distributions	#8	02/05/2014
9	Wed: 02/05/2014	Chap. 4	Dipoles and dielectrics	#9	02/07/2014
10	Fri: 02/07/2014	Chap. 4	Dipoles and dielectrics	#10	02/10/2014
11	Mon: 02/10/2014	Chap. 5	Magnetostatics	#11	02/12/2014

Magnetostatics

Magnetic flux density or magnetic induction field **B**
Steady state (time constant) current density **J**



$$\mathbf{J}(\mathbf{r}) = \sum_i q_i \mathbf{v}_i \delta^3(\mathbf{r} - \mathbf{r}_i)$$

Note that "statics" implies that $\nabla \cdot \mathbf{J} = 0$.

This follows from the continuity equation:

$$\frac{\partial \rho}{\partial t} + \nabla \cdot \mathbf{J} = 0$$

Comparison of electrostatics and magnetostatics

Electrostatic field due to charge density $\rho(\mathbf{r})$:

$$\mathbf{E}(\mathbf{r}) = \frac{1}{4\pi\varepsilon_0}\int d^3r'\rho(\mathbf{r}')\frac{\mathbf{r}-\mathbf{r}'}{|\mathbf{r}-\mathbf{r}'|^3}$$

Magnetostatic field due to current density $\mathbf{J}(\mathbf{r})$:

$$\mathbf{B}(\mathbf{r}) = \frac{\mu_0}{4\pi}\int d^3r'\,\mathbf{J}(\mathbf{r}')\times\frac{\mathbf{r}-\mathbf{r}'}{|\mathbf{r}-\mathbf{r}'|^3}$$

Alternative forms magnetostatic equations

Magnetostatic field due to current density $\mathbf{J}(\mathbf{r})$:

$$\mathbf{B}(\mathbf{r}) = \frac{\mu_0}{4\pi}\int d^3r'\,\mathbf{J}(\mathbf{r}')\times\frac{\mathbf{r}-\mathbf{r}'}{|\mathbf{r}-\mathbf{r}'|^3}$$

Note that: $\nabla\frac{1}{|\mathbf{r}-\mathbf{r}'|} = -\frac{\mathbf{r}-\mathbf{r}'}{|\mathbf{r}-\mathbf{r}'|^3}$

Also note that: $\nabla\times(s(\mathbf{r})\mathbf{V}(\mathbf{r})) = \nabla s(\mathbf{r})\times\mathbf{V}(\mathbf{r}) + s(\mathbf{r})\nabla\times\mathbf{V}(\mathbf{r})$

$$\Rightarrow \mathbf{B}(\mathbf{r}) = \frac{\mu_0}{4\pi}\nabla\times\int d^3r'\,\frac{\mathbf{J}(\mathbf{r}')}{|\mathbf{r}-\mathbf{r}'|}$$

Alternative forms magnetostatic equations -- continued

$$\mathbf{B}(\mathbf{r}) = \frac{\mu_0}{4\pi}\nabla\times\int d^3r'\,\frac{\mathbf{J}(\mathbf{r}')}{|\mathbf{r}-\mathbf{r}'|}$$

$\Rightarrow \nabla\cdot\mathbf{B}(\mathbf{r}) = 0$ No magnetic monopoles

$\Rightarrow \nabla\times\mathbf{B}(\mathbf{r}) = \mu_0\mathbf{J}(\mathbf{r})$ Ampere's law

"Proof" of Ampere's law for magnetostatic system:

$$\nabla\times\mathbf{B}(\mathbf{r}) = \frac{\mu_0}{4\pi}\nabla\times\nabla\times\int d^3r'\,\frac{\mathbf{J}(\mathbf{r}')}{|\mathbf{r}-\mathbf{r}'|}$$

Note that: $\nabla\times\nabla\times\mathbf{V} = \nabla(\nabla\cdot\mathbf{V}) - \nabla^2\mathbf{V}$

Recall that: $\nabla^2\frac{1}{|\mathbf{r}-\mathbf{r}'|} = -4\pi\delta^3(\mathbf{r}-\mathbf{r}')$ and $\nabla\cdot\mathbf{J}(\mathbf{r}) = 0$

Differential forms of magnetostatic equations:

$$\Rightarrow \nabla \cdot \mathbf{B}(\mathbf{r}) = 0 \qquad \text{No magnetic monopoles}$$
$$\Rightarrow \nabla \times \mathbf{B}(\mathbf{r}) = \mu_0 \mathbf{J}(\mathbf{r}) \qquad \text{Ampere's law}$$

