

$\text{Re}(\lambda) = 0$   
(center)

### Notation:

1.  $A$  is called hyperbolic if  $\lambda_1, \lambda_2 \neq 0$  and  $\lambda_1, \lambda_2 \in \mathbb{R}$ .
2. We say that the fixed point  $x=0$  of  $\dot{x} = Ax$  is
  - a.) an attractor if  $\vec{x}(t) \rightarrow 0$  as  $t \rightarrow \infty \Rightarrow \text{Re}(\lambda_1, \lambda_2) < 0$
  - b.) a repeller if  $\vec{x}(t) \rightarrow 0$  as  $t \rightarrow -\infty \Rightarrow \text{Re}(\lambda_1, \lambda_2) > 0$
  - c.) a saddle if  $\lambda_1 < 0 < \lambda_2$
  - d.) non hyperbolic if  $\text{Re}(\lambda_1) = 0$  or  $\text{Re}(\lambda_2) = 0$ .

### Example:

Sketch the phase portrait of

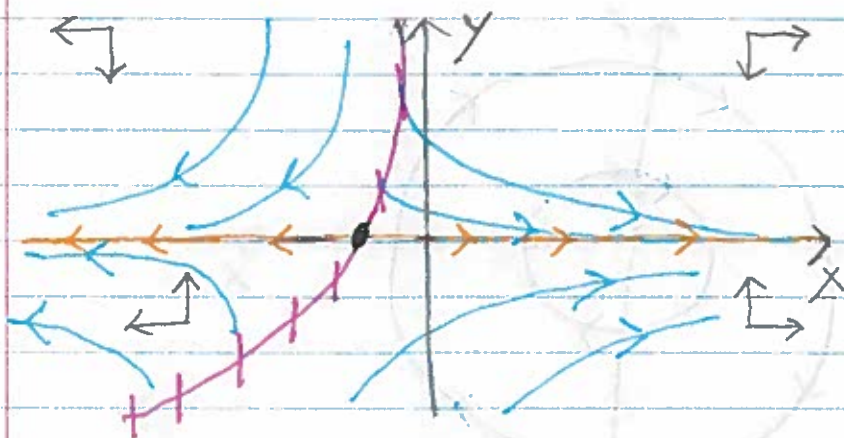
$$\begin{cases} \dot{x} = x + e^y = f(x, y) \\ \dot{y} = -y = g(x, y) \end{cases}$$

Fixed points:  $y=0$  and  $x=-1$ .

### Nullclines:

a.)  $-x = e^y \Rightarrow y = -\ln(-x)$  N1:  $\frac{dx}{dt} = 0$

b.)  $y = 0$  N2:  $\frac{dy}{dt} = 0$



We can also do local analysis;  
linearization about  $(x^*, y^*)$ :

$$f(x, y) = f(x^*) + \nabla f|_{x^*} (\bar{x} - x^*) + \frac{1}{2} (\bar{x} - x^*)^T H(\nabla^2 f)|_{x^*} (\bar{x} - x^*)$$

$$g(x, y) = g(x^*) + \nabla g|_{x^*} (\bar{x} - x^*) + \frac{1}{2} (\bar{x} - x^*)^T H(\nabla^2 g)|_{x^*} (\bar{x} - x^*)$$

Near equilibrium

$$F(\bar{x} - x^*) \approx J(F)|_{x^*} (\bar{x} - x^*)$$

$$J(F) = \begin{pmatrix} f_x & f_y \\ g_x & g_y \end{pmatrix}$$

Changing variables  $\tilde{x} = x - x^*$ ,  $\tilde{y} = y - y^*$

$$F(\tilde{x}) \approx J(F)|_{x^*} \tilde{x}$$

Locally, the solution is

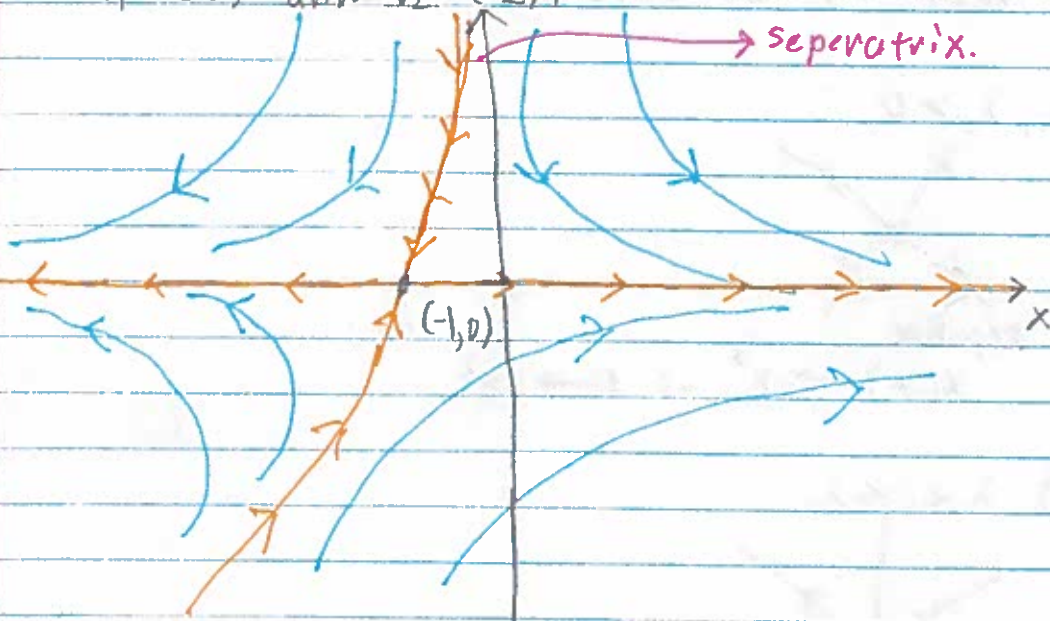
$$\tilde{x} = c_1 e^{\lambda_1 t} \tilde{v}_1 + c_2 e^{\lambda_2 t} \tilde{v}_2$$

where  $\lambda_1, \lambda_2$  are the eigenvalues of  $J|_{x^*}$ .

For our problem:

$$J(F)|_{(-1,0)} = \begin{pmatrix} 1 & -e^{-y} \\ 0 & -1 \end{pmatrix} \Big|_{(-1,0)} = \begin{pmatrix} 1 & -1 \\ 0 & -1 \end{pmatrix}$$

The eigenvalues are  $\lambda_1 = 1$ ,  $\lambda_2 = -1$  with eigenvectors  $\vec{v}_1 = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$  and  $\vec{v}_2 = \begin{pmatrix} 1 \\ 2 \end{pmatrix}$ .



### Phase Plane Summary:

1. Draw Nullclines

2. Draw direction arrows  $\curvearrowright$

3. Find equilibrium

4. Calculate Jacobian

5. Determine eigenvalues and eigenvectors,

$$\lambda_1, \lambda_2 \in \mathbb{R}$$

$\Rightarrow$  hyperbolic points

$$\lambda_1, \lambda_2 \in \mathbb{C} \setminus \mathbb{R}$$

$\Rightarrow$  spirals

\* Determine fast direction.

\* Determine stable or unstable

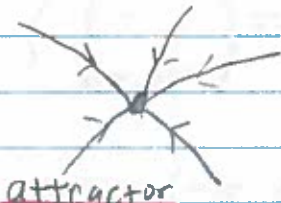
$\lambda_1, \lambda_2$  pure Imaginary

$\Rightarrow$  Need more work.

