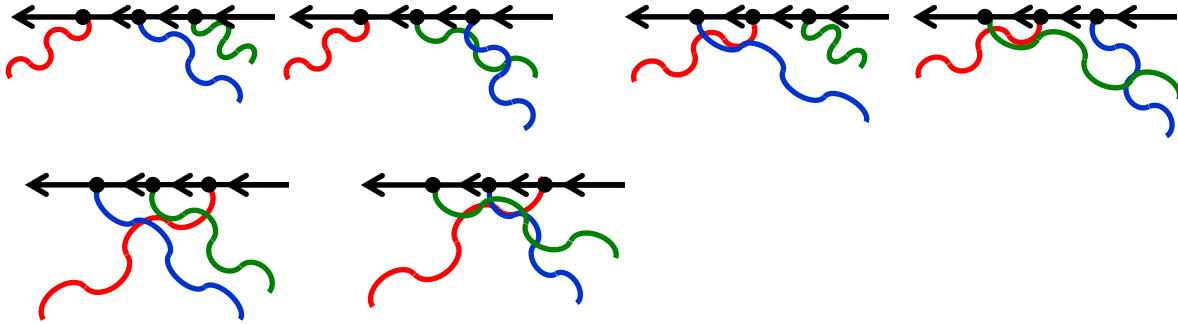


Solutions to Problems 7

3. Draw all six Feynman diagrams for the scattering $e^+\gamma \rightarrow e^+\gamma\gamma$.

This is straightforward, but boring. It is more challenging for me because I am not drawing them by hand. I have color coded the photon lines to try to help keep track of which interaction is which.



5. Calculate the total cross section ratio $\sigma(e^+e^- \rightarrow q\bar{q})/\sigma(e^+e^- \rightarrow \mu^+\mu^-)$, summed over all quarks (multiply each contribution by 3 for colors), treating the quark as massless if $\sqrt{s} > 2m$, and of course ignoring it if $\sqrt{s} < 2m$. Do so in the range $2 \text{ GeV} < \sqrt{s} < 3 \text{ GeV}$ (ignore the c quark), $5 \text{ GeV} < \sqrt{s} < 9 \text{ GeV}$ (include the c quark) and $11 \text{ GeV} < \sqrt{s} < 50 \text{ GeV}$ (include the b quark). Compare to the experimental values illustrated in Fig. 7-10 below. What do you think is going on at the other energies shown?

The requested ratio is just

$$\frac{\sigma(e^+e^- \rightarrow q\bar{q})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)} = 3 \cdot \frac{4\pi\alpha^2 Q^2}{3s} \cdot \frac{3s}{4\pi\alpha^2 (-1)^2} = 3Q^2.$$

The contribution from the up or charm quark would be $3\left(\frac{2}{3}\right)^2 = \frac{4}{3}$. The contribution from the down, strange, or bottom quark would be $3\left(-\frac{1}{3}\right)^2 = \frac{1}{3}$. Depending on the energy range, we then have

$$\frac{\sigma(e^+e^- \rightarrow q\bar{q})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)} = \begin{cases} \frac{4}{3} + 2 \cdot \frac{1}{3} & 2 \text{ GeV} < \sqrt{s} < 3 \text{ GeV} \\ 2 \cdot \frac{4}{3} + 2 \cdot \frac{1}{3} & 5 \text{ GeV} < \sqrt{s} < 9 \text{ GeV} \\ 2 \cdot \frac{4}{3} + 3 \cdot \frac{1}{3} & 11 \text{ GeV} < \sqrt{s} < 50 \text{ GeV} \end{cases} = \begin{cases} 2.00 & 2 \text{ GeV} < \sqrt{s} < 3 \text{ GeV} \\ 3.33 & 5 \text{ GeV} < \sqrt{s} < 9 \text{ GeV} \\ 3.67 & 11 \text{ GeV} < \sqrt{s} < 50 \text{ GeV} \end{cases}$$

I have sketched below dashed lines showing approximately these curve positions. You will note that in the regions listed, they actually work pretty well, though obviously around 40 GeV the Z-resonance (which causes weak interactions) is starting to have some effect. As for other energies, below 2 GeV there is clearly something complicated going on (involving the strong interactions of the light quarks), the spikes for the J/ψ and $\psi(2s)$ are caused by resonances

(bound states of charm/anti-charm quarks), and there are more resonances around 10 GeV (bound states of bottom-anti-bottom quarks).