Magnetostatic vector potential

$$\mathbf{B}(\mathbf{r}) = \nabla \times \mathbf{A}(\mathbf{r})$$
$$\mathbf{B}(\mathbf{r}) = \frac{\mu_0}{4\pi} \nabla \times \int d^3r' \frac{\mathbf{J}(\mathbf{r}')}{|\mathbf{r}-\mathbf{r}'|}$$
$$\Rightarrow \mathbf{A}(\mathbf{r}) = \frac{\mu_0}{4\pi} \int d^3r' \frac{\mathbf{J}(\mathbf{r}')}{|\mathbf{r}-\mathbf{r}'|} + \nabla s(\mathbf{r})$$

Non uniqueness of the magnetostatic vector potential

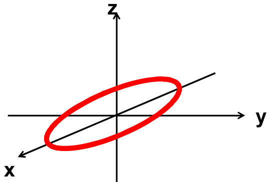
$$\text{Note that: } \quad \mathbf{B}(\mathbf{r}) = \nabla \times \mathbf{A}(\mathbf{r}) = \nabla \times \mathbf{A}'(\mathbf{r})$$
$$\text{if } \quad \mathbf{A}'(\mathbf{r}) = \mathbf{A}(\mathbf{r}) + \nabla s(\mathbf{r})$$

$$\text{Example: } \quad \text{for } \mathbf{B}(\mathbf{r}) = B_0 \hat{\mathbf{z}}$$
$$\mathbf{A}(\mathbf{r}) = \tfrac{1}{2} B_0 (x\hat{\mathbf{y}} - y\hat{\mathbf{x}})$$
$$\text{or} \quad \mathbf{A}(\mathbf{r}) = B_0 x \hat{\mathbf{y}}$$
$$\text{or} \quad \mathbf{A}(\mathbf{r}) = -B_0 y \hat{\mathbf{x}}$$

Differential form of Ampere's law in terms of vector potential:

$$\nabla \times \mathbf{B}(\mathbf{r}) = \nabla \times \nabla \times \mathbf{A}(\mathbf{r}) = \mu_0 \mathbf{J}(\mathbf{r})$$
$$\Rightarrow \nabla(\nabla \cdot \mathbf{A}(\mathbf{r})) - \nabla^2 \mathbf{A}(\mathbf{r}) = \mu_0 \mathbf{J}(\mathbf{r})$$
$$\text{If} \quad \nabla \cdot \mathbf{A}(\mathbf{r}) = 0 \quad \text{(Coulomb gauge)} \quad \Rightarrow \nabla^2 \mathbf{A}(\mathbf{r}) = -\mu_0 \mathbf{J}(\mathbf{r})$$
$$\mathbf{A}(\mathbf{r}) = \frac{\mu_0}{4\pi} \int d^3r' \frac{\mathbf{J}(\mathbf{r}')}{|\mathbf{r}-\mathbf{r}'|}$$

Magnetostatics example: current loop



$$\mathbf{J}(\mathbf{r}') = \frac{I}{a}\sin\theta'\,\delta(\cos\theta')\,\delta(r'-a)\left(-\sin\phi'\,\hat{\mathbf{x}} + \cos\phi'\,\hat{\mathbf{y}}\right)$$

$$\mathbf{A}(\mathbf{r}) = \frac{\mu_0}{4\pi}\int d^3r'\,\frac{\mathbf{J}(\mathbf{r}')}{|\mathbf{r}-\mathbf{r}'|}$$

Magnetostatics example: current loop -- continued

$$\mathbf{J}(\mathbf{r}') = \frac{I}{a}\sin\theta'\,\delta(\cos\theta')\,\delta(r'-a)\left(-\sin\phi'\,\hat{\mathbf{x}} + \cos\phi'\,\hat{\mathbf{y}}\right)$$

$$\mathbf{A}(\mathbf{r}) = \frac{\mu_0}{4\pi}\int d^3r'\,\frac{\mathbf{J}(\mathbf{r}')}{|\mathbf{r}-\mathbf{r}'|}$$

$$\mathbf{A}(\mathbf{r}) = \frac{\mu_0 I}{4\pi a}\int r'^2\,dr'\,d\cos\theta'\,d\phi'\,\frac{\sin\theta'\,\delta(\cos\theta')\,\delta(r'-a)\left(-\sin\phi'\,\hat{\mathbf{x}} + \cos\phi'\,\hat{\mathbf{y}}\right)}{\left(r^2 + r'^2 - 2rr'\left(\cos\theta\cos\theta' + \sin\theta\sin\theta'\cos(\phi-\phi')\right)\right)^{1/2}}$$