$\lambda_1$  or  $\lambda_2 = 0$

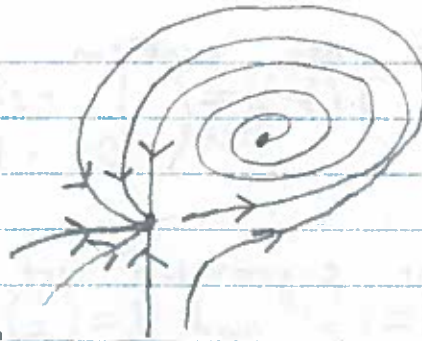
$\Rightarrow$  Need more work

a.)  $\lambda_1, \lambda_2 < 0$



attractor

$X(t) \rightarrow X^*$  as  $t \rightarrow \infty$ .



b.)  $\lambda_1, \lambda_2 > 0$



repeller

$X(t) \rightarrow X^*$  as  $t \rightarrow -\infty$

c.)  $\lambda_1 < 0 < \lambda_2$



Saddle

Example

Predator-Prey

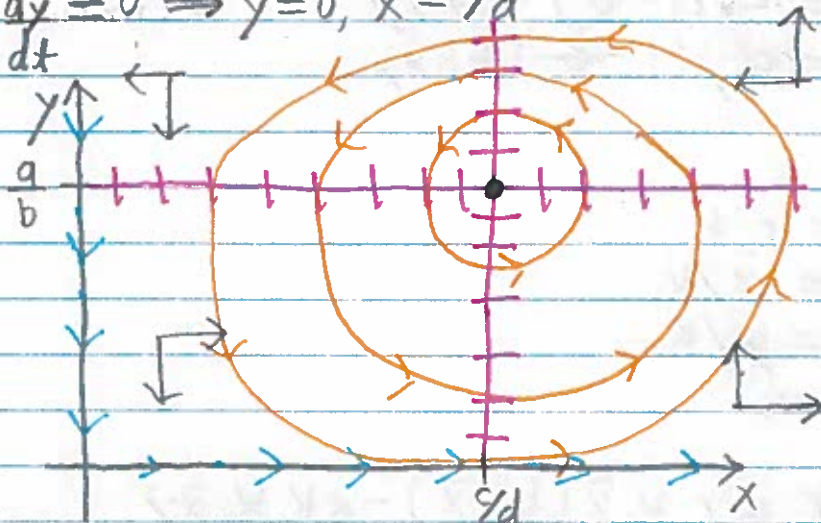
$$\dot{x} = ax - bxy$$

$$\dot{y} = -cy + dxy$$

Null Clines:

$$\frac{dx}{dt} = 0 \Rightarrow x=0, y = a/b$$

$$\frac{dy}{dt} = 0 \Rightarrow y=0, x = c/d$$



Eigenvalues analysis:

$$J = \begin{pmatrix} a-by & -bx \\ dy & -c+dx \end{pmatrix}$$

$$J(0,0) = \begin{pmatrix} a & 0 \\ 0 & -c \end{pmatrix} \Rightarrow \lambda_1 = a, \vec{v}_1 = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \\ \lambda_2 = -c, \vec{v}_2 = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

$$J(c/d, a/b) = \begin{pmatrix} 0 & -bc/d \\ dc/b & 0 \end{pmatrix}$$

The eigenvalues are pure imaginary  $\Rightarrow$  cannot deduce anything. However,

$$\frac{dy}{dx} = \frac{y \cdot (dx-c)}{x \cdot (a-by)}$$

$$\Rightarrow \int \left( \frac{a}{y} - b \right) dy = \int \left( \frac{-c}{x} + d \right) dx$$

$$\Rightarrow a \ln(y) - by = -c \ln(x) + dx + C$$

$$a \ln(y) - by + c \ln(x) - dx = C \\ y^d e^{-by} x^c e^{-dx} = C$$

This describes a family of closed curves  $\Rightarrow$  the solution curves.

Example:

Logistic growth omnivores and prey.

$$\begin{aligned}\dot{x} &= r_1 x \left(1 - \frac{x}{K_1}\right) - \alpha xy && \text{(Invasive Predators)} \\ \dot{y} &= r_2 y \left(1 - \frac{y}{K_2}\right) + \beta xy\end{aligned}$$

Rescale

$$\tau = r_1 t$$

$$\bar{x} = x/K_1$$

$$\bar{y} = y/K_2$$

Therefore,

$$r_1 K_1 \frac{d\bar{x}}{d\tau} = r_1 K_1 \bar{x} (1 - \bar{x}) - \alpha K_1 K_2 \bar{x} \bar{y}$$

$$r_1 K_2 \frac{d\bar{y}}{d\tau} = r_2 K_2 \bar{y} (1 - \bar{y}) + \beta K_1 K_2 \bar{x} \bar{y}$$

$$\Rightarrow \frac{d\bar{x}}{d\tau} = \bar{x} (1 - \bar{x}) - \gamma \bar{x} \bar{y}$$

$$\frac{d\bar{y}}{d\tau} = \eta \bar{y} (1 - \bar{y}) + \delta \bar{x} \bar{y}$$

Equilibrium points:

$$0 = \bar{x} (1 - \bar{x}) - \gamma \bar{x} \bar{y}$$

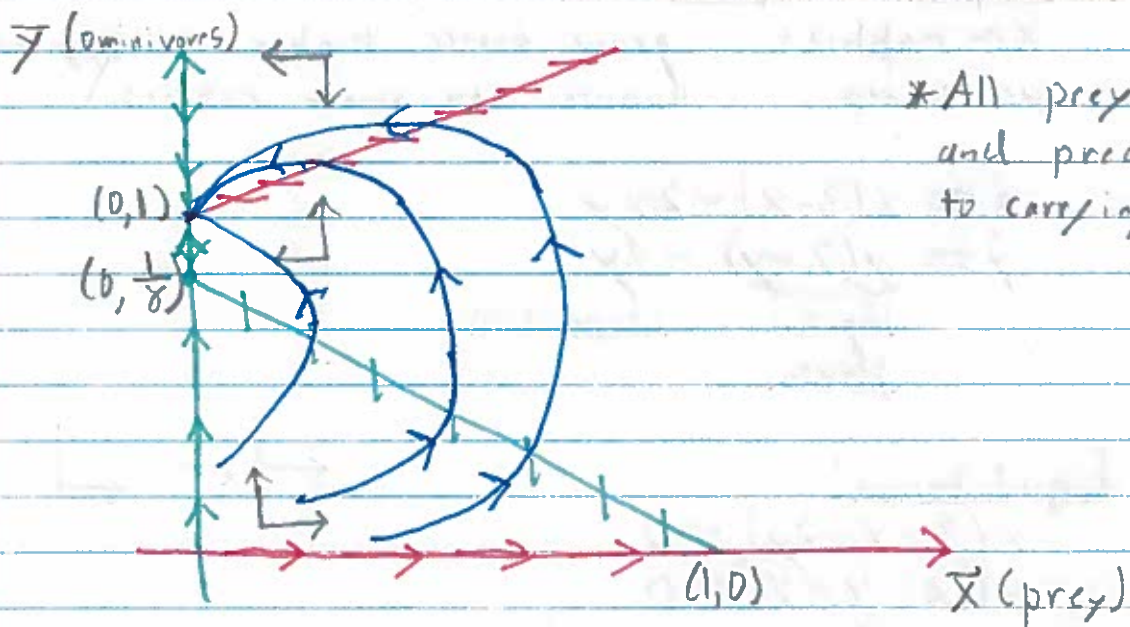
$$0 = \eta \bar{y} (1 - \bar{y}) + \delta \bar{x} \bar{y}$$

$\Rightarrow (\bar{x}, \bar{y}) = (0, 0)$  is one equilibrium. There are others.

