

**PHY 712 Electrodynamics  
9-9:50 AM MWF Olin 105**

**Plan for Lecture 11:**

**Review of Electrostatics - Chapters 1-4**

- Homework problems and related
- Questions

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**Course schedule for Spring 2018**

(Preliminary schedule -- subject to frequent adjustment.)

Lecture date	JDJ Reading	Topic	HW	Due date
Wed: 01/17/2018	No class	Snow		
1 Fri: 01/19/2018	Chap. 1 & Appen.	Introduction, units and Poisson equation	#1	01/26/2018
2 Mon: 01/22/2018	Chap. 1	Electrostatic energy calculations	#2	01/26/2018
3 Wed: 01/24/2018	Chap. 1	Poisson's equation and Green's theorem	#3	01/26/2018
4 Thu: 01/25/2018	Chap. 1 & 2	Poisson's equation in 2 and 3 dimensions		
5 Fri: 01/26/2018	Chap. 1 & 2	Brief introduction to numerical methods	#4	01/29/2018
6 Mon: 01/29/2018	Chap. 2	Method of image charges	#5	01/31/2018
7 Wed: 01/31/2018	Chap. 2 & 3	Cylindrical and spherical geometries	#6	02/02/2018
8 Fri: 02/02/2018	Chap. 3 & 4	Multipole analysis	#7	02/07/2018
9 Mon: 02/05/2018	Chap. 4	Dipoles and Dielectrics	#8	02/09/2018
10 Wed: 02/07/2018	Chap. 4	Dipoles and Dielectrics		
11 Fri: 02/09/2018	Chap. 1-4	Review		
12 Mon: 02/12/2018				
13 Wed: 02/14/2018				
14 Fri: 02/16/2018				
15 Mon: 02/19/2018				

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**HW #4**

1. Consider a two-dimensional charge distribution of the form:

$$\rho(x) = \rho_0 \sin\left(\frac{\pi x}{a}\right) \sin\left(\frac{2\pi y}{a}\right),$$

where  $\rho_0$  represents a density constant and  $a$  represents a length constant. In the problem, you are asked to determine the electrostatic potential  $\Phi(x, y)$  for  $0 \leq x \leq a$  and  $0 \leq y \leq a$ , which satisfies the Poisson equation for the charge density  $\rho(x, y)$ , and satisfies the boundary conditions  $\Phi(0, y) = \Phi(a, y) = \Phi(x, 0) = \Phi(x, a) = 0$ .

- (a) Find the analytic form of the electrostatic potential  $\Phi(x, y)$  for  $0 \leq x \leq a$  and  $0 \leq y \leq a$ .
- (b) Using the finite difference method for the two grids discussed in class, find  $\Phi(x, y)$  on the grid points.
- (c) Using the finite element method for the two grids discussed in class, find  $\Phi(x, y)$  on the grid points.
- (d) Compare the accuracy of the numerical solutions for this example.

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$$\nabla^2 \Phi(x, y) = -\frac{\rho(x, y)}{\epsilon_0}$$
 In this case:  $\rho(x, y) = \rho_0 \sin\left(\frac{\pi x}{a}\right) \sin\left(\frac{2\pi y}{a}\right)$

Solution which vanishes on boundaries:  $x=0, x=a, y=0,$  and  $y=a$

$$\Phi(x, y) = \frac{\rho_0 a^2}{\epsilon_0 5\pi^2} \sin\left(\frac{\pi x}{a}\right) \sin\left(\frac{2\pi y}{a}\right)$$

Solution by finite difference method --

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$$\Phi(x, 4h) = 0$$

$h = \frac{a}{4}$

Algorithm:

$$\Phi(x, y) - \frac{1}{5} S_A - \frac{1}{20} S_B = \frac{3h^2}{10\epsilon_0} \rho(x, y) + \frac{h^4}{40\epsilon_0} \nabla^2 \rho(x, y).$$