Completing integration over r' and θ':

$$\mathbf{A}(\mathbf{r}) = \frac{\mu_0 I a^2}{4\pi a}\int_0^{2\pi} d\phi'\,\frac{\left(-\sin\phi'\,\hat{\mathbf{x}} + \cos\phi'\,\hat{\mathbf{y}}\right)}{\left(r^2 + a^2 - 2ra\left(\sin\theta\cos(\phi-\phi')\right)\right)^{1/2}}$$

Let $\phi - \phi' \equiv \varphi$

$$\sin\phi' = \sin(\phi-\varphi) = \sin\phi\cos\varphi - \cos\phi\sin\varphi$$

$$\cos\phi' = \cos(\phi-\varphi) = \cos\phi\cos\varphi + \sin\phi\sin\varphi$$

Remaining non-trivial terms

$$\mathbf{A}(\mathbf{r}) = \frac{\mu_0 I a}{4\pi}\left(\sin\phi\,\hat{\mathbf{x}} - \cos\phi\,\hat{\mathbf{y}}\right)\int_0^{2\pi} d\varphi\,\frac{\cos\varphi}{\left(r^2 + a^2 - 2ra\left(\sin\theta\cos\varphi\right)\right)^{1/2}}$$

Magnetostatics example: current loop -- continued

$$\mathbf{A}(\mathbf{r}) = \frac{\mu_0 I a}{4\pi}\left(\sin\phi\,\hat{\mathbf{x}} - \cos\phi\,\hat{\mathbf{y}}\right)\int_0^{2\pi} d\varphi\,\frac{\cos\varphi}{\left(r^2 + a^2 - 2ra\left(\sin\theta\cos\varphi\right)\right)^{1/2}}$$

Elliptic integrals:

$$K(m) = \int_0^{\pi/2}\frac{du}{\left(1 - m\sin^2 u\right)^{1/2}}$$

$$E(m) = \int_0^{\pi/2}\left(1 - m\sin^2 u\right)^{1/2} du$$

$$\mathbf{A}(\mathbf{r}) = \frac{\mu_0}{4\pi}\,4Ia\,\frac{\left(\sin\phi\,\hat{\mathbf{x}} - \cos\phi\,\hat{\mathbf{y}}\right)}{\left(r^2 + a^2 + 2ra\sin\theta\right)^{1/2}}\left[\frac{\left(2-k^2\right)K(k) - 2E(k)}{k^2}\right]$$

where: $k^2 \equiv \dfrac{4ar\sin\theta}{r^2 + a^2 + 2ra\sin\theta}$

$$\mathbf{B}(\mathbf{r}) = \nabla\times\mathbf{A}(\mathbf{r})$$

Magnetostatics example: current loop -- continued

$$\mathbf{A}(\mathbf{r}) = \frac{\mu_0}{4\pi} 4Ia \frac{(\sin\phi\hat{\mathbf{x}} - \cos\phi\hat{\mathbf{y}})}{(r^2 + a^2 + 2ra\sin\theta)^{1/2}} \left[\frac{(2-k^2)K(k) - 2E(k)}{k^2}\right]$$

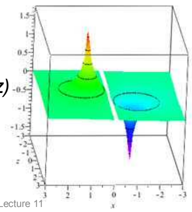
$$\text{where:} \quad k^2 \equiv \frac{4ar\sin\theta}{r^2 + a^2 + 2ra\sin\theta}$$

For $\phi = 0$: $x = r\sin\theta, \; y = 0$

$$\mathbf{A}(\mathbf{r}) = A_y(x,z)\hat{\mathbf{y}} = -\frac{\mu_0}{4\pi} 4Ia\hat{\mathbf{y}} \frac{1}{(x^2 + z^2 + a^2 + 2ax)^{1/2}} \left[\frac{(2-k^2)K(k) - 2E(k)}{k^2}\right]$$

$$\text{where:} \quad k^2 \equiv \frac{4ax}{x^2 + z^2 + a^2 + 2ax}$$

$A_y(x,z)$



Magnetostatics example: current loop -- continued

$$\mathbf{A}(\mathbf{r}) = \frac{\mu_0}{4\pi} 4Ia \frac{(\sin\phi\hat{\mathbf{x}} - \cos\phi\hat{\mathbf{y}})}{(r^2 + a^2 + 2ra\sin\theta)^{1/2}} \left[\frac{(2-k^2)K(k) - 2E(k)}{k^2}\right]$$

$$\text{where:} \quad k \equiv \frac{4ar\sin\theta}{r^2 + a^2 + 2ra\sin\theta}$$

$$\mathbf{B}(\mathbf{r}) = \nabla \times \mathbf{A}(\mathbf{r})$$

Note that for spherical polar coordinates: $\hat{\boldsymbol{\varphi}} = \sin\phi\hat{\mathbf{x}} - \cos\phi\hat{\mathbf{y}}$

