

PHY 712 Electrodynamics
10-10:50 AM MWF Olin 107

Plan for Lecture 19:

Start reading Chap. 8 in Jackson.

A. Examples of waveguides

B. TEM, TE, TM modes

C. Resonant cavities

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Wed 01/29/2014			NAWH out of town - no class		
7 Fri 01/31/2014	Chap. 3	Cylindrical and spherical geometries	#7	02/05/2014	
8 Mon 02/03/2014	Chap. 4	Multipole analysis of charge distributions	#8	02/05/2014	
9 Wed 02/05/2014	Chap. 4	Dipoles and dielectrics	#9	02/07/2014	
10 Fri 02/07/2014	Chap. 4	Dipoles and dielectrics	#10	02/10/2014	
11 Mon 02/10/2014	Chap. 5	Magnetostatics	#11	02/12/2014	
12 Wed 02/12/2014	Chap. 5	Magnetostatics	#12	02/14/2014	
Fri 02/14/2014		Class cancelled because of weather			
13 Mon 02/17/2014	Chap. 5	Magnetostatics	#13	02/28/2014	
14 Mon 02/17/2014	Chap. 6	Maxwell's equations	#14	02/28/2014	
15 Wed 02/19/2014	Chap. 6	Electromagnetic energy and force	#15	02/28/2014	
16 Fri 02/21/2014	Chap. 7	Electromagnetic plane waves	#16	02/28/2014	
17 Fri 02/21/2014	Chap. 7	Dynamic dielectric media and their effects	#17	02/28/2014	
Mon 02/24/2014		No class -- NAWH out of town			
Wed 02/26/2014		No class -- NAWH out of town			
18 Fri 02/28/2014	Chap. 7	Complex dielectrics, TEM modes			
Mon 03/03/2014	APS Meeting	Take-home exam (no class meeting)			
Wed 03/05/2014	APS Meeting	Take-home exam (no class meeting)			
Fri 03/07/2014	APS Meeting	Take-home exam (no class meeting)			
Mon 03/10/2014	Spring Break				
Wed 03/12/2014	Spring Break				
Fri 03/14/2014	Spring Break				
19 Mon 03/17/2014	Chap. 8	Wave guides, Take-home exam due	#18	3/19/2014	
20 Wed 03/19/2014	Chap. 8	Wave guides	#19	3/21/2014	

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Reminder:

- Topic choices for "computational project"
- suggestions available upon request
- due Monday March 24th?

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Fields near the surface on an ideal conductor

Suppose for an isotropic medium: $\mathbf{D} = \epsilon_b \mathbf{E}$ $\mathbf{J} = \sigma \mathbf{E}$

Maxwell's equations in terms of \mathbf{H} and \mathbf{E} :

$$\nabla \cdot \mathbf{E} = 0 \quad \nabla \cdot \mathbf{H} = 0$$

$$\nabla \times \mathbf{E} = -\mu \frac{\partial \mathbf{H}}{\partial t} \quad \nabla \times \mathbf{H} = \sigma \mathbf{E} + \epsilon_b \frac{\partial \mathbf{E}}{\partial t}$$

$$\left(\nabla^2 - \mu \sigma \frac{\partial}{\partial t} - \mu \epsilon_b \frac{\partial^2}{\partial t^2} \right) \mathbf{F} = 0 \quad \mathbf{F} = \mathbf{E}, \mathbf{H}$$

Plane wave form for \mathbf{E} :

$$\mathbf{E}(\mathbf{r}, t) = \Re(\mathbf{E}_0 e^{i\mathbf{k} \cdot \mathbf{r} - i\omega t}) \quad \text{where } \mathbf{k} = (n_R + in_I) \frac{\omega}{c} \hat{\mathbf{k}}$$

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Fields near the surface on an ideal conductor -- continued