Null clines:

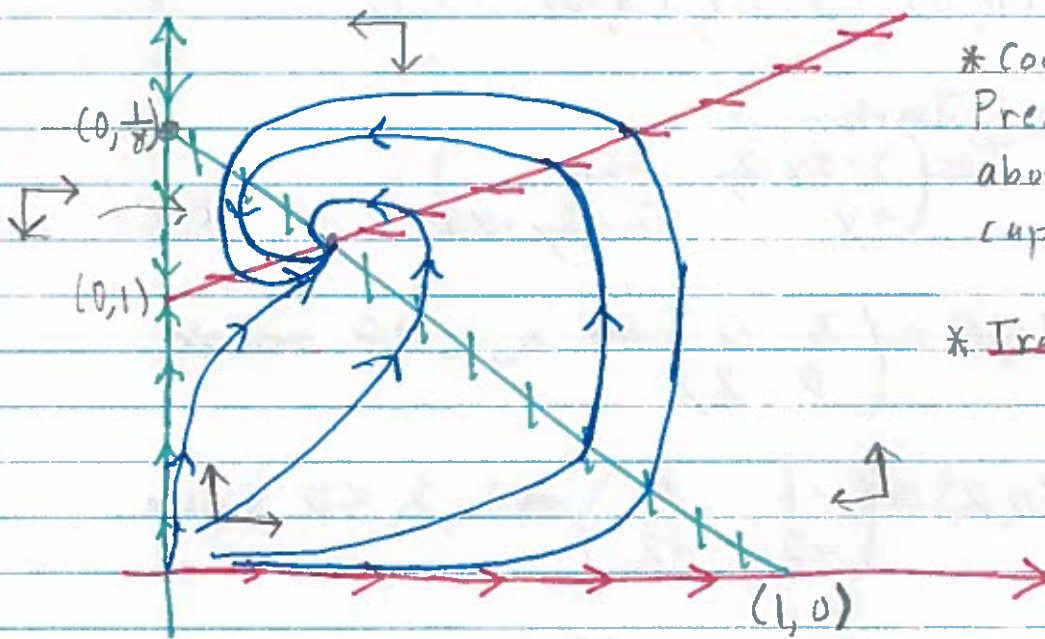
N1:  $\bar{y} = \frac{1 - \bar{x}}{\gamma}$ ,  $\bar{x} = 0$  :  $\frac{d\bar{x}}{d\tau} = 0$ .

N2:  $\bar{x} = \frac{\eta}{\delta} (\bar{y} - 1)$ ,  $\bar{y} = 0$  :  $\frac{d\bar{y}}{d\tau} = 0$



\* All prey die off and predators go to carrying capacity.

Omnivores take over



\* Coexistence. Predators live above carrying capacity.

\* Transcritical bifurcation

Point of intersection:

$$\bar{x}^* = \frac{(1-\gamma)y}{\delta\delta + \gamma} \quad \bar{y}^* = \frac{\delta + \gamma}{\delta\delta + \gamma}$$

The Jacobian is a mess but it is complex valued at the critical point.

⇒ stable spiral

### Example: (Competition)

$x \sim$  rabbits, grow faster higher carrying capacity  
 $y \sim$  sheep, hooves can smash rabbits.

$$\dot{x} = x(3-x) - 2xy$$

$$\dot{y} = \underbrace{y(2-y)}_{\text{logistic growth}} - xy_{\text{Competition}}$$

### Equilibrium:

$$x(3-x-2y) = 0$$

$$y(2-y-x) = 0$$

We get four equilibrium:

$$(0, 0), (0, 2), (3, 0), (1, 1)$$

The Jacobian is

$$J = \begin{pmatrix} 3-2x-2y & -2x \\ -y & 2-2y-x \end{pmatrix}$$

$$J(0, 0) = \begin{pmatrix} 3 & 0 \\ 0 & 2 \end{pmatrix} \Rightarrow \lambda_1, \lambda_2 > 0 \text{ unstable}$$

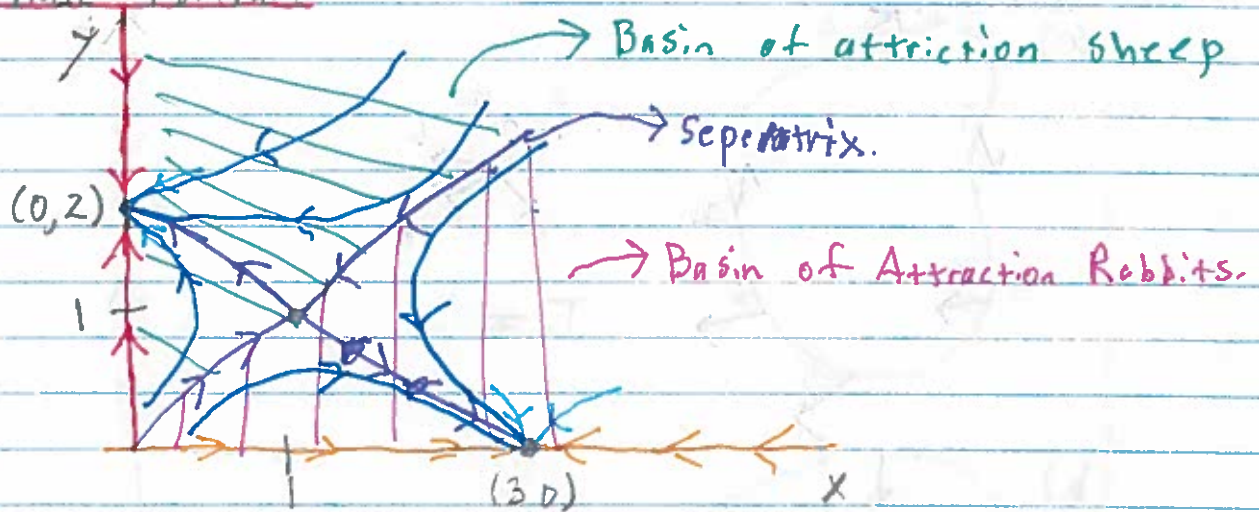
$$J(0, 2) = \begin{pmatrix} -1 & 0 \\ -2 & -2 \end{pmatrix} \Rightarrow \lambda_1, \lambda_2 < 0 \text{ stable}$$

$$J(3, 0) = \begin{pmatrix} -3 & -6 \\ 0 & -1 \end{pmatrix} \Rightarrow \lambda_1, \lambda_2 < 0 \text{ stable}$$

$$J(1, 1) = \begin{pmatrix} -1 & -2 \\ -1 & -1 \end{pmatrix} \Rightarrow \lambda_1, \lambda_2 = -1 \pm \sqrt{2} \text{ saddle}$$



## Phase Portrait



## Index Theory

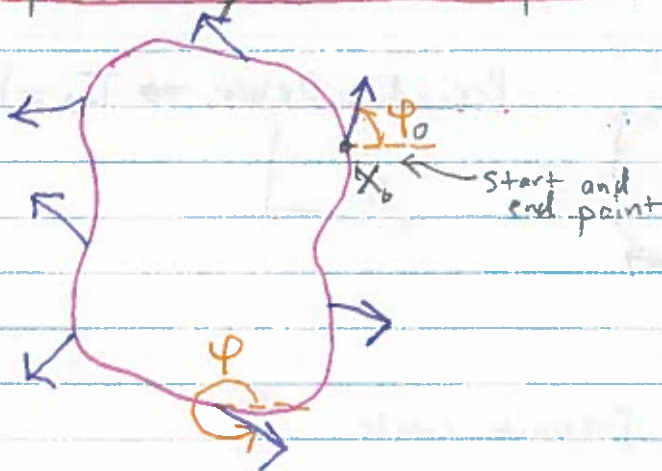
How can we be sure no periodic orbits exist?