$$\mathbf{A}(\mathbf{r}) = A_\phi(\mathbf{r})\hat{\boldsymbol{\varphi}}$$

$$\text{where} \quad A_\phi(\mathbf{r}) = \frac{\mu_0}{4\pi} \frac{4Ia}{(r^2 + a^2 + 2ra\sin\theta)^{1/2}} \left[\frac{(2-k^2)K(k) - 2E(k)}{k^2}\right]$$

$$\mathbf{B}(\mathbf{r}) = \frac{1}{r\sin\theta} \frac{\partial(\sin\theta A_\phi(\mathbf{r}))}{\partial\theta}\hat{\mathbf{r}} - \frac{1}{r}\frac{\partial(rA_\phi(\mathbf{r}))}{\partial r}\hat{\boldsymbol{\theta}}$$

For $r \to \infty$:

$$\mathbf{B}(\mathbf{r}) \approx \frac{\mu_0}{4\pi} \frac{I\pi a^2}{r^3} (2\cos\theta\hat{\mathbf{r}} + \sin\theta\hat{\boldsymbol{\theta}})$$

Other examples of current density sources:

Quantum mechanical expression for current density

for a particle of mass M and charge e and of probability amplitude $\Psi(\mathbf{r})$:

$$\mathbf{J}(\mathbf{r}) = -\frac{e\hbar}{2Mi}\left(\Psi^*(\mathbf{r})\nabla\Psi(\mathbf{r}) - \Psi(\mathbf{r})\nabla\Psi^*(\mathbf{r})\right)$$

For an electron in a spherical potential (such as in an atom):

$$\Psi(\mathbf{r}) \equiv \Psi_{nlm_l}(\mathbf{r}) = R_{nl}(r)Y_{lm_l}(\hat{\mathbf{r}})$$

$$\mathbf{J}(\mathbf{r}) = \frac{e\hbar}{2Mi}|R_{nl}(r)|^2 \frac{1}{r\sin\theta}\left(Y_{lm_l}^*(\hat{\mathbf{r}})\frac{\partial Y_{lm_l}(\hat{\mathbf{r}})}{\partial\phi} - Y_{lm_l}(\hat{\mathbf{r}})\frac{\partial Y_{lm_l}^*(\hat{\mathbf{r}})}{\partial\phi}\right)\hat{\boldsymbol{\varphi}}$$

$$= \frac{e\hbar}{M}\frac{m_l}{r\sin\theta}|\Psi_{nlm_l}(\mathbf{r})|^2\hat{\boldsymbol{\varphi}}$$

Note that: $\hat{\boldsymbol{\varphi}} = -\sin\phi\hat{\mathbf{x}} + \cos\phi\hat{\mathbf{y}} = \dfrac{\hat{\mathbf{z}} \times \mathbf{r}}{r\sin\theta}$

$$\mathbf{J}(\mathbf{r}) = \frac{e\hbar}{M}\frac{m_l}{r^2\sin^2\theta}|\Psi_{nlm_l}(\mathbf{r})|^2(\hat{\mathbf{z}} \times \mathbf{r})$$

Magnetic vector potential for this case:

$$\mathbf{A}(\mathbf{r}) = \frac{\mu_0}{4\pi}\int d^3r' \frac{\mathbf{J}(\mathbf{r}')}{|\mathbf{r}-\mathbf{r}'|}$$

$$\mathbf{J}(\mathbf{r}') = \frac{e\hbar}{M}\frac{m_l}{r'^2\sin^2\theta'}\left|\Psi_{nlm_l}(\mathbf{r}')^2\right|(\hat{\mathbf{z}}\times\mathbf{r}')$$

$$\mathbf{A}(\mathbf{r}) = \frac{\mu_0}{4\pi}\frac{e\hbar m_l}{M}\int d^3r' \frac{(\hat{\mathbf{z}}\times\mathbf{r}')}{|\mathbf{r}-\mathbf{r}'|}\frac{\left|\Psi_{nlm_l}(\mathbf{r}')^2\right|}{r'^2\sin^2\theta'}$$

For example: electron in the $nlm_l = 211$ state of H:

$$\left|\Psi_{211}(\mathbf{r}')^2\right| = \frac{1}{64\pi a^3}\left(\frac{r'}{a}\right)^2 e^{-r'/a}\sin^2\theta'$$

$$\mathbf{A}(\mathbf{r}) = -\frac{\mu_0}{8\pi}\frac{e\hbar}{M}\frac{(\hat{\mathbf{z}}\times\mathbf{r})}{r^3}\left[1-e^{-r/a}\left(1+\frac{r}{a}+\frac{r^2}{2a^2}+\frac{r^3}{8a^3}\right)\right]$$