

PHY 712 Electrodynamics

12-12:50 AM MWF Olin 103

Class notes for Lecture 2:

Reading: Chapter 1 (especially 1.11) in JDJ;

Ewald summation methods

- 1. Motivation**
- 2. Expression to evaluate the electrostatic energy of an extended periodic system**
- 3. Examples**

Comment on HW #1

For Prob. 1.5, what if the potential function had the form:

$$\Phi(r) = \frac{q}{4\pi\epsilon_0} \left[\frac{e^{-\alpha r}}{r} (1 + \alpha r) - \frac{1}{r} \right]$$

Your questions –

From Nick –

- 1) When you say the results should not depend on η , is it that η doesn't matter or that there's something intrinsic about the energy that shouldn't matter what η is? Will you walk through the calculations to show why they are slightly different?
- 2) On the last slide, is this 2 triple integrals? Why do we have d^3r and d^3r' ?

From Gao –

How can the expression of electrostatic energy show its justification in the continuous case as it can not exclude the "self-interaction"?

From Tim –

In the extra notes section on page six you say "there are two kinds of sites...." What do you mean by sites?

PHY 712 Electrodynamics

MWF 10-10:50 PM Online <http://www.wfu.edu/~natalie/s21phy712/>

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Course schedule for Spring 2021

(Preliminary schedule -- subject to frequent adjustment.)

	Lecture date	JDJ Reading	Topic	HW	Due date
1	Wed: 01/27/2021	Chap. 1 & Appen.	Introduction, units and Poisson equation	#1	01/29/2021
2	Fri: 01/29/2021	Chap. 1	Electrostatic energy calculations	#2	02/01/2021
3	Mon: 02/01/2021	Chap. 1	Electrostatic potentials and fields		
4	Wed: 02/03/2021	Chap. 1 - 3	Poisson's equation in 2 and 3 dimensions		
5	Fri: 02/05/2021	Chap. 1 - 3	Brief introduction to numerical methods		
6	Mon: 02/08/2021	Chap. 2 & 3	Image charge constructions		
7	Wed: 02/10/2021	Chap. 2 & 3	Cylindrical and spherical geometries		

Ewald summation methods -- motivation

Consider a collection of point charges $\{q_i\}$ located at points $\{\mathbf{r}_i\}$.

The energy to separate these charges to infinity ($\mathbf{r}_i \rightarrow \infty$) is

$$W = \frac{1}{4\pi\epsilon_0} \sum_{(i,j;i>j)} \frac{q_i q_j}{|\mathbf{r}_i - \mathbf{r}_j|}.$$

Here the summation is over all pairs of (i, j) , excluding $i = j$.

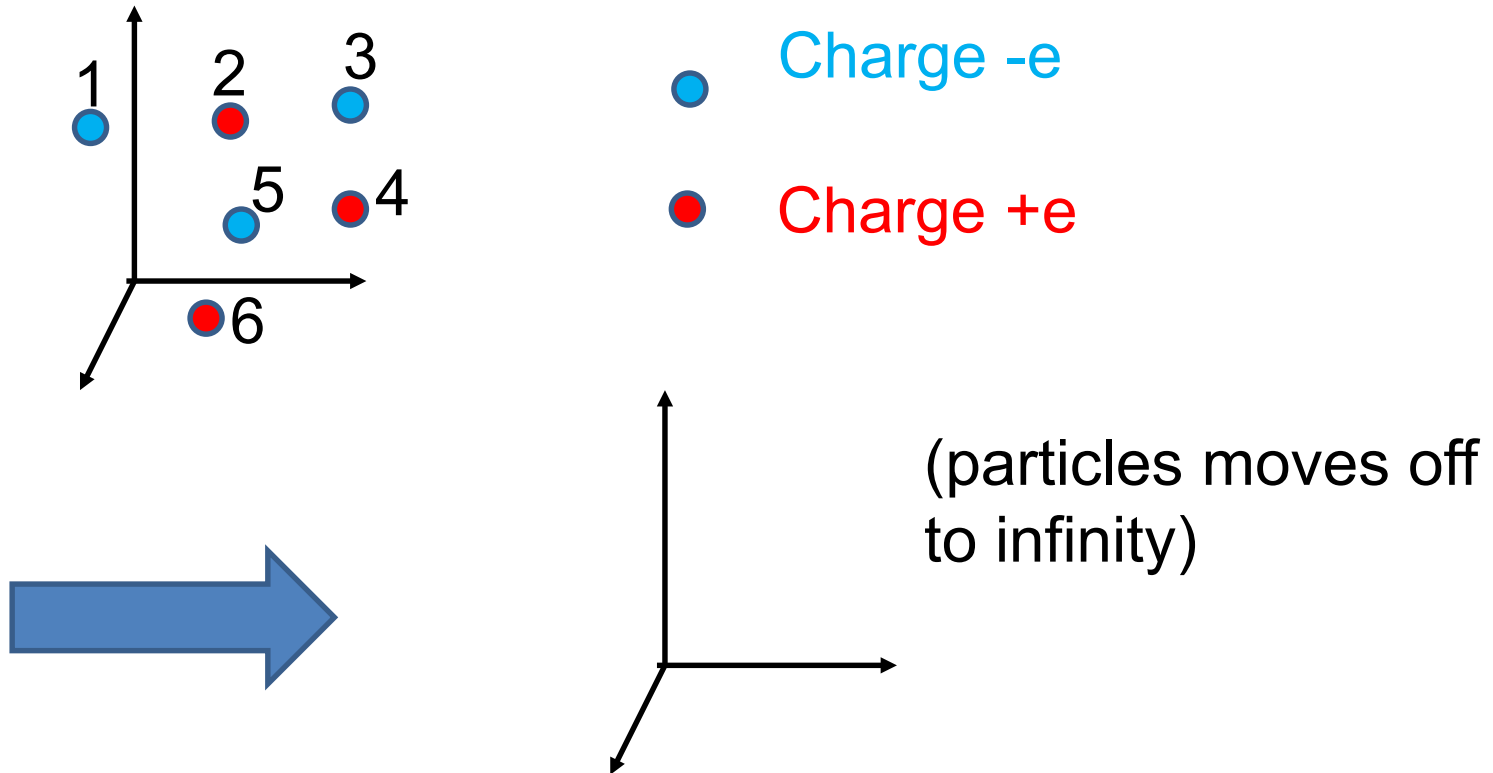
It is convenient to sum over all particles and divide by 2 in order to compensate for the double counting:

$$W = \frac{1}{8\pi\epsilon_0} \sum_{i,j;i \neq j} \frac{q_i q_j}{|\mathbf{r}_i - \mathbf{r}_j|}.$$

Now the summation is over all i and j , excluding $i = j$.

The energy W scales as the number of particles N . As $N \rightarrow \infty$, the ratio W / N remains well-defined in principle, but difficult to calculate in practice.

Example finite charge system for which electrostatic energy W can be calculated in a straightforward way



$$W = W_{12} + W_{13} + W_{14} + W_{15} + W_{16} + W_{23} + W_{24} + W_{25} + W_{26} \\ + W_{34} + W_{35} + W_{36} + W_{45} + W_{46} + W_{56}$$

Evaluation of the electrostatic energy for N point charges:

$$\frac{W}{N} = \frac{1}{8\pi\epsilon_0} \frac{1}{N} \sum_{i,j;i \neq j} \frac{q_i q_j}{|\mathbf{r}_i - \mathbf{r}_j|}.$$

Ewald summation methods – exact results for periodic systems

$$\frac{W}{N} = \sum_{\alpha\beta} \frac{q_\alpha q_\beta}{8\pi\epsilon_0} \left(\frac{4\pi}{\Omega} \sum_{\mathbf{G} \neq 0} \frac{e^{-i\mathbf{G} \cdot \boldsymbol{\tau}_{\alpha\beta}} e^{-G^2/\eta}}{G^2} - \sqrt{\frac{\eta}{\pi}} \delta_{\alpha\beta} + \sum_{\mathbf{T}} \frac{\text{erfc}(\frac{1}{2}\sqrt{\eta} |\boldsymbol{\tau}_{\alpha\beta} + \mathbf{T}|)}{|\boldsymbol{\tau}_{\alpha\beta} + \mathbf{T}|} \right) - \frac{4\pi Q^2}{8\pi\epsilon_0 \Omega \eta}.$$

Note that the results should not depend upon η (assuming that all summations are carried to convergence). In the example of CsCl having a lattice constant a , we show two calculations produce the result:

$$\frac{W}{N} = -\frac{e^2}{8\pi\epsilon_0} \frac{4.070722970}{a} \quad \text{or} \quad \frac{W}{N} = -\frac{e^2}{8\pi\epsilon_0} \frac{4.070723039}{a}$$

See lecture notes for details.

Your question – How is it that the Ewald summation does not depend upon η ?

Comment – This is because of the genius of Ewald (assuming that the summation terms are well converged).

Slight digression:

Comment on electrostatic energy evaluation --

When the discrete charge distribution becomes a continuous charge density: $q_i \rightarrow \rho(\mathbf{r})$, the electrostatic energy becomes

$$W = \frac{1}{8\pi\epsilon_0} \int d^3r \, d^3r' \frac{\rho(\mathbf{r})\rho(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|}.$$

Notice, in this case, it is not possible to exclude the "self-interaction".

Electrostatic energy in terms of $\Phi(\mathbf{r})$ and field $\mathbf{E}(\mathbf{r})$:

Previous expression can be rewritten in terms of the electrostatic potential or field:

$$W = \frac{1}{2} \int d^3r \, \rho(\mathbf{r})\Phi(\mathbf{r}) = -\frac{\epsilon_0}{2} \int d^3r \, (\nabla^2 \Phi(\mathbf{r}))\Phi(\mathbf{r}).$$

$$W = \frac{\epsilon_0}{2} \int d^3r \, |\nabla \Phi(\mathbf{r})|^2 = \frac{\epsilon_0}{2} \int d^3r \, |\mathbf{E}(\mathbf{r})|^2.$$

Your question – What happens to the electrostatic case for continuous charge densities with respect to the self energy contribution.

Comment – It sometimes causes trouble.

Now consider the electrostatic energy of a periodic crystal of CsCl

