

**Physics 712**  
**Solution to Problems 4.5 and 5.1**

5. [5] For problem 4.1, find the total energy if the cylinder has length  $L$ . For problem 4.2, find the total energy. In each case, show that the answer is equivalent to  $W = \frac{1}{2}Q\Delta\Phi$ .

The energy is given for problem 4.1 by

$$\begin{aligned} W &= \frac{1}{2} \int \mathbf{E} \cdot \mathbf{D} d^3\mathbf{x} = \frac{1}{2} \int_0^{2\pi} d\phi \int_0^L dz \left[ \int_a^b \varepsilon \left( \frac{\lambda}{2\pi\varepsilon\rho} \right)^2 \rho d\rho + \int_b^c \varepsilon_0 \left( \frac{\lambda}{2\pi\varepsilon_0\rho} \right)^2 \rho d\rho \right] \\ &= \frac{2\pi\lambda^2 L}{8\pi^2} \left( \frac{1}{\varepsilon} \ln \rho \Big|_a^b + \frac{1}{\varepsilon_0} \ln \rho \Big|_b^c \right) = \frac{\lambda^2 L}{4\pi} \left[ \frac{1}{\varepsilon} \ln \left( \frac{b}{a} \right) + \frac{1}{\varepsilon_0} \ln \left( \frac{c}{b} \right) \right] = \frac{1}{2} \lambda L \Delta\Phi = \frac{1}{2} Q \Delta\Phi, \end{aligned}$$

where at the last step, we interpreted  $\lambda L = Q$  as the total charge. For problem 4.2, we have

$$\begin{aligned} W &= \frac{1}{2} \int \mathbf{E} \cdot \mathbf{D} d^3\mathbf{x} = \frac{1}{2} \int_0^{2\pi} d\phi \int_a^b \left[ \frac{Q}{\pi(3\varepsilon_0 + \varepsilon)r^2} \right]^2 r^2 dr \left( \int_0^{\frac{1}{3}\pi} \varepsilon \sin \theta d\theta + \int_{\frac{1}{3}\pi}^{\pi} \varepsilon_0 \sin \theta d\theta \right) \\ &= \frac{2\pi Q^2}{2\pi^2 (3\varepsilon_0 + \varepsilon)^2} \frac{-1}{r} \Big|_a^b \left( -\varepsilon \cos \theta \Big|_0^{\frac{1}{3}\pi} - \varepsilon_0 \cos \theta \Big|_{\frac{1}{3}\pi}^{\pi} \right) = \frac{Q^2}{\pi (3\varepsilon_0 + \varepsilon)^2} \left( \frac{1}{a} - \frac{1}{b} \right) \left( \frac{1}{2} \varepsilon + \frac{3}{2} \varepsilon_0 \right) \\ &= \frac{Q^2 (b-a)}{2\pi ab (3\varepsilon_0 + \varepsilon)} = \frac{1}{2} Q \Delta\Phi. \end{aligned}$$

1. [10] We are trying to trap a charged particle of mass  $q > 0$  and mass  $m$  by using a combination of magnetic and electric fields given by  $\mathbf{B} = B\hat{\mathbf{z}}$  and  $\mathbf{E} = A(x\hat{\mathbf{x}} + y\hat{\mathbf{y}} - 2z\hat{\mathbf{z}})$ .  
 (a) [1] Obviously,  $\nabla \cdot \mathbf{B} = \nabla \times \mathbf{B} = 0$ . Check that it also satisfies  $\nabla \cdot \mathbf{E} = \nabla \times \mathbf{E} = 0$ .

We simply see that  $\nabla \cdot \mathbf{E} = A + A - 2A = 0$ , and all the terms in  $\nabla \times \mathbf{E}$  vanish.

- (b) [6] Assume the particle has motion given by  $x = R \cos(\omega t)$ ,  $y = R \sin(\omega t)$ . Find an equation for  $\omega$  in terms of  $A$  and  $B$ .

The velocity and acceleration can be found by simply taking derivatives:

$$\mathbf{v} = \dot{\mathbf{x}} = R\omega[-\sin(\omega t)\hat{\mathbf{x}} + \cos(\omega t)\hat{\mathbf{y}}], \quad \mathbf{a} = \dot{\mathbf{v}} = R\omega^2[-\cos(\omega t)\hat{\mathbf{x}} - \sin(\omega t)\hat{\mathbf{y}}]$$

We therefore have

$$\begin{aligned} m\mathbf{a} = \mathbf{F} = (\mathbf{E} + \mathbf{v} \times \mathbf{B}) &= qRA[\cos(\omega t)\hat{\mathbf{x}} + \sin(\omega t)\hat{\mathbf{y}}] + qBR\omega[-\cos(\omega t)\hat{\mathbf{x}} + \sin(\omega t)\hat{\mathbf{y}}] \times \hat{\mathbf{z}}, \\ -mR\omega^2[\cos(\omega t)\hat{\mathbf{x}} + \sin(\omega t)\hat{\mathbf{y}}] &= qRA[\cos(\omega t)\hat{\mathbf{x}} + \sin(\omega t)\hat{\mathbf{y}}] + qBR\omega[\cos(\omega t)\hat{\mathbf{x}} + \sin(\omega t)\hat{\mathbf{y}}], \\ -mR\omega^2 &= qRA + qBR\omega, \\ m\omega^2 + qB\omega + qA &= 0. \end{aligned}$$

We then solve this using the quadratic equation, so

$$\omega = \frac{-qB \pm \sqrt{q^2 B^2 - 4mqA}}{2m}.$$

Until I solved this problem myself, I didn't even realize there were two solutions to this equation.

- (c) [3] Argue that there is a maximum value of  $A$  for which circular motion is possible. Also argue that for  $A > 0$ , the particle will not "wander off" in the  $z$ -direction.

The solution only makes sense if the discriminant is positive, so we must have  $q^2 B^2 \geq 4mqA$ , or  $A \leq qB^2/(4m)$ . Although we have not discussed motion in the  $z$ -direction, it is pretty easy to see that the magnetic field has no influence on it, so the only vertical force is  $F_z = E_z q = -2Azq$ . Such a linear restoring force will result in simple harmonic motion in the  $z$ -direction, so it is stable against motion in the  $z$ -direction.