

Physics 780 – General Relativity
Solution Set T

47. The universe will be finite in size if $\Omega > 1$.

(a) The value given in class for the density parameter is $\Omega_0 = 0.9993 \pm 0.0037$. Taking this literally, for a closed universe, what is the *smallest* possible value for $H_0 a_0$ assuming $1 < \Omega \leq 0.9993 + 0.0037$?

The first Friedmann equation says that

$$\frac{\dot{a}^2}{a^2} = \frac{8\pi G}{3} \rho - \frac{k}{a^2}.$$

The left side evaluated today is Hubble's constant, and the ratio of the first term on the right to the expression on the left is Ω , so evaluating everything today, we have

$$H_0^2 = H_0^2 \Omega_0 - \frac{k}{a_0^2},$$
$$\frac{k}{a_0^2} = H_0^2 (\Omega_0 - 1).$$

We are interested in the case where the universe is closed, $\Omega_0 > 1$, which corresponds to $k = +1$. We see that $\Omega_0 \leq 1.0030$ therefore leads to

$$\frac{1}{a_0^2} \leq H_0^2 (1.0030 - 1),$$
$$a_0^2 H_0^2 \geq \frac{1}{0.0030} = 333,$$
$$a_0 H_0 \geq \sqrt{333} = 18.2.$$

Keeping in mind that we know H_0 actually has units of s^{-1} , and a_0 has units of m, the left hand side has units of velocity. Since we are allowed to multiply or divide by c as necessary, the correct inequality must be $a_0 H_0 \geq 18.2c$.

(b) What is the spatial volume for a closed universe with scale factor a ? You will probably have to use the version of the metric in terms of ψ to get the full range $\psi \in [0, \pi]$ of the whole universe.

The metric when written in terms of ψ , where $r = \sin \psi$, is

$$ds^2 = -dt^2 + a^2 \left[d\psi^2 + \sin^2 \psi (d\theta^2 + \sin^2 \theta d\phi^2) \right].$$

The metric is diagonal, and the space part of the metric is $g_{\psi\psi} = a^2$, $g_{\theta\theta} = a^2 \sin^2 \psi$ and $g_{\phi\phi} = a^2 \sin^2 \psi \sin^2 \theta$. The volume of the universe is just the integral of $\sqrt{g^{(3)}}$, where $g^{(3)}$ is the determinant of the space part of the metric, so we have

$$\begin{aligned} V &= \int \sqrt{a^2 a^2 \sin^2 \psi a^2 \sin^2 \psi \sin^2 \theta} d\psi d\theta d\phi = a^3 \int_0^\pi \sin^2 \psi d\psi \int_0^\pi \sin \theta d\theta \int_0^{2\pi} d\phi \\ &= a^3 \left(\frac{1}{2}\pi\right)(2)(2\pi) = 2\pi^2 a^3. \end{aligned}$$

(c) Write the scale factor a from part (a) in Gpc if $H_0 = 67.7$ km/s/Mpc. Don't forget to add factors of c to get the units right! Then find the minimum volume of the visible universe in Gpc^3 using the result of part (b).

We already put c in as appropriate, so we simply start computing:

$$a_0 \geq \frac{18.2c}{H_0} = \frac{18.2(2.998 \times 10^5 \text{ km/s})}{67.7 \text{ km/s/Mpc}} = 80,810 \text{ Mpc} = 80.81 \text{ Gpc}.$$

We now just put it in the formula from part (b) to find

$$V_0 = 2\pi^2 a_0^3 \geq 2\pi^2 (80.81 \text{ Gpc})^3 = 1.042 \times 10^7 \text{ Gpc}^3.$$

48. We found an integral formula for the current age of the universe times the current Hubble constant $t_0 H_0$ and the current density parameter if there is *only* matter with density Ω_m .

(a) Repeat this exercise and find $t_0 H_0$ if there is *only* radiation with density Ω_r . Perform the integral.

The current density of radiation is given by $\frac{8}{3}\pi G\rho_{r0} = H_0^2 \Omega_r$. Because the density of radiation falls as a^{-4} , the density at other times must be

$$\frac{8\pi}{3} G\rho_r = \frac{8\pi}{3} G\rho_{r0} \left(\frac{a_0}{a}\right)^4 = H_0^2 \Omega_r \left(\frac{a_0}{a}\right)^4$$

We will use the first Friedmann equation, which tells us that

$$\frac{\dot{a}^2}{a^2} = \frac{8\pi}{3} G\rho_r - \frac{k}{a^2}.$$

Evaluating this equation today, we would have

$$H_0^2 = H_0^2 \Omega - \frac{k}{a_0^2},$$

$$\frac{k}{a_0^2} = H_0^2 (\Omega_r - 1)$$

Substituting this back into the first Friedmann equation, we have

$$\frac{\dot{a}^2}{a^2} = \frac{8\pi}{3}G\rho_r - \frac{k}{a^2} = H_0^2\Omega_r \frac{a_0^4}{a^4} - H_0^2(\Omega_r - 1)\frac{a_0^2}{a^2}.$$

Defining $x = a/a_0$, this equation becomes

$$\frac{\dot{x}^2}{x^2} = H_0^2 \left(\frac{\Omega_r}{x^4} + \frac{1-\Omega_r}{x^2} \right).$$

We solving for dx/dt , we find

$$\frac{dx}{dt} = H_0 \sqrt{\Omega_r/x^2 + 1 - \Omega_r}.$$

We want the time, which we will get by rearranging and then integrating. Note that since $x = a/a_0$, we want to integrate a from when the universe was size zero, $x = 0$, to now $x = 1$.

