PHY 712 Electrodynamics 11-11:50 AM MWF Olin 107

Class notes for Lecture 2:

Reading: Chapter 1 (especially 1.11) in JDJ;

- 1. Calculation of the electrostatic energy
- 2. Evaluation of the electrostatic energy of an extended periodic system using Ewald summation methods

1/12/2022

PHY 712 Spring 2022 -- Lecture 2

1

This lecture details some special properties of the electrostatic energy of an extended system. It illustrates some special properties of the long range nature of the Coulomb interaction. Ewald summation methods may or may not be important for your particular field of study. However, it is at least important to be aware of the ideas.

PHY 712 Electrodynamics [MWF 11-11:50 AM Olin 107 Webpage: http://www.wfu.edu/~natalie/s22phy712/ [Instructor: Natalie Holzwarth Office:300 OPL e-mail:natalie@wfu.edu Course schedule for Spring 2022 (Preliminary schedule subject to frequent adjustment.)					
	Lecture date	JDJ Reading Topic		HW	Due date
1	Mon: 01/10/2022	Chap. 1 & Appen.	Introduction, units and Poisson equation	#1	01/14/2022
2	Wed: 01/12/2022	Chap. 1	Electrostatic energy calculations	#2	01/19/2022
3	Fri: 01/14/2022	Chap. 1	Electrostatic energy calculations	#3	01/21/2022
	Mon: 01/17/2022		MLK Holiday no class		
1/12/2022 PHY 712 Spring 2022 Lecture 2					2

The homework problem assigned this time exercises the ideas presented in this lecture.

PHY 712 -- Assignment #2

January 12, 2022

Continue reading Chap. 1 in Jackson.

Calculate the electrostatic energy of the following 5 ion molecule scaled by the factor (1/($4\pi\epsilon_0$)) (q^2 / a). Note that x, y, and zdenote unit vectors in the three Cartesian directions.

- 1. Charge = 4q Position = 0 2. Charge = -q Position = (a/2)(**x**+**y**+**z**) 3. Charge = -q Position = (a/2)(**x**-**y**+**z**) 4. Charge = -q Position = (a/2)(**x**-**y**-**z**) 5. Charge = -q Position = (a/2)(-**x**+**y**-**z**)

1/12/2022

PHY 712 Spring 2022 -- Lecture 2

Comment on HW #1

For Prob. 1.5, what if the potential function had the form:

$$\Phi(r) = \frac{q}{4\pi\epsilon_0} \left[\frac{e^{-2r/a_0}}{r} \left(1 + \frac{r}{a_0} \right) - \frac{1}{r} \right] ?$$



Contribution from only the electron in its ground state

1/12/2022

PHY 712 Spring 2022 -- Lecture 2

Calculation of the electrostatic energy of a system of charges --Consider a collection of point charges $\{q_i\}$ located at points $\{\mathbf{r}_i\}$. The energy to separate these charges to infinity $(\mathbf{r}_i \to \infty)$ is

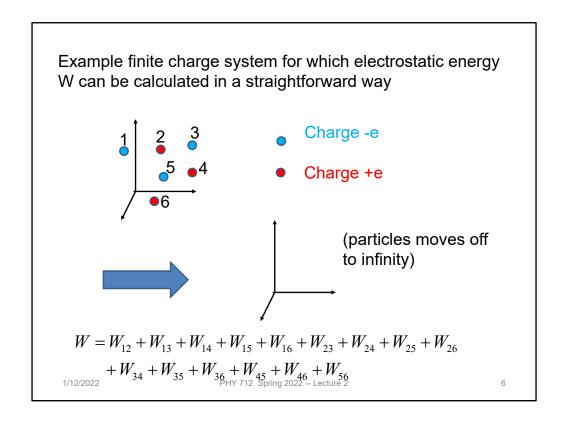
$$W = \frac{1}{4\pi\epsilon_0} \sum_{(i,j;i>j)} \frac{q_i q_j}{|\mathbf{r}_i - \mathbf{r}_j|}.$$

Here the summation is over all pairs of (i, j), excluding i = j. It is convenient to sum over all particles and divide by 2 in order to compensate for the double counting:

$$W = \frac{1}{8\pi\epsilon_0} \sum_{i,j;i\neq j} \frac{q_i q_j}{|\mathbf{r}_i - \mathbf{r}_j|}.$$

Now the summation is over all i and j, excluding i = j. The energy W scales as the number of particles N. As $N \to \infty$, the ratio W / N remains well-defined in principle, but difficult to calculate in practice.