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In our case:

$$\Phi(x, y) - \frac{1}{5} S_A - \frac{1}{20} S_B = \frac{3h^2}{10\epsilon_0} \rho(x, y) + \frac{h^4}{40\epsilon_0} \nabla^2 \rho(x, y).$$

$$\Phi(x_m, y_n) = \Phi\left(\frac{ma}{4}, \frac{na}{4}\right)$$

$$\begin{aligned}
 RHS(x_m, y_n) &= \frac{3h^2}{10\epsilon_0} \rho\left(\frac{ma}{4}, \frac{na}{4}\right) + \frac{h^4}{40\epsilon_0} \nabla^2 \rho\left(\frac{ma}{4}, \frac{na}{4}\right) \\
 &= \frac{a^2 \rho_0}{\epsilon_0} \left( \frac{3}{160} - \frac{5\pi^2}{10240} \right) \sin\left(\frac{\pi m}{4}\right) \sin\left(\frac{2\pi n}{4}\right)
 \end{aligned}$$

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$$\begin{pmatrix} 1 & -1/5 & -1/5 & -1/20 & 0 & 0 \\ -2/5 & 1 & -1/10 & -1/5 & 0 & 0 \\ -1/5 & -1/20 & 1 & -1/5 & -1/5 & -1/20 \\ -1/10 & -1/5 & -2/5 & 1 & -1/10 & -1/5 \\ 0 & 0 & -1/5 & -1/20 & 1 & -1/5 \\ 0 & 0 & -1/10 & -1/5 & -2/5 & 1 \end{pmatrix} \begin{pmatrix} \Phi(h, 3h) \\ \Phi(2h, 3h) \\ \Phi(h, 2h) \\ \Phi(2h, 2h) \\ \Phi(h, h) \\ \Phi(2h, h) \end{pmatrix} = \begin{pmatrix} RHS(h, 3h) \\ RHS(2h, 3h) \\ RHS(h, 2h) \\ RHS(2h, 2h) \\ RHS(h, h) \\ RHS(2h, h) \end{pmatrix}$$

For Maple: `>with(LinearAlgebra)`  
`>evalm(MatrixInverse(LHS).RHS)`

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Solution by finite element method --

$$\Phi(x, y) = \sum_{ij} \psi_{ij} \phi_{ij}(x, y),$$

$$\iint dx dy \phi_{kl}(x, y) (-\nabla^2 \Phi(x, y)) = \frac{1}{\epsilon_0} \iint dx dy \phi_{kl}(x, y) \rho(x, y)$$

$$\sum_{ij} M_{kl, ij} \psi_{ij} = G_{kl},$$

$$M_{kl, ij} \equiv \int dx \int dy \nabla \phi_{kl}(x, y) \cdot \nabla \phi_{ij}(x, y)$$

$$G_{kl} \equiv \int dx \int dy \phi_{kl}(x, y) \rho(x, y) / \epsilon_0.$$

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$$\phi_{ij}(x, y) \equiv \mathcal{X}_i(x) \mathcal{Y}_j(y),$$

$$\mathcal{X}_i(x) \equiv \begin{cases} \left(1 - \frac{|x-x_i|}{h}\right) & \text{for } x_i - h \leq x \leq x_i + h \\ 0 & \text{otherwise} \end{cases}$$

$$G_{kl} = \frac{\rho_0}{\epsilon_0} X_k Y_l$$

where  $X_k = \int_{(k-1)a/4}^{(k+1)a/4} dx \left(1 - \frac{|x-ka/4|}{a/4}\right) \sin\left(\frac{\pi x}{a}\right)$

$$Y_l = \int_{(l-1)a/4}^{(l+1)a/4} dy \left(1 - \frac{|y-la/4|}{a/4}\right) \sin\left(\frac{2\pi y}{a}\right)$$