$$H_0 dt = \frac{dx}{\sqrt{\Omega_r/x^2 + 1 - \Omega_r}},$$

$$H_0 t_0 = \int_0^1 \frac{dx}{\sqrt{\Omega_r/x^2 + 1 - \Omega_r}}.$$

We now have to integrate this, which isn't so bad:

$$H_0 t_0 = \int_0^1 \frac{x dx}{\sqrt{\Omega_r + (1-\Omega_r)x^2}} = \frac{1}{2(1-\Omega_r)} \int_0^1 \frac{d[(1-\Omega_r)x^2]}{\sqrt{\Omega_r + (1-\Omega_r)x^2}} = \frac{\sqrt{\Omega_r + (1-\Omega_r)x^2}}{1-\Omega_r} \Big|_0^1$$

$$= \frac{1 - \sqrt{\Omega_r}}{1-\Omega_r} = \frac{1 - \sqrt{\Omega_r}}{(1-\sqrt{\Omega_r})(1+\sqrt{\Omega_r})} = \frac{1}{1+\sqrt{\Omega_r}}.$$

Technically, this integration was done under the assumption that $\Omega_r \neq 1$, but it is trivial to redo the integral in this case and verify that the final result is still valid.

- (b) Repeat this exercise and find $t_0 H_0$ if there is radiation with density Ω_r and matter with density Ω_m , but the universe is flat, so $\Omega_r + \Omega_m = 1$. Perform the integral. If needed, you can use the fact that $\Omega_r + \Omega_m = 1$ and the fact that they are both positive to conclude they are both smaller than 1.

Since the universe is flat, $k = 0$, so the first Friedmann equation gives

$$\frac{\dot{a}^2}{a^2} = \frac{8\pi}{3}G\rho_r + \frac{8\pi}{3}G\rho_m.$$

The current value of each of the terms on the right are $\frac{8}{3}\pi G\rho_r = H_0^2\Omega_r$ and $\frac{8}{3}\pi G\rho_m = H_0^2\Omega_m$.

Using the fact that $\rho_r \propto a^{-4}$ and $\rho_m \propto a^{-3}$, the Friedmann equation yields

$$\frac{\dot{a}^2}{a^2} = H_0^2 \Omega_r \frac{a_0^4}{a^4} + H_0^2 \Omega_m \frac{a_0^3}{a^3}.$$

As usual, let $x = a/a_0$, then this equation becomes

$$\frac{\dot{x}^2}{x^2} = \frac{H_0^2 \Omega_r}{x^4} + \frac{H_0^2 \Omega_m}{x^3},$$

$$\frac{dx}{dt} = H_0 \sqrt{\frac{\Omega_r}{x^2} + \frac{\Omega_m}{x}}.$$

We rearrange this and integrate it over x to find t_0 :

$$H_0 dt = \frac{dx}{\sqrt{\Omega_r/x^2 + \Omega_m/x}},$$

$$H_0 t_0 = \int_0^1 \frac{dx}{\sqrt{\Omega_r/x^2 + \Omega_m/x}} = \int_0^1 \frac{x dx}{\sqrt{\Omega_r + \Omega_m x}} = \frac{1}{\Omega_m^2} \int_0^1 \frac{(\Omega_m x + \Omega_r - \Omega_r) d(\Omega_m x + \Omega_r)}{\sqrt{\Omega_r + \Omega_m x}}$$

$$= \frac{1}{\Omega_m^2} \int_0^1 \left(\sqrt{\Omega_m x + \Omega_r} - \frac{\Omega_r}{\sqrt{\Omega_m x + \Omega_r}} \right) d(\Omega_m x + \Omega_r)$$

$$= \frac{1}{\Omega_m^2} \left[\frac{2}{3} (\Omega_m x + \Omega_r)^{3/2} - 2\Omega_r \sqrt{\Omega_m x + \Omega_r} \right]_0^1$$

$$= \frac{1}{\Omega_m^2} \left[\frac{2}{3} - 2\Omega_r - \frac{2}{3} \Omega_r^{3/2} + 2\Omega_r^{3/2} \right] = \frac{2}{3\Omega_m^2} (1 - 3\Omega_r + 2\Omega_r^{3/2}),$$

we used $\Omega_r + \Omega_m = 1$ when substituting in the upper limits.

The final expression can be simplified, somewhat, by factoring the numerator as a cubic in $\sqrt{\Omega_r}$, and in the denominator rewriting $\Omega_m = 1 - \Omega_r = (1 - \sqrt{\Omega_r})(1 + \sqrt{\Omega_r})$, so we have

$$H_0 t_0 = \frac{2(1 + 2\sqrt{\Omega_r})(1 - \sqrt{\Omega_r})^2}{3(1 - \sqrt{\Omega_r})^2(1 + \sqrt{\Omega_r})^2} = \frac{2(1 + 2\sqrt{\Omega_r})}{3(1 + \sqrt{\Omega_r})^2}.$$

49. Way back in the previous millennium (i.e., pre-1995) we only knew about matter (and a tiny bit of radiation), and were none too confident about the value of Ω_m . For *this* problem, assume the universe contains matter only.

(a) Show that if $\Omega_m \leq 1$, the universe will never stop growing, i.e., there is no point in the future when $\dot{a} = 0$.

(b) Show that if $\Omega_m > 1$, it is inevitable that the universe will eventually stop growing. Find a formula for the size of the universe compared to now, a/a_0 , when the universe will stop growing as a function of Ω_m .