PHY 712 Spring 2022 - Lecture 2



Summation for a finite number of particles is straightforward, but extension to an infinite system runs into difficulty.

Homework problem #2

Calculate the electrostatic energy of the following 5 ion molecule scaled by the factor $(1/(4\pi\epsilon_0))$ (q^2/a) . Note that x, y, and z denote unit vectors in the three Cartesian directions.

- 1. Charge = 4q Position = 0 2. Charge = -q Position = (a/2)(x+y+z) 3. Charge = -q Position = (a/2)(-x-y+z)
- 4. Charge = -q Position = $(a/2)(x-y-z)^2$
- 5. Charge = -q Position = (a/2)(-x+y-z)



1/12/2022

PHY 712 Spring 2022 -- Lecture 2

Now consider the case of an infinite periodic system -- Consider a collection of point charges $\{q_i\}$ located at points $\{\mathbf{r}_i\}$. The energy to separate these charges to infinity $(\mathbf{r}_i \to \infty)$ is

$$W = \frac{1}{4\pi\epsilon_0} \sum_{(i,j;i>j)} \frac{q_i q_j}{|\mathbf{r}_i - \mathbf{r}_j|}.$$

Here the summation is over all pairs of (i, j), excluding i = j. It is convenient to sum over all particles and divide by 2 in order to compensate for the double counting:

$$W = \frac{1}{8\pi\epsilon_0} \sum_{i,j;i\neq j} \frac{q_i q_j}{|\mathbf{r}_i - \mathbf{r}_j|}.$$

Now the summation is over all i and j, excluding i = j. The energy W scales as the number of particles N. As $N \to \infty$, the ratio W / N remains well-defined in principle, but difficult to calculate in practice.

Evaluation of the electrostatic energy for *N* point charges:

$$\frac{W}{N} = \frac{1}{8\pi\epsilon_0} \frac{1}{N} \sum_{i,j;i\neq j} \frac{q_i q_j}{|\mathbf{r}_i - \mathbf{r}_j|}.$$

Ewald summation methods – exact results for periodic systems
$$\frac{W}{N} = \sum_{\alpha\beta} \frac{q_{\alpha}q_{\beta}}{8\pi\varepsilon_{0}} \left(\frac{4\pi}{\Omega} \sum_{\mathbf{G}\neq\mathbf{0}} \frac{e^{-i\mathbf{G}\cdot\boldsymbol{\tau}_{\alpha\beta}} \ e^{-G^{2}/\eta}}{G^{2}} - \sqrt{\frac{\eta}{\pi}} \delta_{\alpha\beta} + \sum_{\mathbf{T}} \frac{\operatorname{erfc}(\frac{1}{2}\sqrt{\eta} \mid \boldsymbol{\tau}_{\alpha\beta} + \mathbf{T} \mid)}{\mid \boldsymbol{\tau}_{\alpha\beta} + \mathbf{T} \mid} \right) - \frac{4\pi Q^{2}}{8\pi\varepsilon_{0}\Omega\eta}.$$

Note that the results should not depend upon η (assuming that all summations are carried to convergence). In the example of CsCl having a lattice constant a, we show two calculations produce the result:

$$\frac{W}{N} = -\frac{e^2}{8\pi\epsilon_0} \frac{4.070722970}{a}$$
 or $\frac{W}{N} = -\frac{e^2}{8\pi\epsilon_0} \frac{4.070723039}{a}$

Details to follow --

1/12/2022

PHY 712 Spring 2022 -- Lecture 2

9

In the Extra lecture notes this expression will be derived and illustrated.