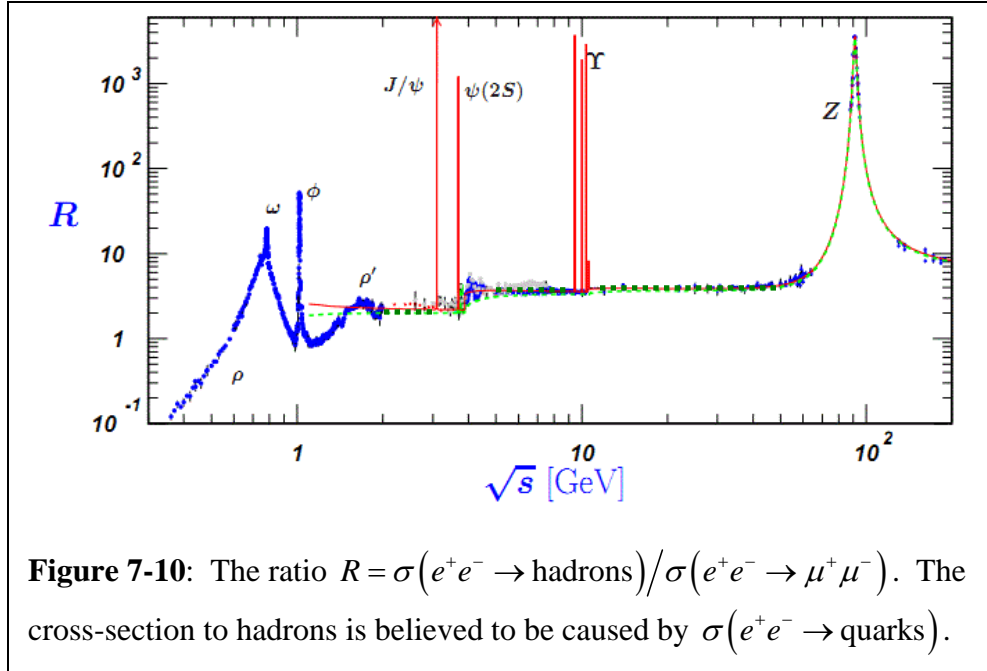


Figure 7-10: The ratio $R = \sigma(e^+e^- \rightarrow \text{hadrons}) / \sigma(e^+e^- \rightarrow \mu^+\mu^-)$. The cross-section to hadrons is believed to be caused by $\sigma(e^+e^- \rightarrow \text{quarks})$.

6. Calculate the differential and total cross-section for $e^+e^- \rightarrow f\bar{f}$, treating the electron mass as zero, but not ignoring the fermion mass m . Sketch the result for $3\sigma E^2 / \pi Q^2 \alpha^2$ as a function of m/E . It should be 1 if $m = 0$.