Consider

$$\dot{\vec{x}} = F(\vec{x})$$

with  $F: \mathbb{R}^2 \rightarrow \mathbb{R}^2$  continuously differentiable.

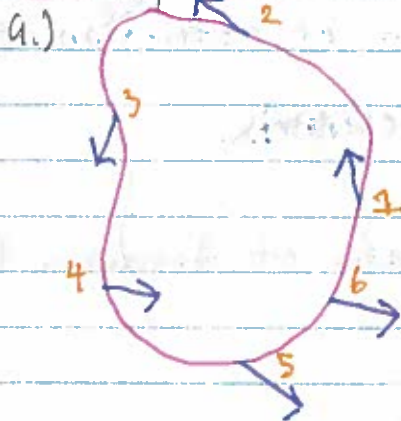
Take a closed curve  $\Gamma$  with no self intersections, that does not pass through a fixed point



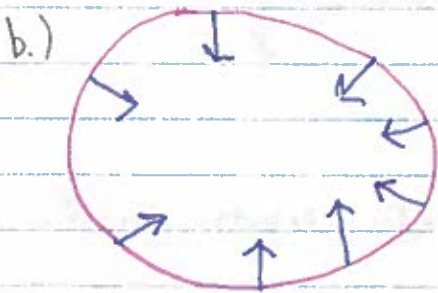
1. Start at  $x_0$ , traverse  $\Gamma$  counterclockwise and take angle  $\varphi$  of  $F(\vec{x}) \rightarrow$  this angle changes continuously as  $\Gamma$  is traversed
2. After one pass we again end up at  $x_0$  with an angle  $\varphi_1 = \varphi_0 + 2\pi n$ ;  $n \in \mathbb{Z}$

$$I_{\Gamma} = \frac{1}{2\pi} (\varphi_1 - \varphi_0)$$

# Examples



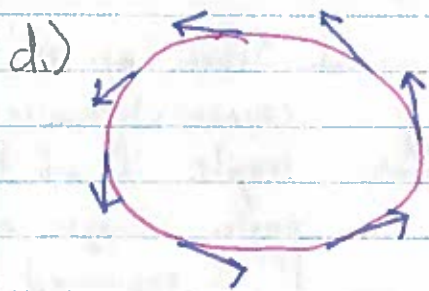
$$I_r = 1$$



$$I_r = 1$$

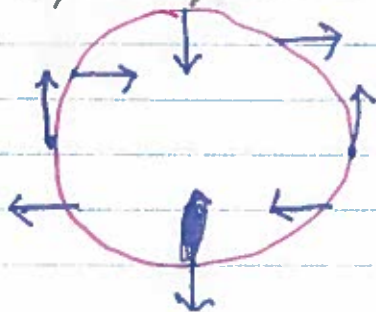


$$I_r = -1$$



Periodic Orbit  $\Rightarrow I_r = 1$

e.)  $\begin{cases} \dot{x} = x^2 y^2 \\ \dot{y} = x - y^2 \end{cases}$ ,  $\Gamma = \text{unit circle}$



$$I_r = 0$$

## Properties of the Index

1. If  $\Gamma$  can be deformed continuously into  $\tilde{\Gamma}$  without passing through any equilibrium points then
- $$I_{\Gamma} = I_{\tilde{\Gamma}}$$

proof:

$I_{\Gamma}$  varies continuously as  $\Gamma$  is deformed, but  $I_{\Gamma}$  is integer valued.

2. If  $\Gamma$  does not contain any fixed points then  $I_{\Gamma} = 0$

proof:

Property 1 implies we can shrink  $\Gamma$  to a point without changing the index.

3. If we replace  $F(\vec{x})$  by  $F(+\vec{x})$  the index is not changed.

proof:

Each angle is replaced by  $\varphi + \pi$ , hence  $\varphi_1 - \varphi_0$  is the same.

4. The index of a periodic orbit is one.

5. If  $F(\vec{x})$  is deformed continuously without creating any fixed points on  $\Gamma$ ,  $I_{\Gamma}$  stays the same.

Theorem - Assume  $F$  is continuously differentiable Inside each periodic orbit, there is at least one equilibrium.

proof:

Follows from items 2 and 4.

Index of isolated fixed point - Let  $\vec{x}^*$  be an isolated fixed point of  $\vec{x} = F(\vec{x})$ . Define

$I(\vec{x}^*) =$  index of simple closed curve that encloses  $x^*$  and no other fixed points

$I(\vec{x}^*)$  is well defined by property 4.

### Consequences

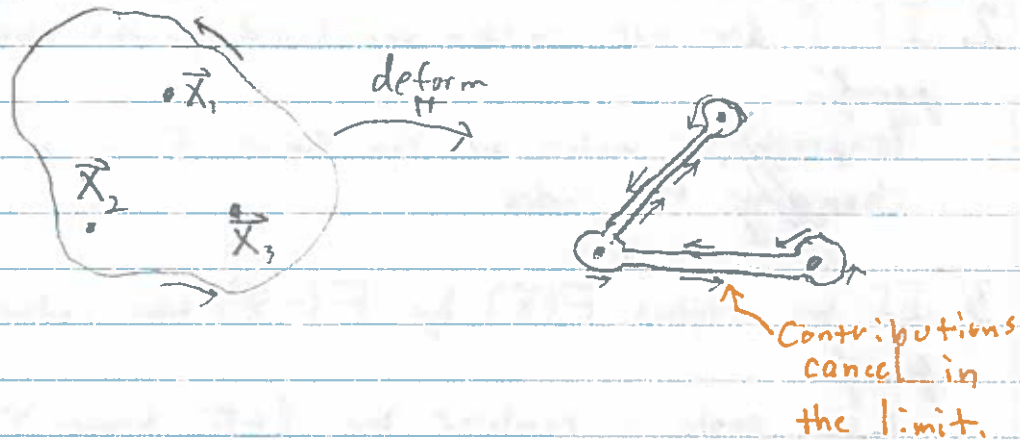
1. If  $\bar{x}^*$  is an attractor or repeller then  $I(\bar{x}^*) = 1$
2. If  $\bar{x}^*$  is a saddle point then  $I(\bar{x}^*) = -1$ .

proof:

Follows from examples b and c and properties 1, 3, 5.