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$$\begin{pmatrix} 8/3 & -1/3 & -1/3 & -1/3 & 0 & 0 \\ -2/3 & 8/3 & -2/3 & -1/3 & 0 & 0 \\ -1/3 & -1/3 & 8/3 & -1/3 & -1/3 & -1/3 \\ -2/3 & -1/3 & -2/3 & 8/3 & -2/3 & -1/3 \\ 0 & 0 & -1/3 & -1/3 & 8/3 & -1/3 \\ 0 & 0 & -2/3 & -1/3 & -2/3 & 8/3 \end{pmatrix} \begin{pmatrix} \Phi(h, 3h) \\ \Phi(2h, 3h) \\ \Phi(h, 2h) \\ \Phi(2h, 2h) \\ \Phi(h, h) \\ \Phi(2h, h) \end{pmatrix} = \begin{pmatrix} G(h, 3h) \\ G(2h, 3h) \\ G(h, 2h) \\ G(2h, 2h) \\ G(h, h) \\ G(2h, h) \end{pmatrix}$$

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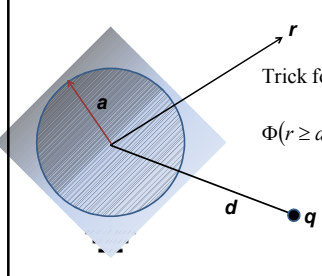
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HW #5 Find total surface charge for  $d > a$  and for  $d < a$   
 A grounded metal sphere of radius  $a$ , in the presence of a point charge  $q$  at a distance  $d$  from its center.



Trick for  $r \geq a$ :

$$\Phi(r \geq a) = \frac{1}{4\pi\epsilon_0} \left( \frac{q}{|\mathbf{r} - \mathbf{d}|} - \frac{q}{\frac{d}{a} |\mathbf{r} - \mathbf{d} \frac{a^2}{d^2}|} \right)$$

Interpreted as  
 Image charge of  $q' = -q \frac{a}{d}$   
 Located along  $\hat{\mathbf{d}}$  at  $\hat{\mathbf{d}} \frac{a}{d}$

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For  $d > a$

Surface charge density:

$$\sigma(\hat{\mathbf{r}}) = -\epsilon_0 \frac{\partial \Phi}{\partial r} \Big|_{r=a} = -\frac{q}{4\pi a^2} \frac{a}{d} \frac{(1 - \frac{a^2}{d^2})}{(1 + \frac{a^2}{d^2} - 2\frac{a}{d} \hat{\mathbf{r}} \cdot \hat{\mathbf{d}})^{3/2}}$$

Let  $\hat{\mathbf{r}} \cdot \hat{\mathbf{d}} \equiv \cos \gamma$

$$Q = 2\pi a^2 \int_{-1}^1 d\cos \gamma \left( -\frac{q}{4\pi a^2} \frac{a}{d} \frac{(1 - \frac{a^2}{d^2})}{(1 + \frac{a^2}{d^2} - 2\frac{a}{d} \cos \gamma)^{3/2}} \right)$$

$$= -\frac{qa}{d}$$

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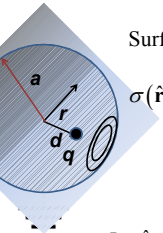
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For  $d < a$

Surface charge density:



$$\sigma(\hat{\mathbf{r}}) = \epsilon_0 \left. \frac{\partial \Phi}{\partial r} \right|_{r=a} = \frac{q}{4\pi a^2} \frac{a}{d} \frac{(1 - \frac{a^2}{d^2})}{(1 + \frac{a^2}{d^2} - 2\frac{a}{d} \hat{\mathbf{r}} \cdot \hat{\mathbf{d}})^{3/2}}$$