We start with eq. (7.19), but we keep the masses. We have

$$\begin{aligned}
 \frac{1}{4} \sum_{\text{spins}} |i\mathcal{M}|^2 &= \frac{e^4 Q^2}{4s^2} \text{Tr}(\not{p}_2 \gamma^\mu \not{p}_1 \gamma^\nu) \text{Tr}[(\not{p}_3 + m) \gamma_\mu (\not{p}_4 - m) \gamma_\nu] \\
 &= \frac{e^4 Q^2}{4s^2} \text{Tr}(\not{p}_2 \gamma^\mu \not{p}_1 \gamma^\nu) \text{Tr}(\not{p}_3 \gamma_\mu \not{p}_4 \gamma_\nu - m^2 \gamma_\mu \gamma_\nu) \\
 &= \frac{4e^4 Q^2}{s^2} (p_2^\mu p_1^\nu + p_1^\mu p_2^\nu - g^{\mu\nu} p_2 \cdot p_1) (p_{3\mu} p_{4\nu} + p_{3\nu} p_{4\mu} - g_{\mu\nu} p_3 \cdot p_4 - m^2 g_{\mu\nu}) \\
 &= \frac{4e^4 Q^2}{s^2} \left[2(p_2 \cdot p_3)(p_1 \cdot p_4) + 2(p_2 \cdot p_4)(p_1 \cdot p_3) \right. \\
 &\quad \left. + (p_2 \cdot p_1)(p_3 \cdot p_4)(-1-1-1-1+4) + m^2(p_1 \cdot p_2)(-1-1+4) \right] \\
 &= \frac{8e^4 Q^2}{s^2} [(p_2 \cdot p_3)(p_1 \cdot p_4) + (p_2 \cdot p_4)(p_1 \cdot p_3) + m^2(p_1 \cdot p_2)]
 \end{aligned}$$

In the center of mass frame, the electron's energy E is the same as their momenta, and we have total energy $2E$. The final state particles will split this energy between them, and will therefore each have energy E as well, but their momenta will be reduced to $p = \sqrt{E^2 - m^2}$. Hence the momenta will be

$$\begin{aligned}
 p_1^\mu &= (E, 0, 0, E), & p_3^\mu &= (E, p \sin \theta \cos \phi, p \sin \theta \sin \phi, p \cos \theta), \\
 p_2^\mu &= (E, 0, 0, -E), & p_4^\mu &= (E, -p \sin \theta \cos \phi, -p \sin \theta \sin \phi, -p \cos \theta).
 \end{aligned}$$

We then get all the relevant dot products:

$$p_1 \cdot p_4 = p_2 \cdot p_3 = E^2 + Ep \cos \theta, \quad p_1 \cdot p_3 = p_2 \cdot p_4 = E^2 - Ep \cos \theta, \quad p_1 \cdot p_2 = 2E^2.$$

Substituting this in, along with $s = 4E^2$, we have

$$\frac{1}{4} \sum_{\text{spins}} |i\mathcal{M}|^2 = \frac{8e^4 Q^2}{(4E^2)^2} \left[(E^2 + Ep \cos \theta)^2 + (E^2 - Ep \cos \theta)^2 + 2m^2 E^2 \right] = \frac{e^4 Q^2}{E^2} (E^2 + p^2 \cos^2 \theta + m^2).$$

We now proceed to the cross-section in the usual way.

$$\begin{aligned} \sigma &= \frac{D}{4|E_2 \mathbf{p}_1 - E_1 \mathbf{p}_2|} = \frac{1}{8E^2} \frac{p}{16\pi^2 E_{cm}} \int d\Omega \frac{1}{4} \sum_{\text{spins}} |i\mathcal{M}|^2 = \frac{pe^4 Q^2}{256\pi^2 E^3 E^2} \int d\Omega (E^2 + p^2 \cos^2 \theta + m^2) \\ &= \frac{\alpha^2 Q^2 p}{16E^5} 2\pi \int_{-1}^1 (E^2 + p^2 \cos^2 \theta + m^2) d \cos \theta = \frac{\pi \alpha^2 Q^2 p}{8E^5} (2E^2 + \frac{2}{3} p^2 + 2m^2) \\ &= \frac{\pi \alpha^2 Q^2}{4E^5} \sqrt{E^2 - m^2} (E^2 + \frac{1}{3} E^2 - \frac{1}{3} m^2 + m^2) = \frac{\pi Q^2 \alpha^2}{3E^2} \left\{ \left(1 + \frac{m^2}{2E^2} \right) \sqrt{1 - \frac{m^2}{E^2}} \right\}. \end{aligned}$$

The factor in $\{ \}$'s is what we were asked to graph. A sketch of the result is given at right. Note that for even modestly small values of m/E , it is very close to one; for example, at $m/E = 0.5$ it is already 97.4% of its maximum value.