For our system:

$$\frac{\omega}{c} n_R = \omega \sqrt{\frac{\mu \epsilon_b}{2} \left(\sqrt{1 + \left(\frac{\sigma}{\omega \epsilon_b} \right)^2} + 1 \right)^{1/2}}$$

$$\frac{\omega}{c} n_I = \omega \sqrt{\frac{\mu \epsilon_b}{2} \left(\sqrt{1 + \left(\frac{\sigma}{\omega \epsilon_b} \right)^2} - 1 \right)^{1/2}}$$

$$\text{For } \frac{\sigma}{\omega} \gg 1 \quad \frac{\omega}{c} n_R \approx \frac{\omega}{c} n_I \approx \sqrt{\frac{\mu \sigma \omega}{2}} \equiv \frac{1}{\delta}$$

$$\Rightarrow \mathbf{E}(\mathbf{r}, t) \approx e^{-\hat{\mathbf{k}} \cdot \mathbf{r} / \delta} \Re(\mathbf{E}_0 e^{i\hat{\mathbf{k}} \cdot \mathbf{r} - i\omega t})$$

$$\Rightarrow \mathbf{H}(\mathbf{r}, t) \approx \Re\left(\frac{n}{c\mu} \hat{\mathbf{k}} \times \mathbf{E}(\mathbf{r}, t) \right) \approx \Re\left(\frac{1+i}{\delta \mu \omega} \hat{\mathbf{k}} \times \mathbf{E}(\mathbf{r}, t) \right)$$

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Fields near the surface on an ideal conductor -- continued

$$\text{For } \frac{\sigma}{\omega} \gg 1 \quad \frac{\omega}{c} n_R \approx \frac{\omega}{c} n_I \approx \sqrt{\frac{\mu \sigma \omega}{2}} \equiv \frac{1}{\delta}$$

$$\text{In this limit, } \sqrt{\frac{\mu \epsilon}{\mu_0 \epsilon_0}} = c \sqrt{\mu \epsilon} = n_R + in_I = \frac{c}{\omega} \frac{1+i}{\delta}$$

$$\mathbf{E}(\mathbf{r}, t) = e^{-\hat{\mathbf{k}} \cdot \mathbf{r} / \delta} \Re(\mathbf{E}_0 e^{i\hat{\mathbf{k}} \cdot \mathbf{r} - i\omega t})$$

$$\mathbf{D}(\mathbf{r}, t) = \epsilon \mathbf{E}(\mathbf{r}, t) = \frac{i\sigma}{\omega} \mathbf{E}(\mathbf{r}, t)$$

$$\mathbf{H}(\mathbf{r}, t) = \frac{n}{c\mu} \hat{\mathbf{k}} \times \mathbf{E}(\mathbf{r}, t) = \frac{1+i}{\delta \mu \omega} \hat{\mathbf{k}} \times \mathbf{E}(\mathbf{r}, t)$$

$$\mathbf{B}(\mathbf{r}, t) = \mu \mathbf{H}(\mathbf{r}, t) = \frac{n}{c} \hat{\mathbf{k}} \times \mathbf{E}(\mathbf{r}, t) = \frac{1+i}{\delta \omega} \hat{\mathbf{k}} \times \mathbf{E}(\mathbf{r}, t)$$

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Fields near the surface on an ideal conductor -- continued

$$\mathbf{E}(\mathbf{r}, t) = e^{-\hat{\mathbf{k}} \cdot \mathbf{r} / \delta} \Re(\mathbf{E}_0 e^{i\hat{\mathbf{k}} \cdot \mathbf{r} / \delta - i\omega t})$$

$$\mathbf{D}(\mathbf{r}, t) = \epsilon \mathbf{E}(\mathbf{r}, t) = \frac{i\sigma}{\omega} \mathbf{E}(\mathbf{r}, t)$$

$$\mathbf{H}(\mathbf{r}, t) = \frac{n}{c\mu} \hat{\mathbf{k}} \times \mathbf{E}(\mathbf{r}, t) = \frac{1+i}{\delta\mu\omega} \hat{\mathbf{k}} \times \mathbf{E}(\mathbf{r}, t)$$

$$\mathbf{B}(\mathbf{r}, t) = \mu \mathbf{H}(\mathbf{r}, t) = \frac{n}{c} \hat{\mathbf{k}} \times \mathbf{E}(\mathbf{r}, t) = \frac{1+i}{\delta\omega} \hat{\mathbf{k}} \times \mathbf{E}(\mathbf{r}, t)$$

