

Physics 712 Solutions to Chapter 11 Problems

5. Consider a line of charge with linear charge density λ arranged, in a primed frame, along the y' -axis at rest. Write the electric field at all points in Cartesian coordinates in the primed frame. Now, consider a line of charge with the same linear charge density, parallel to the y -axis, but this time moving in the $+x$ direction at velocity v . Find the electric and magnetic fields everywhere in the unprimed frame.

For a line of charge along the y' -axis, we can draw a cylinder of radius r' and length L around the linear charge density. The charge enclosed will be λL . Symmetry argues that the electric field will point directly out of the cylinder on the lateral surface, and will depend only on the distance away, so that $\mathbf{E}' = E'(r)\hat{\mathbf{r}}$, where $\hat{\mathbf{r}}$ is a unit vector pointing away from the y' -axis.

We then use Gauss's Law to conclude that the electric field everywhere is

$$\frac{\lambda L}{\epsilon_0} = \int_s \mathbf{E}' \cdot \hat{\mathbf{n}} da = 2\pi r' L E'(r'), \quad \text{so} \quad E'(r') = \frac{\lambda}{2\pi\epsilon_0 r'}.$$

We therefore have

$$\mathbf{E} = E'(r)\hat{\mathbf{r}} = \frac{\lambda\hat{\mathbf{r}}}{2\pi\epsilon_0 r'} = \frac{\lambda\mathbf{r}'}{2\pi\epsilon_0 r'^2} = \frac{\lambda(x'\hat{\mathbf{x}} + z'\hat{\mathbf{z}})}{2\pi\epsilon_0 (x'^2 + z'^2)},$$

where we recall that in this context r' is the distance of the point from the y' -axis. Of course, in the primed frame, there is no magnetic field at all.

To solve the "harder" problem, we now simply perform a Lorentz boost by speed $-v$ in the x -direction. There is one (apparent) subtlety here – are we sure the linear charge density λ is the same in both frames? We know that charge is Lorentz-invariant, and a boost in the x -direction does not affect distances in the y -direction, and since linear charge density is the charge per unit length (in the y -direction), the linear charge density should be unchanged.

The Lorentz transformations for the fields for this Lorentz boost will be

$$\begin{aligned} \mathbf{E}_{\parallel} = \mathbf{E}'_{\parallel} &= \frac{\lambda x'\hat{\mathbf{x}}}{2\pi\epsilon_0 (x'^2 + z'^2)}, & \mathbf{E}_{\perp} &= \gamma(\mathbf{E}'_{\perp} - \mathbf{v} \times \mathbf{B}') = \frac{\gamma\lambda z'\hat{\mathbf{z}}}{2\pi\epsilon_0 (x'^2 + z'^2)}, \\ \mathbf{B}_{\parallel} = \mathbf{B}'_{\parallel} &= 0, & \mathbf{B}_{\perp} &= \gamma(\mathbf{B}'_{\perp} + \mathbf{v} \times \mathbf{E}'/c^2) = \frac{\gamma v \hat{\mathbf{x}} \times \lambda z'\hat{\mathbf{z}}}{2\pi\epsilon_0 c^2 (x'^2 + z'^2)} = \frac{-\gamma\mu_0 \lambda z'\hat{\mathbf{y}}}{2\pi (x'^2 + z'^2)}. \end{aligned}$$

The coordinates are related by

$$t' = \gamma(t - vx/c^2), \quad x' = \gamma(x - vt), \quad y' = y, \quad z' = z.$$

Substituting this into the previous expressions, we have

$$\mathbf{E} = \frac{\lambda\gamma[(x-vt)\hat{\mathbf{x}} + z\hat{\mathbf{z}}]}{2\pi\epsilon_0 [\gamma^2(x-vt)^2 + z^2]}, \quad \mathbf{B} = \frac{-\gamma\mu_0 \lambda z\hat{\mathbf{y}}}{2\pi [\gamma^2(x-vt)^2 + z^2]}, \quad \text{where} \quad \gamma = \frac{1}{\sqrt{1-v^2/c^2}}.$$