Slight digression:

Comment on electrostatic energy evaluation --

When the discrete charge distribution becomes a continuous charge density: $q_i \rightarrow \rho(\mathbf{r})$, the electrostatic energy becomes

 $W = \frac{1}{8\pi\epsilon_0} \int d^3r \ d^3r' \frac{\rho(\mathbf{r})\rho(\mathbf{r'})}{|\mathbf{r} - \mathbf{r'}|}.$

Notice, in this case, it is not possible to exclude the ``self-interaction".

Electrostatic energy in terms of $\Phi(\mathbf{r})$ and field $\mathbf{E}(\mathbf{r})$:

Previous expression can be rewritten in terms of the electrostatic potential or field:

potential or field:

$$W = \frac{1}{2} \int d^3 r \ \rho(\mathbf{r}) \Phi(\mathbf{r}) = -\frac{\epsilon_0}{2} \int d^3 r \left(\nabla^2 \Phi(\mathbf{r}) \right) \Phi(\mathbf{r}).$$

$$W = \frac{\epsilon_0}{2} \int d^3r \left| \nabla \Phi(\mathbf{r}) \right|^2 = \frac{\epsilon_0}{2} \int d^3r \left| \mathbf{E}(\mathbf{r}) \right|^2.$$
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We will see and use these expressions throughout the course. However, the so-called self energy often leads to difficulties....

Some details --

Electrostatic potential

$$\Phi(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \int d^3r' \frac{\rho(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|}$$

Electrostatic field

$$\mathbf{E}(\mathbf{r}) = -\nabla \Phi(\mathbf{r})$$

Poisson equation

$$\nabla^2 \Phi(\mathbf{r}) = -\frac{\rho(\mathbf{r})}{\varepsilon_0}$$

1/12/2022

PHY 712 Spring 2022 -- Lecture 2

Summary --Electrostatic energy

$$W = \frac{1}{8\pi\epsilon_0} \int d^3r \ d^3r' \frac{\rho(\mathbf{r})\rho(\mathbf{r'})}{|\mathbf{r} - \mathbf{r'}|}.$$

Evaluation of electrostatic energy in terms of potential $\Phi(\mathbf{r})$ and field $\mathbf{E}(\mathbf{r})$:

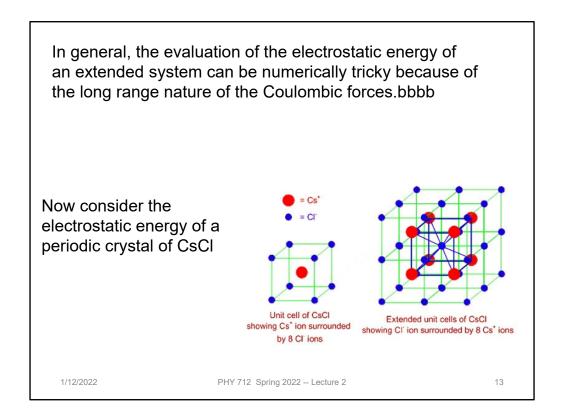
$$W = \frac{1}{2} \int d^3 r \, \rho(\mathbf{r}) \Phi(\mathbf{r}) = -\frac{\epsilon_0}{2} \int d^3 r \left(\nabla^2 \Phi(\mathbf{r}) \right) \Phi(\mathbf{r}).$$

$$W = \frac{\epsilon_0}{2} \int d^3 r \left| \nabla \Phi(\mathbf{r}) \right|^2 = \frac{\epsilon_0}{2} \int d^3 r \left| \mathbf{E}(\mathbf{r}) \right|^2.$$

$$W = \frac{\epsilon_0}{2} \int d^3r \left| \nabla \Phi(\mathbf{r}) \right|^2 = \frac{\epsilon_0}{2} \int d^3r \left| \mathbf{E}(\mathbf{r}) \right|^2.$$

1/12/2022

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In general, the evaluation of the electrostatic energy of an extended system can be numerically tricky because of the long range nature of the Coulombic forces.

However, thanks to very clever mathematicians, it is possible to perform this sort of calculation for periodic systems.

Ewald, Paul Peter, 1888-1985

American crystallographer, emigrated from Germany



1/12/2022

PHY 712 Spring 2022 -- Lecture 2