Theorem - If  $\Gamma$  is a closed simple curve that contains  $n$  isolated fixed points  $\bar{x}_1, \dots, \bar{x}_n$  then  $I_\Gamma = I(\bar{x}_1) + \dots + I(\bar{x}_n)$

proof:



Corollary: A periodic orbit must enclose fixed points whose indices sum to +1.

### Omnivore example

The index of all the fixed points is  $-2$ ,

### Sheep and Rabbits

The index of all the fixed points is  $0$ .

Example:

$$\dot{x} = y$$

$$\dot{y} = -x + (1 - x^2 - y^2)y$$

Null-clines:

N1:  $\dot{x} = 0, y = 0$

N2:  $\dot{y} = 0, y(1 - x^2 - y^2) = 0$

$$\begin{aligned} & y - yx^2 - y^3 = 0 \\ & y - y^3 = x^2 + y^2 \\ & \Rightarrow \frac{y - y^3}{1 + y} = x^2 + y^2 \rightarrow \text{Not very informative.} \end{aligned}$$

Idea:

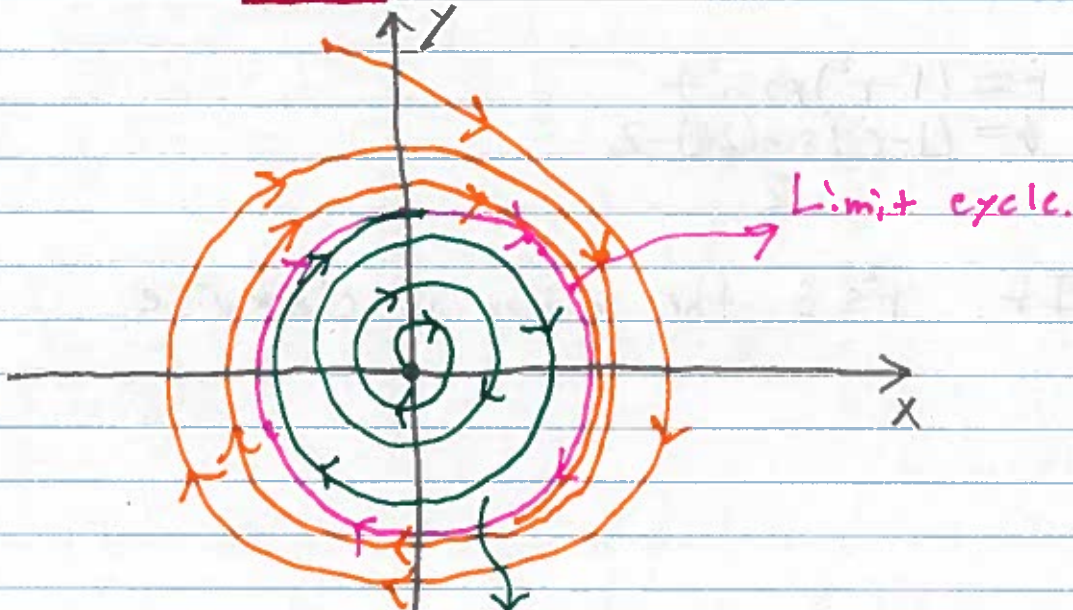
$$\dot{x} + x = (1 - x^2 - y^2)y$$

$$\Rightarrow \dot{x}\ddot{x} + x\dot{x} = (1 - x^2 - y^2)\dot{x} \cdot y$$

$$\Rightarrow \frac{1}{2} \frac{d}{dt} (\dot{x}^2 + x^2) = (1 - x^2 - y^2)\dot{x}^2$$

$$\Rightarrow \frac{1}{2} \frac{d}{dt} (y^2 + x^2) = (1 - x^2 - y^2)y^2$$

radius



radius

growing since  $x^2 + y^2 < 1$ .

Convert to polar coordinates:

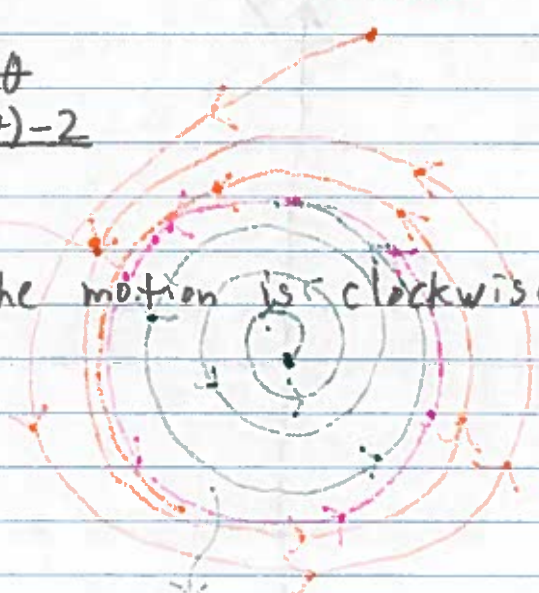
$$\begin{aligned}
 r^2 &= x^2 + y^2 & x &= r \cos \theta \\
 2r\dot{r} &= 2x\dot{x} + 2y\dot{y} & y &= r \sin \theta \\
 \Rightarrow \dot{r} &= (x\dot{x} + y\dot{y})/r \\
 &= \frac{x\dot{y} - y\dot{x}}{r} \\
 &= (1-r^2) \cdot r \sin^2 \theta
 \end{aligned}$$

$$\begin{aligned}
 \dot{\theta} &= \tan^{-1}\left(\frac{y}{x}\right) \\
 \dot{\theta} &= \frac{1}{1 + \frac{y^2}{x^2}} \left( \frac{x\dot{y} - y\dot{x}}{x^2} \right) \\
 &= \frac{x\dot{y} - y\dot{x}}{r^2} \\
 &= \frac{-x^2 + (1-r^2)xy - y^2}{r^2} \\
 &= -1 + (1-r^2)\sin\theta\cos\theta \\
 &= -1 + \frac{(1-r^2)\sin(2\theta)}{2} \\
 &= \frac{(1-r^2)\sin 2\theta - 2}{2}
 \end{aligned}$$

$$\begin{aligned}
 \dot{r} &= (1-r^2)r \sin^2 \theta \\
 \dot{\theta} &= \frac{(1-r^2)\sin(2\theta) - 2}{2}
 \end{aligned}$$

Since  $r \leq 3$

If  $r \leq 3$  the motion is clockwise.



Since  $r \leq 3$  the motion is clockwise.

