Electrodynamics – PHY712 Lecture 2 – Ewald summation methods Reference: Chap. 1.11 in J. D. Jackson's textbook.

- 1. Motivation
- 2. Expression to calculate electrostatic energy of an extended periodic system and its derivation
- 3. Examples



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Lecture 2 – Ewald summation methods – Motivation

Consider a collection of point charges $\{q_i\}$ located at points $\{\mathbf{r}_i\}$. The energy to separate these charges to infinity ($\{\mathbf{r}_i \to \infty\}$ is

$$W = \frac{1}{4\pi\epsilon_0} \sum_{(i,j;i>j)} \frac{q_i q_j}{|\mathbf{r}_i - \mathbf{r}_j|}.$$
(1)

Here the summation is over all pairs of (i, j), excluding i = j. It is convenient to sum over all particles and divide by 2 to compensate for the double counting:

$$W = \frac{1}{8\pi\epsilon_0} \sum_{i,j;i\neq j} \frac{q_i q_j}{|\mathbf{r}_i - \mathbf{r}_j|}.$$
(2)

Here the summation is over all pairs of i, j, excluding i = j. The energy W scales as the number of particles N. As $N \to \infty$, the ratio W/N remains well-defined in principle, but difficult to calculate in practice.





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Lecture 2 – Ewald summation methods – slight digression

When the discrete charge distribution becomes a continuous charge density: $q_i \rightarrow \rho(\mathbf{r})$, the electrostatic energy becomes

$$W = \frac{1}{8\pi\epsilon_0} \int d^3r \ d^3r' \ \frac{\rho(\mathbf{r})\rho(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|}.$$
(3)

Notice, in this case, it is not possible to exclude the "self-interaction". This expression can be written in terms of the electrostatic potential $\Phi(\mathbf{r})$ and field $\mathbf{E}(\mathbf{r})$:

$$W = \frac{1}{2} \int d^3 r \ \rho(\mathbf{r}) \Phi(\mathbf{r}) = -\frac{\epsilon_0}{2} \int d^3 r \left(\nabla^2 \Phi(\mathbf{r})\right) \Phi(\mathbf{r}). \tag{4}$$

$$W = \frac{\epsilon_0}{2} \int d^3 r \left| \nabla \Phi(\mathbf{r}) \right|^2 = \frac{\epsilon_0}{2} \int d^3 r \left| \mathbf{E}(\mathbf{r}) \right|^2.$$
(5)



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Lecture 2 – Ewald summation methods – exact result for periodic system

$$\frac{W}{N} = \sum_{\alpha\beta} \frac{q_{\alpha}q_{\beta}}{8\pi\varepsilon_0} \left(\frac{4\pi}{\Omega} \sum_{\mathbf{G}\neq\mathbf{0}} \frac{e^{-i\mathbf{G}\cdot\tau_{\alpha\beta}} e^{-G^2/\eta}}{G^2} - \sqrt{\frac{\eta}{\pi}} \delta_{\alpha\beta} + \sum_{\mathbf{T}}' \frac{\operatorname{erfc}(\frac{1}{2}\sqrt{\eta}|\tau_{\alpha\beta}+\mathbf{T}|)}{|\tau_{\alpha\beta}+\mathbf{T}|} \right) - \frac{4\pi Q^2}{8\pi\varepsilon_0 \Omega\eta}.$$
(6)



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