Let  $\hat{\mathbf{r}} \cdot \hat{\mathbf{d}} \equiv \cos \gamma$

$$Q = 2\pi a^2 \int_{-1}^1 d \cos \gamma \left( \frac{q}{4\pi a^2} \frac{a}{d} \frac{(1 - \frac{a^2}{d^2})}{(1 + \frac{a^2}{d^2} - 2\frac{a}{d} \cos \gamma)^{3/2}} \right)$$

$$= -q$$

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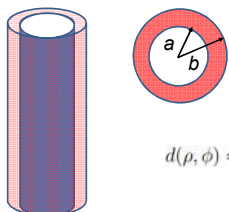
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HW#6

top view



$$d(\rho, \phi) = \begin{cases} 0 & \text{for } 0 \leq \rho < a \\ d_0 & \text{for } a \leq \rho \leq b \\ 0 & \text{for } \rho > b \end{cases}$$

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Green's function appropriate for this geometry with boundary conditions at  $\rho = 0$  and  $\rho = \infty$ :

$$\left( \frac{1}{\rho} \frac{\partial}{\partial \rho} \left( \rho \frac{\partial}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2}{\partial \phi^2} \right) G(\rho, \rho', \phi, \phi') = -4\pi \frac{\delta(\rho - \rho')}{\rho} \delta(\phi - \phi')$$

$$G(\rho, \rho', \phi, \phi') = -\ln(\rho_>^2) + 2 \sum_{m=1}^{\infty} \frac{1}{m} \left( \frac{\rho_<}{\rho_>} \right)^m \cos(m(\phi - \phi'))$$

$$\Phi(\rho, \phi) = \frac{1}{4\pi\epsilon_0} \int_0^{2\pi} d\phi' \int_0^{\infty} \rho' d\rho' G(\rho, \rho', \phi, \phi') d(\rho', \phi')$$

$$= -\frac{1}{\epsilon_0} \int_0^{\infty} \rho' d\rho' \ln(\rho_>) d(\rho')$$

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HW #7

PHY 712 – Assignment #7

February 2, 2018

Complete reading Chapter 3 and start Chapter 4 in Jackson.

1. Consider the charge density of an electron bound to a proton in a hydrogen atom –  $\rho(r) = (1/\pi a_0^3) e^{-2r/a_0}$ , where  $a_0$  denotes the Bohr radius. Find the electrostatic potential  $\Phi(r)$  associated with  $\rho(r)$ . Compare your result to HW#1.

Example for isolated charge density  $\rho(\mathbf{r})$  with electrostatic potential vanishing for  $r \rightarrow \infty$ :

$$\begin{aligned} \Phi(\mathbf{r}) &= \frac{1}{4\pi\epsilon_0} \int d^3r' \frac{\rho(\mathbf{r}')}{|\mathbf{r}-\mathbf{r}'|} \\ &= \frac{1}{4\pi\epsilon_0} \int d^3r' \rho(\mathbf{r}') \left( \sum_{lm} \frac{4\pi}{2l+1} \frac{r_{<}^l}{r_{>}^{l+1}} Y_{lm}(\theta, \varphi) Y_{lm}^*(\theta', \varphi') \right) \end{aligned}$$

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In our case,  $\rho(r) = \frac{1}{\pi a_0^3} e^{-2r/a_0}$

$$\Phi(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \int d^3r' \rho(\mathbf{r}') \left( \sum_{lm} \frac{4\pi}{2l+1} \frac{r_{<}^l}{r_{>}^{l+1}} Y_{lm}(\theta, \varphi) Y_{lm}^*(\theta', \varphi') \right)$$

Integral over solid angles

$$\int d\Omega Y_{lm}^*(\theta', \varphi') = \sqrt{4\pi} \delta_{l,0} \delta_{m,0}$$

Non-trivial radial integral

$$\begin{aligned} \Phi(\mathbf{r}) &= \frac{1}{\epsilon_0} \int_0^\infty r'^2 dr' \rho(r') \frac{1}{r_{>}} \\ &= \frac{1}{\epsilon_0} \left( \int_0^r r'^2 dr' \rho(r') + \int_r^\infty r' dr' \rho(r') \right) \end{aligned}$$

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