## Conservative Systems

Inertial Systems of the form:

$$\ddot{x} = F(x)$$

A first integral can be found as follows:

$$\dot{x} \ddot{x} = \dot{x} F(x)$$

$$\Rightarrow \frac{1}{2} \frac{d(\dot{x}^2)}{dt} = \frac{dx}{dt} \left( \frac{-dV}{dx} \right), \quad (\text{For any solution curve})$$

where  $V(x) = -\int_x F(x) dx$  ( $x$  can be chosen

$$\Rightarrow \frac{d}{dt} \left( \frac{\dot{x}(t)^2}{2} + V \right) = 0 \quad (\text{arbitrarily})$$

For any solution curve there is a constant E such that

$$\boxed{\frac{\dot{x}(t)^2}{2} + V(x(t)) = E}$$

We can also write as a system:

$$\dot{x} = v$$

$$\dot{v} = F(x)$$

$$\frac{v^2}{2} + V(x) = E$$

phase portrait



contour plot

Theorem - A conservative system cannot have any attractor or repellers

proof:

Suppose there exists  $(x^*, v^*)$  that is an attracting point with a basin of attraction  $A$ . Then, for all  $(x_1, v_1), (x_2, v_2) \in A$  it follows that  $E(x_1, v_1) = E(x_2, v_2)$ . Since

$$E(x_1, v_1) = \lim_{t \rightarrow \infty} E(x_1(t), v_1(t))$$

$$= E(x^*, v^*(t))$$

$$= \lim_{t \rightarrow \infty} E(x_2(t), v_2(t))$$

$$= E(x_2, v_2)$$

Therefore,  $E$  must be constant in entire basin of attraction which we preclude by definition.  $\blacksquare$

### Example.

$$\ddot{x} + \sin(x) = 0$$

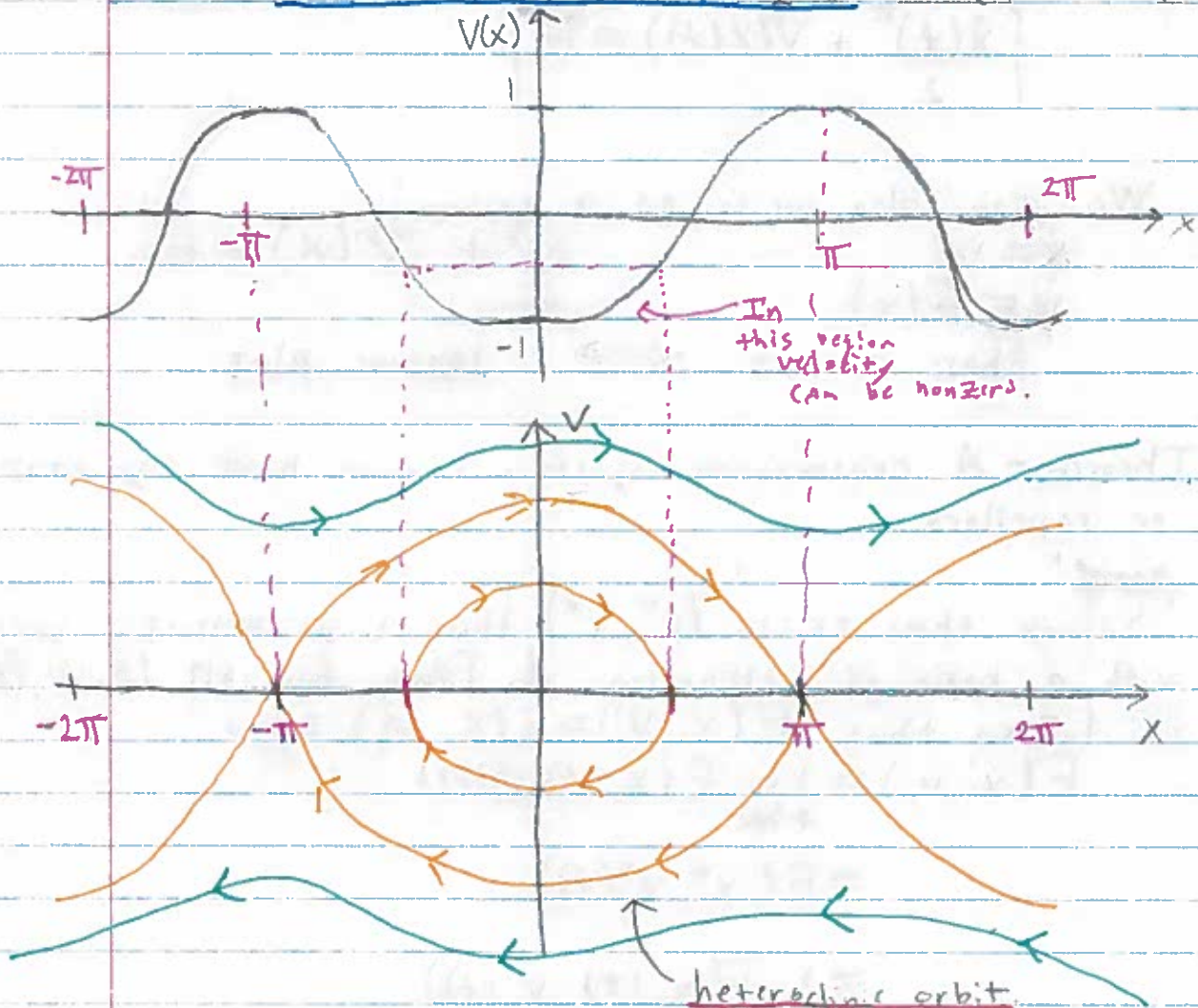
$$V(x) = -\cos(x).$$

$$E = \frac{1}{2} v^2 - \cos(x)$$

$$\Rightarrow v = \pm \sqrt{2E + 2\cos(x)}$$

$$v = \pm \sqrt{2(\cos(x) - \cos(x_0))} \Leftrightarrow (\text{If } v(0) = 0)$$

$$v = \pm \sqrt{2(\cos(x_0) - \cos(x)) + \frac{1}{2} v_0^2} \Leftrightarrow (\text{If } v(0) \neq 0)$$