Note that we can express the results in terms of the \mathbf{H} field:

$$\mathbf{H}(\mathbf{r}, t) = e^{-\hat{\mathbf{k}} \cdot \mathbf{r} / \delta} \Re(\mathbf{H}_0 e^{i\hat{\mathbf{k}} \cdot \mathbf{r} / \delta - i\omega t})$$

$$\mathbf{E}(\mathbf{r}, t) = \delta\mu\omega \Re\left(\frac{1-i}{2} \hat{\mathbf{k}} \times \mathbf{H}(\mathbf{r}, t)\right)$$

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Boundary values for ideal conductor

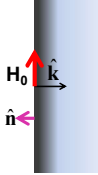
Inside the conductor:

$$\mathbf{H}(\mathbf{r}, t) = e^{-\hat{\mathbf{k}} \cdot \mathbf{r} / \delta} \Re(\mathbf{H}_0 e^{i\hat{\mathbf{k}} \cdot \mathbf{r} / \delta - i\omega t})$$

$$\mathbf{E}(\mathbf{r}, t) = \delta\mu\omega \frac{1-i}{2} \hat{\mathbf{k}} \times \mathbf{H}(\mathbf{r}, t)$$

At the boundary of an ideal conductor, the \mathbf{E} and \mathbf{H} fields decay in the direction normal to the interface.

Ideal conductor boundary conditions:

$$\hat{\mathbf{n}} \times \mathbf{E}|_S = 0 \quad \hat{\mathbf{n}} \cdot \mathbf{H}|_S = 0$$


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Waveguide terminology

- TEM: transverse electric and magnetic (both E and H fields are perpendicular to wave propagation direction)
- TM: transverse magnetic (H field is perpendicular to wave propagation direction)
- TE: transverse electric (E field is perpendicular to wave propagation direction)

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Wave guides

Top view:

Inside medium, μ, ϵ assumed to be real

Coaxial cable TEM modes

(following problem 8.2 in Jackson's text)

Maxwell's equations inside medium: for $a \leq \rho \leq b$

$$\nabla \times \mathbf{E} = i\omega \mathbf{B} \quad \nabla \cdot \mathbf{E} = 0$$

$$\nabla \times \mathbf{B} = -i\omega \mu \epsilon \mathbf{E} \quad \nabla \cdot \mathbf{B} = 0$$

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Electromagnetic waves in a coaxial cable -- continued

Top view: Example solution for $a \leq \rho \leq b$

$$\mathbf{E} = \hat{\rho} \Re \left(\frac{E_0 a}{\rho} e^{ikz - i\omega t} \right) \quad \text{Find:}$$

$$\mathbf{B} = \hat{\phi} \Re \left(\frac{B_0 a}{\rho} e^{ikz - i\omega t} \right) \quad k = \omega \sqrt{\mu \epsilon}$$

$$E_0 = \frac{B_0}{\sqrt{\mu \epsilon}}$$

$$\hat{\rho} = \cos \phi \hat{x} + \sin \phi \hat{y}$$

$$\hat{\phi} = -\sin \phi \hat{x} + \cos \phi \hat{y}$$

Poynting vector within cable medium (with μ, ϵ):

$$\langle \mathbf{S} \rangle_{\text{avg}} = \frac{1}{2\mu} \Re(\mathbf{E} \times \mathbf{B}^*) = \frac{|B_0|^2}{2\mu\sqrt{\mu\epsilon}} \left(\frac{a}{\rho} \right)^2 \hat{z}$$

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Electromagnetic waves in a coaxial cable -- continued

Top view:

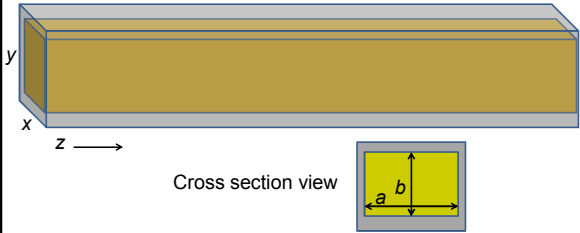
Time averaged power in cable material:

$$\int_0^{2\pi} d\phi \int_a^b \rho d\rho \langle \mathbf{S} \rangle_{\text{avg}} \cdot \hat{z} = \frac{|B_0|^2 \pi a^2}{\mu \sqrt{\mu \epsilon}} \ln \left(\frac{b}{a} \right)$$

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Analysis of rectangular waveguide

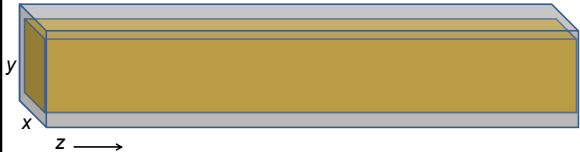
Boundary conditions at surface of waveguide:
 $E_{\text{tangential}}=0, \mathbf{B}_{\text{normal}}=0$



Cross section view

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Analysis of rectangular waveguide



$\mathbf{B} = \Re\{B_x(x, y)\hat{x} + B_y(x, y)\hat{y} + B_z(x, y)\hat{z}\}e^{ikz-i\omega t}$

$\mathbf{E} = \Re\{E_x(x, y)\hat{x} + E_y(x, y)\hat{y} + E_z(x, y)\hat{z}\}e^{ikz-i\omega t}$

$$k^2 = \mu\epsilon\omega^2 - \left(\frac{m\pi}{a}\right)^2 - \left(\frac{n\pi}{b}\right)^2$$

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Maxwell's equations within the pipe:

$$\frac{\partial B_x}{\partial x} + \frac{\partial B_y}{\partial y} + ikB_z = 0.$$

$$\frac{\partial E_x}{\partial x} + \frac{\partial E_y}{\partial y} + ikE_z = 0.$$

$$\frac{\partial E_z}{\partial y} - ikE_y = i\omega B_x, \quad \frac{\partial B_z}{\partial y} - ikB_y = -i\mu\epsilon\omega E_x.$$

$$ikE_x - \frac{\partial E_z}{\partial x} = i\omega B_y, \quad ikB_x - \frac{\partial B_z}{\partial x} = -i\mu\epsilon\omega E_y.$$

$$\frac{\partial E_y}{\partial x} - \frac{\partial E_x}{\partial y} = i\omega B_z, \quad \frac{\partial B_y}{\partial x} - \frac{\partial B_x}{\partial y} = -i\mu\epsilon\omega E_z.$$

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Solution of Maxwell's equations within the pipe:

Combining Faraday's Law and Ampere's Law, we find that each field component must satisfy a two-dimensional Helmholtz equation:

$$\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} - k^2 + \mu\epsilon\omega^2\right) E_x(x, y) = 0.$$

For the rectangular wave guide discussed in Section 8.4 of your text a solution for a TE mode can have:

$$E_z(x, y) \equiv 0 \quad \text{and} \quad B_z(x, y) = B_0 \cos\left(\frac{m\pi x}{a}\right) \cos\left(\frac{n\pi y}{b}\right),$$

$$\text{with } k^2 \equiv k_{mn}^2 = \mu\epsilon\omega^2 - \left[\left(\frac{m\pi}{a}\right)^2 + \left(\frac{n\pi}{b}\right)^2\right]$$

Some of the other field components are:

$$B_x = -\frac{k}{\omega} E_y \quad \text{and} \quad B_y = \frac{k}{\omega} E_x.$$

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TE modes for rectangular wave guide continued:

$$E_z(x, y) \equiv 0 \quad \text{and} \quad B_z(x, y) = B_0 \cos\left(\frac{m\pi x}{a}\right) \cos\left(\frac{n\pi y}{b}\right),$$

$$E_x = \frac{\omega}{k} B_y = \frac{-i\omega}{k^2 - \mu\epsilon\omega^2} \frac{\partial B_z}{\partial y} = \frac{-i\omega}{\left[\left(\frac{m\pi}{a}\right)^2 + \left(\frac{n\pi}{b}\right)^2\right]} \frac{n\pi}{b} B_0 \cos\left(\frac{m\pi x}{a}\right) \sin\left(\frac{n\pi y}{b}\right),$$

$$E_y = -\frac{\omega}{k} B_x = \frac{i\omega}{k^2 - \mu\epsilon\omega^2} \frac{\partial B_z}{\partial x} = \frac{i\omega}{\left[\left(\frac{m\pi}{a}\right)^2 + \left(\frac{n\pi}{b}\right)^2\right]} \frac{m\pi}{a} B_0 \sin\left(\frac{m\pi x}{a}\right) \cos\left(\frac{n\pi y}{b}\right).$$

Check boundary conditions:

$$\mathbf{E}_{\text{tangential}} = 0 \quad \text{because:} \quad E_x(x, 0) = E_x(x, b) = 0$$

$$\text{and } E_y(0, y) = E_y(a, y) = 0.$$

$$\mathbf{B}_{\text{normal}} = 0$$

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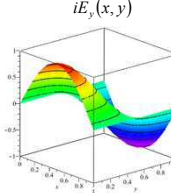
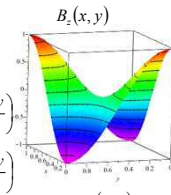
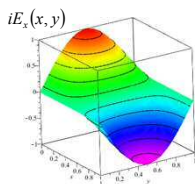
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Solution for m=n=1

$$B_z(x, y) = B_0 \cos\left(\frac{m\pi x}{a}\right) \cos\left(\frac{n\pi y}{b}\right)$$

$$iE_x(x, y) = B_0 \frac{\omega m\pi / b}{\left(\frac{m\pi}{a}\right)^2 + \left(\frac{n\pi}{b}\right)^2} \cos\left(\frac{m\pi x}{a}\right) \sin\left(\frac{n\pi y}{b}\right)$$

$$iE_y(x, y) = B_0 \frac{-\omega m\pi / a}{\left(\frac{m\pi}{a}\right)^2 + \left(\frac{n\pi}{b}\right)^2} \sin\left(\frac{m\pi x}{a}\right) \cos\left(\frac{n\pi y}{b}\right)$$



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Solution for $m=n=1$

$$k^2 \equiv k_{mn}^2 = \mu\epsilon\omega^2 - \left[\left(\frac{m\pi}{a} \right)^2 + \left(\frac{n\pi}{b} \right)^2 \right]$$

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Resonant cavity

$$0 \leq x \leq a$$

$$0 \leq y \leq b$$

$$0 \leq z \leq d$$

$$\mathbf{B} = \Re\{B_x(x, y, z)\hat{x} + B_y(x, y, z)\hat{y} + B_z(x, y, z)\hat{z}\}e^{-i\omega t}$$

$$\mathbf{E} = \Re\{E_x(x, y, z)\hat{x} + E_y(x, y, z)\hat{y} + E_z(x, y, z)\hat{z}\}e^{-i\omega t}$$

In general: $E_i(x, y, z) = E_i(x, y)\sin(kz)$ or $E_i(x, y)\cos(kz)$
 $B_i(x, y, z) = B_i(x, y)\sin(kz)$ or $B_i(x, y)\cos(kz)$

$$\Rightarrow k = \frac{p\pi}{d}$$

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Resonant cavity

$$0 \leq x \leq a$$

$$0 \leq y \leq b$$

$$0 \leq z \leq d$$

$$k^2 = \left(\frac{p\pi}{d} \right)^2 = \mu\epsilon\omega^2 - \left(\frac{m\pi}{a} \right)^2 - \left(\frac{n\pi}{b} \right)^2$$

$$\Rightarrow \omega^2 = \frac{1}{\mu\epsilon} \left[\left(\frac{m\pi}{a} \right)^2 + \left(\frac{n\pi}{b} \right)^2 + \left(\frac{p\pi}{d} \right)^2 \right]$$

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