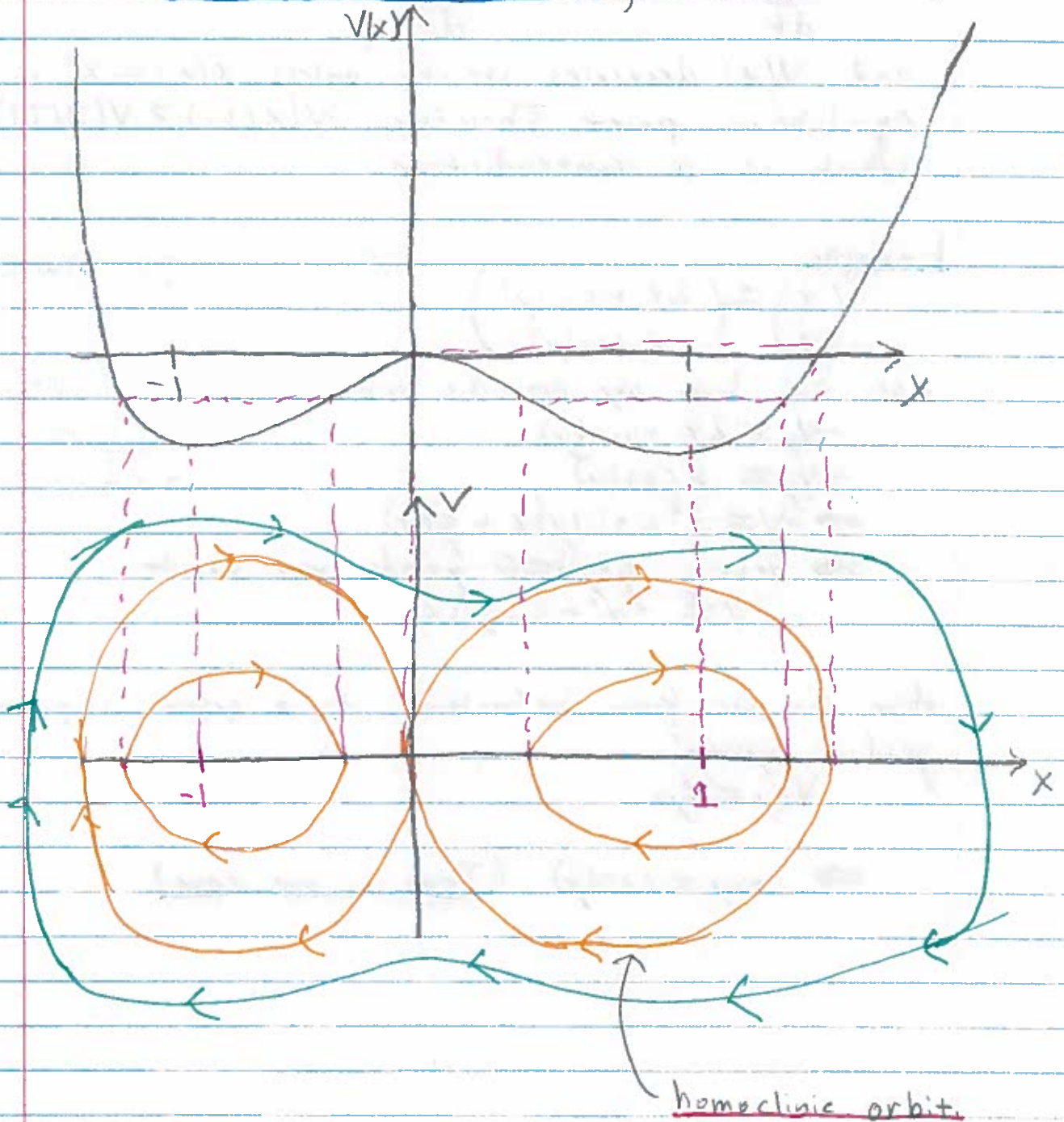
Example

$$\ddot{x} = x - x^3$$

$$V(x) = -\frac{x^2}{2} + \frac{x^4}{4}$$

$$E = \frac{1}{2}v^2 - \frac{x^2}{2} + \frac{x^4}{4}$$

$$\Rightarrow v = \pm \sqrt{x^2 - \frac{x^4}{2} - x_0^2 + \frac{x_0^4}{2}}, \quad (\text{If } v(0) = 0)$$



## G Gradient Systems.

$$\dot{\vec{x}} = -\nabla V, \text{ where } V: \mathbb{R}^2 \rightarrow \mathbb{R}$$

Lemma - Gradient systems cannot have closed orbits

proof:

Let  $x(t)$  be a closed orbit with period  $T$ . Then,

$$\frac{d}{dt} V(x(t)) = \nabla V \cdot \frac{dx}{dt} = -|\nabla V|^2 \leq 0.$$

and  $V(x)$  decreases strictly unless  $x(t) = x^*$  is an equilibrium point. Therefore,  $V(x(0)) > V(x(T))$  which is a contradiction.

Example

$$\begin{pmatrix} \dot{x} \\ \dot{y} \end{pmatrix} = \begin{pmatrix} 2x + \sin(y) \\ x \cos(y) \end{pmatrix}$$

does not have any periodic orbits.

$$-V_x = 2x + \sin(y)$$

$$-V_y = x \cos(y)$$

$$\Rightarrow -V = x^2 + \sin(y)x + g(y)$$

$\Rightarrow$  Setting  $g(y) = 0$  yields the result

$$V = x^2 - \sin(y)x$$

How can we know beforehand if a system is potentially a gradient system?

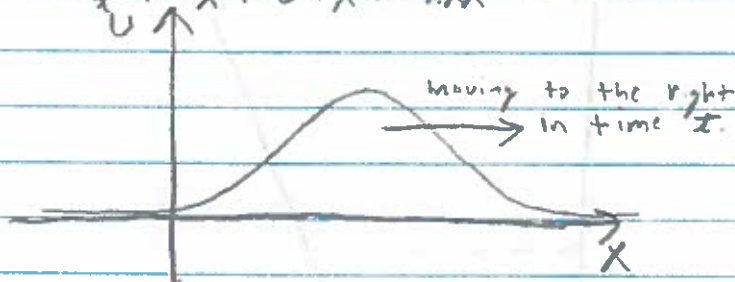
$$V_{xy} = V_{yx}$$

$$\Rightarrow \cos(y) = \cos(y) \quad (\text{True in our case})$$

## Solitons

Shallow water waves in a narrow canal

$$u_t + uv_x + u_x v + u_{xxx} = 0$$



$$\text{Let } z = x - ct$$

$$\frac{\partial}{\partial x} = \frac{\partial z}{\partial x} \frac{d}{dz} = \frac{d}{dz}$$

$$\frac{\partial}{\partial t} = \frac{\partial z}{\partial t} \frac{d}{dz} = -c \frac{d}{dz}$$

$$\Rightarrow (1-c) \frac{dv}{dz} + \frac{1}{2} \frac{d}{dz} (v^2) + u_{zzz} = 0$$

$$\Rightarrow (1-c) \dot{v} + \frac{1}{2} v^2 + u_{zz} = \mu$$

$$\Rightarrow v_{zz} = -(1-c)v - \frac{1}{2} v^2 + \mu$$

This is a conservative system with potential:

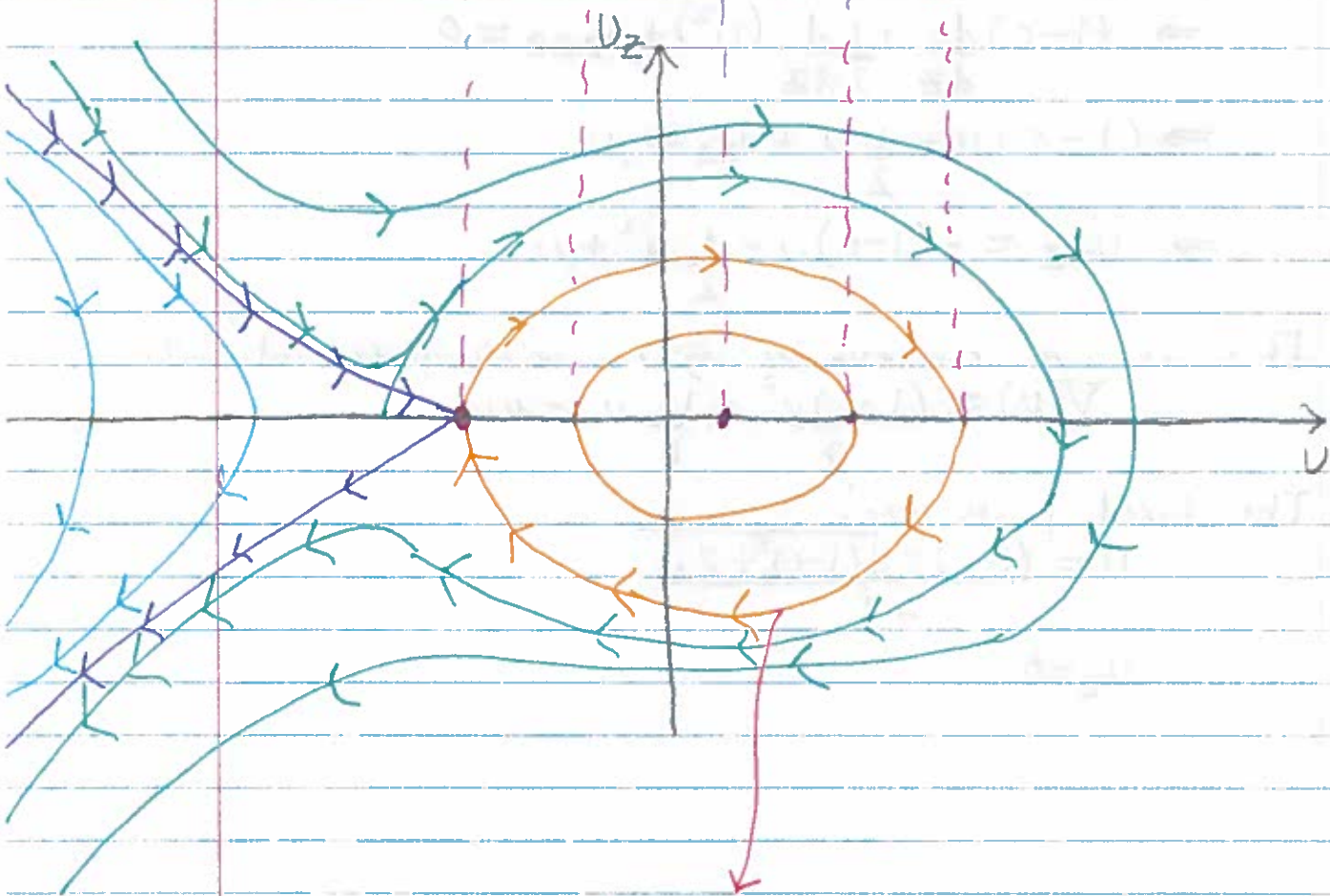
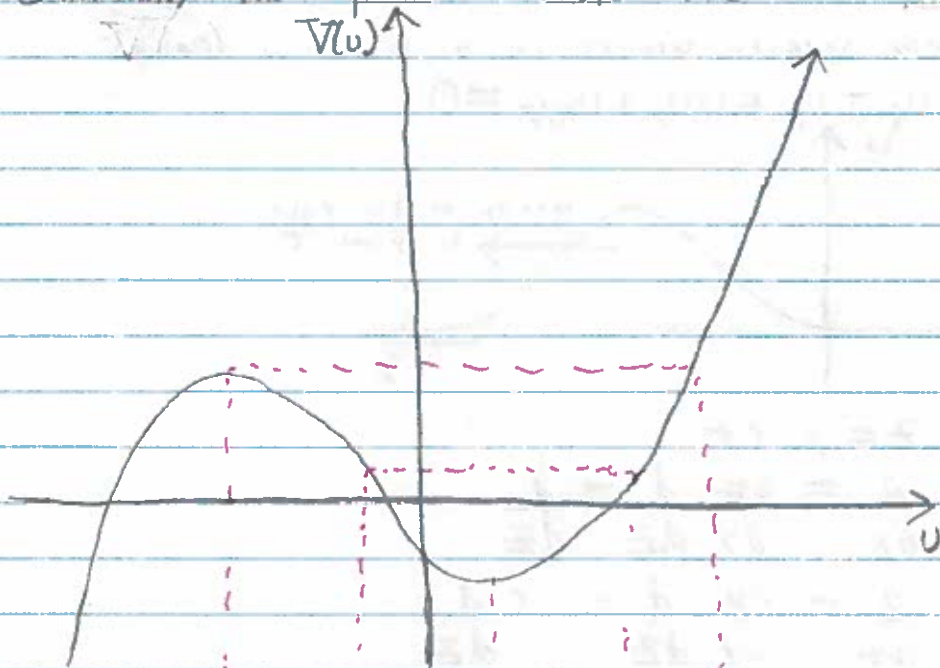
$$V(v) = \frac{(1-c)v^2}{2} + \frac{1}{6} v^3 - \mu v$$

The fixed points are:

$$v = \frac{(1-c) \pm \sqrt{(1-c)^2 + 2\mu}}{-1}$$

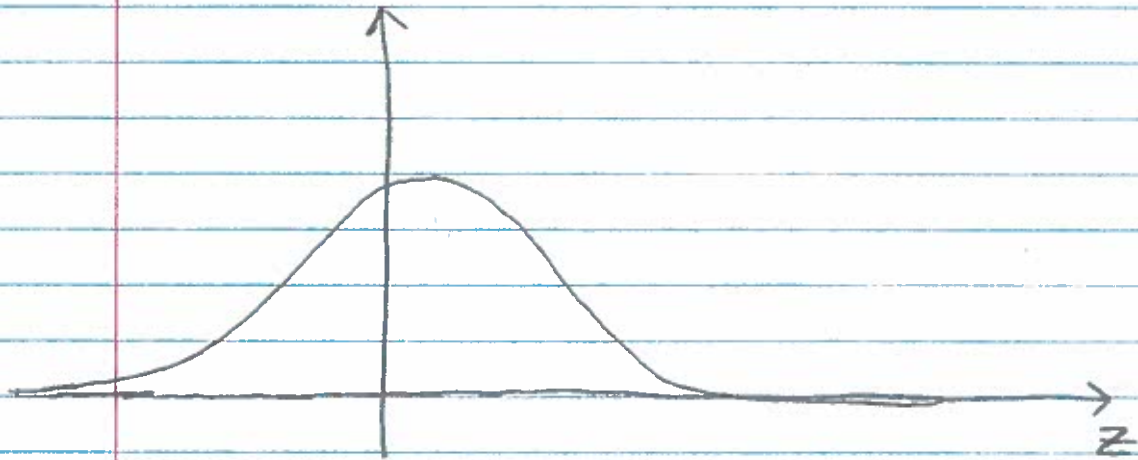
$$\bar{v}_z = 0$$

Generically the potential looks like



Separatrix/homoclinic orbit

The travelling wave corresponds to the separatrix / homoclinic



Sketch of the solution.

## Lyapunov Functions

$$\dot{\vec{x}} = F(\vec{x})$$

A continuously differentiable function  $L: \mathbb{R}^2 \rightarrow \mathbb{R}$  is called a Lyapunov function if  $L(x(t))$  strictly decreases along each solution of  $\dot{\vec{x}} = F(\vec{x})$  that is not an equilibrium.

Lemma - If  $\dot{\vec{x}} = F(\vec{x})$  admits a Lyapunov function, then it cannot have any periodic orbits.

### Example

$$\ddot{x} + \alpha \dot{x} = g(x), \quad \alpha > 0,$$

Let  $V(x) = -\int_{x_0}^x g(x) dx$ . Then,

$$\frac{1}{2} \frac{d(\dot{x}^2)}{dt} + \alpha \dot{x}^2 = -\frac{d}{dt}(V(x(t)))$$

$$\Rightarrow \frac{d}{dt} \left( \frac{1}{2} \dot{x}^2 + V(x(t)) \right) = -\alpha \dot{x}^2 < 0.$$

The function  $L(v, x) = \frac{1}{2}v^2 + V$  is a Lyapunov function.

### Summary:

1.  $\ddot{x} = -\frac{dV}{dx} \rightarrow$  conservative  $E(x, \dot{x})$  is conserved } Many periodic orbits.
2.  $\begin{pmatrix} \dot{x} \\ \dot{y} \end{pmatrix} = -\nabla V \rightarrow$  gradient system,  $V$  decreases along solutions } No periodic solutions.
3.  $\ddot{x} + \alpha \dot{x} = -\frac{dV}{dx} \rightarrow E(x, \dot{x})$  decreases